

Impact of Non-Pointing Photon Searches at ATLAS for Dark Matter Models



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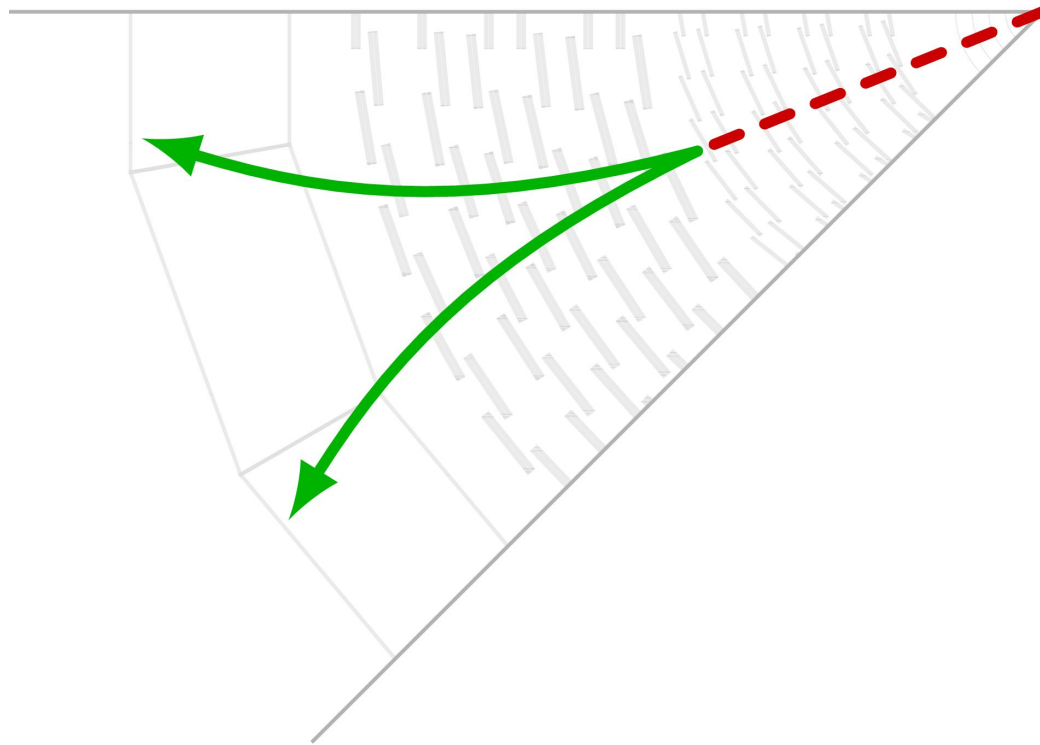
Based on the following work:
P. Arias, B. Díaz, L. Duarte, J. Jones-Perez, WR, D. Zegarra (2507.15930)

Motivation

- Dark matter remains unexplained → suggests hidden sectors.

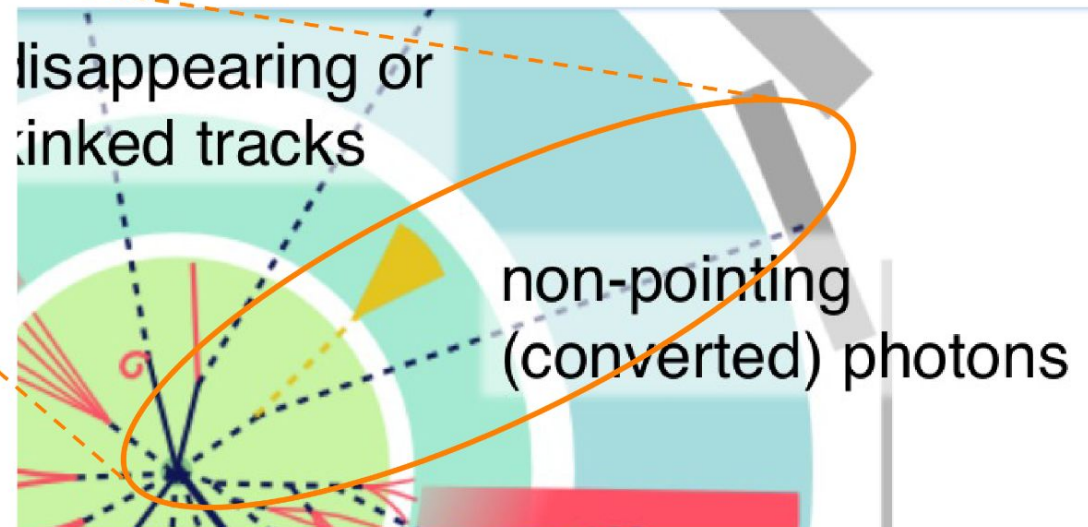
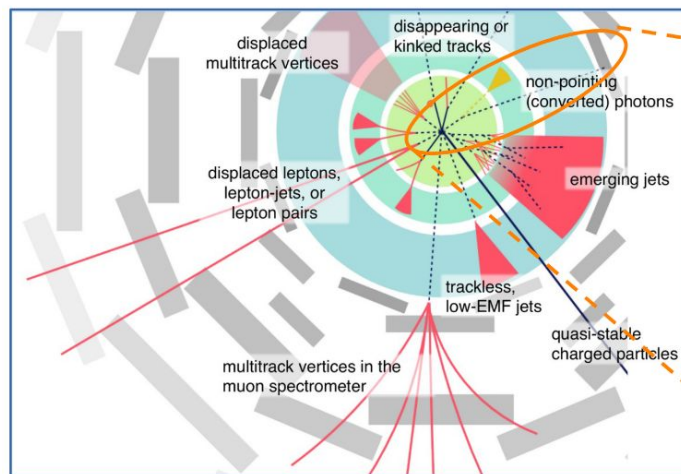
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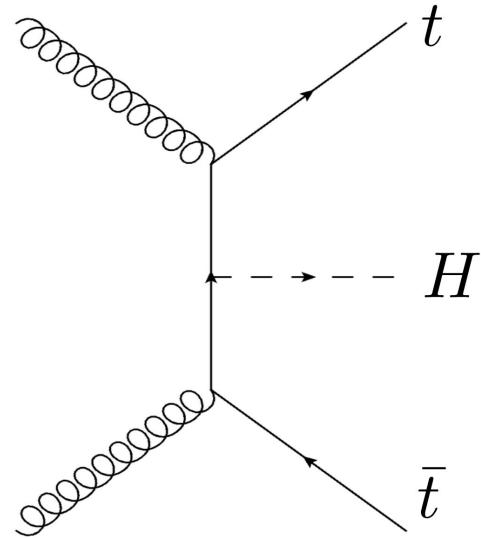
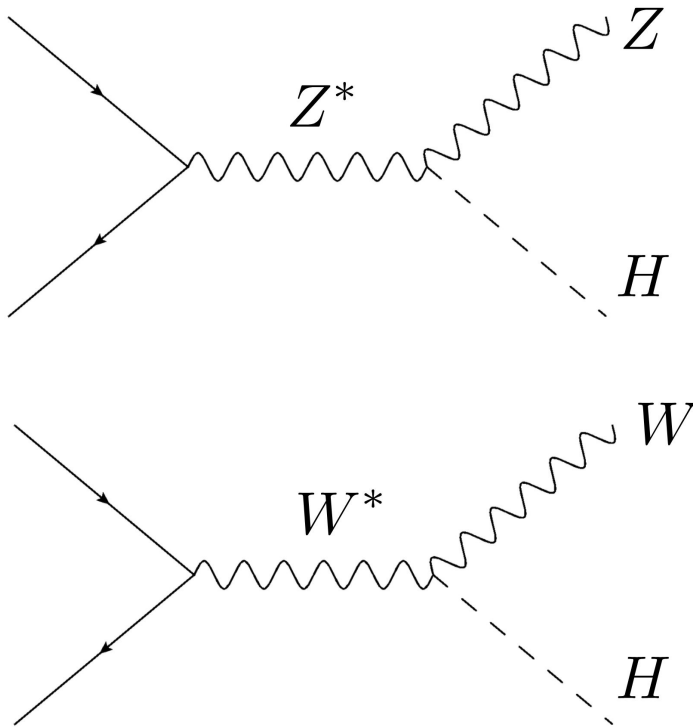
Design by H. Russel

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- Many models predict long-lived particles (LLPs), including in Higgs decays.
- LLPs can give displaced photons, a clean and striking LHC signature.
- This work recasts ATLAS displaced-photon searches to probe these scenarios using refined collider searches.

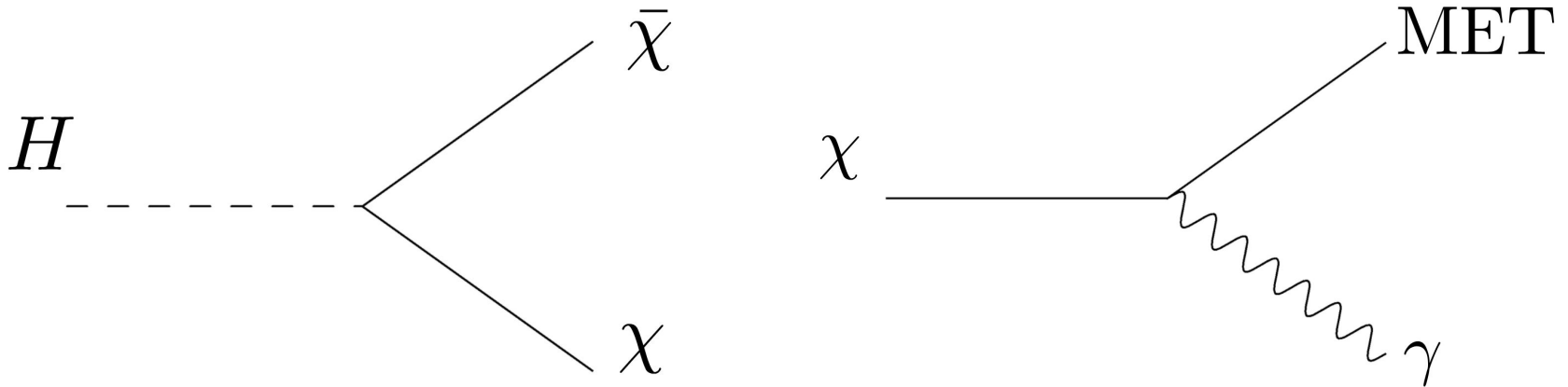
Recasting

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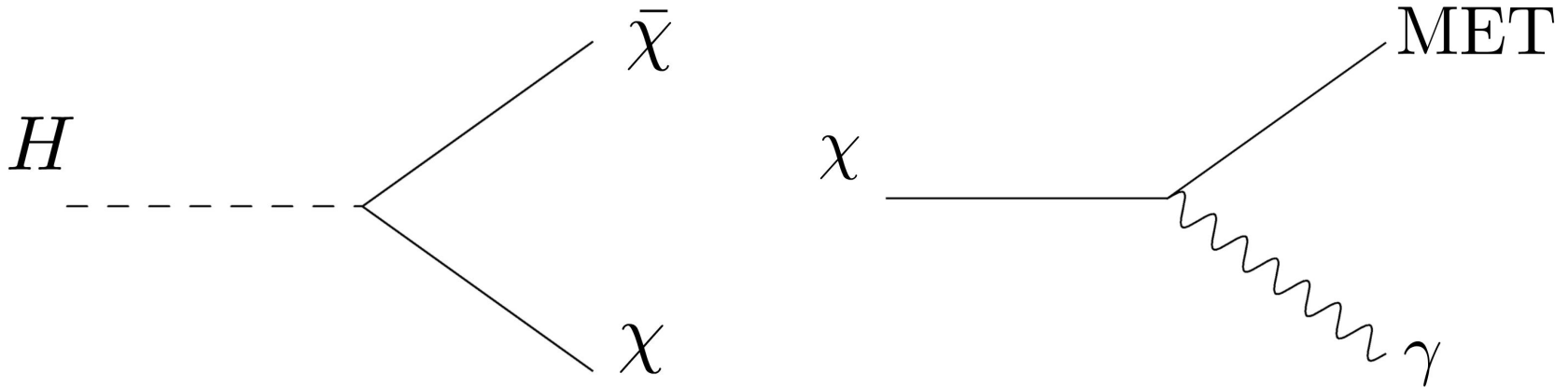
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- Place cuts on t_γ and $|\Delta z_\gamma|$. Distinguish single and multi-photon samples.

$$(1) \quad 1.5 \text{ ns} < t_\gamma < 12 \text{ ns} \qquad 1 \text{ ns} < t_\gamma < 12 \text{ ns} \quad (2+)$$

$$|\Delta z_\gamma| > 300 \text{ mm}$$

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- Place cuts on t_γ and $|\Delta z_\gamma|$. Distinguish single and multi-photon samples.
- Results obtained for SUSY:

	1	2+	1+
Expected	3.8 ± 1.6	0.28 ± 0.04	4.1 ± 1.7
Observed	4	0	4

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- The model is endowed with \mathbb{Z}_2 symmetry:
 - Dark fields A'_μ and ϕ are odd under \mathbb{Z}_2
 - Standard Model (SM) fields are even under \mathbb{Z}_2

$$\mathcal{L} \supset -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{1}{2}m_{\gamma'}^2 A'^2_\mu + \frac{1}{2}(\partial_\mu\phi)^2 \\ - \frac{1}{2}\tilde{m}_\phi^2 \phi^2 - \lambda_\phi \phi^4 - \lambda_{HS} \phi^2 |H|^2$$

Y. Farzan, A. R. Akbarieh (1207.4272 [hep-ph])

P. Arias, A. Arza, J. Jaeckel, D. Vargas-Arancibia (2007.12585 [hep-ph])

B. Diaz (2405.06113 [hep-ph])

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 - Dark fields A'_μ and ϕ are odd under \mathbb{Z}_2
 - Standard Model (SM) fields are even under \mathbb{Z}_2
- A 5-dimension EFT operator involving A'_μ and preserving \mathbb{Z}_2 and $U'(1)$ symmetries is used to include the decoupled sector.

$$\mathcal{L}_5 = \frac{g_D}{2} \phi F'_{\mu\nu} \tilde{B}^{\mu\nu}$$

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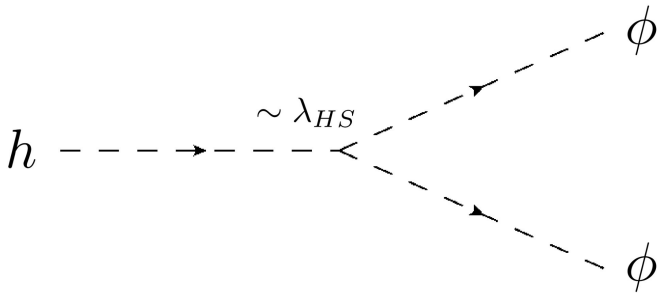
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- We want to reinterpret the collider search results to apply them to our new theoretical scenario.

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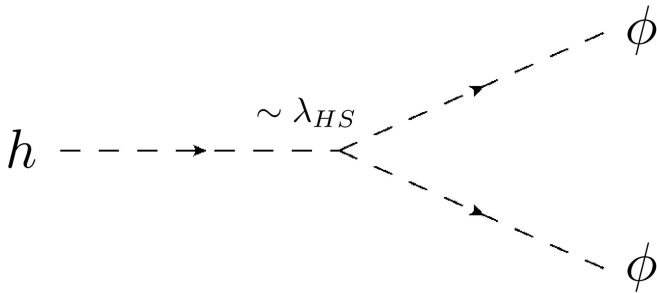
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$$\lambda_{HS} \phi^2 |H|^2$$


The diagram illustrates the decay of a Higgs boson h into two particles ϕ . The decay is mediated by a loop with a coupling proportional to λ_{HS} . The diagram shows a dashed line for the Higgs boson h entering from the left, which then splits into two dashed lines for the particles ϕ exiting to the right. The coupling is indicated by a dashed line with an arrow pointing to the vertex, labeled $\sim \lambda_{HS}$.

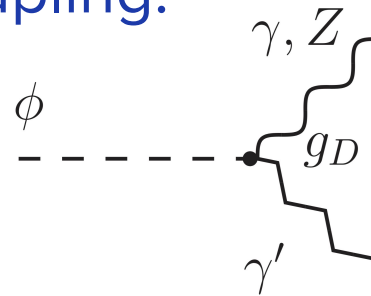
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- Meanwhile, the displaced photon and the MET will be generated via the effective coupling.

$$\mathcal{L}_5 = \frac{g_D}{2} \phi F'_{\mu\nu} \tilde{B}^{\mu\nu}$$



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$$\Gamma_\phi(\phi \rightarrow \gamma\gamma') = \frac{g_D^2 c_W^2}{32\pi} m_\phi^3 \left(1 - \frac{m_{\gamma'}^2}{m_\phi^2}\right)^3$$

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- We require a sufficient number of scalars produced from the Higgs.

$$\Gamma(h \rightarrow \phi\phi) = \frac{\lambda_{HS}^2 v^2}{8\pi m_h} \sqrt{1 - 4 \frac{m_\phi^2}{m_h^2}}$$

this calls for $\mathcal{O}(10^{-5}) \lesssim \lambda_{HS} \lesssim \mathcal{O}(1)$



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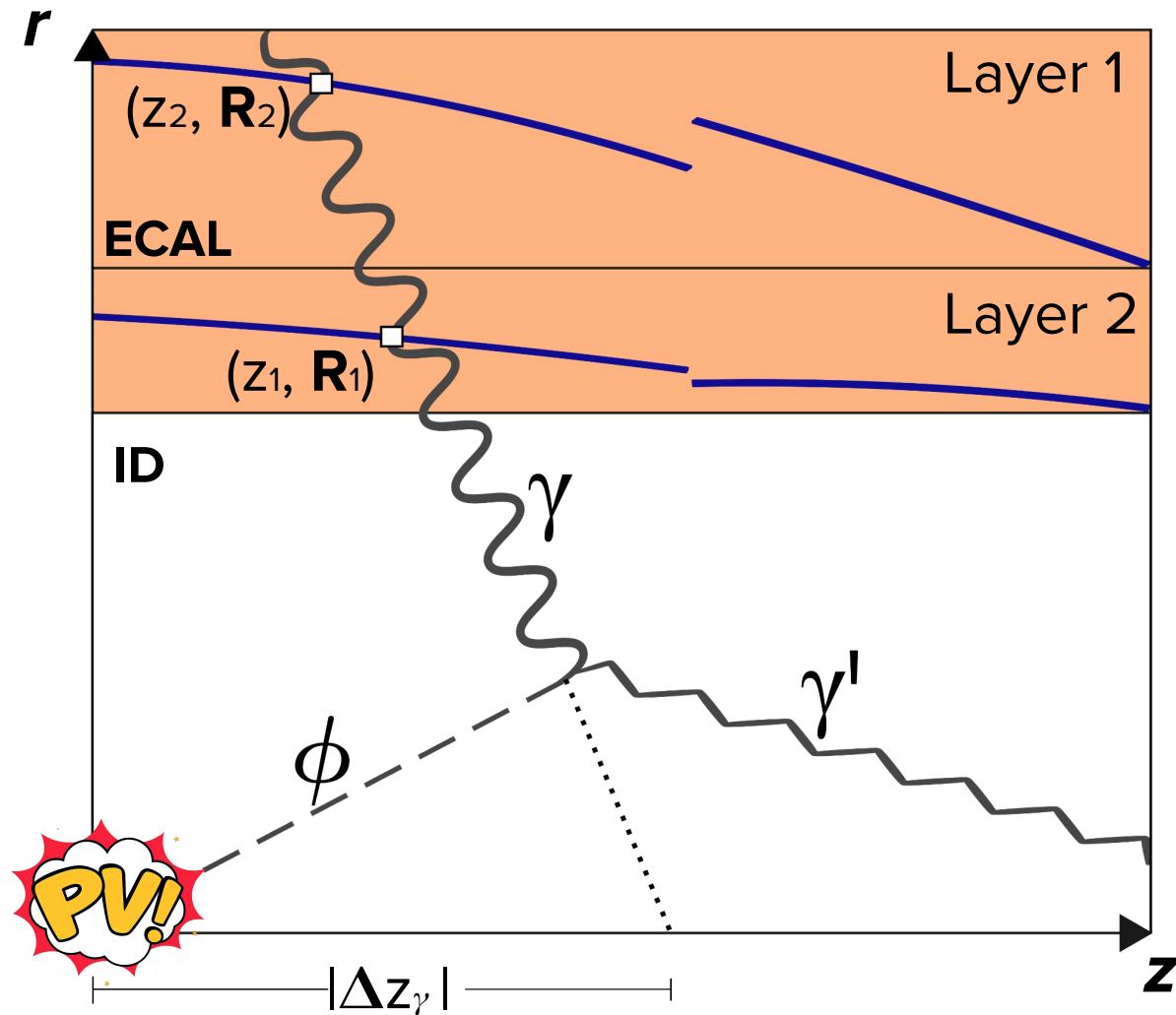
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- This dark matter scenario also reproduces the correct DM relic abundance and satisfies cosmological constraints.

Non Pointing Variable: ATLAS



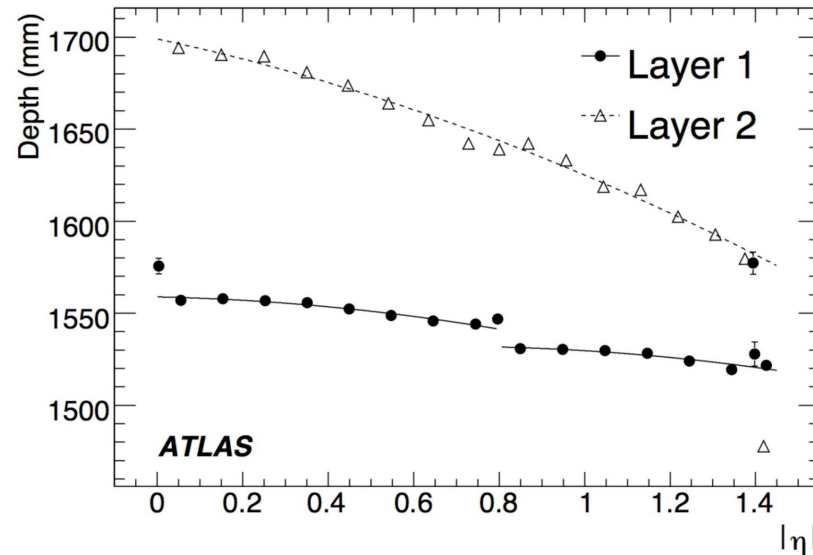
Non Pointing Variables: Δz_γ

- With this, at first instance, one would use the following expression to find the non pointing parameter:

$$\Delta z_\gamma = \frac{z_1 R_2 - z_2 R_1}{R_2 - R_1}$$

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$$\Delta z_\gamma = r_{\phi z} - \frac{p_{\gamma z}}{p_{\gamma T}^2} (r_{\phi x} p_{\gamma x} + r_{\phi y} p_{\gamma y}) \\ + \frac{p_{\gamma z}}{p_{\gamma T}} \left(\frac{R_1 R_2}{R_2 - R_1} \right) \left\{ \left(1 - \frac{d_0^2}{R_1^2} \right)^{1/2} - \left(1 - \frac{d_0^2}{R_2^2} \right)^{1/2} \right\}$$

with d_0 : photon
impact parameter

$$d_0 = (r_{\phi x} p_{\gamma y} - r_{\phi y} p_{\gamma x}) / p_{\gamma T}$$



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- Note that for these calculations, the truth level of Pythia is used.

Simulation Refinement

- ATLAS Delphes Card was modified to adapt them to match the performance more accurately.

Higgs Decay to Invisible & Undetected

- Invisible decays: Long-lived scalars ϕ that decay outside detector contribute to the Higgs invisible BR, limited by ATLAS **$B_{inv} < 0.107$**

$$B_{inv} = \text{BR}(h \rightarrow \phi\phi) \exp \left[-\frac{2L_{\text{det}}}{\gamma_{\text{rel}} \beta_{\text{rel}} c \tau_{\phi}} \right] < 0.107$$

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- Assuming a **spherical detector** of $L_{det} = 1.97 \text{ m}$ radius, the probability of ϕ decaying beyond the detector is known

$$P(L_{det}, \infty) = \exp \left[- \frac{L_{det}}{\gamma_{rel} \beta_{rel} c \tau_{\phi}} \right]$$



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- Undetected signals: Global fits to Higgs data (ATLAS/CMS) constrain additional exotic decays not covered by direct searches, yielding **$B_{und} < 0.12$**

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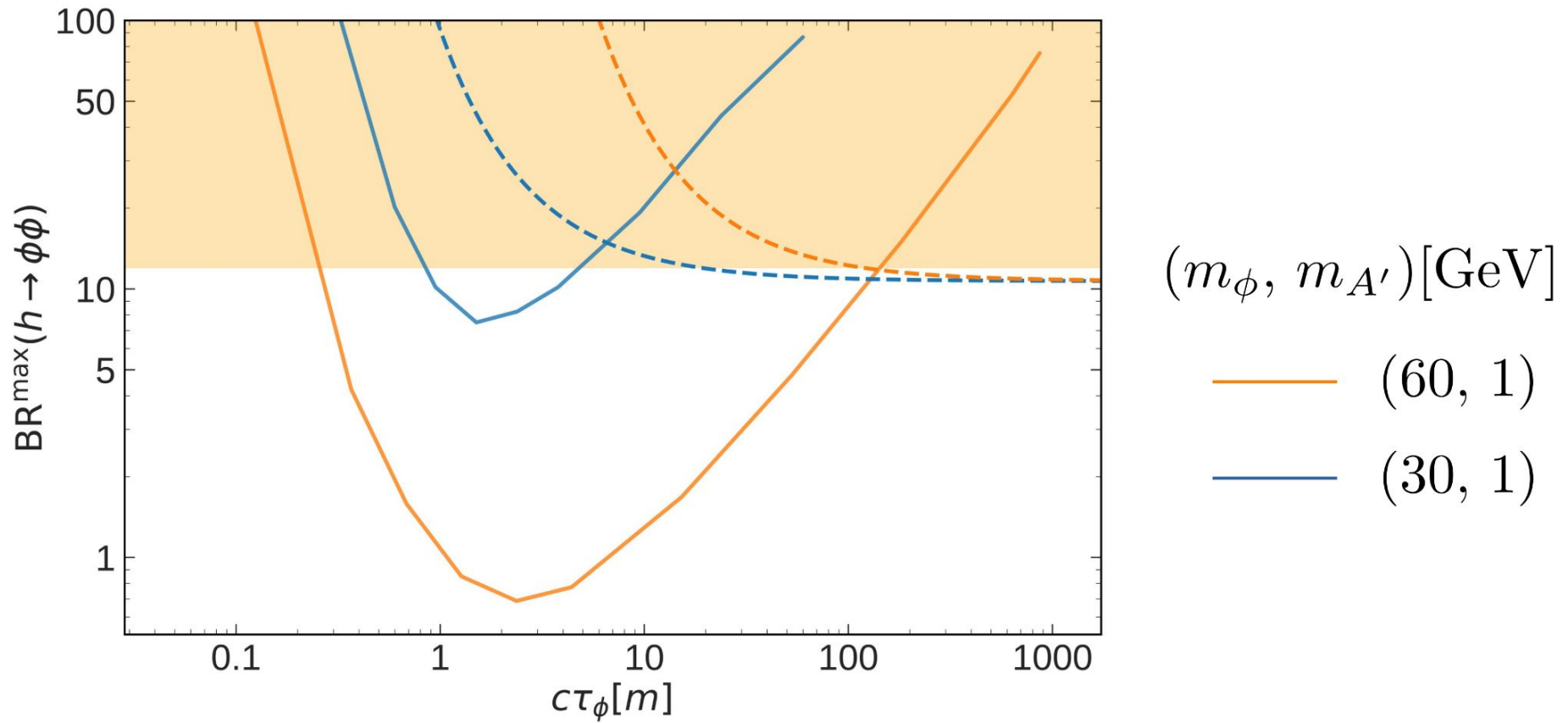
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- No excess observed → bounds on $\text{BR}(h \rightarrow \phi\phi)$ and LLP parameter space.

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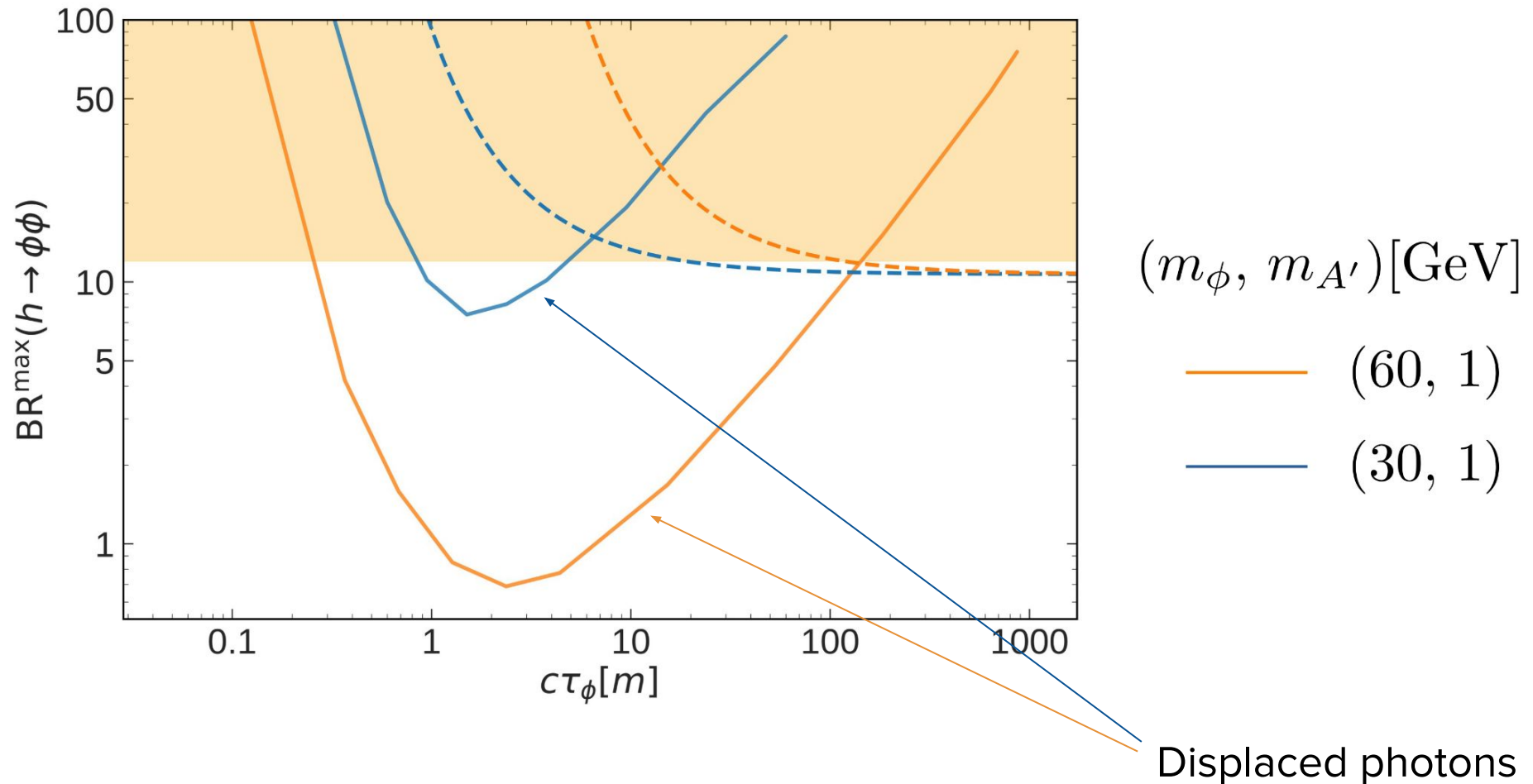
- Efficiency (in %) for $m_\phi = 60$ GeV, $m_{\gamma'} = 1$ GeV. We combine the three production processes. Results shown for $gD = 5.7 \times 10^{-10}$ and 1.6×10^{-10} GeV⁻¹, which respectively imply decay lengths $c\tau_\phi = 0.69$ and 4.4 m.

Cut	$c\tau_\phi = 0.69$ m		$c\tau_\phi = 4.4$ m	
Trigger, $p_{\gamma T} > 10$ GeV, Acceptance	71		44	
Isolation, Efficiencies, Z-veto	39		18	
	Channel			
	1	2+	1	2+
$E_{\text{cell}} > 10$ GeV	15	20	11	3.2
MET > 50 GeV	6.2	8.3	5.4	1.3
$ \Delta z_\gamma > 300$ mm	2.2	1.8	2.6	0.46
$t_\gamma > 1$ (1.5) ns	0.20	0.45	0.89	0.27

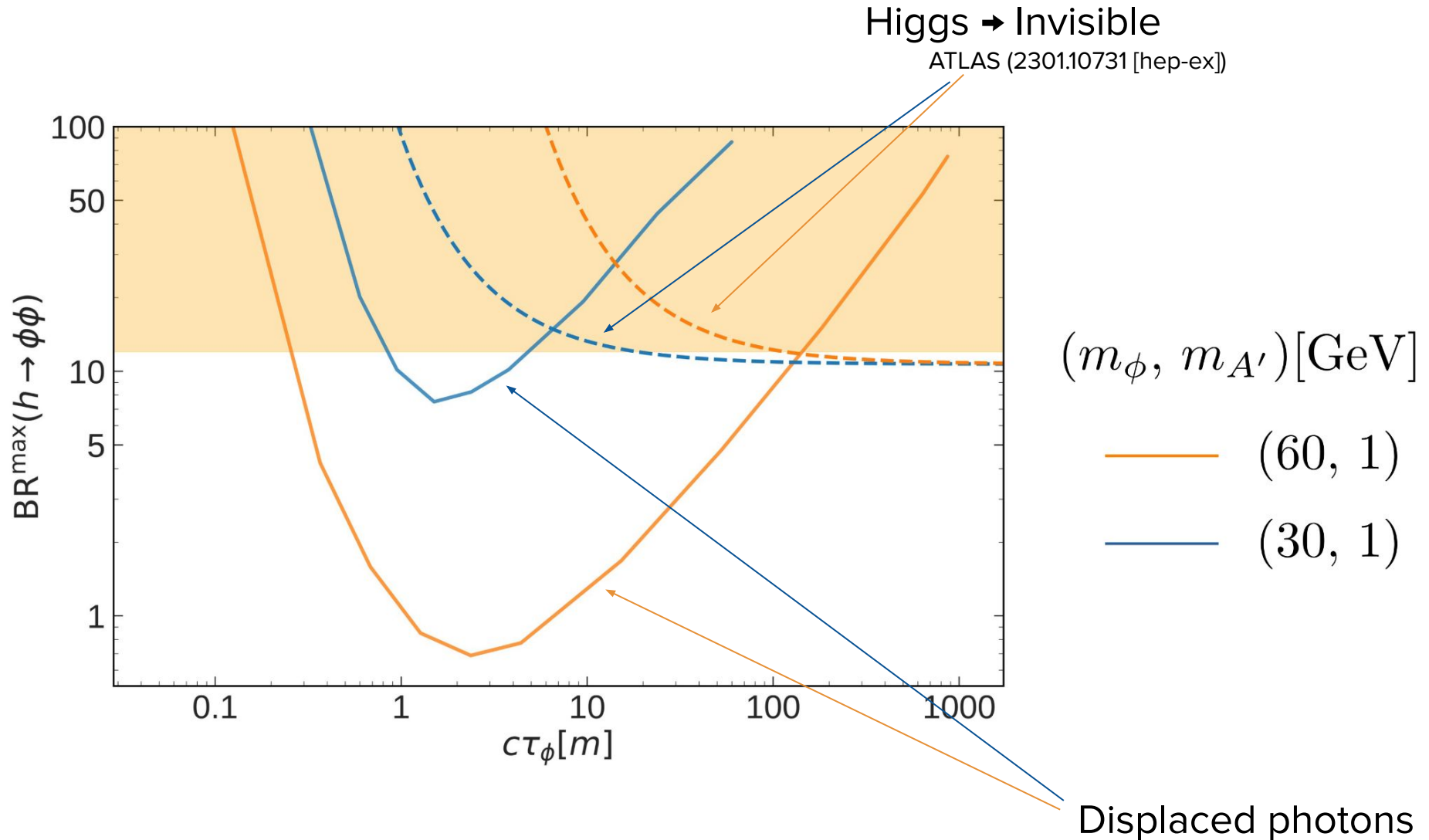
Results



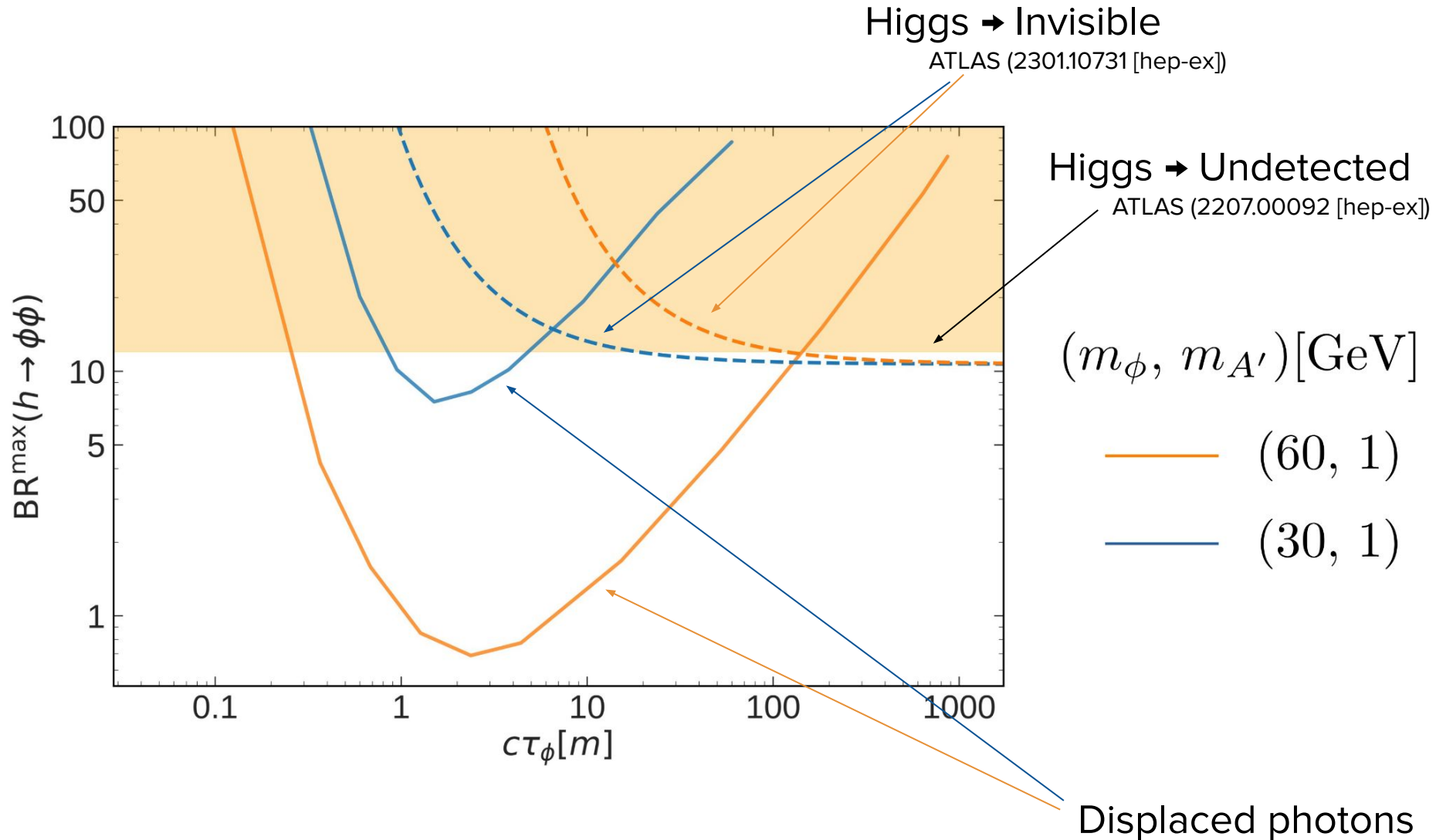
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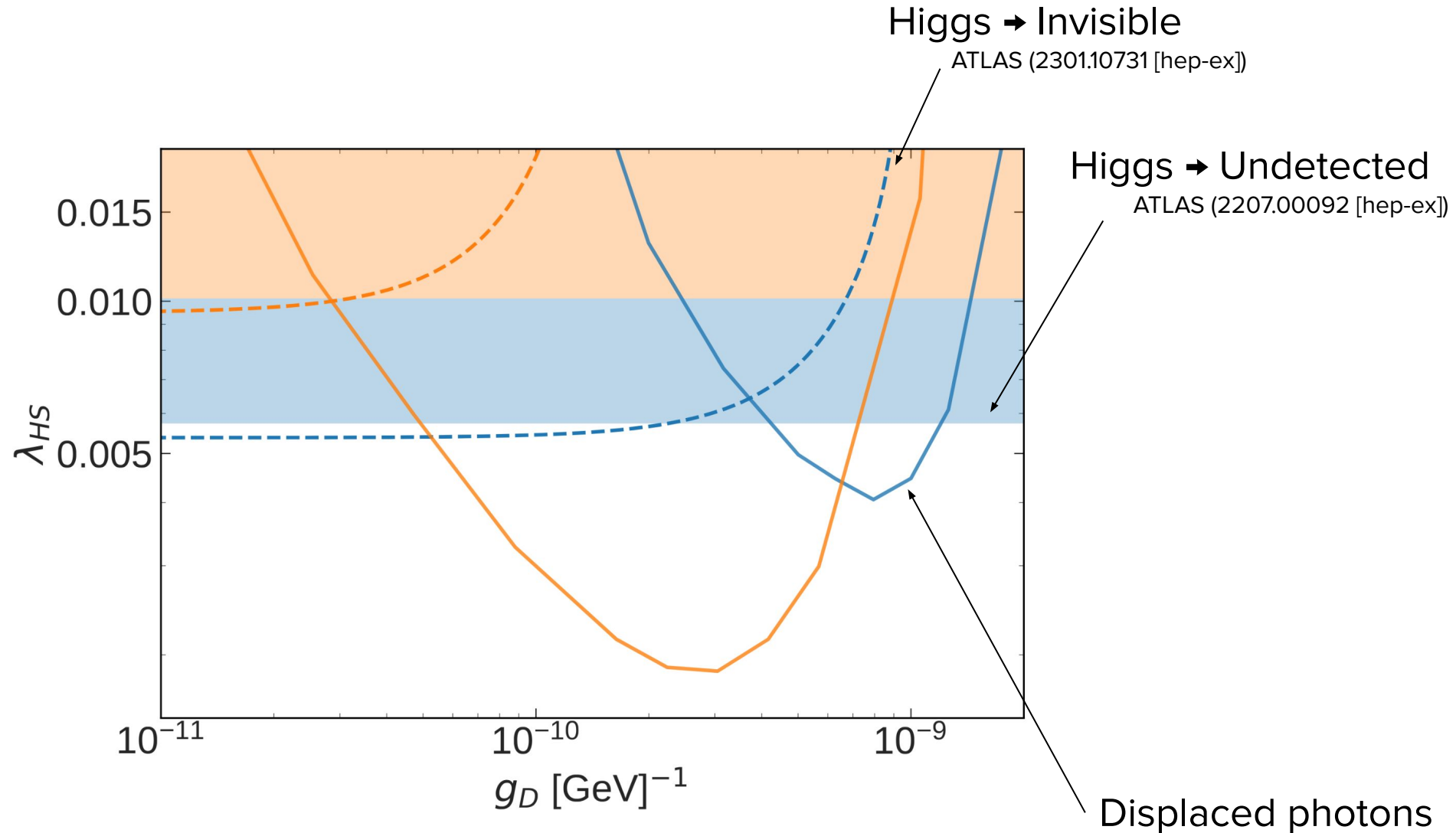
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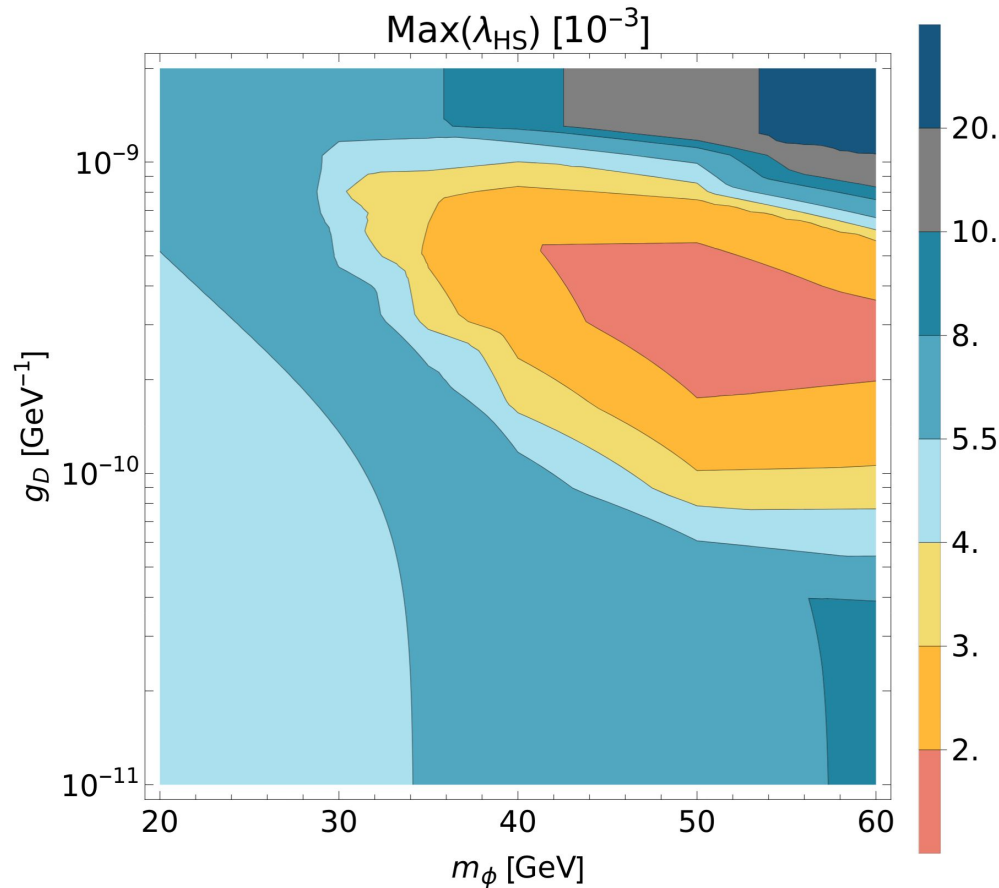
Results



Results



Results: Combined bounds



- Large coupling implies ϕ decays promptly.
- Heavier masses for ϕ are favoured: ϕ moves slowly so t_γ is larger.
- For large g_D , constraints from fit are strongest. For small g_D , constraints from invisible decays are strongest.
- Small couplings implies ϕ decays outside detector.

- $m_{\gamma'} = 1$ GeV.

Conclusions

- Searches for displaced photons can place bounds on models with LLPs decaying into photons and MET.
- These searches have largest sensitivity for $c\tau \sim \mathcal{O}(1 \text{ m})$
- For lower (higher) lifetimes, Higgs \rightarrow undetected (invisible) searches present better bounds.
- Recast of search in the context of dark photon with scalar portal places bounds $\lambda_{HS} \sim \mathcal{O}(10^{-3})$

Gracias

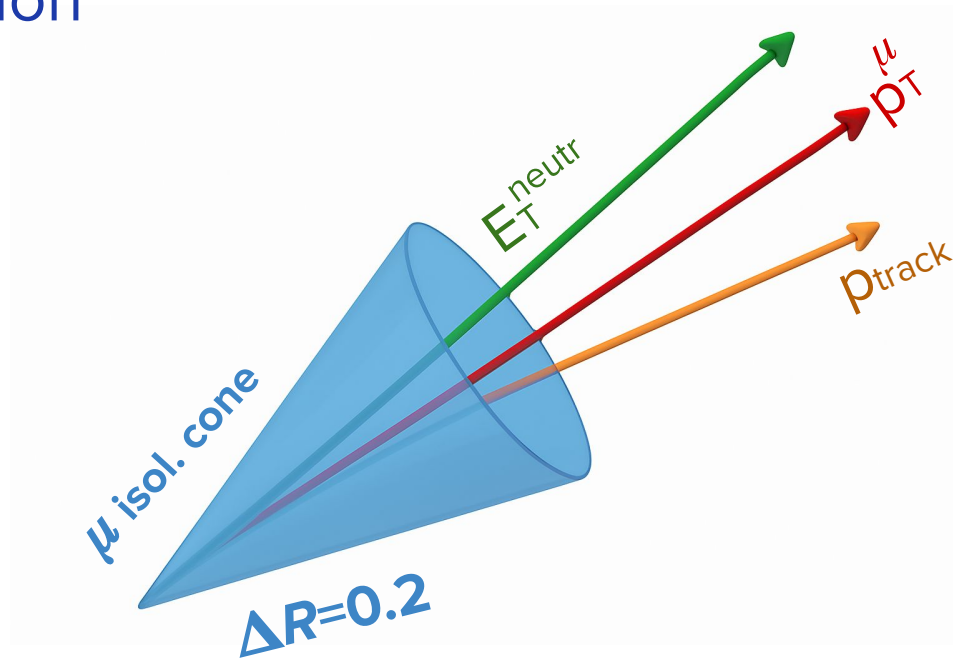
Backup Slides



PUCP

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- Muon Isolation

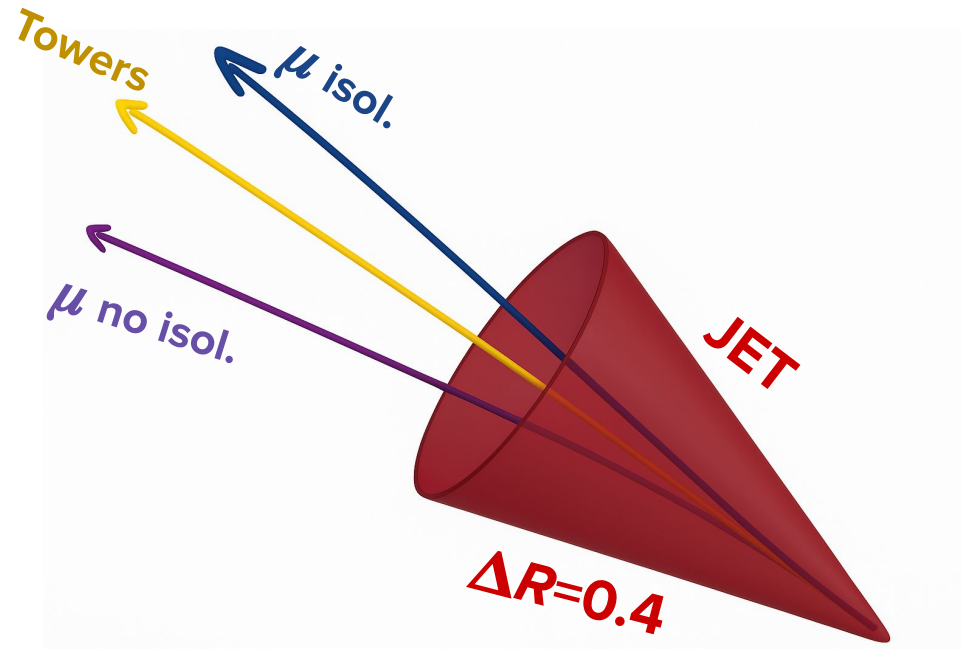


$$\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$$

$$(p_{\text{trk}} + 0.4E_T^{\text{neut}}) < 0.16 p_T^\mu$$

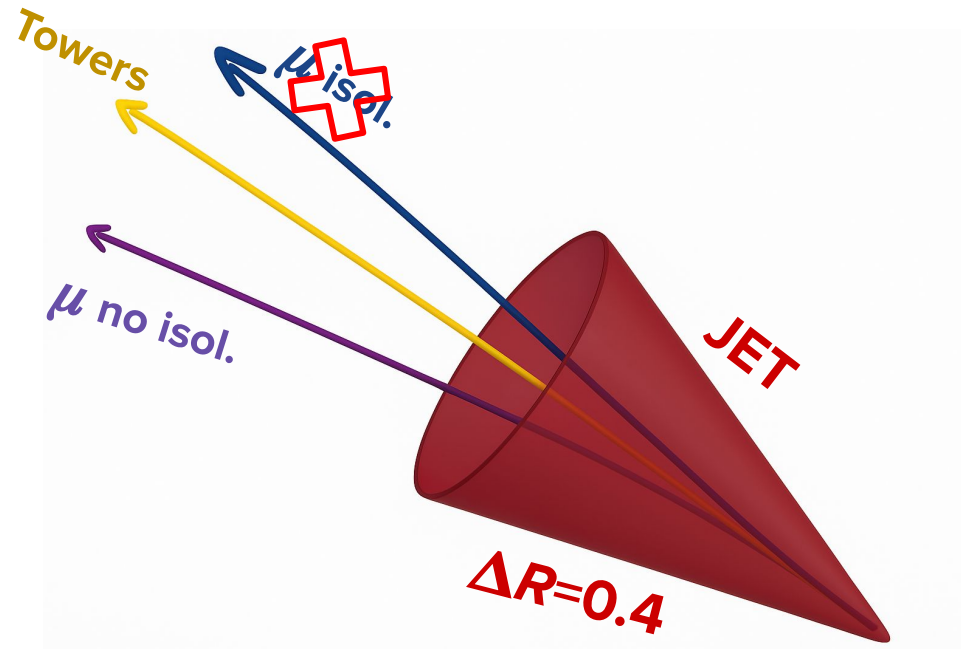
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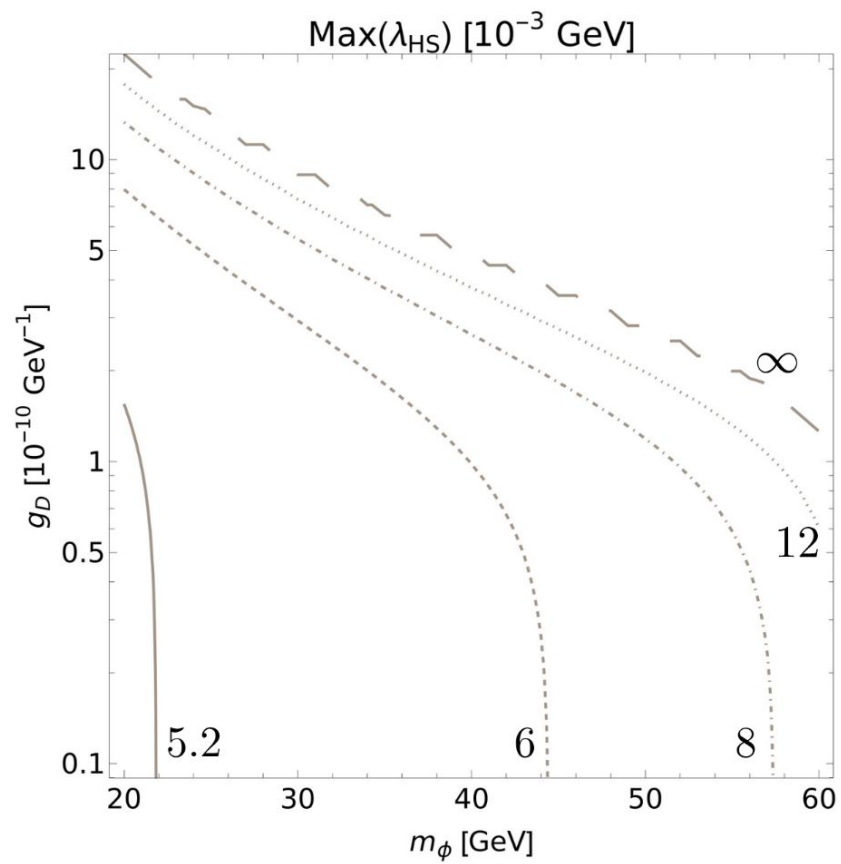


Delphes

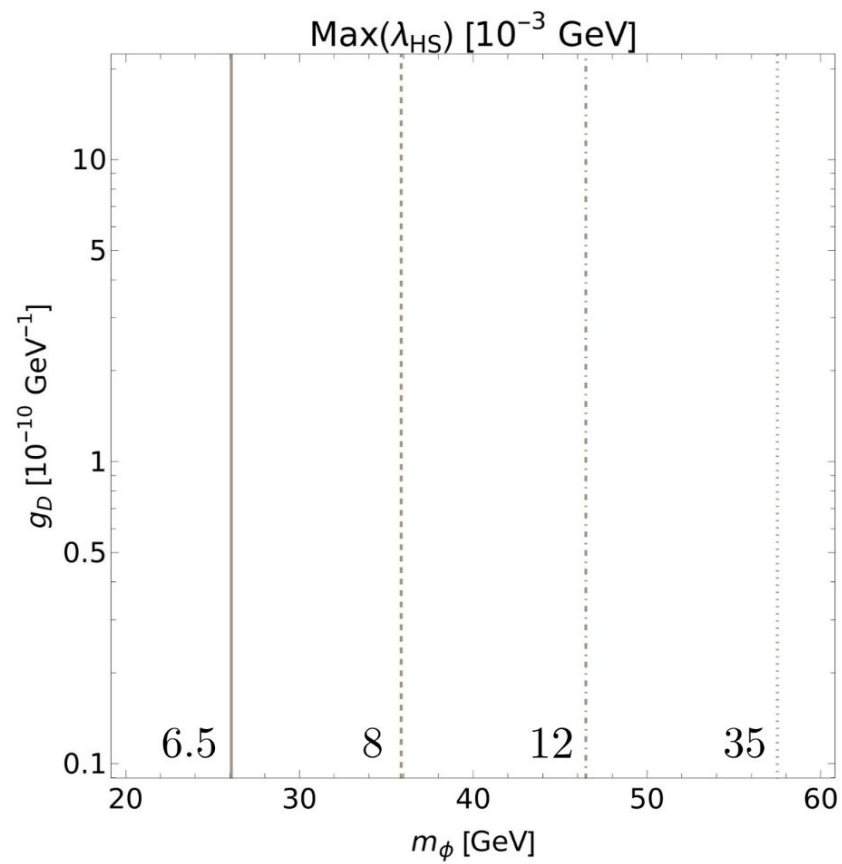
- Adapted to read HepMC with non-pointing Monte Carlo information.
- The photon and electron efficiencies are applied after Delphes.
- Muon isolation uses track and calorimeter information.
- Jets only include non-isolated muons, which later helps overlap removal to favour jets over muons.

Post-Delphes cuts

- Apply gaussian smear on delta Z (later apply on t gamma)
- Apply momentum and eta cuts on photons and separate (later apply these on electrons, muons and jets)
- Apply electron efficiencies and ID
- Implement overlap removal for photons, electrons, jets and muons.
- Assign signal region if $\text{MET} > 50 \text{ GeV}$



Higgs \rightarrow invisible



Higgs \rightarrow undetected



