

Alignment limit and strong first-order electroweak phase transition in 2HDM

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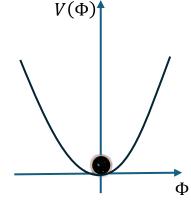


SM explains mass generation : EWPT

- A single scalar SU(2) doublet $\Phi = \begin{pmatrix} \phi^+ \\ \phi^o \end{pmatrix}$,
- With $\mu^2 < 0, \lambda > 0$,

$$V(\Phi) = -\mu^2 (\Phi^{\dagger} \Phi) + \frac{\lambda}{4} (\Phi^{\dagger} \Phi)^2$$

BEH mechanism





EW symmetry spontaneously broken to $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$

➤ Mass to particles (Fermions and Gauge Bosons) and scalar Higgs boson

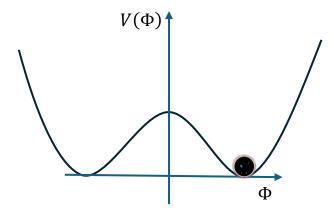


Fig 1. Smooth transition to the EW minimum







At finite-temperature:

$$V_{eff} = \underbrace{V_{tree} + V_{CW} + V_{CT}}_{\text{T-independent}} + \underbrace{V_{T} + V_{daisy}}_{\text{thermal corrections}}$$

 V_{tree} : Tree-level scalar potential

 V_{CW} : One-loop Coleman Weinberg potential (Ref.2)

 V_{CT} : UV-finite counter-term potential

 V_T : One-loop thermal corrections

 V_{daisy} : Resummation of the daisy diagrams

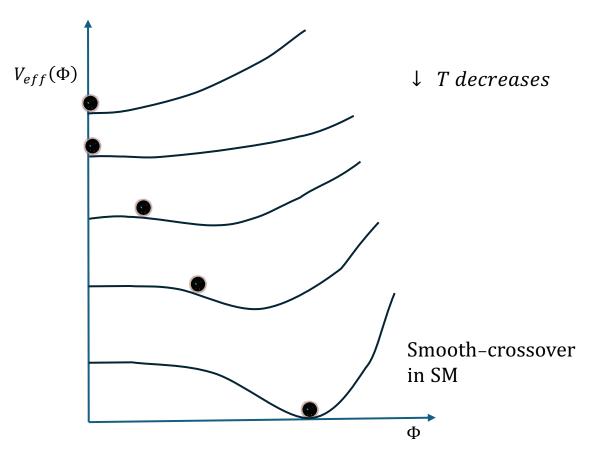


Fig 2.1: Finite-temperature potential: thermal evolution of vacuum as a smooth cross-over

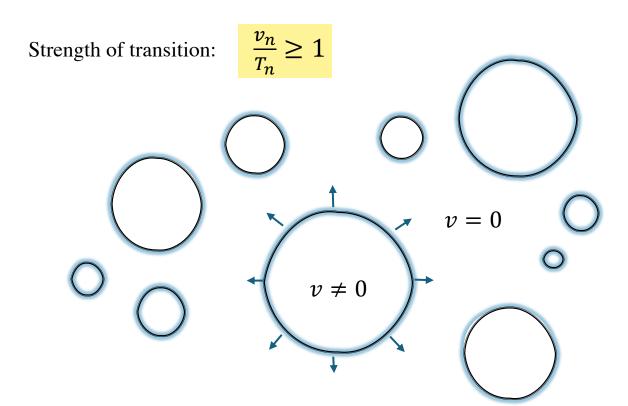


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First-order Phase Transition

- At very high temperatures:
 EW symmetry preserved at v = 0
- Temperature decreased :
 - -> Universe cooled -> develops non-zero vev
 - -> EW symmetry breaks at T_n -> sudden change of vev
 - -> Nucleation of vacuum bubbles



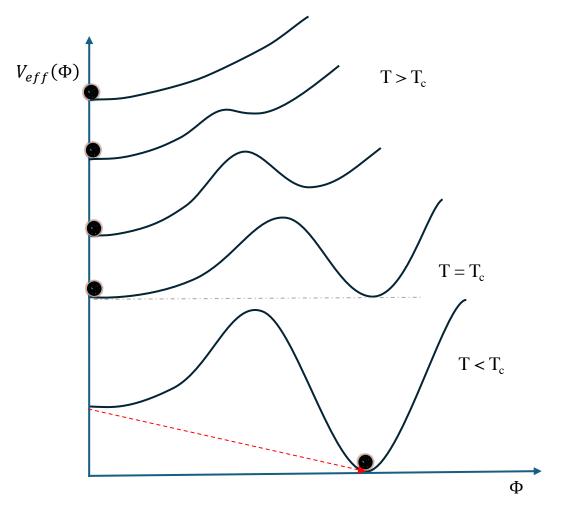


Fig 2.2: First-order phase transition of the universe









- ☐ Matter-Antimatter asymmetry
 - -> *SM prediction*: 6 . 10^{-19} ; Observed: 6 . 10^{-10}
 - -> Sakharov conditions

(B, C and CP violation + Out-of-Thermal-Equilibrium Process)

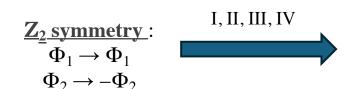
- ☐ EW baryogenesis -> Strong first-order phase transition
- ✓ BSM (extended) Higgs sector required!



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Two-Higgs-Doublet Model (2HDM)



$$V = m_{11}^{2}(\Phi_{1}^{\dagger}\Phi_{1}) + m_{22}^{2}(\Phi_{2}^{\dagger}\Phi_{2}) - m_{12}^{2}(\Phi_{1}^{\dagger}\Phi_{2} + \Phi_{2}^{\dagger}\Phi_{1}) + \frac{\lambda_{1}}{2}(\Phi_{1}^{\dagger}\Phi_{1})^{2} + \frac{\lambda_{2}}{2}(\Phi_{2}^{\dagger}\Phi_{2})^{2} + \lambda_{3}(\Phi_{1}^{\dagger}\Phi_{1})(\Phi_{2}^{\dagger}\Phi_{2}) + \lambda_{4}(\Phi_{1}^{\dagger}\Phi_{2})(\Phi_{2}^{\dagger}\Phi_{1}) + \frac{\lambda_{5}}{2}[(\Phi_{1}^{\dagger}\Phi_{2})^{2} + (\Phi_{2}^{\dagger}\Phi_{1})^{2}].$$

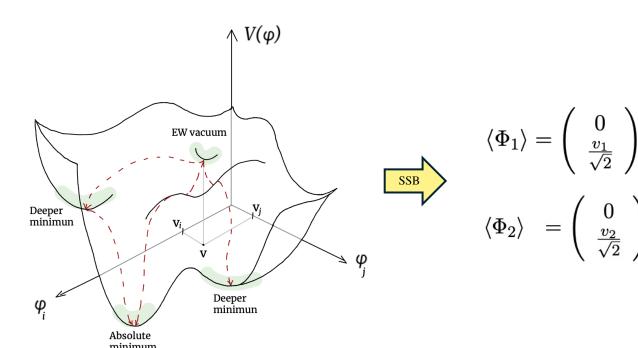


Fig 3: Representation of minima structure of 2HDM (or BSM) scalar potential

In mass basis:

- \rightarrow 5 physical scalar fields (h, H, A, H[±]) diagonalised by α and β,
- \rightarrow 3 would-be-Goldstone bosons (G^0 , G^{\pm})
- Non-trivial structure of the scalar potential







Constraints on 2HDM Parameter Space

Free parameters:

 $an eta, \cos (eta - lpha), \ v, \ m_h$, m_H , m_{H^\pm}, m_A , m_{12}^2 Enters in self-couplings -> barrier shape

Alignment limit:

- $\cos(\beta \alpha) \rightarrow 0$
- Light CP-even Higgs h has SM-like couplings of h_{125} at the lowest order
- Possible in both decoupling and non-decoupling regime

Theoretical constraints:

- Vacuum stability and boundedness-frombelow
- Perturbative unitarity

Experimental constraints:

- Electroweak precision observables
- Constraints are from the limits from searches for additional Higgs bosons
- Measurements of the properties of h_{125}
- Flavour Constraints







Mass splittings

Key Regions:

- 1. $m_A = m_{H^{\pm}}$
- 2. $m_H = m_{H^{\pm}}$
- Larger mass splittings $m_A m_H$ driven by the quartic couplings in the scalar potential correlate to stronger first-order phase transitions.
- For region 1: mass splitting range: 100 200 GeV
- For region 2: mass splitting range: 150 400 GeV
- This dependence is constrained by perturbative unitarity.

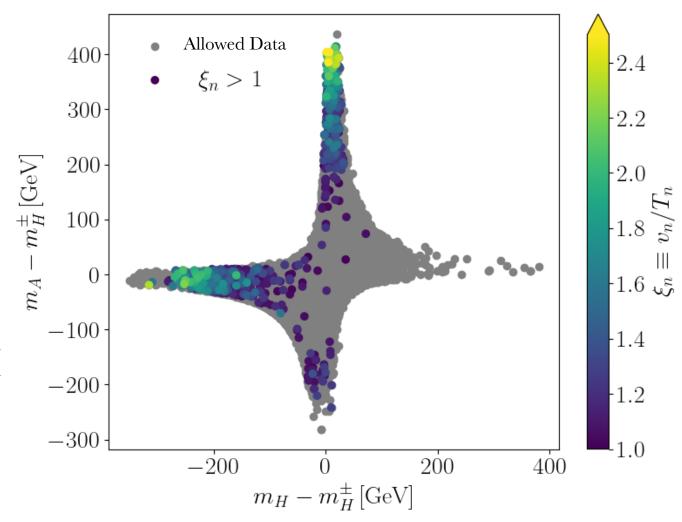


Fig 4: SFOEWPT strength depends on the mass splittings, here the free parameters vary as $1 \le \tan \beta \le 50$ and $150 \le m_H \le 1400$ GeV.







SFOEWPT in $(m_H, \cos(\beta-\alpha))$ plane

For low mass values m_H around 200 GeV:

- SFOEWPT occurs for wide range of $\cos(\beta-\alpha)$ values
- Both CP-even neutral Higgs bosons h and H take part in the EW phase transition → enhanced thermal corrections → relatively higher strengths

For higher mass values m_H :

• SFOEWPT occurs for $cos(\beta-\alpha) \rightarrow 0$, i.e. occurs close to the alignment limit $\rightarrow h$ drives the phase transition

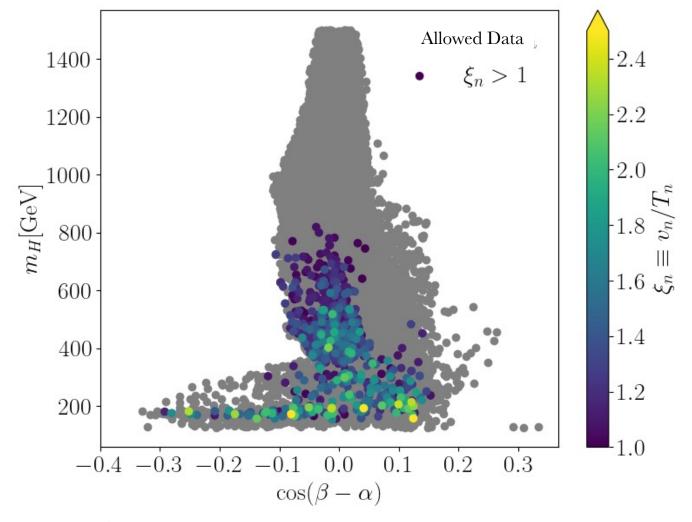


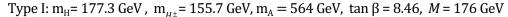
Fig 5: SFOEWPT for CP-even heavy Higgs mass values, here the free parameters vary as $1 \le \tan \beta \le 50$

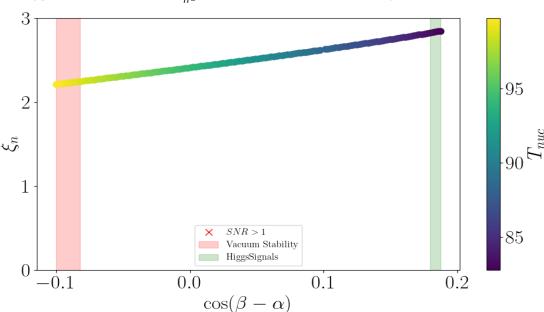


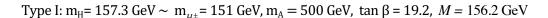




SFOEWPT occurs for wide range of $\cos (\beta - \alpha)$, strength higher far from the alignment limit







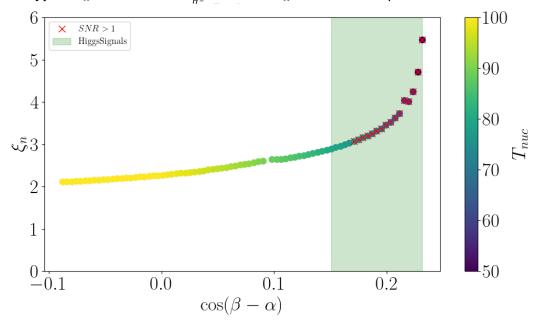


Fig 6a and b: Variation of strength with $cos(\beta-\alpha)$ for lower CP-even Higgs masses to correlate the SFOEWPT parameter space close to as well as away from alignment limit

SFOEWPT vs cos (β - α) for higher m_H





SFOEWPT occurs strongest for $\cos (\beta - \alpha) \rightarrow 0$, i.e. highest strength at the alignment limit

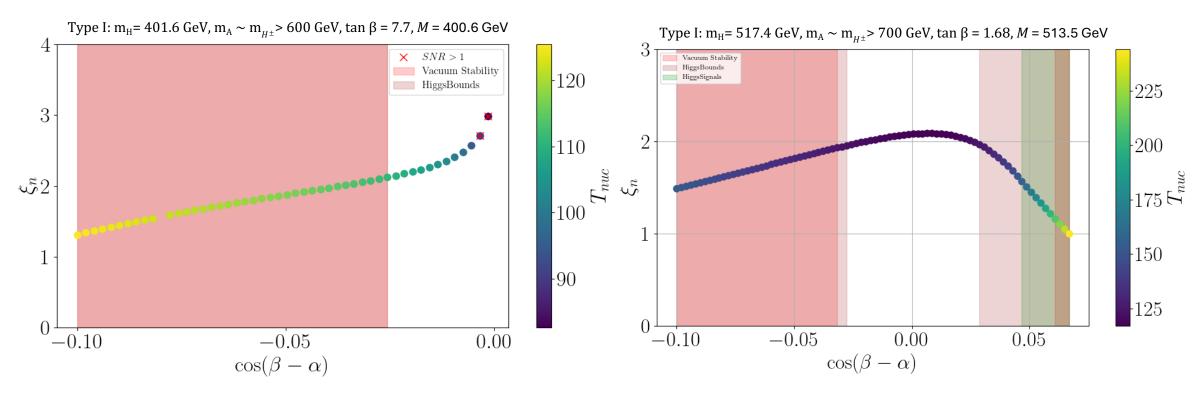


Fig 7a and b: Variation of strength with $cos(\beta-\alpha)$ for higher CP-even Higgs masses to correlate the SFOEWPT parameter space close to alignment limit





Conclusions

- Extended Higgs sectors can feature a SFOEWPT, necessary for EW baryogenesis, can potentially lead to detectable GW signals
- The parameter region giving rise to a SFOEWPT in the 2HDM has a mass splitting between m_A and m_H of about 200 GeV
- For high values of m_H (up to 800 GeV): SFOEWPT occurs close to the alignment limit
- For low values of m_H (~ 200 GeV): SFOEWPT possible for larger deviations from the alignment limit
- Qualitatively, for SFOEWPT parameter points with higher m_H , the strength of phase transition increases with $\cos{(\beta-\alpha)}$, and is highest near the alignment limit, whereas for the case of lower m_H , the strength is higher as we go away from the alignment limit.





Next Steps:

- Extend the current results for the types II, III and IV of 2HDM
- Check for κ_{λ} variation for the above parameter regions showing SFOEWPT in correlation to the alignment limit
- Investigate the prospects for possible GW signals

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