

Fundamentals of X-ray interaction with matter

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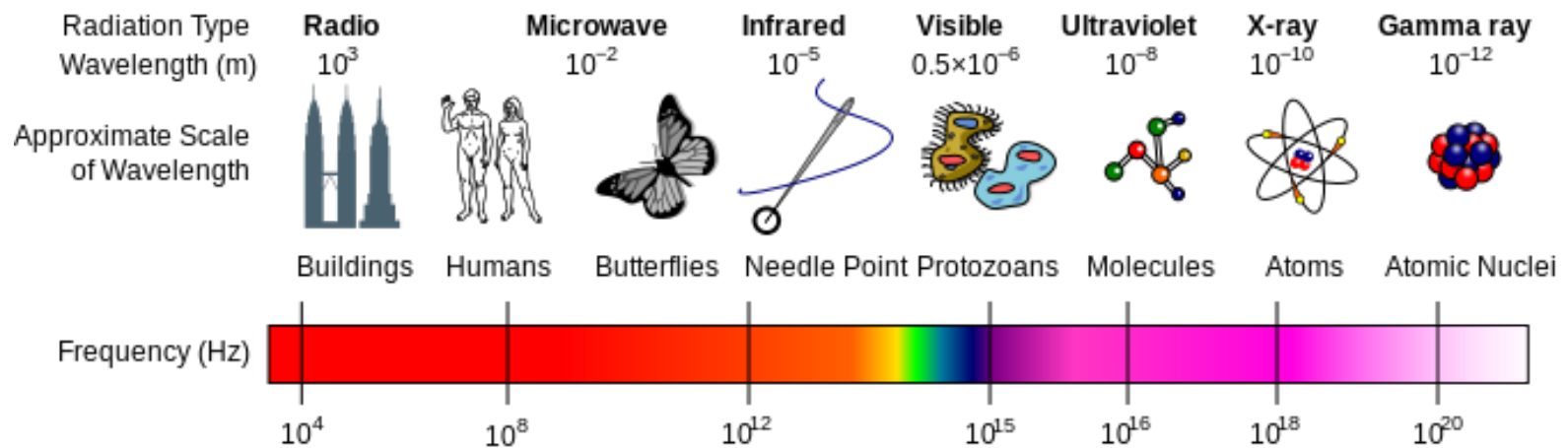
Outline

- > Introduction
- > Theory of X-ray–matter interaction
- > X-ray absorption
- > X-ray induced relaxation: X-ray fluorescence and Auger-Meitner decay
- > X-ray scattering
- > [Exercises: computer simulations with XATOM]

What is X-ray?

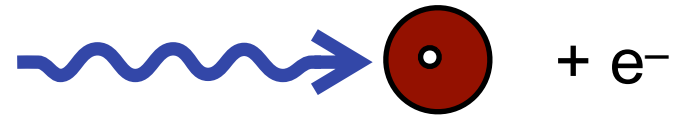
X-ray is...

- > discovered by Wilhelm Röntgen in 1895
- > photon energy: 100 eV ~ 100 keV
- > wavelength: 0.01 nm to 10 nm
- > penetrability
- > element selectivity



taken from https://en.wikipedia.org/wiki/Electromagnetic_spectrum

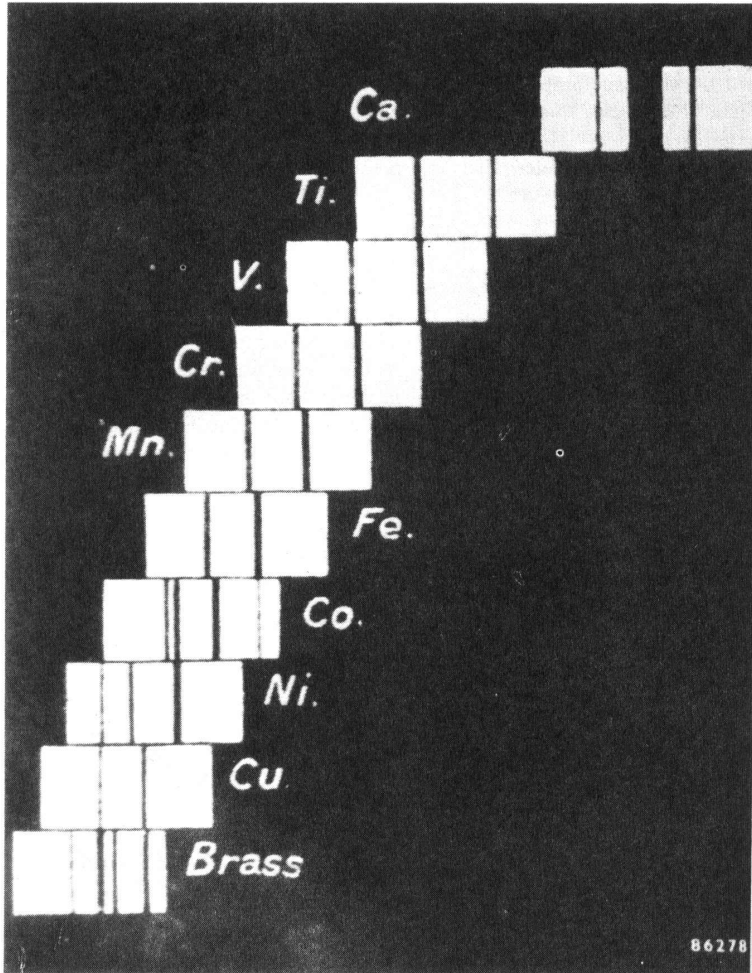
X-ray absorption



Radiography: see inside due to X-ray absorption contrast

Hand with rings: Wilhelm Röntgen's first medical X-ray in 1895 (Nobel prize in 1901)
taken from https://en.wikipedia.org/wiki/Wilhelm_Röntgen

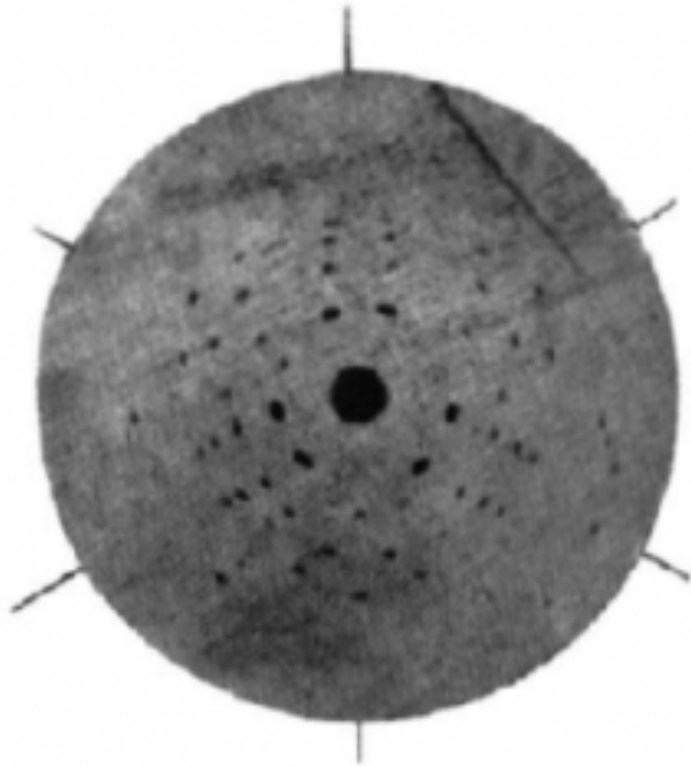
X-ray fluorescence



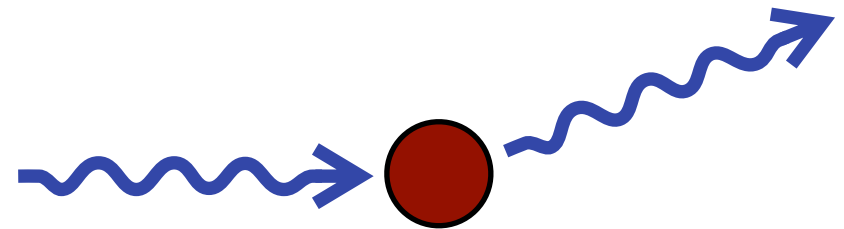
Different chemical compositions identified by X-ray fluorescence

Moseley's Staircase: First qualitative X-ray chemical analysis by Henry Moseley, *Phil. Mag.* **26**, 1024 (1913)

X-ray scattering / diffraction



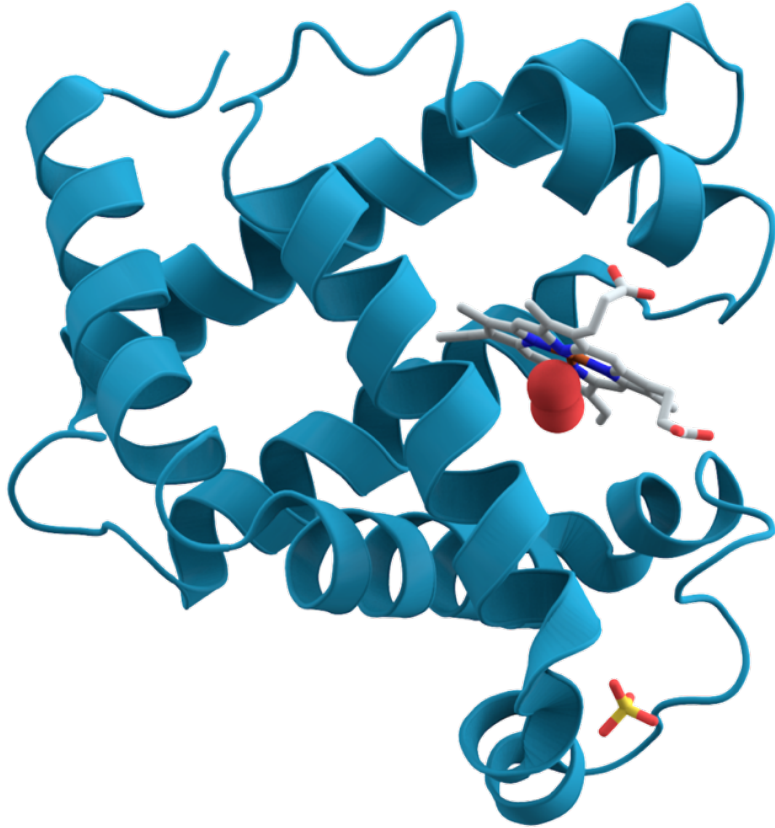
Zinc blend: First published diffraction pattern of a crystal by Max von Laue in 1912 (Nobel prize in 1914)



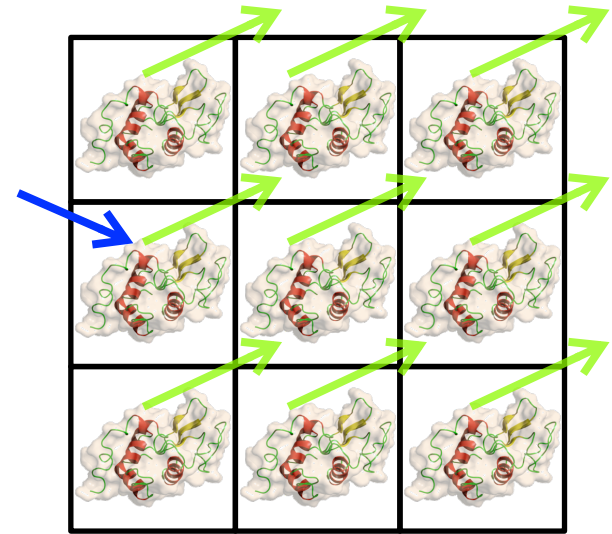
Elastic X-ray scatterings from a crystal make a diffraction pattern

Further reading: M. Eckert, *Ann. Phys.* **524**, A83 (2012)

X-ray crystallography

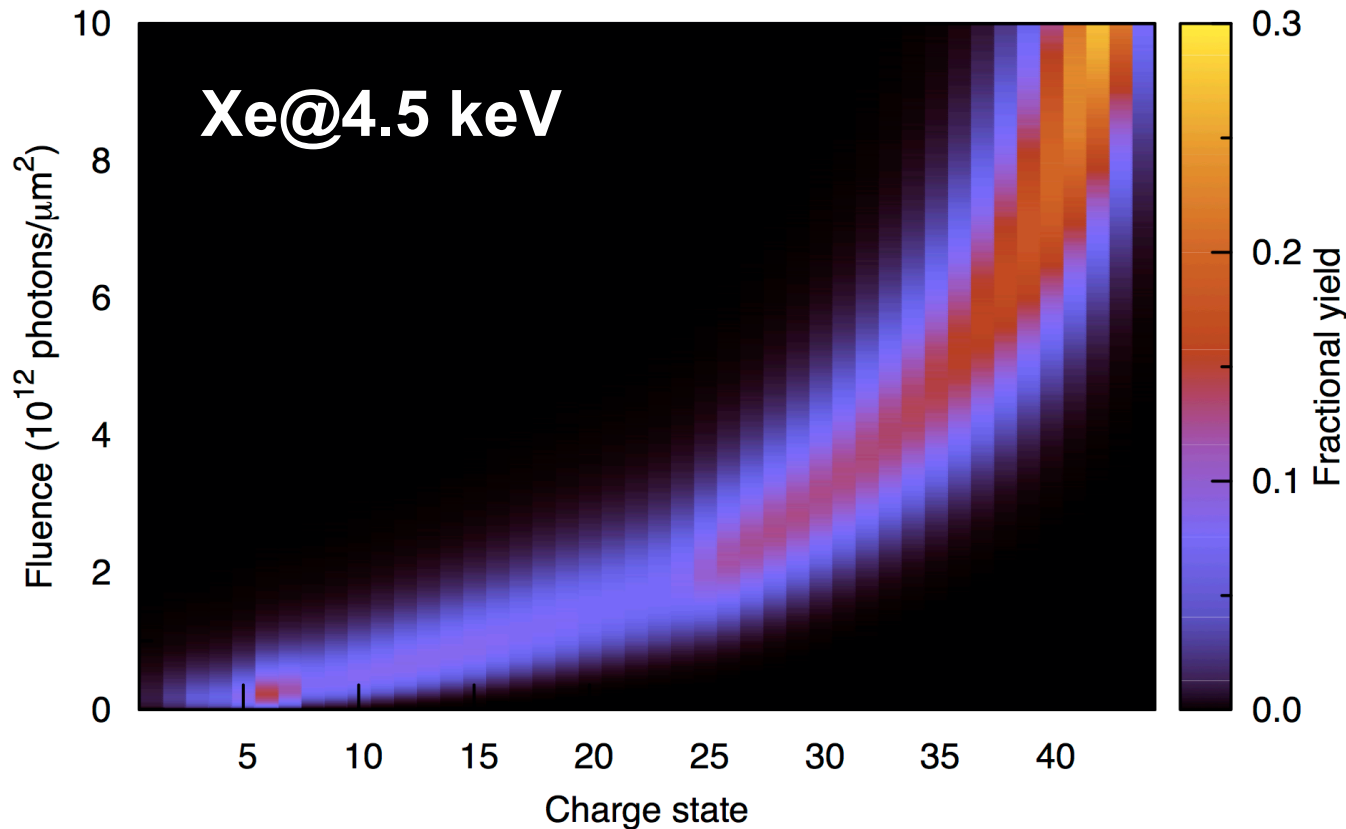


Myoglobin: first protein structure revealed by X-ray crystallography by John Kendrew in 1958 (Nobel prize in 1962 with Max Perutz)



X-ray crystallography can reveal atomic structure of biological macromolecules.

High X-ray intensity

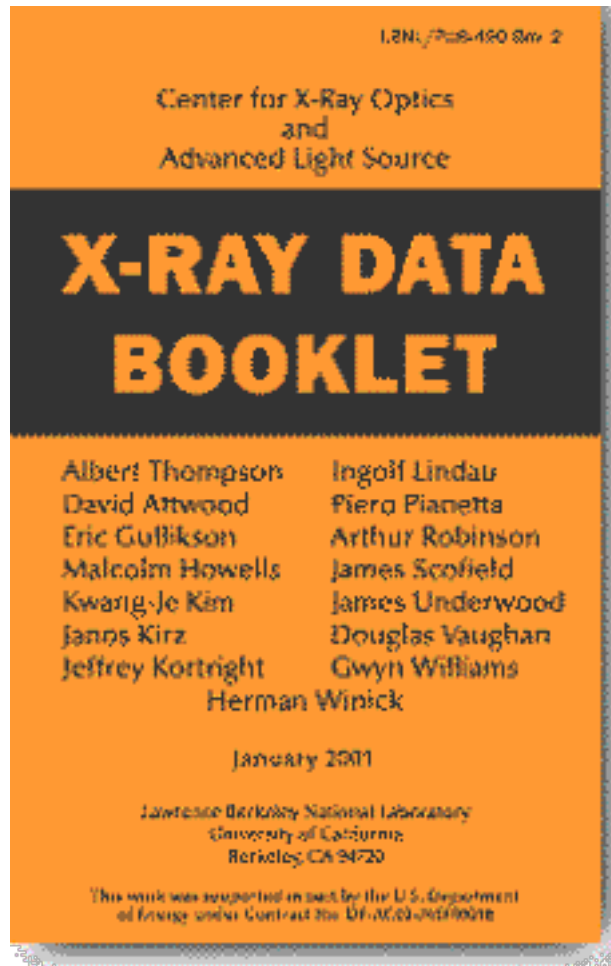


Son and Santra,
Phys. Rev. A **85**,
063415 (2012)

- Here we will discuss only one-photon absorption limit.
- Later you will see more interesting physics with intense x-ray pulses.

Main references

<http://xdb.lbl.gov>



R. Santra, *J. Phys. B* **42**, 023001 (2010).

IOP PUBLISHING
J. Phys. B: At. Mol. Opt. Phys. **42** (2009) 023001 (16pp)
doi:10.1088/0953-4075/42/2/023001

PHD TUTORIAL

Concepts in x-ray physics

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Abstract

A basic introduction to the theory underlying x-ray processes is provided. After general remarks on the practical advantages of using x-rays for probing matter, the derivation of the minimal-coupling Hamiltonian within nonrelativistic quantum electrodynamics is outlined. Perturbation theory is reviewed and applied to describe x-ray-induced processes. In connection with x-ray absorption, inner-shell binding energies and the photon energy dependence of the x-ray absorption cross section are discussed. In the context of x-ray scattering, atomic and molecular scattering factors are introduced, the complex index of refraction is derived, and the nonrelativistic theory of Compton scattering is described. The final topic is x-ray fluorescence and Auger decay of inner-shell-excited systems.

1. Introductory remarks

Since their discovery in the year 1895 by Wilhelm Conrad Röntgen [1], x-rays have become an indispensable tool for studying the structure and electronic properties of matter. Equally important is the role x-rays have come to play in medicine, archaeology, art, security, astronomy and other applications. To date, 19 Nobel prizes have been awarded for x-ray-related research (W C Röntgen 1901, M von Laue 1914, W H Bragg and W L Bragg 1915, C G Barkla 1917, K M G Siegbahn 1924, A H Compton 1927, P J W Debye 1936, H J Muller 1946, M F Perutz and J C Kendrew 1962, F H C Crick, J D Watson, and M H F Wilkins 1962, D Crowfoot Hodgkin 1964, W N Lipscomb 1976, A M Cormack and G N Hounsfield 1979, K M Siegbahn 1981, H A Hauptman and J Karle 1985, J Deisenhofer, R Huber and H Michel 1988, P D Boyer and J E Walker 1997, P Agre and R MacKinnon 2003, R Kornberg 2006).

As a new generation of x-ray sources—so-called x-ray free-electron lasers [2, 3]—is about to come online [4–6], it is timely to familiarize newcomers, experimentalists and theorists alike, with some of the basic properties that make x-rays such a powerful tool. This is attempted in this tutorial. The emphasis is on the development of a consistent theoretical framework that may be employed to describe a variety of x-ray processes. Throughout, the x-rays are assumed to be used as

a weak, essentially nonperturbative *probe*. Basic applications of x-rays will be covered. However, since this is a tutorial, there is not a single topic that is discussed in great depth, and some topics, e.g., x-ray sources and x-ray optics [7, 8], are not discussed at all.

This tutorial is structured as follows. In section 2, the derivation of the Hamiltonian underlying nonrelativistic quantum electrodynamics is sketched. This Hamiltonian, in combination with perturbation theory (section 3), allows one to describe all basic x-ray processes. X-ray absorption is the topic of section 4. In section 5, x-ray scattering processes are discussed. Finally, in section 6, relaxation processes following the excitation of an inner-shell electron are treated.

Atomic units are employed, i.e., $m_e = 1$, $|e| = 1$, $\hbar = 1$ and $c = 1/\alpha$, where m_e is the electron mass, e is the electron charge, \hbar is Planck's constant divided by 2π , c is the speed of light in vacuum and $\alpha = \frac{e^2}{\hbar c} \approx 1/137$ is the fine-structure constant. The atomic unit of length is the bohr, $a_0 = \frac{\hbar^2}{m_e e^2} \approx 0.529 \text{ \AA}$. The atomic unit of cross section is $a_0^2 \approx 28.0 \text{ Mb}$, where a *barn* (b) equals 10^{-28} m^2 . The atomic unit of energy is the hartree, $E_h = m_e c^2 \alpha^2 \approx 27.2 \text{ eV}$.

Depending on the photon energy, we typically distinguish between the following two x-ray regimes [7, 8]. Photons with an energy between $\sim 10 \text{ eV}$ ($\sim 300 \text{ eV}$) and $\sim 100 \text{ eV}$ (a few keV) are called *soft x-rays*. Soft x-rays cover, roughly, the 1s binding energies for elements ranging from carbon

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How to treat interaction between X-ray and matter?

Total Hamiltonian

$$\hat{H} = \hat{H}_{\text{EM}} + \hat{H}_{\text{mol}} + \hat{H}_{\text{int}}$$

- > X-rays interact with the electrons only
- > Non-relativistic quantum electrodynamics
- > Principle of Minimal Coupling
- > Time-dependent perturbation theory

Atomic units used:

$$m_e = 1$$

$$|e| = 1$$

$$\hbar = 1$$

$$1/4\pi\epsilon_0 = 1$$

$$(c = 1/\alpha, a_0 = 1)$$

Hamiltonian for free electromagnetic field

Hamiltonian for the free electromagnetic field (harmonic oscillator)

$$\hat{H}_{\text{EM}} = \sum_{\mathbf{k}, \lambda} \omega_{\mathbf{k}} \hat{a}_{\mathbf{k}, \lambda}^{\dagger} \hat{a}_{\mathbf{k}, \lambda}, \quad \omega_{\mathbf{k}} = |\mathbf{k}|/\alpha$$

Photon mode: wave vector \mathbf{k} and polarization vector $\boldsymbol{\epsilon}_{\mathbf{k}, \lambda}$
$$\begin{cases} \mathbf{k} \cdot \boldsymbol{\epsilon}_{\mathbf{k}, \lambda} = 0, \\ \boldsymbol{\epsilon}_{\mathbf{k}, 1}^* \cdot \boldsymbol{\epsilon}_{\mathbf{k}, 2} = 0 \end{cases}$$

Field operator: annihilates or creates a photon in the mode (\mathbf{k}, λ) .

$$\hat{a}_{\mathbf{k}, \lambda} |n_{\mathbf{k}, \lambda}\rangle = \sqrt{n_{\mathbf{k}, \lambda}} |n_{\mathbf{k}, \lambda} - 1\rangle$$

$$\hat{a}_{\mathbf{k}, \lambda}^{\dagger} |n_{\mathbf{k}, \lambda}\rangle = \sqrt{n_{\mathbf{k}, \lambda} + 1} |n_{\mathbf{k}, \lambda} + 1\rangle$$

Commutator relations for
photon mode operators

$$[\hat{a}_{\mathbf{k}, \lambda}, \hat{a}_{\mathbf{k}', \lambda'}] = 0$$

$$[\hat{a}_{\mathbf{k}, \lambda}, \hat{a}_{\mathbf{k}', \lambda'}^{\dagger}] = \delta_{\mathbf{k}, \mathbf{k}'} \delta_{\lambda, \lambda'}$$

$$[\hat{a}_{\mathbf{k}, \lambda}^{\dagger}, \hat{a}_{\mathbf{k}', \lambda'}^{\dagger}] = 0$$

$$[a, b] = ab - ba$$

Molecular Hamiltonian

Molecular Hamiltonian $\hat{H}_{\text{mol}} = \hat{T}_{\text{N}} + \hat{V}_{\text{NN}} + \hat{H}_{\text{el}}$

$$\hat{T}_{\text{N}} = - \sum_a \frac{\nabla_a^2}{2M_a} \quad \hat{V}_{\text{NN}} = \sum_{a < a'} \frac{Z_a Z_{a'}}{|\mathbf{R}_a - \mathbf{R}_{a'}|}$$

Electronic Hamiltonian using the second quantization

$$\begin{aligned} \hat{H}_{\text{el}} = & \int d^3x \, \hat{\psi}^\dagger(\mathbf{x}) \left\{ -\frac{\nabla^2}{2} - \sum_a \frac{Z_a}{|\mathbf{x} - \mathbf{R}_a|} \right\} \hat{\psi}(\mathbf{x}) \\ & + \frac{1}{2} \int d^3x \int d^3x' \, \hat{\psi}^\dagger(\mathbf{x}) \hat{\psi}^\dagger(\mathbf{x}') \frac{1}{|\mathbf{x} - \mathbf{x}'|} \hat{\psi}(\mathbf{x}') \hat{\psi}(\mathbf{x}) \end{aligned}$$

Anticommutator relations
for electron field operators

$$\begin{aligned} \{a, b\} &= ab + ba \\ \{\hat{\psi}_\sigma(\mathbf{x}), \hat{\psi}_{\sigma'}(\mathbf{x}')\} &= 0, \\ \{\hat{\psi}_\sigma(\mathbf{x}), \hat{\psi}_{\sigma'}^\dagger(\mathbf{x}')\} &= \delta_{\sigma, \sigma'} \delta^{(3)}(\mathbf{x} - \mathbf{x}'), \\ \{\hat{\psi}_\sigma^\dagger(\mathbf{x}), \hat{\psi}_{\sigma'}^\dagger(\mathbf{x}')\} &= 0. \end{aligned}$$

Interaction Hamiltonian

Interaction between electrons and photons
(minimal coupling, Coulomb gauge)

$\mathbf{p} \cdot \mathbf{A}$ and \mathbf{A}^2

$$\hat{H}_{\text{int}} = \alpha \int d^3x \hat{\psi}^\dagger(\mathbf{x}) \left[\hat{\mathbf{A}}(\mathbf{x}) \cdot \frac{\nabla}{i} \right] \hat{\psi}(\mathbf{x}) + \frac{\alpha^2}{2} \int d^3x \hat{\psi}^\dagger(\mathbf{x}) \hat{A}^2(\mathbf{x}) \hat{\psi}(\mathbf{x})$$

Mode expansion of
vector potential

$$\hat{\mathbf{A}}(\mathbf{x}) = \sum_{\mathbf{k}, \lambda} \sqrt{\frac{2\pi}{V\omega_{\mathbf{k}}\alpha^2}} \left\{ \hat{a}_{\mathbf{k}, \lambda} \boldsymbol{\epsilon}_{\mathbf{k}, \lambda} e^{i\mathbf{k} \cdot \mathbf{x}} + \hat{a}_{\mathbf{k}, \lambda}^\dagger \boldsymbol{\epsilon}_{\mathbf{k}, \lambda}^* e^{-i\mathbf{k} \cdot \mathbf{x}} \right\}$$

$\mathbf{p} \cdot \mathbf{A}$: one-photon absorption in first order; one-photon scattering in second order

\mathbf{A}^2 : one-photon scattering in first order

Time-dependent perturbation theory

Full Hamiltonian and unperturbed part $\hat{H} = \underbrace{\hat{H}_{\text{mol}} + \hat{H}_{\text{EM}}}_{\hat{H}_0} + \hat{H}_{\text{int}}$

Initial state: $|I\rangle = |\Psi_0^{N_{\text{el}}}\rangle |N_{\text{EM}}\rangle$

Interaction-picture state vector to second order in perturbation

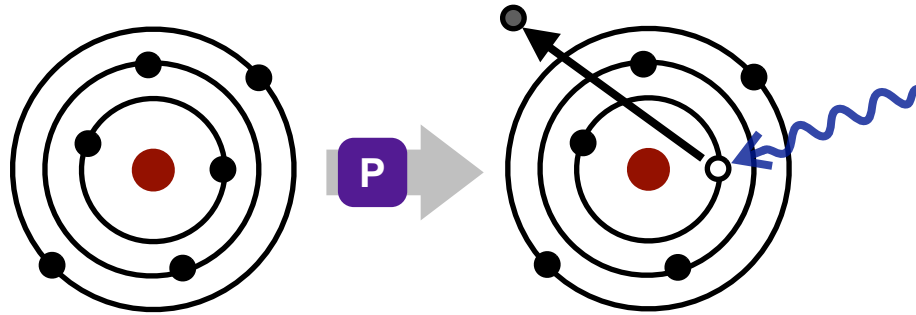
$$\begin{aligned} |\Psi, t\rangle_{\text{int}} = & |I\rangle - i \int_{-\infty}^t dt' e^{i\hat{H}_0 t'} \hat{H}_{\text{int}} e^{-\epsilon|t'|} e^{-i\hat{H}_0 t'} |I\rangle \\ & - \int_{-\infty}^t dt' e^{i\hat{H}_0 t'} \hat{H}_{\text{int}} e^{-\epsilon|t'|} e^{-i\hat{H}_0 t'} \int_{-\infty}^{t'} dt'' e^{i\hat{H}_0 t''} \hat{H}_{\text{int}} e^{-\epsilon|t''|} e^{-i\hat{H}_0 t''} |I\rangle + \dots \end{aligned}$$

Transition rate to second order in perturbation

$$\Gamma_{FI} = 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle + \sum_M \frac{\langle F | \hat{H}_{\text{int}} | M \rangle \langle M | \hat{H}_{\text{int}} | I \rangle}{E_I - E_M - i\epsilon} + \dots \right|^2$$

X-ray absorption

X-ray absorption



Initial state: $|I\rangle = |\Psi_0^{N_{\text{el}}}\rangle |N_{\text{EM}}\rangle$

Final state: $|F\rangle = |\Psi_F^{N_{\text{el}}}\rangle |N_{\text{EM}} - 1\rangle$

Absorption rate: $\Gamma_{FI} = 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2$

X-ray absorption cross section

X-ray
absorption rate

$$\Gamma_{FI} = \frac{4\pi^2}{\omega_{\text{in}}} \frac{N_{\text{EM}}}{V} \delta(E_F^{N_{\text{el}}} - E_0^{N_{\text{el}}} - \omega_{\text{in}}) \times \left| \langle \Psi_F^{N_{\text{el}}} | \int d^3x \hat{\psi}^\dagger(\mathbf{x}) e^{i\mathbf{k}_{\text{in}} \cdot \mathbf{x}} \boldsymbol{\epsilon}_{\mathbf{k}_{\text{in}}, \lambda_{\text{in}}} \cdot \frac{\nabla}{i} \hat{\psi}(\mathbf{x}) | \Psi_0^{N_{\text{el}}} \rangle \right|^2$$

X-ray photon flux

$$J_{\text{EM}} = \frac{1}{\alpha} \frac{N_{\text{EM}}}{V}$$

X-ray absorption cross section

$$\sigma = \Gamma_{FI} / J_{\text{EM}}$$

$$\sigma_F(\mathbf{k}_{\text{in}}, \lambda_{\text{in}}) = \frac{4\pi^2}{\omega_{\text{in}}} \alpha \delta(E_F^{N_{\text{el}}} - E_0^{N_{\text{el}}} - \omega_{\text{in}}) \times \left| \langle \underline{\Psi_F^{N_{\text{el}}}} | \int d^3x \hat{\psi}^\dagger(\mathbf{x}) e^{i\mathbf{k}_{\text{in}} \cdot \mathbf{x}} \boldsymbol{\epsilon}_{\mathbf{k}_{\text{in}}, \lambda_{\text{in}}} \cdot \frac{\nabla}{i} \hat{\psi}(\mathbf{x}) | \underline{\Psi_0^{N_{\text{el}}}} \rangle \right|^2$$

how to treat many-electron systems?

Electronic many-body problem

Introducing a set of orthonormal spin orbitals $\{\varphi_p(\mathbf{x})\}$

which are eigenstates of an one-body Hamiltonian $\hat{F}|\varphi_p\rangle = \varepsilon_p|\varphi_p\rangle$

An electronic one-body Hamiltonian
$$\hat{F} = \sum_p \varepsilon_p \hat{c}_p^\dagger \hat{c}_p$$

Creation and annihilation
operators for $|\varphi_p\rangle$: \hat{c}_p^\dagger and \hat{c}_p

$$\begin{aligned}\{\hat{c}_p, \hat{c}_q\} &= 0, \\ \{\hat{c}_p, \hat{c}_q^\dagger\} &= \delta_{p,q}, \\ \{\hat{c}_p^\dagger, \hat{c}_q^\dagger\} &= 0.\end{aligned}$$

Connecting with the field operators

$$\begin{aligned}\hat{\psi}(\mathbf{x}) &= \sum_p \varphi_p(\mathbf{x}) \hat{c}_p, \\ \hat{\psi}^\dagger(\mathbf{x}) &= \sum_p \varphi_p^\dagger(\mathbf{x}) \hat{c}_p^\dagger\end{aligned}$$

Mean-field model

Independent-electron model: $\hat{H}_{\text{el}} \approx \hat{F}$

Mean-field ground state in terms of a single Slater determinant:

$$|\Psi_0^{N_{\text{el}}}\rangle \approx |\Phi_0^{N_{\text{el}}}\rangle \equiv \prod_{i=1}^{N_{\text{el}}} \hat{c}_i^\dagger |\text{vacuum}\rangle$$

Mean-field ground-state energy:

$$E_0^{N_{\text{el}}} \equiv \langle \Psi_0^{N_{\text{el}}} | \hat{H}_{\text{el}} | \Psi_0^{N_{\text{el}}} \rangle \approx \langle \Phi_0^{N_{\text{el}}} | \hat{F} | \Phi_0^{N_{\text{el}}} \rangle = \sum_{i=1}^{N_{\text{el}}} \varepsilon_i$$

Hartree-Fock (HF) method: minimizing $\langle \Phi_0^{N_{\text{el}}} | \hat{H}_{\text{el}} - \hat{F} | \Phi_0^{N_{\text{el}}} \rangle$

the best choice among the mean-field models

Koopmans' theorem

One-hole state $|\Phi_i^{N_{\text{el}}-1}\rangle \equiv \hat{c}_i |\Phi_0^{N_{\text{el}}}\rangle$

Energy of one-hole state $E_i^{N_{\text{el}}-1} \approx \langle \Phi_i^{N_{\text{el}}-1} | \hat{F} | \Phi_i^{N_{\text{el}}-1} \rangle = \left\{ \sum_{j=1}^{N_{\text{el}}} \varepsilon_j \right\} - \varepsilon_i$

Ionization potential $I_i \equiv E_i^{N_{\text{el}}-1} - E_0^{N_{\text{el}}} \approx -\varepsilon_i$

Koopmans' theorem: valid for the HF solution

$$I_i = \langle \Phi_i^{\text{HF}, N_{\text{el}}-1} | \hat{H}_{\text{el}} | \Phi_i^{\text{HF}, N_{\text{el}}-1} \rangle - \langle \Phi_0^{\text{HF}, N_{\text{el}}} | \hat{H}_{\text{el}} | \Phi_0^{\text{HF}, N_{\text{el}}} \rangle = -\varepsilon_i$$

Note: Koopmans received a Nobel prize in Economics in 1975, of course, not for this theorem.

For Neon, Koopmans	EXP
$-\varepsilon_{2p} = 23.1 \text{ eV}$	21.6 eV
$-\varepsilon_{2s} = 52.5 \text{ eV}$	48.5 eV
$-\varepsilon_{1s} = 892 \text{ eV}$	870 eV

X-ray absorption within mean-field model

Particle-hole state $|\Phi_i^a\rangle \equiv \hat{c}_a^\dagger |\Phi_i^{N_{\text{el}}-1}\rangle = \hat{c}_a^\dagger \hat{c}_i |\Phi_0^{N_{\text{el}}}\rangle$

Initial state $|I\rangle = |\Phi_0^{N_{\text{el}}}\rangle |N_{\text{EM}}\rangle$

Final state $|F\rangle = \hat{c}_a^\dagger \hat{c}_i |\Phi_0^{N_{\text{el}}}\rangle |N_{\text{EM}}-1\rangle$

Cross section for particle-hole transition

$$\sigma_i^a(\mathbf{k}_{\text{in}}, \lambda_{\text{in}}) = \frac{4\pi^2}{\omega_{\text{in}}} \alpha \delta(\varepsilon_a + I_i - \omega_{\text{in}}) \left| \langle \varphi_a | e^{i\mathbf{k}_{\text{in}} \cdot \mathbf{x}} \boldsymbol{\epsilon}_{\mathbf{k}_{\text{in}}, \lambda_{\text{in}}} \cdot \frac{\nabla}{i} | \varphi_i \rangle \right|^2$$

Subshell x-ray absorption cross section

$$\sigma_i(\mathbf{k}_{\text{in}}, \lambda_{\text{in}}) = \sum_a \sigma_i^a(\mathbf{k}_{\text{in}}, \lambda_{\text{in}})$$

Energy of excited electron

$$\varepsilon_a = \omega_{\text{in}} - I_i > 0 \quad \text{“Photoelectric effect”}$$

Note: Einstein received a Nobel prize in 1921 for the photoelectric effect, not for theory of relativity.

X-ray absorption cross section

- > After employing the dipole approximation and atomic orbitals, X-ray absorption cross section of i -th orbital at ω_{in} is calculated as

$$\sigma(i, \omega_{\text{in}}) = \frac{4}{3} \alpha \pi^2 \omega_{\text{in}} N_i \sum_{l_j=|l_i-1|}^{l_i+1} \frac{l_{>}}{2l_i + 1} \left| \int_0^\infty \underbrace{u_{n_i l_i}(r)}_{\text{bound-state orbital}} \underbrace{u_{\epsilon l_j}(r)}_{\text{continuum-state orbital}} r dr \right|^2$$

bound-state
orbital

continuum-state
orbital

K-shell photoionization: a simple model

- > Simple model of an atomic species of atomic number Z

- 1s orbital \sim H-like 1s orbital of Z ,

$$u_{1s}(r) \sim Z^{3/2} r e^{-Zr}, \quad I_{1s} = Z^2/2$$

- photoelectron wave function \sim plane wave (assuming $\omega_{\text{in}} \gg I_{1s}$)

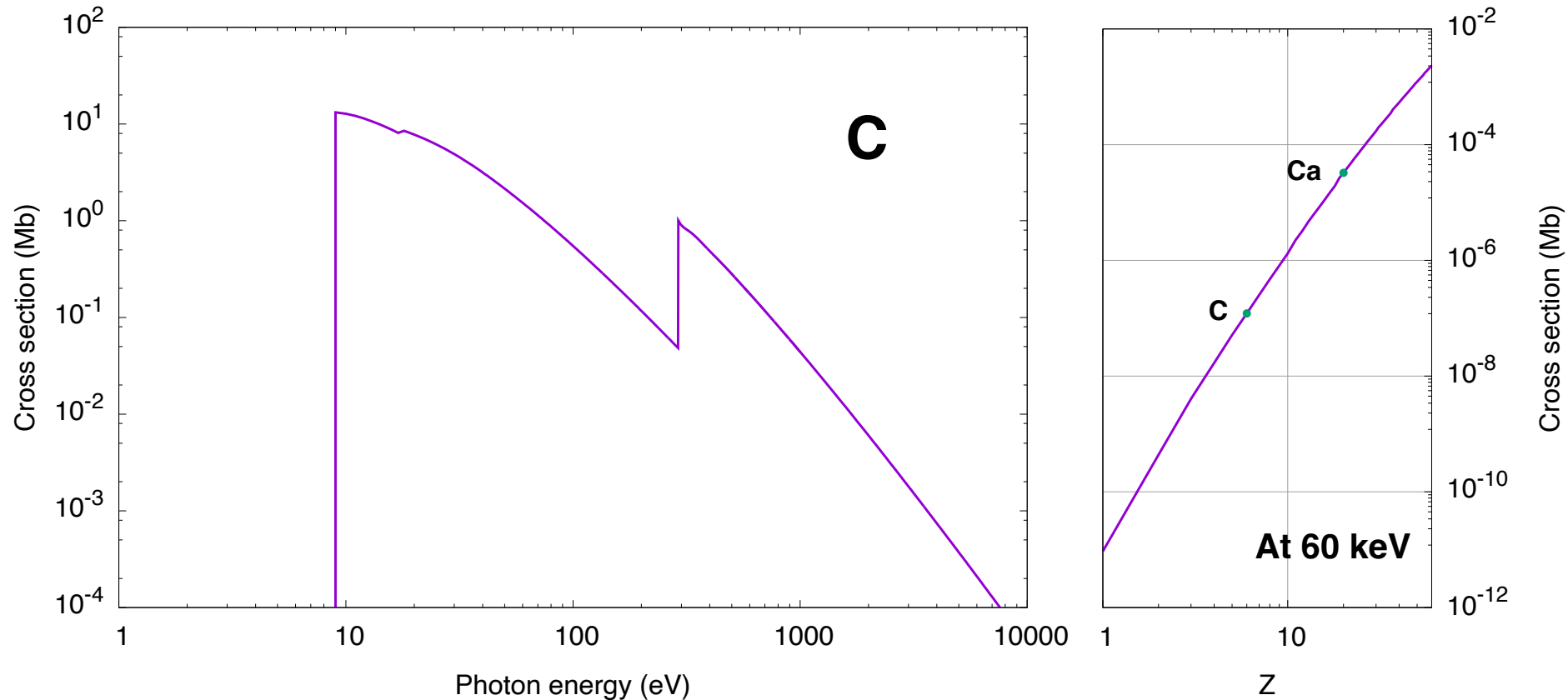
- > At a fixed photon energy ω_{in} , the calculated cross section: $\sigma \sim Z^5$

Compare σ_{C} ($Z=6$) and σ_{Ca} ($Z=20$): $(20/6)^5 \sim 400$ times

- > At a fixed Z , the calculated cross section: $\sigma \sim \omega_{\text{in}}^{-7/2}$

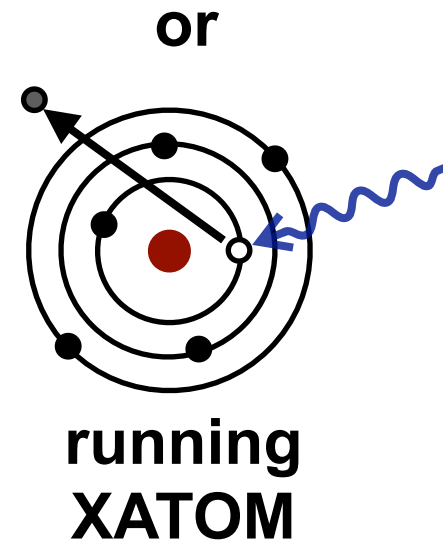
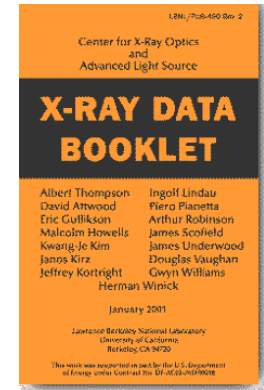
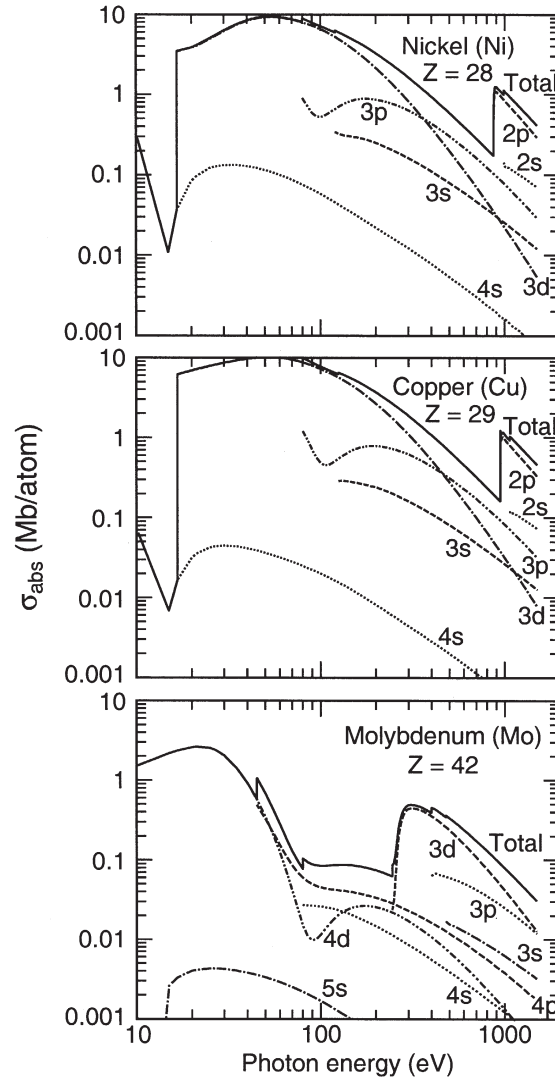
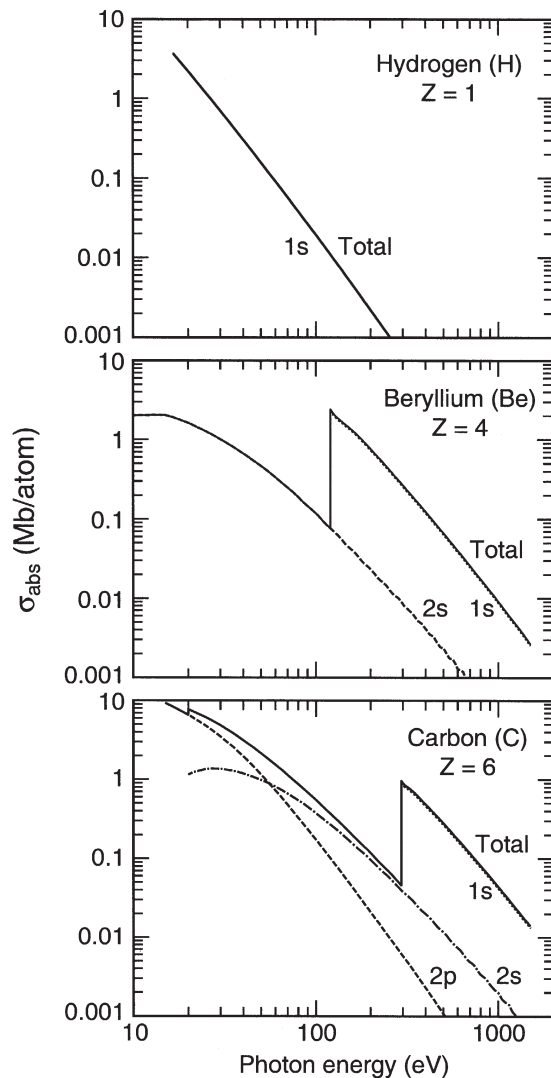
For a given atom, the x-ray absorption cross section above the 1s threshold decreases with increasing photon energy.

Photoionization cross section calculation



➤ PCS changes by varying ω_{in} and Z ; edge structures

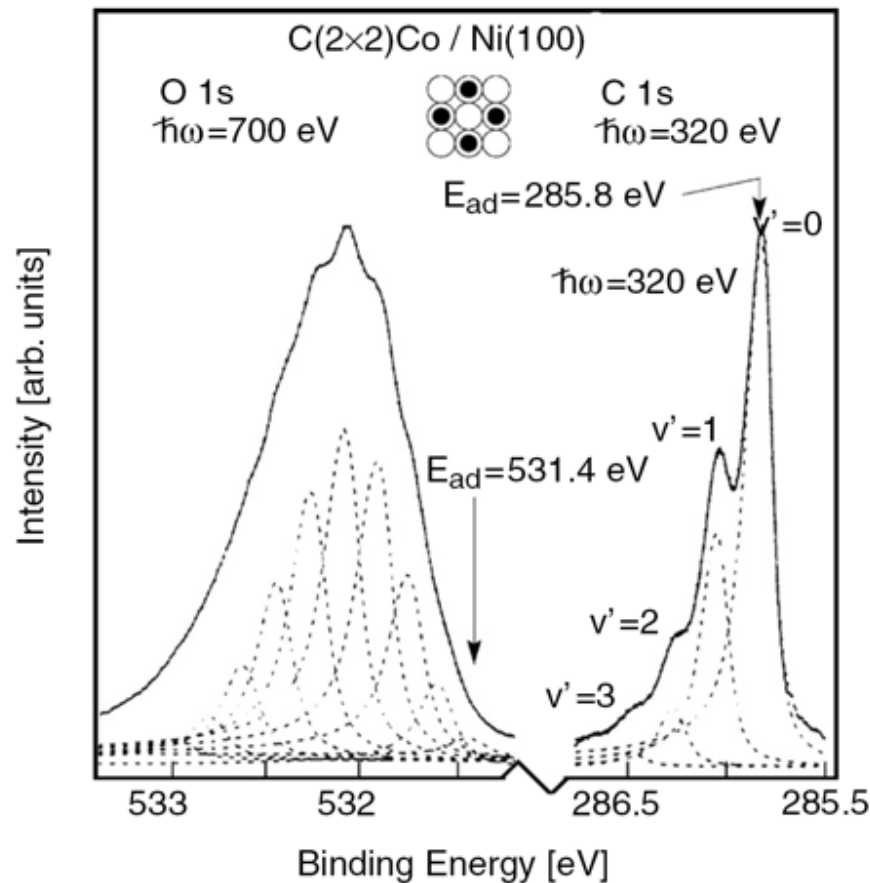
Table of photoionization cross sections



Applications of X-ray absorption

- X-ray photoelectron spectroscopy (XPS): measuring *photoelectron energy*
 - PES: Photo-Electron Spectroscopy or Photo-Emission Spectroscopy
 - ESCA: Electron Spectroscopy for Chemical Analysis
- X-ray absorption spectroscopy (XAS): measuring *absorption cross section* as a function of *photon energy*
 - XANES: X-ray Absorption Near-Edge Structure
 - NEXAFS: Near Edge X-Ray Absorption Fine Structure
 - EXAFS: Extended X-ray Absorption Fine Structure

XPS: X-ray Photoelectron Spectroscopy



PES: Photo-Emission
(or Photo-Electron)
Spectroscopy

$$\varepsilon_a = \omega_{in} - I_i$$

↑ measured ↑ fixed

Figure 8. High-resolution XPS spectra of C1s and O1s from a $c(2 \times 2)$ CO/Ni(100) [38]. The fine structure is due to intramolecular stretch vibrations of the adsorbed CO molecules.

ESCA: Electron Spectroscopy for Chemical Analysis

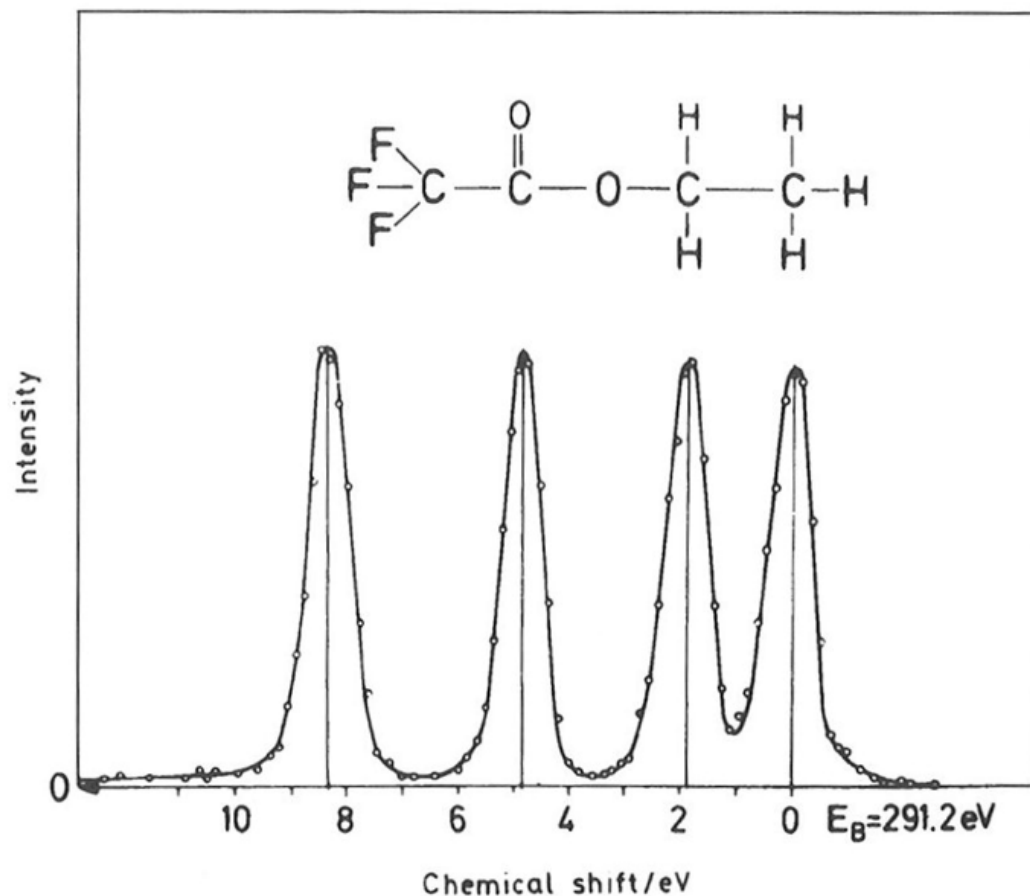


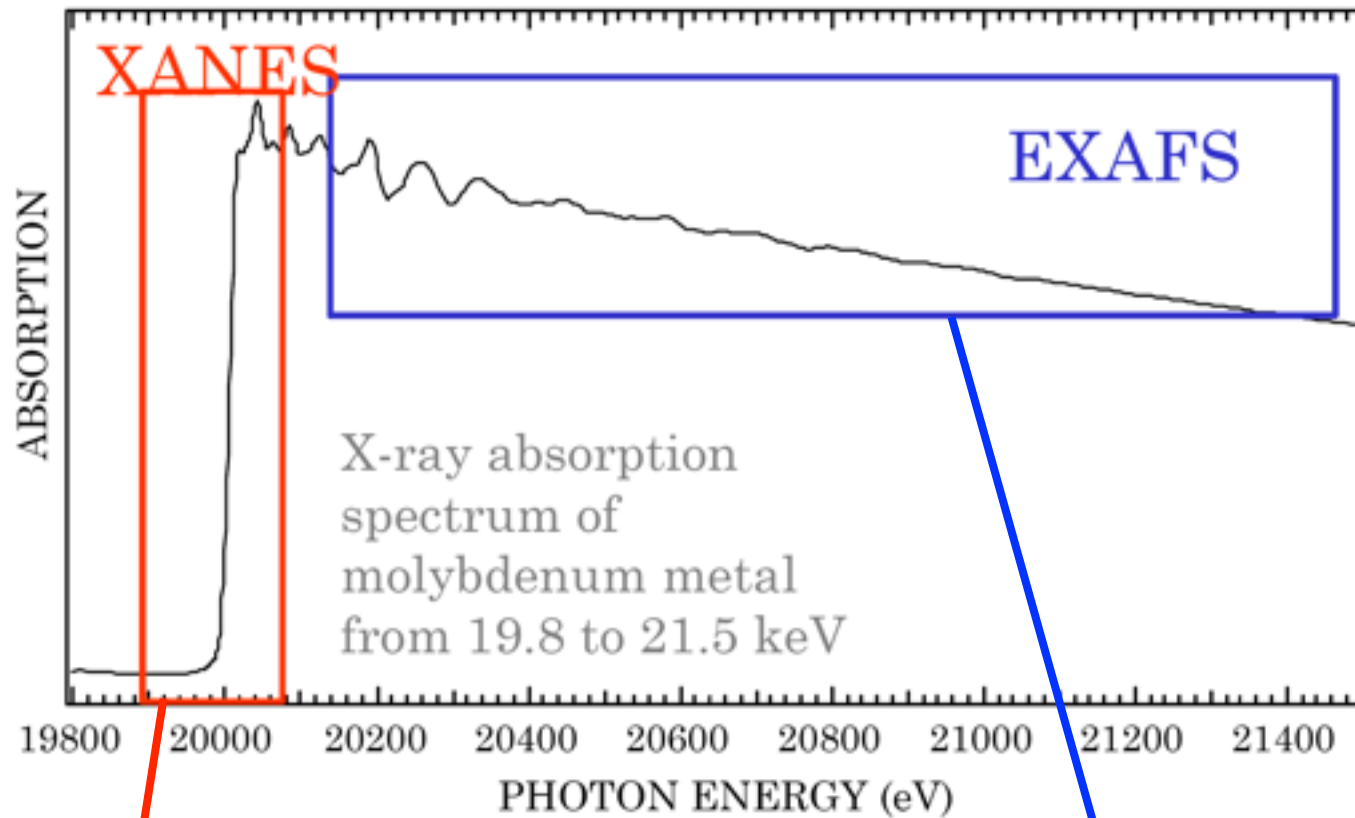
Figure 8.14 The monochromatized AlK α carbon 1s XPS spectrum of ethyltrifluoroacetate showing the chemical shifts relative to an ionization energy of 291.2 eV. (Reproduced, with permission, from Gelius, U., Basilier, E., Svensson, S., Bergmark, T., and Siegbahn, K., *J. Electron Spectrosc.*, 2, 405, 1974)

$$\varepsilon_a = \omega_{\text{in}} - I_i$$

↑ measured ↑ fixed

ESCA molecule:
first molecule used for
ESCA by Kai Siegbahn
(Nobel prize in 1981)

XAS: X-ray Absorption Spectroscopy



$$\varepsilon_a = \omega_{\text{in}} - I_i$$

↑
varied

solid target:
transmission

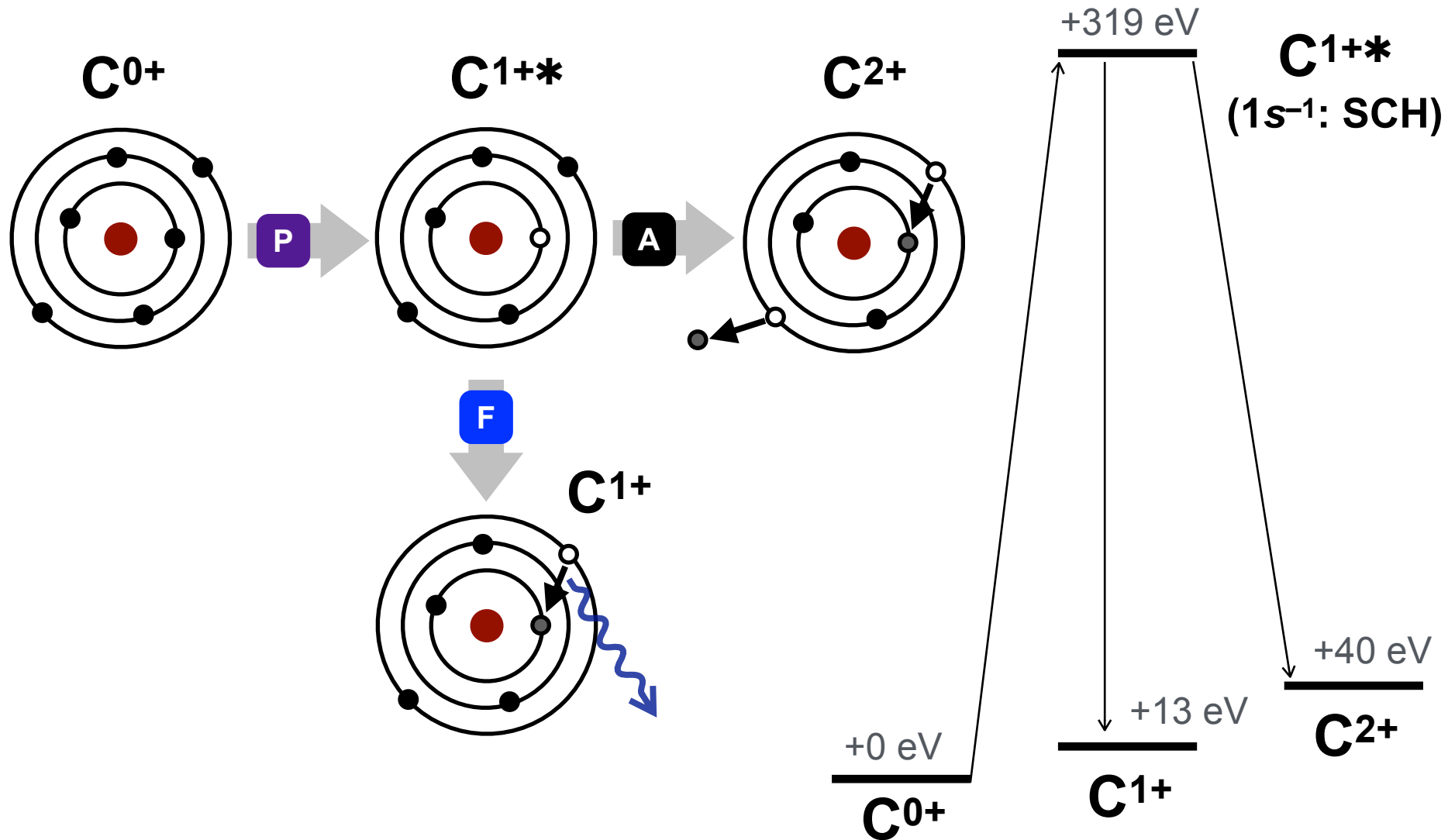
gas phase or very
dilute target:
fluorescence

NE: depends on local coordination
environment and oxidation state

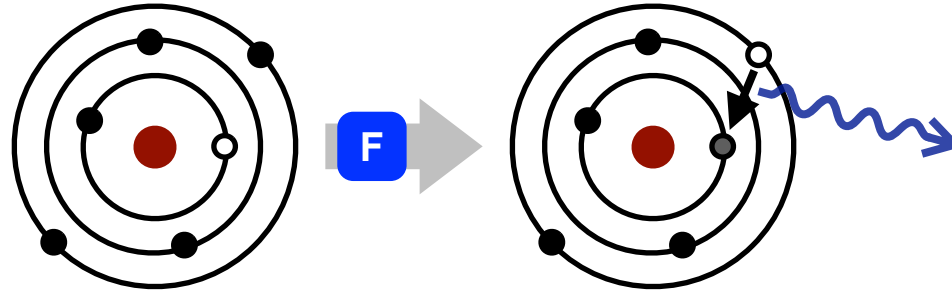
E: depends on nearest-neighbor
geometry

Relaxation processes

Decay of inner-shell-excited systems



X-ray fluorescence within mean-field model



Initial state: $|I\rangle = \hat{c}_i |\Phi_0^{N_{\text{el}}}\rangle |0\rangle$

Final state: $|F\rangle = \hat{c}_{i'} |\Phi_0^{N_{\text{el}}}\rangle \hat{a}_{\mathbf{k}_F, \lambda_F}^\dagger |0\rangle$

Fluorescence rate: $\Gamma_{FI} = 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2$

X-ray fluorescence rate and energy

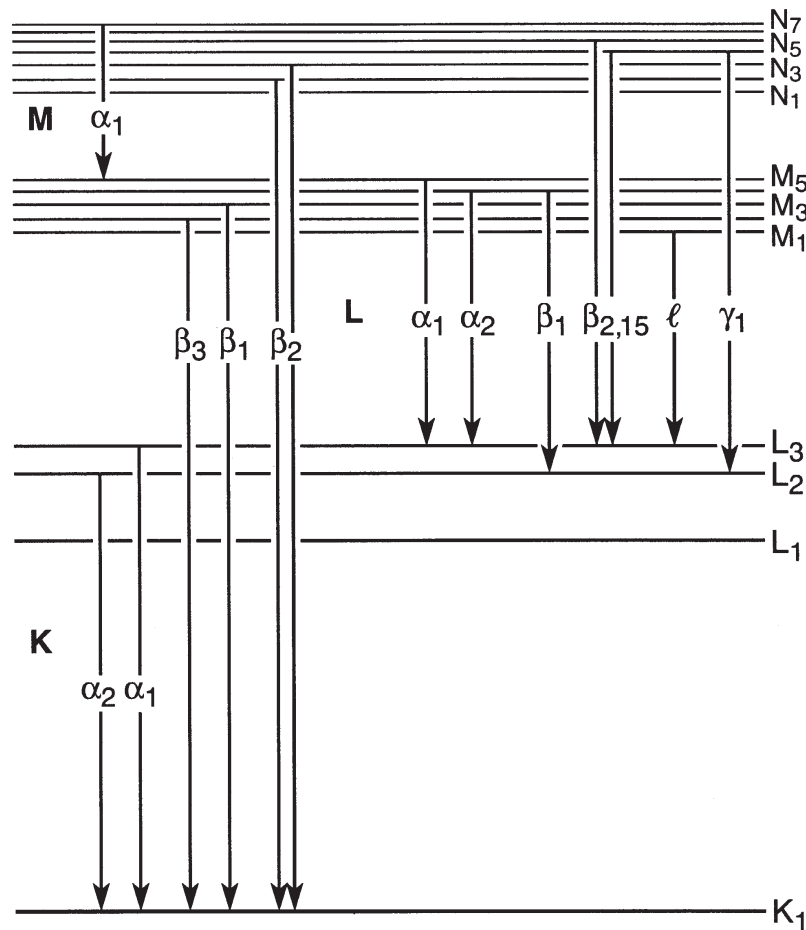
Partial X-ray fluorescence rate

$$\begin{aligned}\Gamma_{FI} &= 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2 \\ &= \frac{4\pi^2}{V\omega_F} \delta(I_{i'} + \omega_F - I_i) \left| \langle \varphi_i | e^{-i\mathbf{k}_F \cdot \mathbf{x}} \boldsymbol{\epsilon}_{\mathbf{k}_F, \lambda_F}^* \cdot \frac{\nabla}{i} | \varphi_{i'} \rangle \right|^2\end{aligned}$$

$$\Gamma_{i,i'} = \sum_{\mathbf{k}_F, \lambda_F} \Gamma_{FI} = \sum_{\lambda_F} \frac{V}{(2\pi)^3} \int d\Omega_F \int_0^\infty dk_F k_F^2 \Gamma_{FI}$$

X-ray fluorescence energy $\omega_F = I_i - I_{i'}$

X-ray fluorescence transition lines



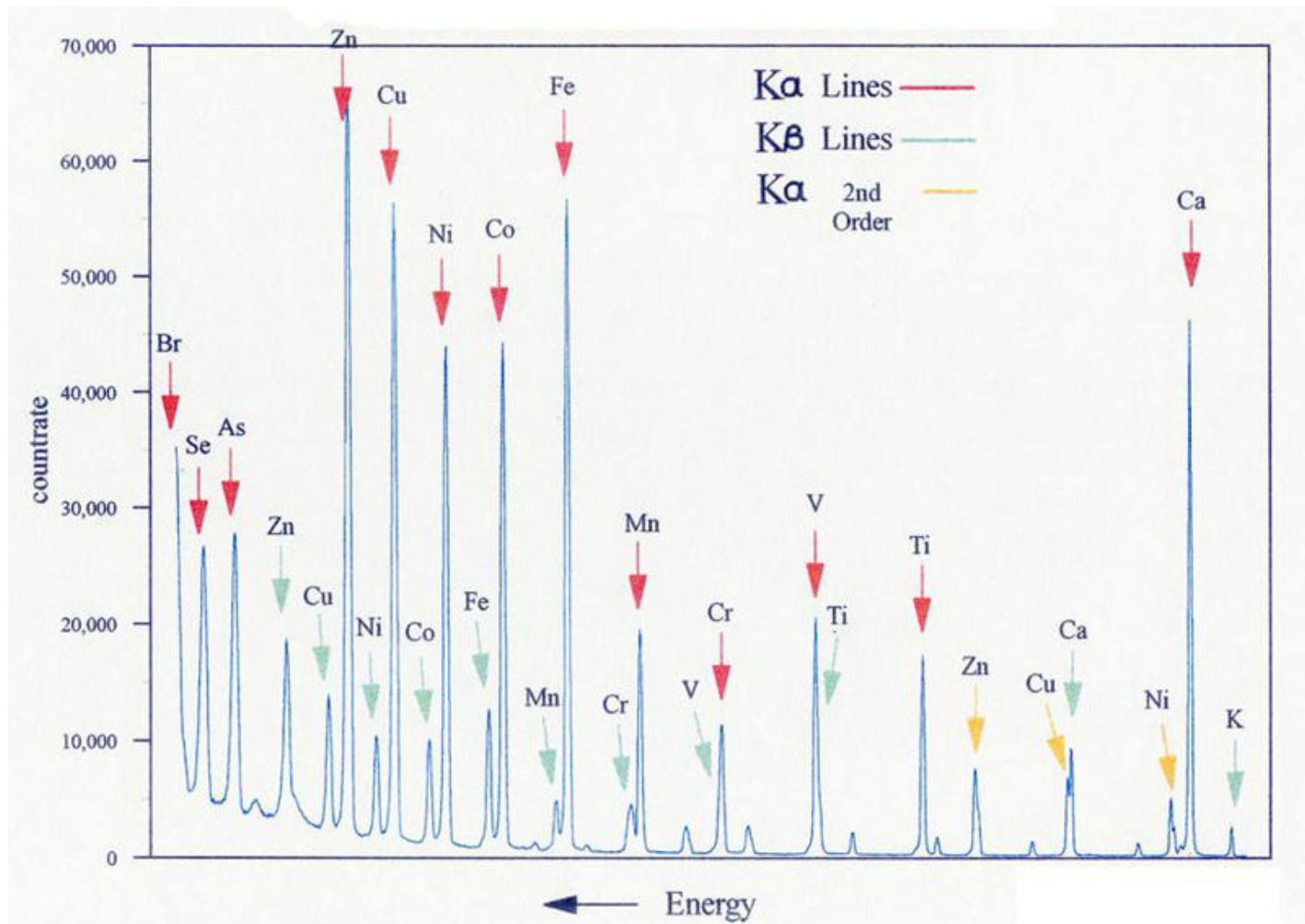
Characteristic transition energies
→ fingerprint of elements

Energy (eV)	Element	Line	Relative intensity
54.3	3 Li	$K\alpha_{1,2}$	150
108.5	4 Be	$K\alpha_{1,2}$	150
183.3	5 B	$K\alpha_{1,2}$	151
277	6 C	$K\alpha_{1,2}$	147
348.3	21 Sc	Ll	21
392.4	7 N	$K\alpha_{1,2}$	150
395.3	22 Ti	Ll	46
395.4	21 Sc	$L\alpha_{1,2}$	111
399.6	21 Sc	$L\beta_1$	77
446.5	23 V	Ll	28
452.2	22 Ti	$L\alpha_{1,2}$	111
458.4	22 Ti	$L\beta_1$	79
500.3	24 Cr	Ll	17
511.3	23 V	$L\alpha_{1,2}$	111
519.2	23 V	$L\beta_1$	80

524.9	8 O	$K\alpha_{1,2}$	151
556.3	25 Mn	Ll	15
572.8	24 Cr	$L\alpha_{1,2}$	111
582.8	24 Cr	$L\beta_1$	79
615.2	26 Fe	Ll	10
637.4	25 Mn	$L\alpha_{1,2}$	111
648.8	25 Mn	$L\beta_1$	77
676.8	9 F	$K\alpha_{1,2}$	148
677.8	27 Co	Ll	10
705.0	26 Fe	$L\alpha_{1,2}$	111
718.5	26 Fe	$L\beta_1$	66
742.7	28 Ni	Ll	9
776.2	27 Co	$L\alpha_{1,2}$	111
791.4	27 Co	$L\beta_1$	76
811.1	29 Cu	Ll	8
833	57 La	$M\alpha_1$	100
848.6	10 Ne	$K\alpha_{1,2}$	150

Fig. 1-1. Transitions that give rise to the emission lines in Table 1-3.

X-ray fluorescence transition lines (cont.)



K α fluorescence: a simple model

> Simple model of an atomic species of atomic number Z

- 1s orbital ~ H-like 1s orbital of Z
- 2p orbital ~ H-like 2p orbital of Z

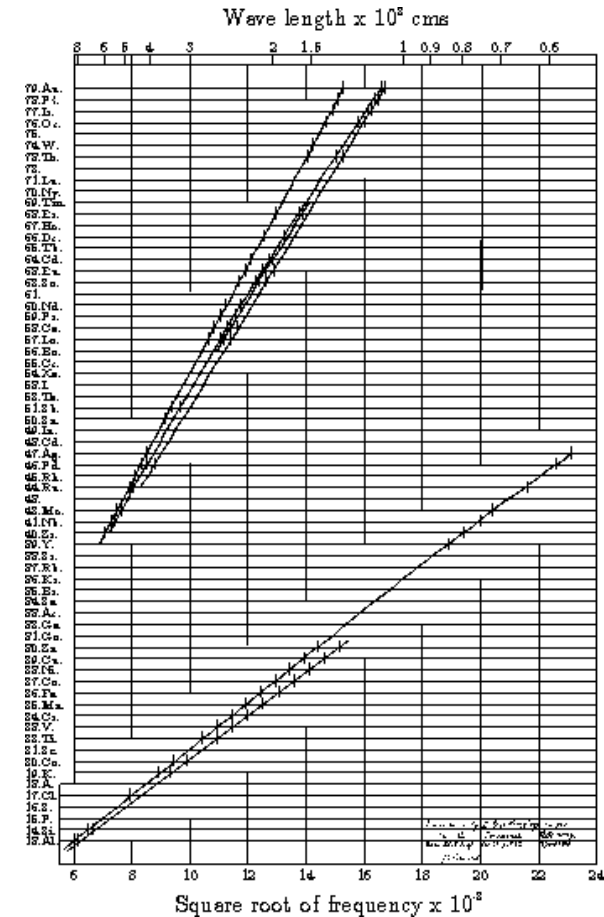
> K α photon energy: $\omega_{K\alpha} \sim Z^2$

> K α fluorescence rate: $\Gamma_{K\alpha} \sim Z^4$

Moseley's law

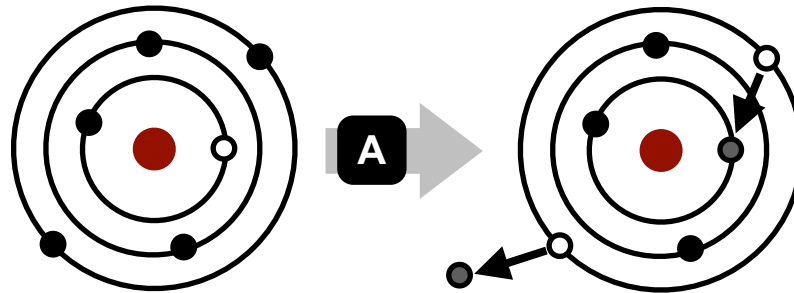
$$\sqrt{\omega} \propto Z$$

Phil. Mag. 27, 703 (1914).



Auger-Meitner within mean-field model

A renaming proposal to Auger-Meitner:
D. Matsakis *et al.*, *Phys. Today* **72**, 10 (2019).



Initial state: $|I\rangle = \hat{c}_i |\Phi_0^{N_{\text{el}}}\rangle$

Final state: $|F\rangle = \hat{c}_a^\dagger \hat{c}_j \hat{c}_{j'} |\Phi_0^{N_{\text{el}}}\rangle, \quad j, j' \neq i$

Auger rate: $\Gamma_{FI} = 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2$

Auger-Meitner rate and energy

Perturbation theory $\hat{H}_0 = \hat{F}$, $\hat{H}_{\text{int}} = \frac{1}{2} \sum_{pqrs} v_{pqrs} \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r - \sum_{pq} \sum_i \{v_{piqi} - v_{pii q}\} \hat{c}_p^\dagger \hat{c}_q$

Coulomb matrix element $v_{pqrs} = \int d^3x \int d^3x' \varphi_p^\dagger(\mathbf{x}) \varphi_q^\dagger(\mathbf{x}') \frac{1}{|\mathbf{x} - \mathbf{x}'|} \varphi_r(\mathbf{x}) \varphi_s(\mathbf{x}')$

Partial Auger-Meitner rate

$$\Gamma_{i,jj'} = 2\pi \sum_a |v_{aijj'} - v_{aij'j}|^2 \delta(\varepsilon_a + I_j + I_{j'} - I_i)$$

Auger-Meitner electron energy $\varepsilon_a = I_i - I_j - I_{j'}$

AES: Auger-Meitner Electron Spectrum

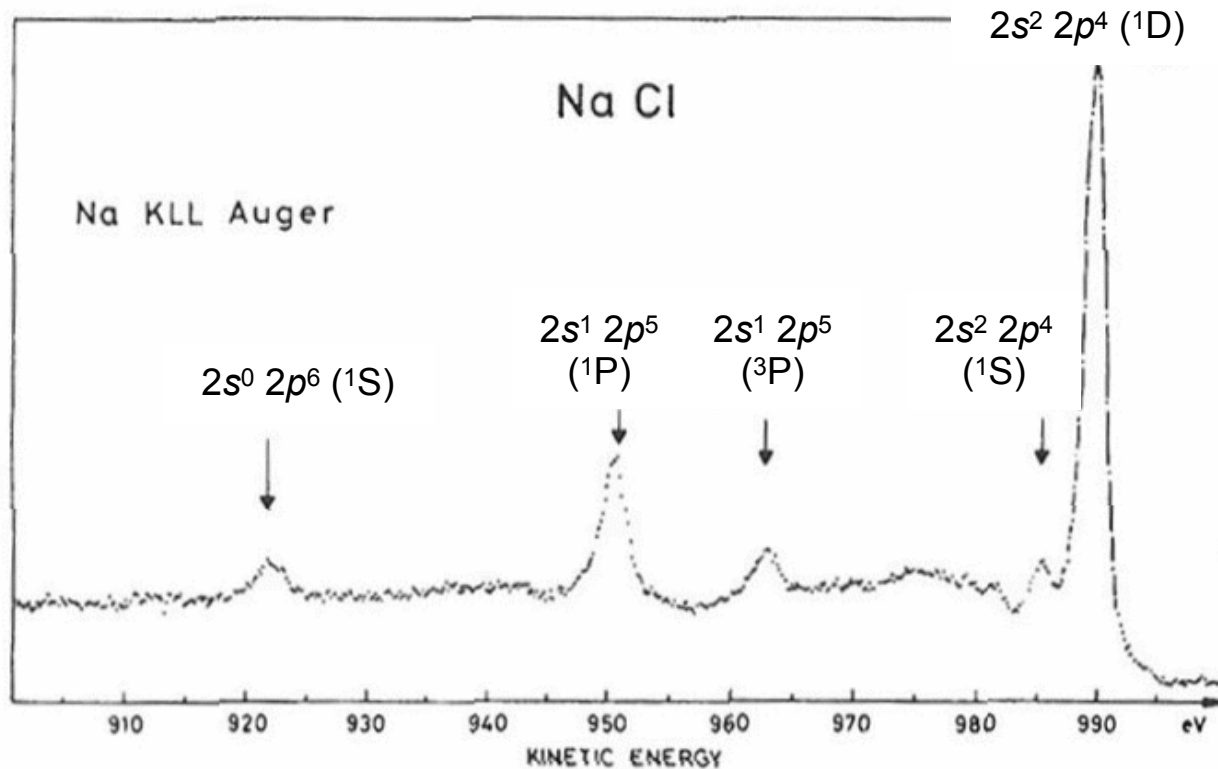


Figure 8.26 The $KL_{II,III}L_{II,III}$ (1D_2 and 1S_0) Auger spectrum of sodium in solid NaCl. (Reproduced, with permission, from Siegbahn, K. (June 1976), *Publication No. UUIP-940*, p. 81, Uppsala University Institute of Physics)

$$\varepsilon_a = I_i - I_j - I_{j'}$$

Auger-Meitner channels

$1s - 2s\ 2s : 2s^0 2p^6\ ^1S$

$1s - 2s\ 2p : 2s^1 2p^5\ ^1P$

3P

$1s - 2p\ 2p : 2s^2 2p^4\ ^1S$

~~3P~~

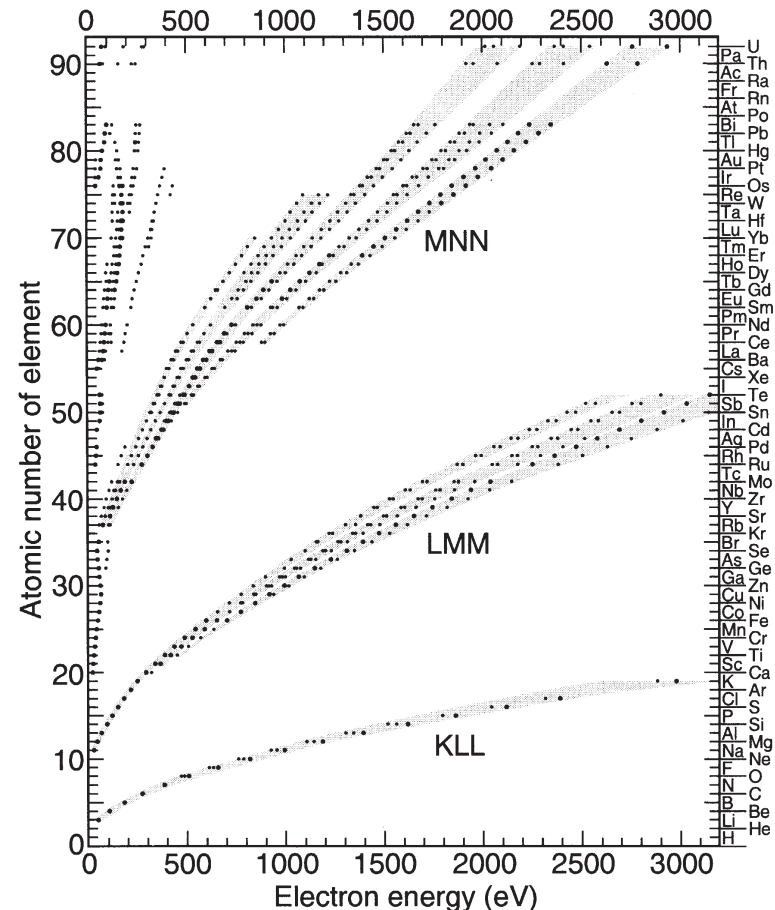
1D

position: AM energy

intensity: AM rate

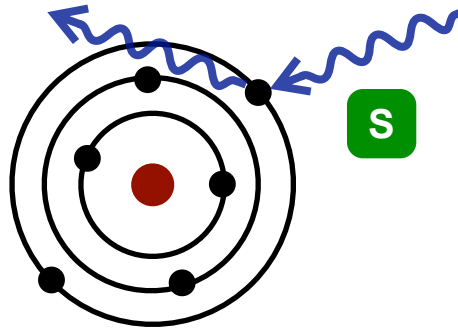
KLL Auger-Meitner decay: a simple model

- > Simple model of an atomic species of atomic number Z
 - 1s orbital ~ H-like 1s orbital of Z
 - 2p orbital ~ H-like 2p orbital of Z
 - Auger-Meitner electron wave function ~ plane wave
- > Focus on the channel with two 2p holes in the final state: *KLL*
- > *KLL* AM electron E : $\varepsilon_{KLL} \sim Z^2$
- > *KLL* AM decay rate: $\Gamma_{KLL} \sim Z^0$



X-ray scattering

Elastic X-ray scattering



Initial state: $|I\rangle = |\Psi_0^{N_{\text{el}}}\rangle |N_{\text{EM}}\rangle$

Final state: $|F\rangle = |\Psi_0^{N_{\text{el}}}\rangle \hat{a}_{\mathbf{k}_F, \lambda_F}^\dagger |N_{\text{EM}} - 1\rangle$

Scattering rate: $\Gamma_{FI} = 2\pi\delta(E_F - E_I) \left| \langle F | \hat{H}_{\text{int}} | I \rangle \right|^2$

Elastic X-ray scattering form factor

Transition rate

$$\Gamma_{FI} = \frac{(2\pi)^3}{V} \frac{N_{\text{EM}}}{V} \frac{1}{\omega_F \omega_{\text{in}}} \left| \boldsymbol{\epsilon}_{\mathbf{k}_F, \lambda_F}^* \cdot \boldsymbol{\epsilon}_{\mathbf{k}_{\text{in}}, \lambda_{\text{in}}} \right|^2 \left| f^0(\mathbf{k}_{\text{in}} - \mathbf{k}_F) \right|^2 \delta(\omega_F - \omega_{\text{in}})$$

Scattering vector
(or momentum transfer)

$$\mathbf{Q} = \mathbf{k}_{\text{in}} - \mathbf{k}_F$$

Scattering (or form) factor

$$f^0(\mathbf{Q}) = \int d^3x \rho^{(\text{el})}(\mathbf{x}) e^{i\mathbf{Q} \cdot \mathbf{x}}$$

Electron density

$$\rho^{(\text{el})}(\mathbf{x}) = \langle \Psi_0^{N_{\text{el}}} | \hat{\psi}^\dagger(\mathbf{x}) \hat{\psi}(\mathbf{x}) | \Psi_0^{N_{\text{el}}} \rangle$$

Elastic x-ray scattering cross section

Differential cross section

$$d\sigma = \sum_{\mathbf{k}_F, \lambda_F} \frac{\Gamma_{FI}}{J_{EM}} = \frac{V}{(2\pi)^3} \alpha^3 d\Omega \sum_{\lambda_F} \int_0^\infty d\omega_F \omega_F^2 \frac{\Gamma_{FI}}{J_{EM}}$$

$$\frac{d\sigma(\mathbf{Q})}{d\Omega} = \alpha^4 |f^0(\mathbf{Q})|^2 \sum_{\lambda_F} |\boldsymbol{\epsilon}_{\mathbf{k}_F, \lambda_F}^* \cdot \boldsymbol{\epsilon}_{\mathbf{k}_{in}, \lambda_{in}}|^2 = \left(\frac{d\sigma}{d\Omega} \right)_T |f^0(\mathbf{Q})|^2$$

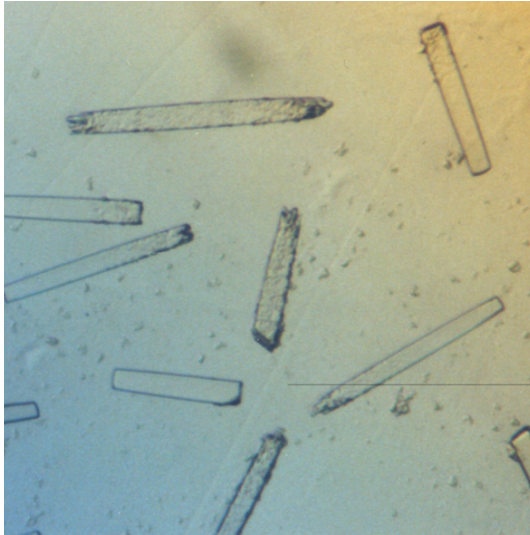
Total form factor
decomposed in terms of
atomic contributions

$$f^0(\mathbf{Q}) \approx \int d^3x \sum_n \rho_n^{(el)}(\mathbf{x} - \mathbf{R}_n) e^{i\mathbf{Q} \cdot \mathbf{x}}$$

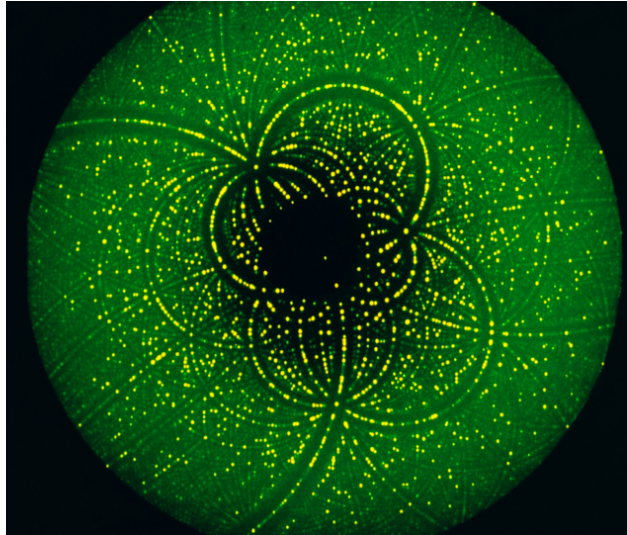
$$= \sum_n \left\{ \int d^3x' \rho_n^{(el)}(\mathbf{x}') e^{i\mathbf{Q} \cdot \mathbf{x}'} \right\} e^{i\mathbf{Q} \cdot \mathbf{R}_n}$$

$$\equiv \sum_n f_n^0(\mathbf{Q}) e^{i\mathbf{Q} \cdot \mathbf{R}_n}$$

X-ray crystallography

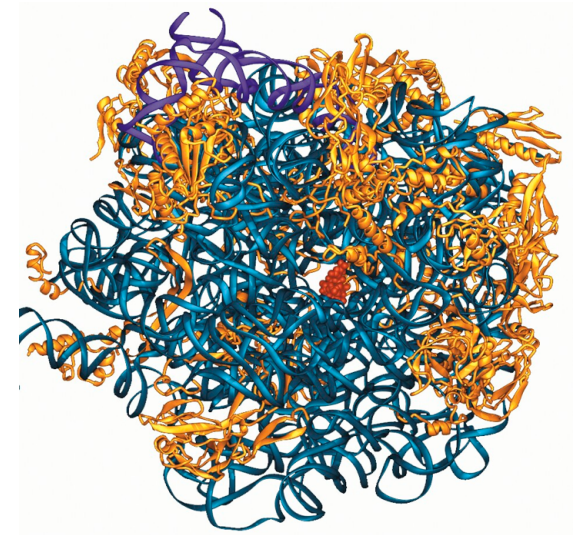


Protein crystal



X-ray diffraction pattern

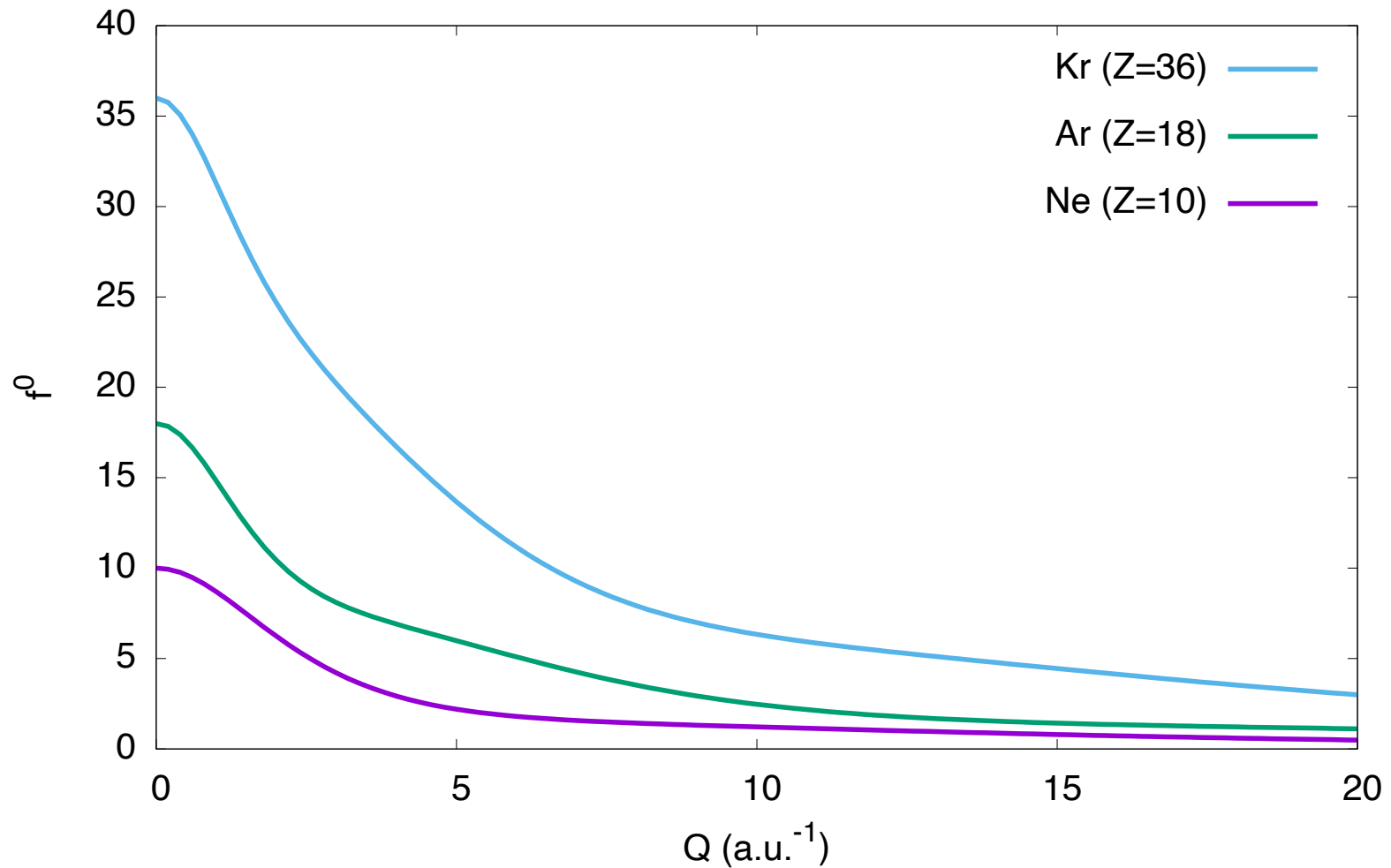
$$I \sim |f^0(\mathbf{Q})|^2$$



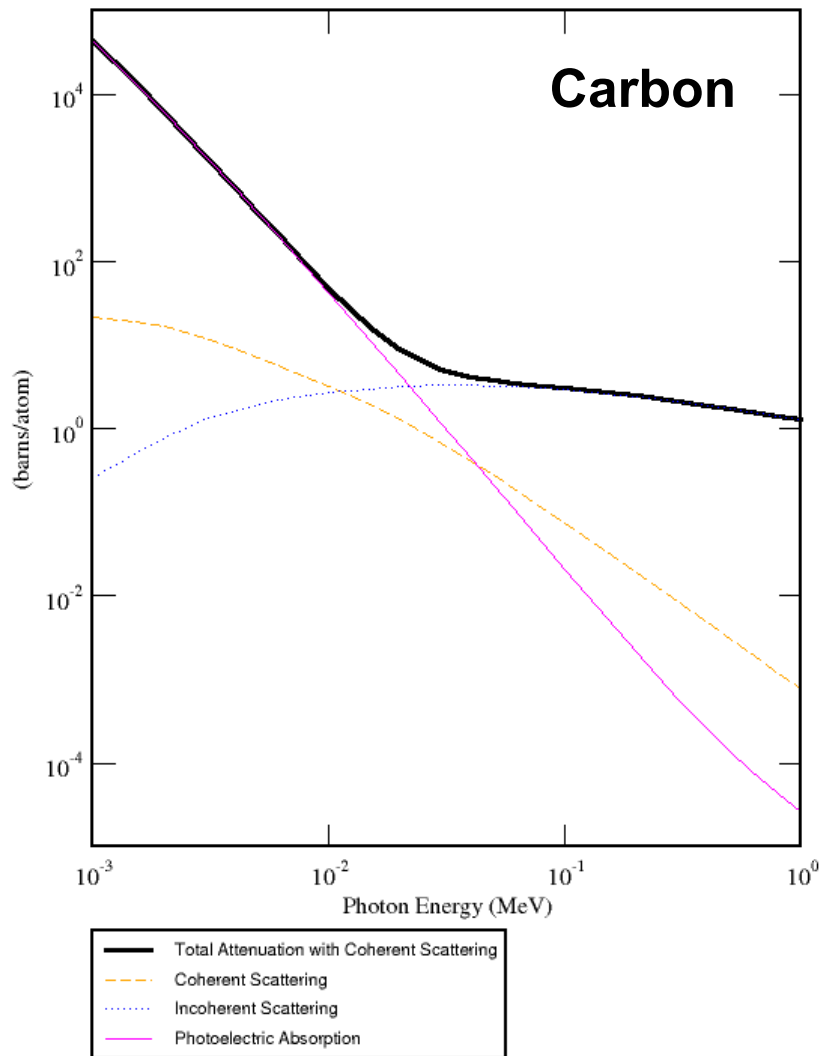
Molecular structure
of a protein

$$f^0(\mathbf{Q}) \underset{\text{FT}}{\overset{\text{IFT}}{\rightleftharpoons}} \rho^{(\text{el})}(\mathbf{x})$$

Atomic form factor



Absorption vs. scattering



C@10 keV

$\sigma_{\text{absorption}}=41$ barn

$\sigma_{\text{el. scattering}}=3.2$ barn

$\sigma_{\text{inel. scattering}}=2.7$ barn

$$\frac{\sigma_{\text{el. sc}}}{\sigma_{\text{abs}}} \sim \frac{1}{13}$$

XCOM, NIST: <https://doi.org/10.18434/T48G6X>

X-ray scattering

elastic vs. inelastic

$$|\Psi_I^{\text{el}}\rangle = |\Psi_F^{\text{el}}\rangle \quad |\Psi_I^{\text{el}}\rangle \neq |\Psi_F^{\text{el}}\rangle$$

without M vs. with M

$$(A^2)$$

$$(p \cdot A)$$

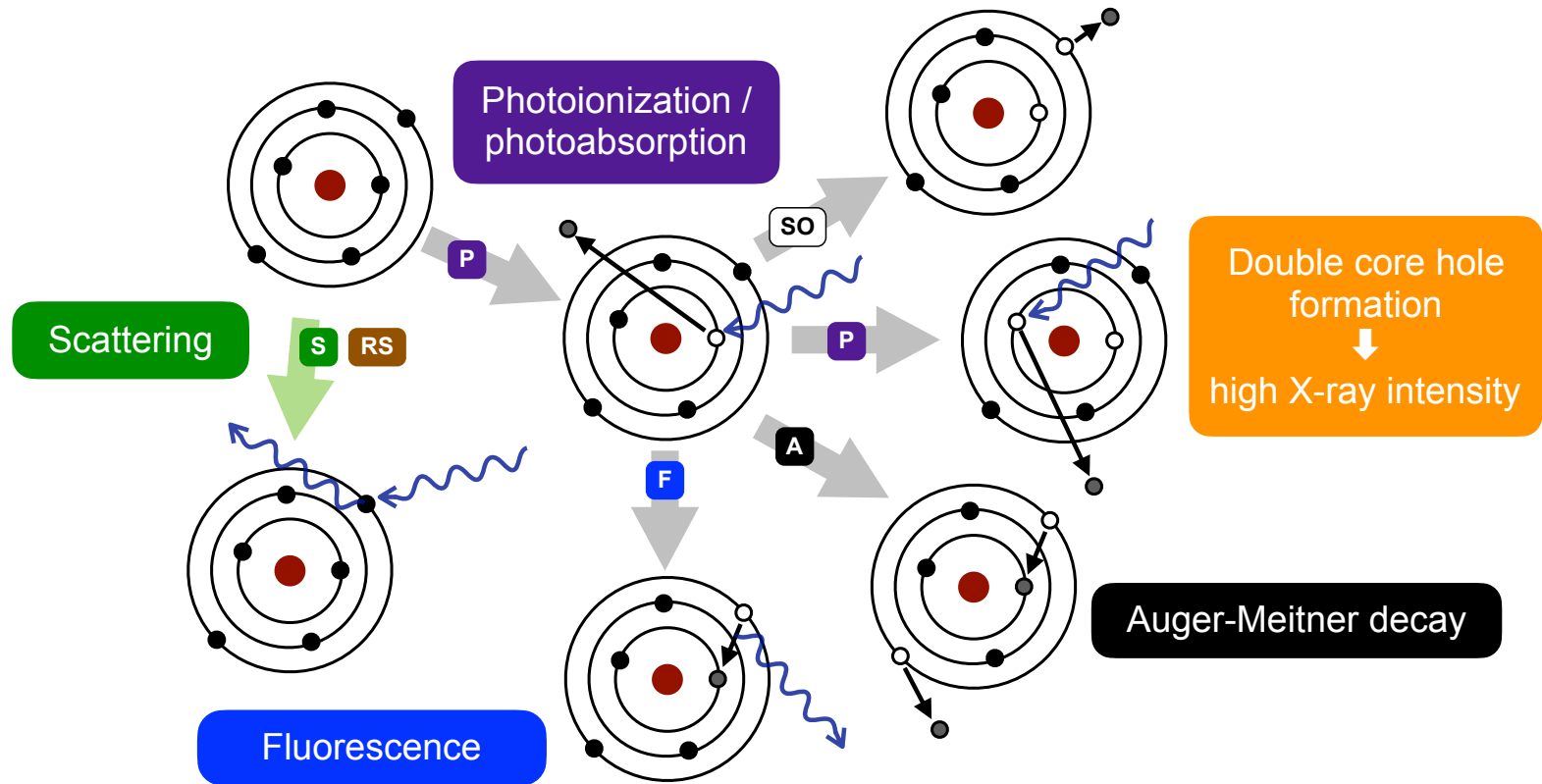
> Elastic X-ray scattering

- Non-resonant elastic X-ray scattering via A^2
- Resonant elastic X-ray scattering via $p \cdot A$ (dispersion correction)

> Inelastic X-ray scattering

- Non-resonant inelastic X-ray scattering via A^2 (Compton scattering)
- Resonant inelastic X-ray scattering via $p \cdot A$ (RIXS or resonant X-ray Raman)

Summary



- > Theoretical aspects of X-ray-induced processes
- > Many-electron system: mean-field approach
- > σ and Γ calculated within the same framework (QED and TDPT)

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