

Magnetic Characterization of PEALD Coated Thin Films for SRF Cavity Research

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Motivation



SRF cavities approach thermodynamic limit of niobium

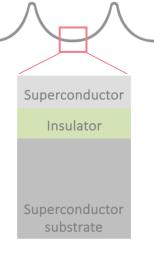
Goals

- To achieve lower
 - \triangleright Coat inside since R_{BCS}
- To enhance acce

Magnetic Characterization Study

⇒ comprehensively investigate the magnetic properties of SIS layers

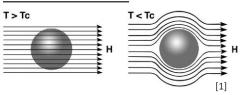
Surpass field limit H_{sh} by coating thin Superconductor-Insulator-Superconductor (SIS) multilayers





Superconductivity in Particle Accelerators

Meissner effect



BCS and GL theory

Microscopic electron-latticeinteraction















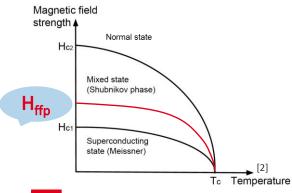






Type I SC

Type II SCs



Niobium as material of choice for SRF cavities

Highest known $T_{\rm C}$ of **9.27 K**

Highest known H_{C1} of **180 mT**

⇒ Upper limit for rf vortex penetration even higher

Superheating field

 $H_{\rm sh} \sim 200 \text{ to } 250 \text{ mT}$

Limiting Performance Factors

Surface Resistance

$$R_{\rm S} = A \frac{\omega^2}{T} \exp\left(-\frac{\Delta}{k_{\rm B}T}\right) + R_{\rm res}$$

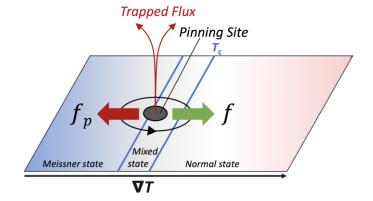
 $R_{\rm BCS}$ vanishes for $T \rightarrow 0 \, {\rm K}$

Everything except R_{BCS}

- Weakly coupled grain boundaries, surface oxides etc.
- Trapped magnetic flux

 Q_0 improves with lower R_S

- ... but is limited by various factors!
- \rightarrow Magnetic Quenching when $H_{applied} > H_{C1}$
- → Dissipation becomes unmanagable!

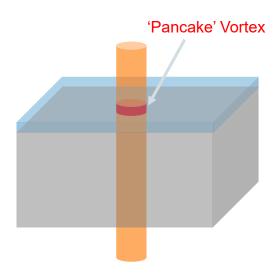


E_{acc} is limited by dissipation because of high surface fields!

Thin Films and Screening Currents

A. Gurevich, APL 88 (2006), T. Kubo, SUST 30 (2017)

- Gurevich: Vortices are energetically supressed in SC thin film
 - $\rightarrow d < \lambda_1$
 - Pancake vortices
- SIS multilayer idea: insulator as vortex barrier
 - → global vortex penetration is prevented
- Kubo's new argument: Vortex penetration shifts to H_{ffp}>H_{C1} through "counter current" at interfaces
 - Layer must be thin enough so that rf field sees counter currents



Vortex dissipation becomes energetically unfavourable



SIS Multilayer Theory

- Top superconductor sees majority of rf field
 - → Surface resistance improves
- Insulating layer creates more interfaces
 - → More counter currents
- Layers must be thinner than $\lambda_{L,Top}$
 - → RF field is affected by counter currents

Theoretical maximum surface field a SIS coated cavity can withstand:

$$H_{\text{max}} = \min\{\tilde{\gamma_1}^{-1}\underline{H_{\text{sh,S}}}, \tilde{\gamma_2}^{-1}\underline{H_{\text{sh,sub}}}\}$$

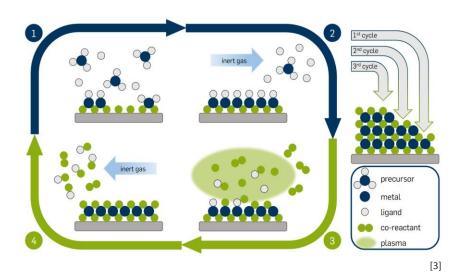
$$\frac{\tilde{\gamma}_{1}}{\tilde{\gamma}_{1}} = \frac{\sinh\frac{d_{S}}{A_{S}} + \frac{\lambda_{sub} + d_{I}}{A_{S}}\cosh\frac{d_{S}}{A_{S}}}{\cosh\frac{d_{S}}{A_{S}} + \frac{\lambda_{sub} + d_{I}}{A_{S}}\sinh\frac{d_{S}}{A_{S}}}$$

$$\frac{\tilde{\gamma}_{2}}{\tilde{\gamma}_{2}} = \frac{1}{1}$$



Sample Preparation

Plasma-Enhanced Atomic Layer Deposition (PEALD)



- Uniform layer-by-layer thin film coatings
- Various advantages
 - Low-temperature deposition
 - ➤ High quality, high conformality
 - Vast choice of precursors
- Self-regulating process
- Convenient for coating the inside of an SRF cavity



Sample Preparation

PEALD Process Parameters

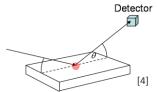
Precursor	Target Material
TMA	AIN
TDMAT	TiN
TBTDEN	NbN

- Plasma mix of H₂ and N₂
- Deposition temperature 250 °C



Characterization Measurements

X-Ray Reflectivity (XRR)



- Determine thin film thickness
- Reflectivity curve shows interference oscillations
 → periodicity related to layer thickness

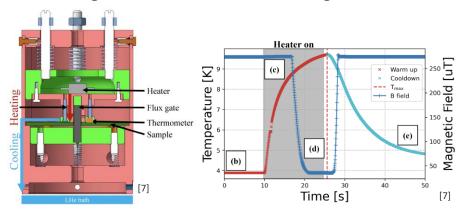
Oscillation width =



Characterization Measurements

Magnetic Flux Expulsion Lens (MFEL)

 SC/NC sample transition with pulsed heating under controlled cooling conditions







PEALD Growth per Cycle for NbN, TiN and NbTiN

1 supercycle NbTiN = 3 x NbN + 1 x TiN

Two-fold results for NbTiN:

- GPC calculated from measured XRR NbTiN thickness
- GPC calculated from individual GPCs of NbN and TiN

Results within the expected range

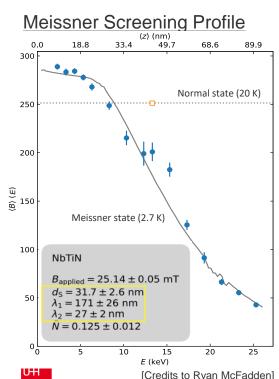
Coating	Cycles/	Estimated	Measured	GPC	GPC
	Supercycles	thickness	thickness (nm)	measured	calculated
	(#)	(nm)		(nm/(sup)cyc.)	(nm/supcyc.)
NbN	1000	~ 60	52.23 ± 2.91	0.052 ± 0.003	-
TiN	200	15.95	15.34 ± 0.85	0.077 ± 0.004	-
NbTiN	243	~ 60	56.74 ± 2.91	0.233 ± 0.012	0.233 ± 0.005



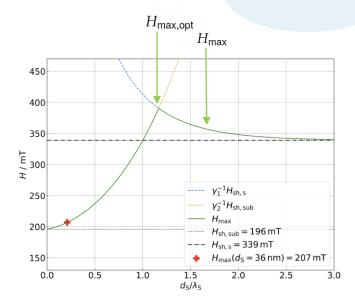
matching NbTiN GPCs



λ_{L} and the Estimation of H_{ffp}



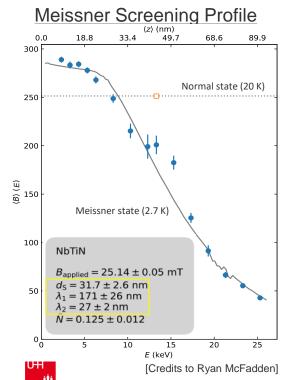
- Nb/NbTiN with $d \sim 36$ nm
- Kubo's counter current model used for fit
- Expected decrease in B with increasing E and z
- H_{max} estimated for Nb/NbTiN SS bi-layer



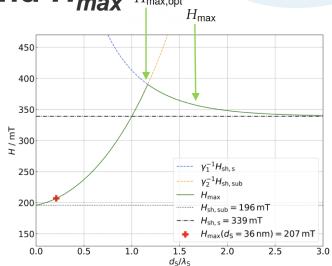
Material	λ (nm)	λ_{L} (nm)	ξ_0 (nm)	κ (nm)	H_{c1} (mT)	H_{c} (mT)	H_{sh} (mT)
NbTiN	171	150 [65]	5 [65]	43	24	452	339
Nb	27	28 [56]	35 [65]	0.82	(no calc.)	200 [65]	196

Clean limit of λ_{L}

Discussion on Meissner Profile and H_{max} $H_{max, opt}$

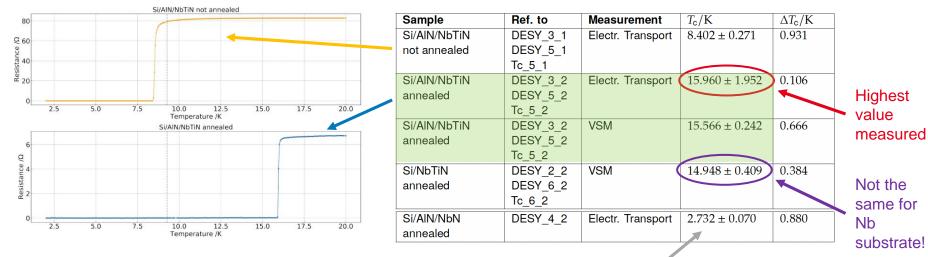


- Kubo's counter current model suitable to determine λ_L
- H_{max} estimate visualizesKubo's prediction
 - $\rightarrow H_{\text{max}} > H_{\text{sh,Nb}}$
 - SIS multilayers needed!



- Estimate based on Meissner Fit → improvable!
- Illustrates principle of field enhancement in an SS bi-layer

PPMS T_C Measurements on Si

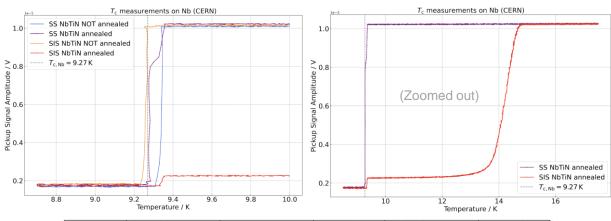


- $T_{\rm C}$ can be clearly assigned to the SC thin film
- Highest T_C for **Si/AIN/NbTiN** in both measurement modes
- Measurement of an SIS Si/AIN/NbN sample \rightarrow very low $T_{\rm C}$



Contactless Inductive $T_{\rm C}$ Measurements on Nb

- As-deposited: T_C not higher than 7 to 8 K
- SS with NbTiN behaves similar as unannealed samples
 - $\succ T_{\rm C} \sim T_{\rm C,Nb}$
 - > Two transitions?
- Only SIS with NbTiN
 - → shows clear stepwise transition
 - \rightarrow reaches high T_c



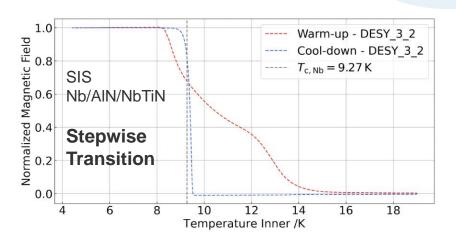
Sample	Structure	$T_{c(1)}/K$	$\Delta T_{c(1)}/K$	$T_{c(2)}/K$	$\Delta T_{c(2)}/K$
Tc_6_1	Nb/NbTiN	9.339 ± 6.890	0.032	-	. :
	not annealed				
Tc_5_1	Nb/AIN/NbTiN	9.259 ± 6.998	0.023	=	-
	not annealed				
Tc_6_2	Nb/NbTiN	9.276 ± 10.663^{8}	0.044	9.340 ± 0.326	0.033
	annealed				
Tc_5_2	Nb/AIN/NbTiN	9.343 ± 1.407	0.0144	14.196 ± 0.014	0.835
	annealed				



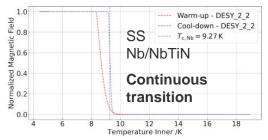


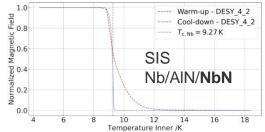
MFEL T_C Measurements

- Annealed Nb SIS and SS samples
- Only SIS with NbTiN
 - → shows stepwise transition
 - \rightarrow reaches high T_c
- No double transition or significant increase in T_c for
 - → SS Nb/NbTiN
 - → SIS Nb/AIN/NbN
- All T_c measurements deliver matching results!













Magnetic Flux Expulsion

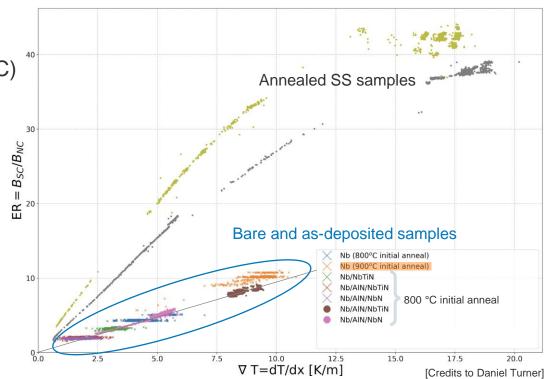
As-deposited SIS and SS thin films

ER(Nb@900 °C) > ER(Nb@800 °C)

Limited spatial thermal gradient

Annealed SS thin films

- Much higher flux expulsion
- Greater expulsion for 900 °C initial anneal

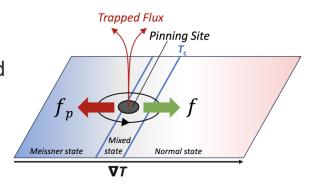






Discussion on $T_{\rm C}$ and Flux Expulsion

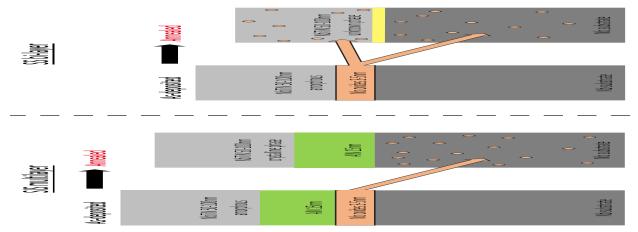
- Better flux expulsion for 900 °C initial anneal
 - Increase in grain size with higher T_{anneal}
 - Reduction in pinning sites
- Post-deposition annealing enhances flux expulsion
 - Even for SS Nb/NbTiN even if no high T_C is achieved
 - Assumption: Pancake effect adding additional force that counteracts pinning force





Discussion on $T_{\rm C}$ and Flux Expulsion

- High-T annealing required to activate the thin film and achieve high T_C
- Insulating layer on Nb required to ensure increase in T_C
 - Barrier layer prevents oxygen from diffusing into NbTiN layer
 - \triangleright XRD confirms NbTiN δ -phase formation for SS sample





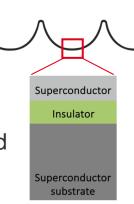
Conclusion

- Unclear if reduction of T_C implies reduction of H_C
 - Experimental determination required
 - Useful comparison with estimation of H_{max}
- Findings on oxygen diffusion allow new considerations on μSR data analysis
 - Intermediate layer?
 - Electron Microscopy probably shows oxygen-enriched phase with varying stoichiometry
- 900 °C annealing of SIS Nb/AIN/NbN
 - \triangleright Does **not** form intended high- $T_{\rm C} \delta$ -phase
 - \triangleright Ti as stabiliser of the cubic high- $T_{\rm C} \delta$ -phase
 - NbN excluded from SIS studies



Summary

- Further improvement of Nb SRF cavity performance neccessary...
- ... possible with PEALD coated SIS multilayer!
 - Coating thin high-T_C superconducting and insulating layers
 - Pushing the field of first flux penetration H_{ffp}
- More characterization and experimental testing of the SIS theory required



SIS is a promising approach towards new technologies and improved future applications in SRF research!



Thank you! Gracias! Danke!

Isabel, Ricardo, Rezvan, Jonas, Chris, Getnet, Marco, Artem, Chriag, Leon, Michel, Anton, Thomas, Cem, Marc, Rakshith, Cornelius, Wolfgang, Robert, Thorsten, Detlef, Dimitri, Alexey, Dan, Alick, Ryan, Andreas



References

- [1] Irwin Yousept. "Optimal Control of Non-Smooth Hyperbolic Evolution Maxwell Equations in Type-II Superconductivity". In: SIAM Journal on Control and Optimization 55 (Jan. 2017), pp. 2305–2332. COI: 10.1137/16M1074229
- [2] Irwin Yousept. "Optimal Control of Non-Smooth Hyperbolic Evolution Maxwell Equations in Type-II Superconductivity". In: SIAM Journal on Control and Optimization 55 (Jan. 2017), pp. 2305–2332. COI: 10.1137/16M1074229
- [3] Lukas Mai. "Investigation of amino-alkyl coordinated complexes as new precursor class for atomic layer deposition of aluminum, tin and zinc oxide thin films and their application". doctoralthesis. Ruhr-Universität Bochum, Universitätsbibliothek, 2020. DOI: 10.13154/294-7658
- [4] https://de.wikipedia.org/wiki/Röntgenreflektometrie
- [5] https://de.wikipedia.org/wiki/Larmorpräzession
- [6] https://de.wikipedia.org/wiki/Myon
- [7] Daniel Turner et al. "Flux Expulsion Lens: Concept and Measurements". In: JACoW SRF2023 (2023), MOPMB003. DOI: 10.18429/JACoW-SRF2023-MOPMB003



Backup

PEALD System at Universität Hamburg

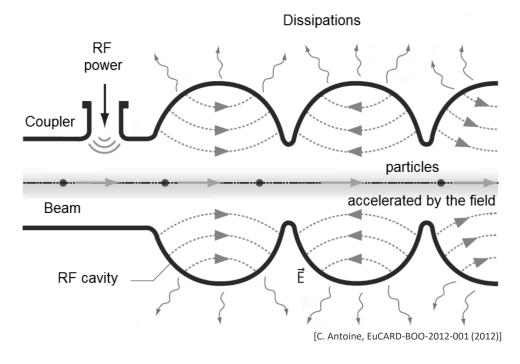


UHV Furnace at Universität Hamburg

XERION.



Cavity Resonator



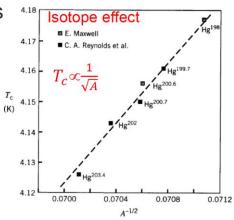


Theory of Superconductivity

$\Delta(0) = 1.76k_{\rm B} \cdot T_{\rm c}$

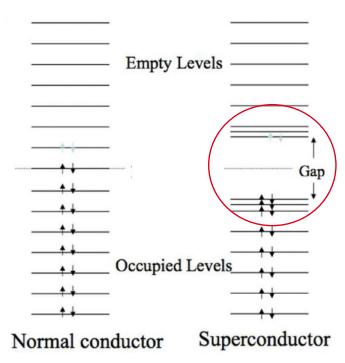
Isotope effect

First clue that SC is linked to atomic lattice, not just electrons



SC Gap

- Sc electrons
 (bound state) only
 conduction
 electrons that
 interact with lattice
 at finite phonon
 frequency
- Energy gap aroundE_F

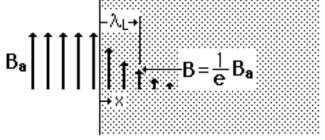




Theory of Superconductivity

London Penetration depth

$$\lambda_{L}(T) = \frac{\lambda_{L}(T = 0 \text{ K})}{\sqrt{1 - \left(\frac{T}{T_{c}}\right)^{4}}}$$



$$B_{inside} = B_a e^{-x/\lambda_L}$$

The London penetration depth is the distance required to fall to 1/e times the externally applied field $B_{\mathbf{a}}$.



Surface Resistance

$$R_{\rm S} = A \frac{\omega^2}{T} \exp\left(-\frac{\Delta}{k_{\rm B}T}\right) + R_{\rm res}$$

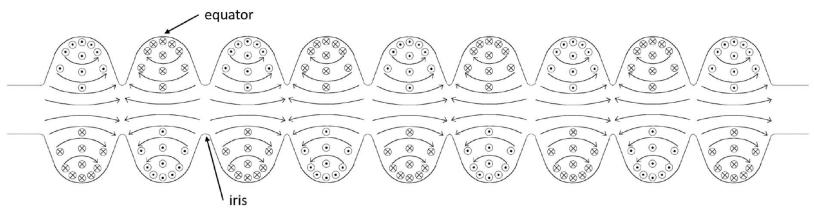
 $R_{\rm BCS}$ vanishes for $T \rightarrow 0 \text{ K}$

- R_{BCS} decreases exponentially with respect to the temperature (higher T more electron will participate in the losses)
- Proportional to the square of the frequency (higher frequency means less effective shielding)
- Non-vanishing AC resistance
- Cooper pairs have inertia
 - Can not follow the RF field instantly
 - Do not shield RF field perfectly
- 'Normal' electrons are accelerated and dissipate power



9-cell TESLA-type niobium cavity

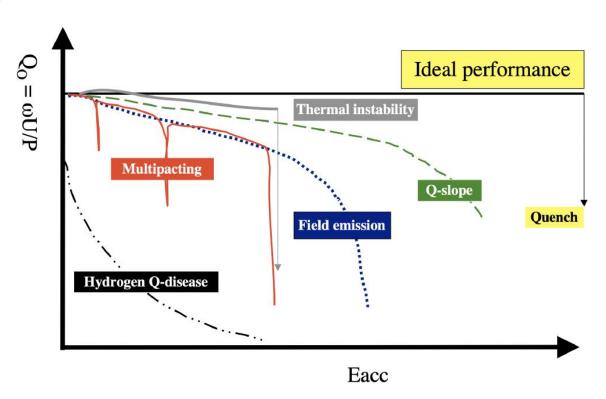
$$E_{\text{acc}} = \frac{1}{L} \int_{-L/2}^{L/2} E(z) \cos\left(\frac{\omega z}{v}\right) dz$$





Quality Factor Q_0

$$Q_0 = \frac{\omega_0 U}{P_{\text{diss}}} \equiv \frac{G}{R_{\text{S}}}$$





Thin Films and Pancake Vortices (Gurevich)

- Nucleation of parallel magnetic vortices is energetically suppressed in SC thin films
- Yield of H_{C1} : $H_{\text{C1}} = \frac{2\phi_0}{\pi d^2} \ln \frac{d}{\tilde{\xi}}$
- A coating of N superconducting thin layers then screens the interface field H_i on the Nb surface down to

$$H_{\rm i} = H_0 \exp\left(\frac{-Nd}{\lambda}\right)$$

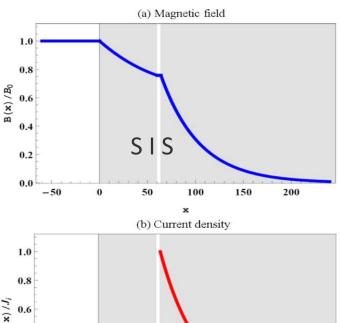
Kubo's Theory

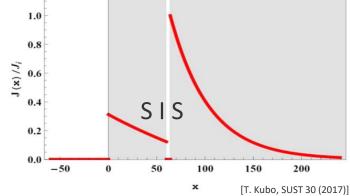
Theoretical maximum surface field a SIS coated cavity can withstand:

$$H_{\text{max}} = \min\{\tilde{\gamma_1}^{-1}\underline{H_{\text{sh,S}}}, \tilde{\gamma_2}^{-1}\underline{H_{\text{sh,sub}}}\}$$

$$\frac{\tilde{\gamma_1}}{\cos h} = \frac{\sinh \frac{d_S}{A_S} + \frac{\lambda_{sub} + d_I}{A_S} \cosh \frac{d_S}{A_S}}{\cosh \frac{d_S}{A_S} + \frac{\lambda_{sub} + d_I}{A_S} \sinh \frac{d_S}{A_S}}$$

$$\frac{\tilde{\gamma_2}}{\cosh\frac{d_S}{\lambda_S} + \frac{\lambda_{sub} + d_I}{\lambda_S}} \sinh\frac{d_S}{\lambda_S}$$







Prediction of maximum surface field H_{max}

$$\kappa = \frac{\lambda}{\xi_{\rm GL}} = \frac{2\sqrt{3}}{\pi} \frac{\lambda^2}{\xi_0 \lambda_{\rm L}}$$

$$H_{\rm sh} \approx 0.75 \, H_{\rm c} \quad \text{for} \quad \kappa \gg 1$$

$$H_{\rm c} = \frac{\sqrt{2}\kappa H_{\rm c1}}{\ln \kappa}$$

$$H_{\rm c1} = \frac{\Phi_0}{4\pi \lambda^2} (\ln(\kappa) + 0.497)$$

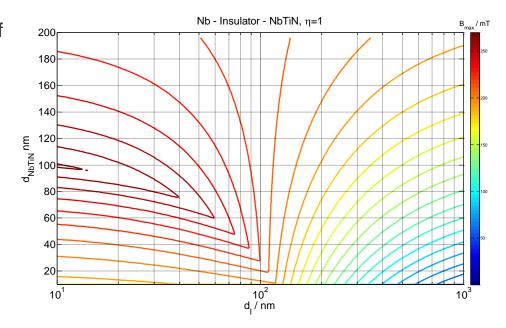
Material	λ (nm)	λ_{L} (nm)	ξ_0 (nm)	κ (nm)	H_{c1} (mT)	H_{c} (mT)	$H_{sh} \; (mT)$
NbTiN	171	150 [65]	5 [65]	43	24	452	339
Nb	27	28 [56]	35 [65]	0.82	(no calc.)	200 [65]	196

Clean limit of λ_L



Prediction of maximum surface field H_{max}

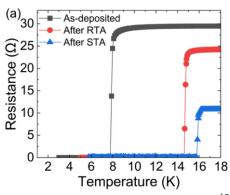
- Field enhancement depends on H_{C1}, λ_L of both sc layer and substrate and η
- Assume Nb Insulator NbTiN
 - $H_{C1.Nb}$ 180 mT
 - $-H_{C1, NbTiN}$ 213 mT
 - $-\eta = 1$
- ► Isolator 15 nm, NbTiN 100 nm
 - $H_{C1} = 270 \text{ mT oder } 63 \text{ MV/m}$
 - $-\eta = 0.7$
- ► Isolator 10 nm, NbTiN 60 nm
 - $H_{C1} = 224 \text{ mT oder } 52 \text{ MV/m}$
 - $-\eta = 0.7$





Post-Deposition Annealing

\rightarrow Improvement of $T_{\rm C}$ and $H_{\rm C1}$



$$\mu_0 H_{C1} = 15 \text{ mT}$$

 $\mu_0 H_{C1} = 81 \text{ mT}$
 $\mu_0 H_{C1} = 98 \text{ mT}$

$$\mu_0 H_{\text{C1,bulk}} \approx 33 \text{ mT}$$

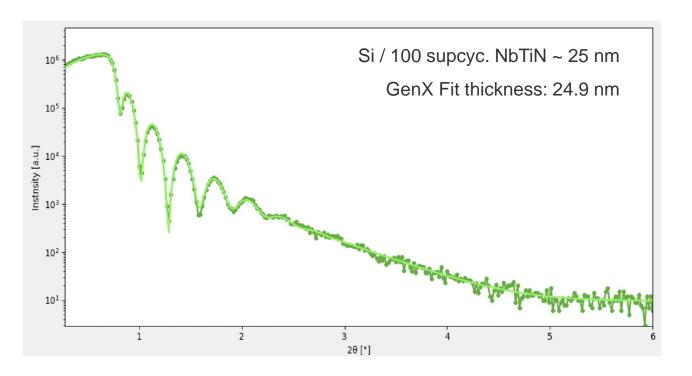
[Gonzalez.I. et al., J. Appl. Phys. 134, 159902 (2023)] [Gonzalez.I. et al., J. Appl. Phys. 134, 035301 (2023)]

900 °C annealing for 1 hour

- → heating rate 3 °C/min
- → controlled cooling rate ~ 1.45 °C/min for 8 hours to 200 °C



XRR



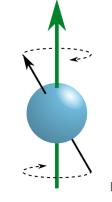


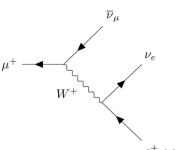
How to measure the Meissner Profile with λ_L

- Positively charged, spin-polarized muons are implanted into the material sample to be examined
- passage of a muon through the muon detector determines the time zero of the implantation



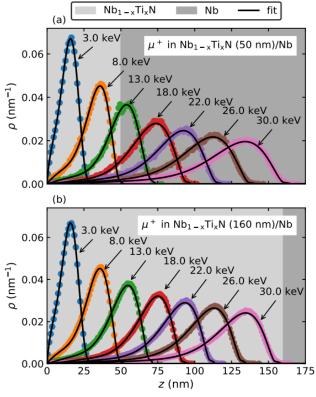
- Interaction of the muon spin with the magnetic moments of the host lattice leads to a characteristic temporal evolution of the spin polarization → different precession frequencies because of field distribution
- The original phase relationship of the individual spins is increasingly lost
 → rotation signal therefore exhibits a temporal attenuation







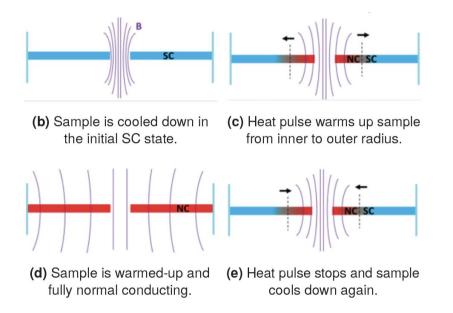
Stopping Profiles

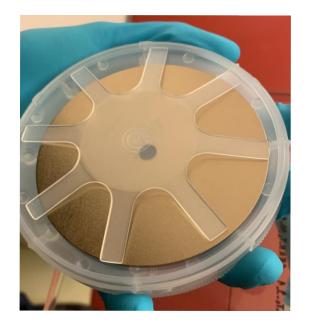






Flux Expulsion Measurement









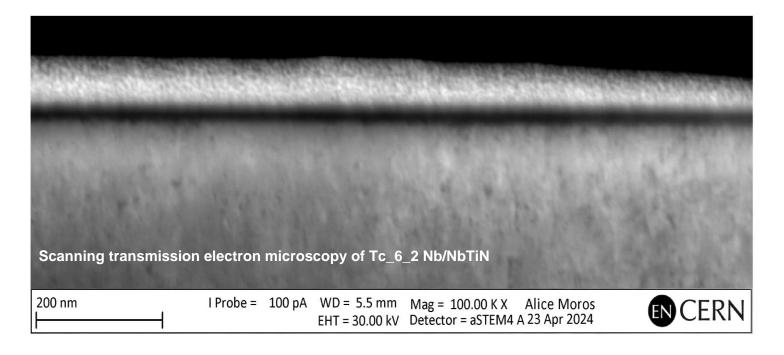
Flux Expulsion Measurement

Sample	Initial annealing T (°C)	Coated structure	I (~ 15 nm)	S (~ 60 nm)
DESY_0	800	-	-	-
DESY_1	900	-	-	-
DESY_2_1	800	SS	-	NbTiN
DESY_2_2				
DESY_3_1	800	SIS	AIN	NbTiN
DESY_3_2				
DESY_4_1	800	SIS	AIN	NbN
DESY_4_2				
DESY_5_1	900	SIS	AIN	NbTiN
DESY_5_2				
DESY_6_1	900	SS	-	NbTiN
DESY_6_2				



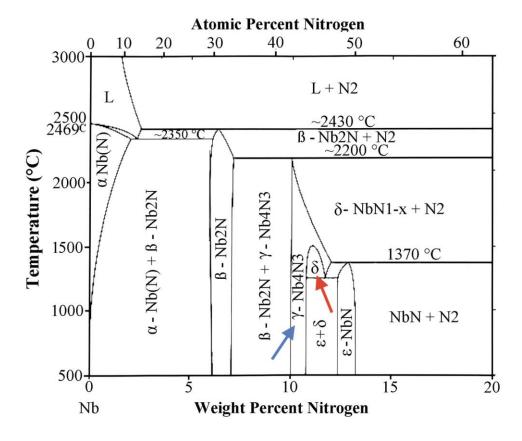


STEM Measurements at CERN



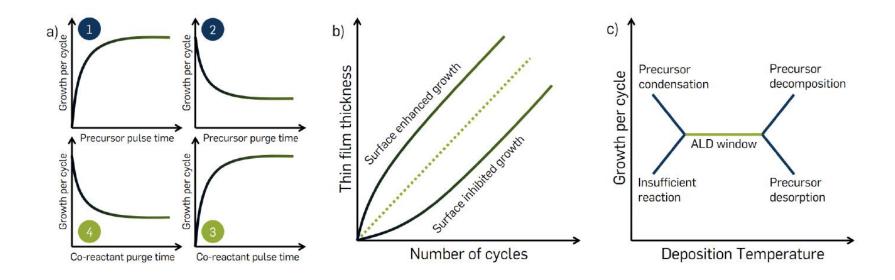


NbN Phase Diagram





PEALD





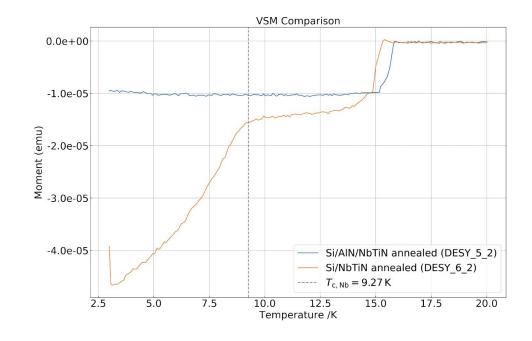
Deposition specifications

Precursor	TMA trimethyl-	TDMAT tetrakis	TBTDEN (t-butylimido)
	aluminum	(dimethylamino)	tris(diethylamino)
		titanium(IV)	niobium(V)
Target Material	AIN	TiN	NbN
Precursor Pulse Duration (s)	0.021	0.5	0.5
Precursor Temperature (°C)	room temperature	70	90
1st Purge Duration (s)	30	40	40
Plasma Exposure Duration (s)	10	60	90
Plasma Flow of H ₂ , N ₂ (SCCM)	30,10	45,20	45,12
2nd Purge Duration (s)	100	120	120



VSM T_c Measurements Results

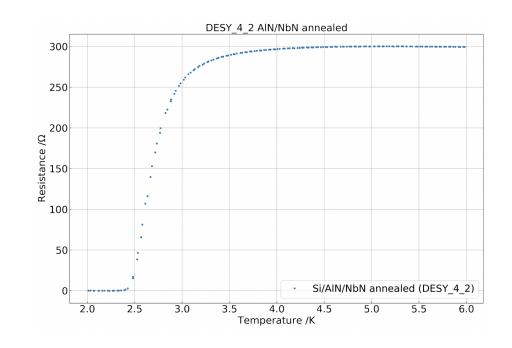
text





Electrical-transport T_c Measurement Result for NbN

- 900 °C annealing of SIS Nb/AIN/NbN
 - ightharpoonup intended high- $T_{\rm C}$ δ-phase is **not** formed
 - Ti as stabiliser of the cubic high- $T_{\rm C}$ δ -phase
 - NbN excluded from SIS studies



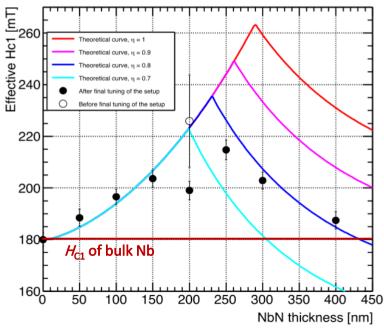


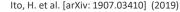
*H*_{C1} measurement at KEK

- Determination of $H_{C1}(d)$ for SIS mandatory
- Third-harmonic voltage measurement with local magnetometry at KEK in Japan
 - SIS coated Nb samples with varying layer thicknesses

Compare results with theory!

Measured H_{C1} of Nb/SiO₂/NbN samples vs. NbN thickness:







SC