Introduction

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based on a series of works with Pietro Benetti Genolini, Chris Couzens, Jerome Gauntlett, Yusheng Jiao, and James Sparks

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Motivation

Introduction

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CFT observables: free energy, central charges, fluxes, ...

Integrals over (subspaces of) the internal space

$$\mathcal{F} \sim \int_{M_q} \operatorname{vol}_{M_q}$$

Introduction

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SUGRA SCFT
$$Y_d \times \tilde{M}_q \qquad W_{d-1} = \partial Y_d$$

- Observables: on-shell action, free energy, entropy, fluxes, ...
- Typically reduce on the internal space (truncate) \longrightarrow d-dim supergravity

$$I \sim \int_{Y_d} \operatorname{vol}_{Y_d}(R + \dots)$$

Motivation

Introduction

$$\mathcal{F} \sim \int_{M} \operatorname{vol}_{M}, \qquad I \sim \int_{Y} \operatorname{vol}_{Y}(R + \dots)$$

- Looks like we need to know an explicit metric
- This would involve solving supergravity EOM → Difficult!
- Is there a simpler way to get this result?

Motivation

Introduction

$$\mathcal{F} \sim \int_{M} \operatorname{vol}_{M}, \qquad I \sim \int_{Y} \operatorname{vol}_{Y}(R + \dots)$$

- Looks like we need to know an explicit metric
- This would involve solving supergravity EOM → Difficult!
- Is there a simpler way to get this result?
 - → Yes! Equivariant localization

[Benetti Genolini, Gauntlett, Sparks '23]

Introduction

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- M-theory example

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- M-theory example

Equivariant Cohomology

Introduction

- U(1) action with associated vector field ξ
- Equivariant exterior derivative

$$d_{\xi} \equiv d - \xi \rfloor$$

$$+1 \nearrow \qquad \nwarrow$$

acts on polyforms

$$\Phi = \Phi_n + \Phi_{n-2} + \cdots + \Phi_0$$

• Φ is equivariantly closed if $d_{\varepsilon}\Phi = 0$

$$\mathrm{d}\Phi_n = 0$$
 $\xi \, \lrcorner \, \Phi_n = \mathrm{d}\Phi_{n-2}$... $\xi \, \lrcorner \, \Phi_2 = \mathrm{d}\Phi_0$

Introduction

BVAB Theorem: [Berline, Vergne '82, Atiyah, Bott '84] The integral of an equivariantly closed form *localizes* to fixed point contributions

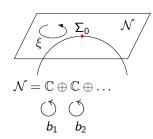
$$\int_{M_n} \Phi = \sum_{\Sigma_0} \frac{(2\pi)^k}{\prod_{i=1}^k b_i} \Phi_0 \Big|_{\Sigma_0}, \qquad n = 2k$$

• $\Sigma_0 \subset M_n$ fixed points set:

$$\xi|_{\Sigma_0}=0$$

• b_i weights of U(1) action:

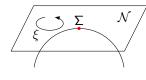
$$\xi = \sum_{i=1}^k b_i \partial_{\varphi_i}$$



BVAB Theorem: [Berline, Vergne '82, Atiyah, Bott '84] The integral of an equivariantly closed form *localizes* to fixed point contributions

$$\int_{M_n} \Phi = \sum_{\Sigma_0} \frac{(2\pi)^k}{\prod_{i=1}^k b_i} \Phi_0 \Big|_{\Sigma_0} + \sum_{\Sigma_2} \frac{(2\pi)^{k-1}}{\prod_{i=1}^{k-1} b_i} \int_{\Sigma_2} \Big[\Phi_2 - \Phi_0 \sum_j \frac{2\pi}{b_j} c_1(L_j) \Big] + \dots$$

$$\xi = \sum_{i=1}^k b_i \partial_{\varphi_i}, \quad \mathcal{N} = L_1 \oplus L_2 \oplus \dots$$



Introduction

BVAB Theorem: [Berline, Vergne '82, Atiyah, Bott '84] The integral of an equivariantly closed form *localizes* to fixed point contributions

$$\int_{M_n} \Phi_n = \sum_{\Sigma_0}$$
 "nut contrib." $+ \sum_{\Sigma_2}$ "bolt contrib."

 \implies Can be evaluated using Φ_0 and Φ_2

It circumvents the need to build an explicit solution to the EOM!

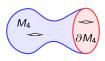
- 2 4d truncation example
- M-theory example

Introduction

Minimal 4d gauged supergravity ←→ Einstein–Maxwell

$$I = -rac{1}{16\pi G_4} \int_{M_4} \left(R + 6 - F^2\right) \mathrm{vol}_4 + \frac{1}{16\pi G_4} \int_{M_4} \left(R + 6 - F^2\right) \mathrm{vol}_4$$

- Can add matter → scalar & gauge fields Interactions specified by a prepotential function \mathcal{F}
- Dual to 3d SCFT on ∂M_4 : $I = -\log Z$



Action

Introduction

$$I = \frac{1}{8\pi G_4} \int_{M_4} \Phi_4$$

Equivariantly closed form

$$\Phi = \Phi_4 + \Phi_2 + \Phi_0 \,, \qquad \mathrm{d}_{\mathcal{E}} \Phi = 0$$

Equivariant localization

$$I = \frac{\pi}{2G_4} \left\{ \sum_{\text{puts}} \frac{\Phi_0}{b_1 b_2} + \sum_{\text{bolts}} \int_{\Sigma} \frac{\Phi_2}{2\pi b} - \frac{\Phi_0 c_1(L)}{b^2} \right\}$$

Introduction

On-shell action of minimal 4d gauged supergravity

[Benetti Genolini, Perez Ipiña, Sparks '19]

$$I = -rac{\pi}{8G_4} \left[\sum_{ ext{nuts}} rac{(b_1 - b_2)^2}{b_1 b_2} + \sum_{ ext{bolts}} \left(\mathfrak{p} - \int_{\Sigma} c_1(L)
ight) \right]$$

- Magnetic charges: $\mathfrak{p} = \frac{1}{2\pi} \int_{\Sigma} F$
- ullet Adding matter straightforward: just introduces prepotential ${\cal F}$

Main Result

Introduction

On-shell action of 4d gauged supergravity with arbitrary matter [Benetti Genolini, Gauntlett, Jiao, AL, Sparks '24]

$$I = -rac{\pi}{G_4} \left[\sum_{
m nuts} rac{(b_1 - b_2)^2}{b_1 b_2} \mathrm{i} \mathcal{F}(u^J)
ight.
onumber \ \left. + \sum_{
m bolts} \left(\partial_I \mathrm{i} \mathcal{F}(u^J) \mathfrak{p}^I - \mathrm{i} \mathcal{F}(u^J) \int_{\Sigma} c_1(L)
ight)
ight]$$

- F general prepotential (specifies the theory)
- u^J combinations of the scalar fields at the fixed points

e.g.
$$\mathcal{F}_{STU}(u^I) = -2i\sqrt{u^0u^1u^2u^3}$$

STU Result

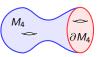
Introduction

On-shell action of 4d U(1)⁴ supergravity

[Benetti Genolini, Gauntlett, Jiao, AL, Sparks '24]

$$I = -\frac{2\pi}{G_4} \left[\sum_{\text{nuts}} \frac{(b_1 - b_2)^2}{b_1 b_2} \sqrt{u^0 u^1 u^2 u^3} + \sum_{\text{bolts}} \sqrt{u^0 u^1 u^2 u^3} \left(\sum_{I=0}^3 \frac{\mathfrak{p}^I}{u^I} - \int_{\Sigma} c_1(L) \right) \right]$$

- STU on $M_4 \longleftrightarrow \text{dual to ABJM on } \partial M_4$
- Pick a topology for M_4 : nuts and/or bolts



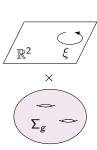
Euclidean Black Hole

$$M_4 = \mathbb{R}^2 imes \Sigma_g$$

Fixed pt set: Σ_g at the origin of \mathbb{R}^2

$$\implies I = \text{bolt contrib.}$$

$$= -\frac{\pi}{G_4} \sqrt{u^0 u^1 u^2 u^3} \sum_{l=0}^{3} \frac{\mathfrak{p}^l}{u^l}$$



Euclidean Black Hole

Introduction

$$M_4 = \mathbb{R}^2 imes \Sigma_g$$

Fixed pt set: Σ_{φ} at the origin of \mathbb{R}^2

$$\implies I = \text{bolt contrib.}$$

$$= -\frac{\pi}{G_4} \sqrt{u^0 u^1 u^2 u^3} \sum_{l=0}^{3} \frac{\mathfrak{p}^l}{u^l}$$

$$\begin{array}{c|c}
 & \searrow \\
 & & \xi \\
 & \times \\
 & & \times \\
 & & \Sigma_g
\end{array}$$

Rotating Black Hole

Introduction

$$M_4 = \mathcal{O}(p)
ightarrow \Sigma_g$$

Fixed pt set: Σ_g at the origin of \mathbb{R}^2

$$\implies I = \text{bolt contrib. with } c_1 \neq 0$$

$$= -\frac{\pi}{G_4} \sqrt{u^0 u^1 u^2 u^3} \Big(\sum_{l=0}^3 \frac{\mathfrak{p}^l}{u^l} - 2\mathfrak{p} \Big)$$

Rotating black saddles Solution unknown!

$$\longrightarrow$$
 ABJM on $S^1 \ltimes \Sigma_g$
[Toldo, Willett '17]

Accelerating Black Hole

Introduction

$$M_4 = \mathbb{R}^2 \times \text{spindle}$$

Fixed pt set: 2 poles of the spindle

$$\begin{array}{c|c}
 & \downarrow \\
 & \xi \\
 & \times
\end{array}$$

$$\implies I = \text{north nut} + \text{south nut}$$

$$= -\frac{2\pi}{G_4 b_0} \left[\sqrt{y_N^0 y_N^1 y_N^2 y_N^3} \pm \sqrt{y_S^0 y_S^1 y_S^2 y_S^3} \right]$$

$$\pm \sqrt{y_S^0 y_S^1 y_S^2 y_S^3} \bigg]$$



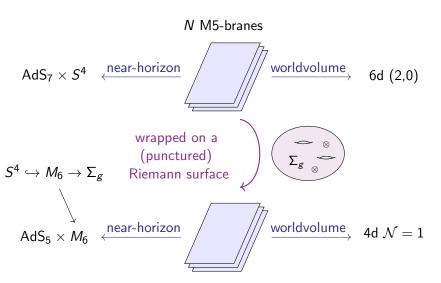
Accelerating black saddles [Crisafio, Fontanarossa, Martelli '24] came after our result!

ABJM on $S^1 \times \text{spindle}$ [Colombo, Hosseini, Martelli, Pittelli, Zaffaroni '24]

2 4d truncation example

M-theory example

Motivation



M-theory Setup

11d supergravity on AdS₅ \times M_6 [Gauntlett, Martelli, Sparks, Waldram '04]

- ullet Dual to 4d $\mathcal{N}=1$ SCFT
- Central charge

$$a = \frac{1}{2(2\pi)^6 \ell_{\rm p}^9} \int_{M_6} {\rm e}^{9\lambda} {
m vol}_{M_6}$$

Introduction

Central charge

$$a = \frac{1}{2(2\pi)^6 \ell_{\rm p}^9} \int_{M_6} \Phi_6$$

Equivariantly closed form

$$\Phi = \Phi_6 + \Phi_4 + \Phi_2 + \Phi_0 \,, \qquad \mathrm{d}_{\xi} \Phi = 0$$

Equivariant localization

$$\begin{aligned} a &= \frac{1}{2(2\pi)^3 \ell_{\rm p}^9} \bigg\{ \sum_{\rm nuts} \frac{\Phi_0}{\epsilon_1 \epsilon_2 \epsilon_3} \\ &+ \sum_{\rm bolts} \frac{1}{\epsilon_1 \epsilon_2} \int_{\Sigma} \frac{\Phi_2}{2\pi} - \Phi_0 \bigg(\frac{c_1(L_1)}{\epsilon_1} + \frac{c_1(L_2)}{\epsilon_2} \bigg) \bigg\} \end{aligned}$$

Smooth Riemann surface

Introduction

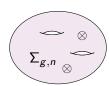
[Benetti Genolini, Gauntlett, Sparks '23]

$$M_6 = S^4 \xrightarrow[\mathcal{O}(p_1) \oplus \mathcal{O}(p_2)]{} \Sigma_g \,, \quad p_1 + p_2 = 2(g-1)$$
 $\xi = b_1 \partial_{\varphi_1} + b_2 \partial_{\varphi_2} \,, \qquad b_1 + b_2 = 1$ Fixed pt set: Σ_g at the poles of S^4

$$\Rightarrow a^{\Sigma_g} = ext{bolt contrib.}$$
 $= -rac{9}{8}b_1b_2(b_1p_2 + b_2p_1)N^3$ [Bah, Beem, Bobev, Wecht '12]

Adding Punctures

Introduction



$$a^{\sum_{g,n}} = a^{\sum_g} - \sum_{l=1}^n \delta a^l$$

M-theory example 00000●

Compute puncture contrib. δa with equivariant localization

→ Completely **new** results! [Couzens, **AL**, Sparks '25]

Summary

Introduction

Lower dimensions

- 4d [Benetti Genolini, Gauntlett, Jiao, AL, Sparks '24]
- 5d [Cassani, Ruiperez, Turetta '24, Colombo, Dimitrov, Martelli, Zaffaroni '25, Benetti Genolini, Gauntlett, Jiao, Park, Sparks '25]
- 6d [Couzens, Gregory, Muniz, Sieper, Sparks '25]
- near-horizon [Benetti Genolini, Gauntlett, Jiao, AL, Sparks '24, Suh '24]

Internal spaces

- M-theory on AdS₅ \times M_6 [Benetti Genolini, Gauntlett, Sparks '23]
- Massive type IIA on $AdS_4 \times M_6$ [Couzens, AL '24]
- Type IIB AdS₃ \times M_7 [Couzens, AL, Sparks '25]

Orbifold geometries

- Punctured Riemann surface [Couzens, AL, Sparks '25 + to appear]
- Other punctured manifolds
- Defects

Outlook

Take home message

Equivariant localization is a powerful tool to compute holographic observables in supergravity without needing the explicit solution.

Many interesting open questions

- Full 11d spaces
- No supersymmetry
- Higher-derivative corrections

Outlook

Introduction

Take home message

Equivariant localization is a powerful tool to compute holographic observables in supergravity without needing the explicit solution.

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Thank you! Questions?



Appendix

Supergravity Applications

Supergravity EOM ←→ Killing spinor equations

$$(\nabla + \dots)\epsilon = 0$$

• Spinor bilinears \overrightarrow{f} dual to Killing vect. ξ

$$S = \bar{\epsilon}\epsilon$$
, $P = \bar{\epsilon}\gamma_*\epsilon$, $\xi^{\flat} = \bar{\epsilon}\gamma_{\mu}\gamma_*\epsilon dx^{\mu}$

- → follow constraints coming from KSE
- Tool to build equiv. closed forms $(d \xi)\Phi = 0$

e.g.
$$dP = \xi \, \lrcorner \, F \implies \Phi = F - P$$
 equiv. closed

• Game: build such forms whose top form part is an obs.

Supergravity Applications

Applications to supergravity [Benetti Genolini, Gauntlett, Sparks '23]

- Supersymmetry implies the existence of a U(1) symmetry
- Use Killing spinor equations and spinor bilinears to build equiv. closed forms → once
- Localize on various topologies → different sets of fixed points

Obs. =
$$\int_{M_n} \Phi_n = \sum_{\substack{\text{fixed} \\ \text{points}}} \Phi_0$$