

QCD

Part 4

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HELMHOLTZ



And now for something completely different

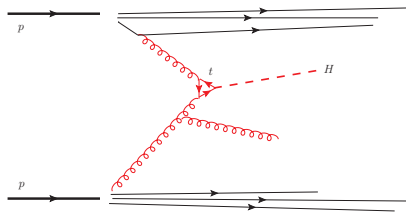
a few words about general-purpose event generators

e.g. Herwig, Pythia, Sherpa

note: Many other generators exist, often with a specialised scope and approach. Not all of them fit the description given in the following.

Monte Carlo generators e.g. Herwig, Pythia, Sherpa

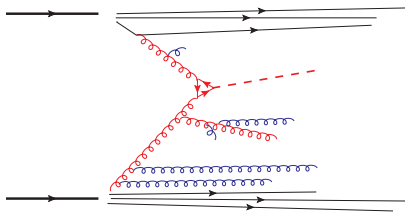
- ▶ build on structure of factorisation formulae e.g. for $pp \rightarrow H + g + X$
- ▶ but compute fully specified events, i.e. no “+X” schematically:



- ▶ ingredients:
 - parton densities and hard-scattering matrix elements

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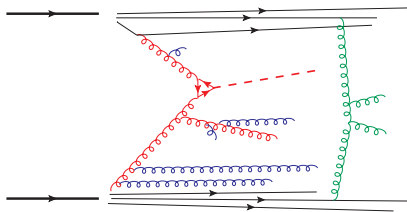
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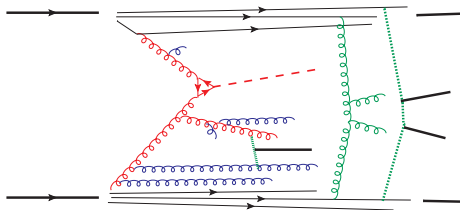
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 - models for multiparton interactions

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 - parton showers: collinear and soft radiation from partons in initial and final state (in perturbative region)
 - models for multiparton interactions and hadronisation

Parton density fits

Parton densities involve QCD at low momentum scales $\leftrightarrow$ large coupling

- ▶ can compute $\partial f(x, \mu)/\partial \mu$ in perturbation theory, but not $f(x, \mu)$
- ▶ ongoing effort to compute with non-perturbative methods
e.g. in lattice QCD
- ▶ in practice: determined from experimental data

Principle of PDF determinations:

- ▶ data for observables with factorisation formulae
most important: DIS ($ep \rightarrow e + X$), Drell-Yan ($pp \rightarrow \ell^+ \ell^- + X$,
 $pp \rightarrow \ell \nu + X$), jets in ep and pp , $t\bar{t}$ production in pp , ...
- ▶ parameterise PDFs at “starting” scale μ_0
use DGLAP eqs. to evolve to scales μ needed in fact. formulae
- ▶ determine PDF parameters by fit to data

Uncertainties on extracted PDFs

“PDF errors”

- ▶ errors (stat. and syst.) of fitted data propagated to PDF parameters

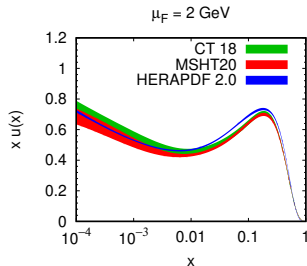
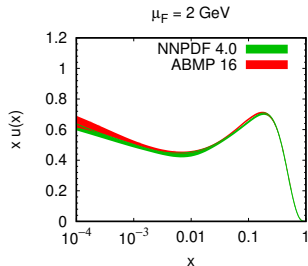
“systematic theory uncertainties”

- ▶ selection of data sets and kinematics
- ▶ perturbative order of evolution and hard-scattering cross sections
- ▶ values of α_s and m_c, m_b and possibly other constants if they are taken as external parameters rather than fitted
- ▶ fine details of perturbative calculations
e.g. treatment of heavy quarks, resummation of large logarithms
- ▶ power corrections (try to avoid by using data with $Q > Q_{\min}$)

ongoing effort: include uncertainties from higher orders in PDF errors

Harland-Lang, Thorne 2018; Khalek et al. 2019; McGowan et al. 2023

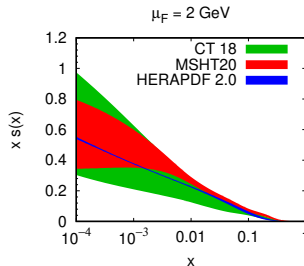
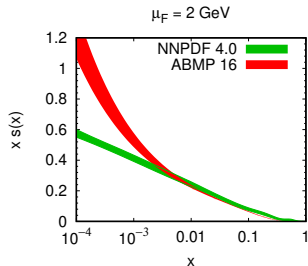
Illustration of PDF sets and their errors



all PDFs at NNLO, error bands for 68% CL

- ▶ spread between different parameterisation often larger than error bands of single parameterisation
- ▶ error bands propagate uncertainties of fitted data into PDFs but do **not** reflect “systematic theory uncertainties” of extraction

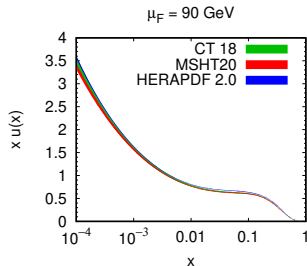
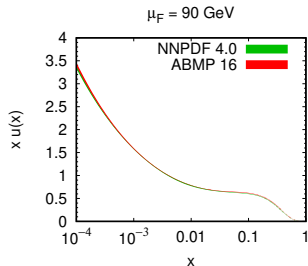
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- ▶ strangeness distribution remains poorly known
sometimes assume $s(x) \propto \bar{u}(x) + \bar{d}(x)$ or $s(x) \propto \bar{d}(x)$ at $\mu = \mu_0$
 $\rightsquigarrow$ small errors in fit

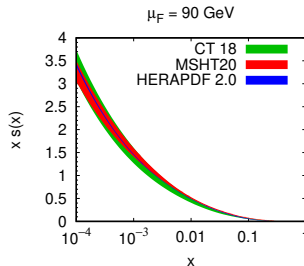
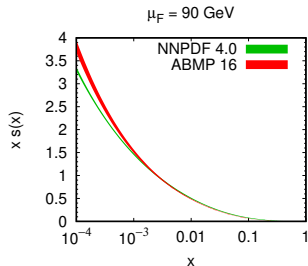
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- evolution to higher scales $\rightsquigarrow q\bar{q}$ pairs at low x

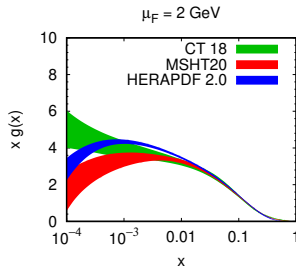
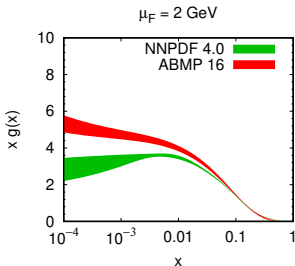
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- ▶ evolution to higher scales $\rightsquigarrow q\bar{q}$ pairs at low x
- ▶ all $q(x)$ and $\bar{q}(x)$ become similar at high scales and low x
- ▶ relative uncertainties shrink

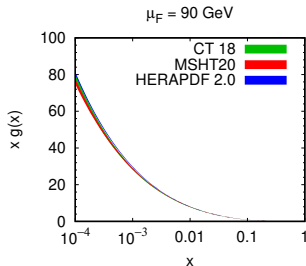
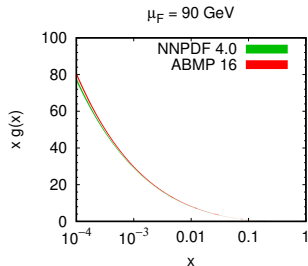
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all PDFs at NNLO, error bands for 68% CL

- ▶ $g(x) \gg q(x)$ for x below 0.1
- ▶ at low scale and low x gluon known very poorly

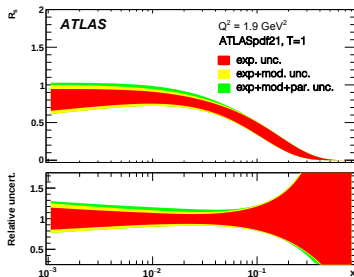
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all PDFs at NNLO, error bands for 68% CL

- ▶ $g(x) \gg q(x)$ for x below 0.1
- ▶ evolution for $g(x)$ even stronger than for $q(x)$

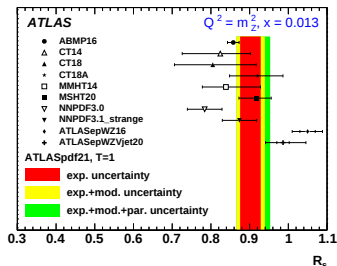
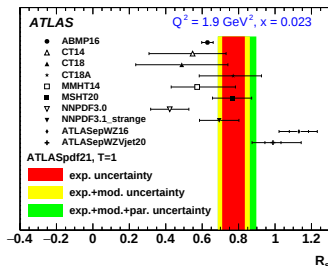
Strange quarks: recent results



ratio of strange to
non-strange sea quarks:

$$R_s(x) = \frac{s(x) + \bar{s}(x)}{\bar{u}(x) + \bar{d}(x)}$$

ATLAS, arXiv:2112.11266



Standard Model Production Cross Section Measurements

