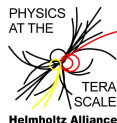


Monte Carlo — Lecture III

Stefan Gieseke

Institut für Theoretische Physik
KIT

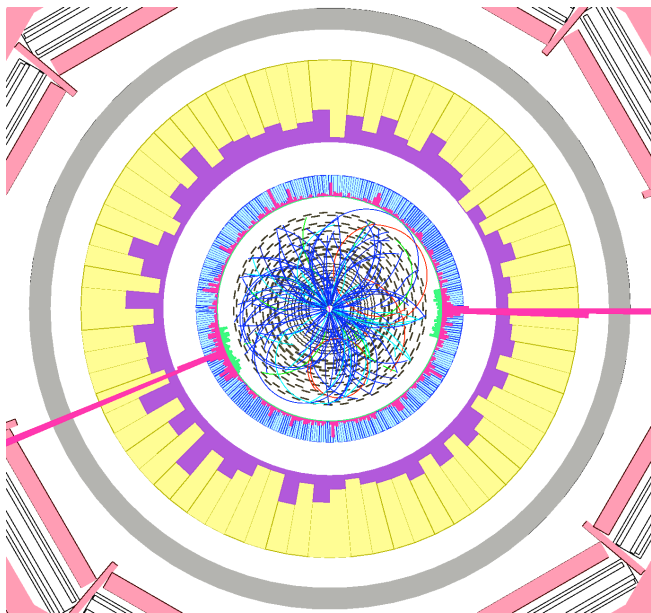
4th Terascale MC School, 12-15 Mar 2012

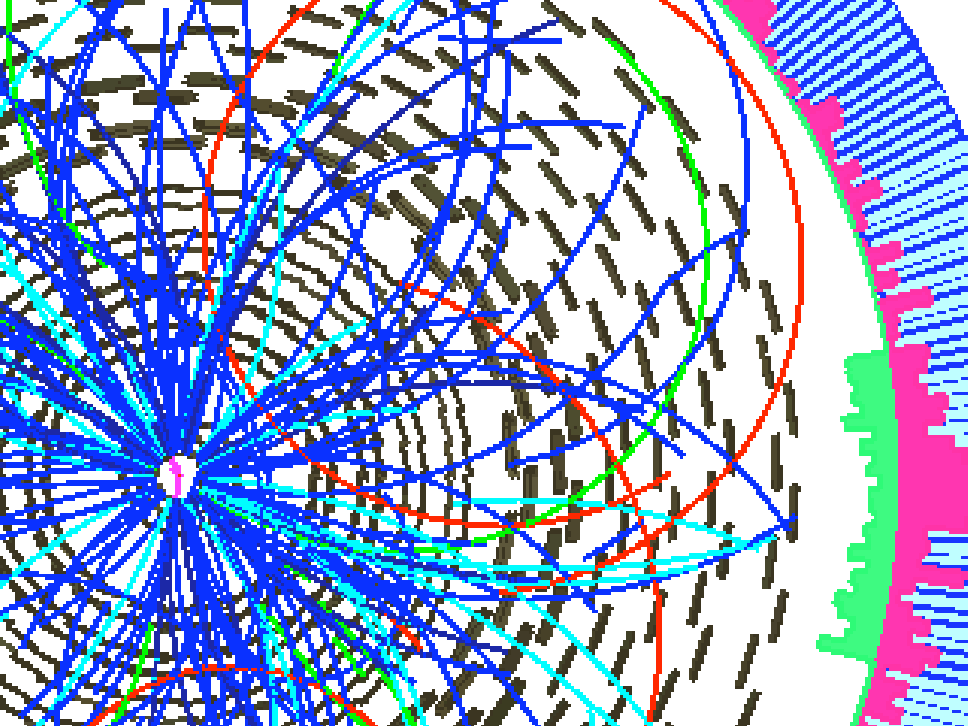


- ▶ Lecture I — Basics
 - ▶ Introduction
 - ▶ Monte Carlo techniques
- ▶ Lecture II — Perturbative physics
 - ▶ Hard scattering
 - ▶ Parton showers
- ▶ Lecture III — Non-perturbative physics
 - ▶ Hadronization
 - ▶ Hadronic decays
 - ▶ Comparison to data
- ▶ Lecture IV — Multiple Partonic Interactions
 - ▶ Minimum Bias/Underlying Event in data
 - ▶ Modelling

- ▶ Hadronization
 - ▶ Confinement
 - ▶ Flux tube model
 - ▶ Independent fragmentation
 - ▶ Lund string model
 - ▶ Cluster hadronization model
- ▶ Hadronic decays
- ▶ Comparison to data

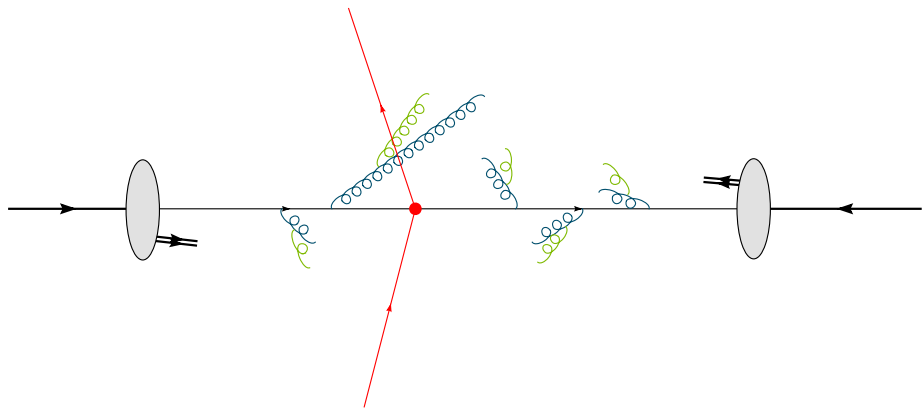
Can you spot the Higgs?



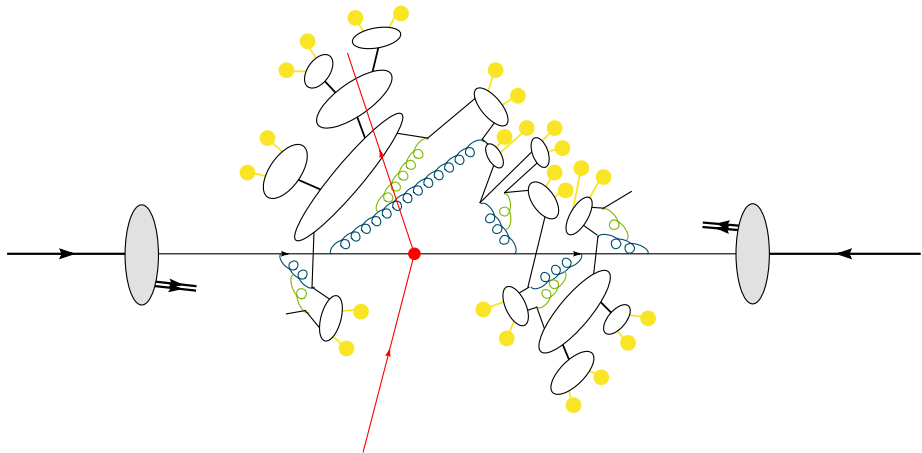


Hadronization

Parton shower



Parton shower \longrightarrow hadrons



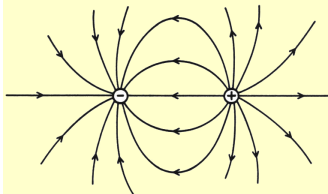
Parton shower \longrightarrow hadrons

- ▶ Parton shower terminated at $t_0 =$ lower end of PT.
- ▶ Can't measure quarks and gluons.
- ▶ Degrees of freedom in the detector are **hadrons**.
- ▶ Need a description of **confinement**.

Physical input

Self coupling of gluons
↔ “attractive field lines”

QED FIELD LINES

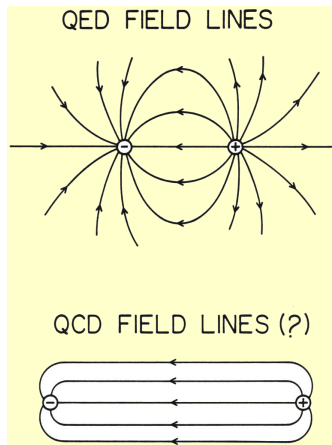


QCD FIELD LINES (?)

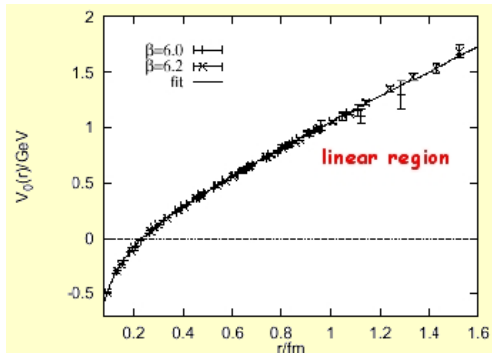


Physical input

Self coupling of gluons
 \leftrightarrow “attractive field lines”



Linear static potential $V(r) \approx \kappa r$.



Supported by lattice QCD,
hadron spectroscopy.

Older models:

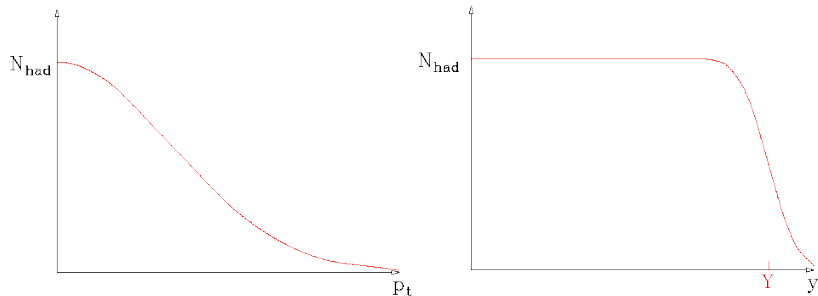
- ▶ Flux tube model.
- ▶ Independent fragmentation.

Today's models.

- ▶ Lund string model (Pythia).
- ▶ Cluster model (Herwig).

Flux tube model

$e^+e^- \rightarrow$ hadrons: hadrons in “flux tube”.
Limited p_\perp , flat in y .



Suggests simple model for jet mass (Feynman, '72).

Estimate hadronization corrections to perturbative quantities.
Jet Energy and momentum:

$$E = \int_0^Y dy d^2p_{\perp} \rho(p_{\perp}) p_{\perp} \cosh y = \lambda \sinh Y$$

$$P = \int_0^Y dy d^2p_{\perp} \rho(p_{\perp}) p_{\perp} \sinh y = \lambda (\cosh Y - 1) \approx E - \lambda$$

$$\langle p_{\perp} \rangle = \int d^2p_{\perp} \rho(p_{\perp}) p_{\perp} = \lambda .$$

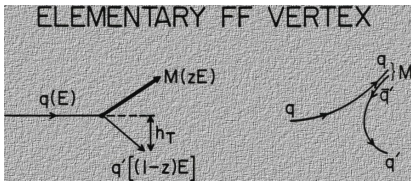
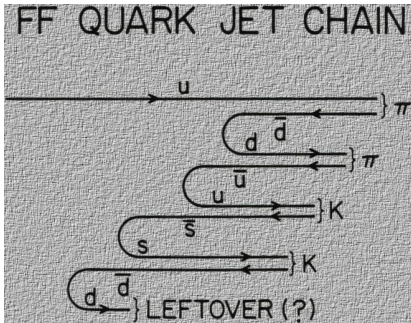
Motion inside confined hadrons, estimate

$$\langle p_{\perp} \rangle \sim 1/R_{\text{had}} \sim m_{\text{had}} \sim 1 \text{ GeV}.$$

\implies Jets acquire non-perturbative mass $M^2 \sim E^2 - P^2 \sim 2\lambda E$.

Large hadronization correction, $O(10 \text{ GeV})$ for $O(100 \text{ GeV})$ jet.

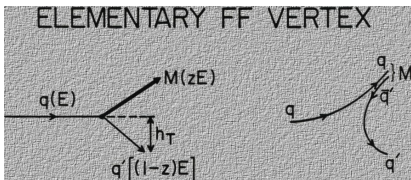
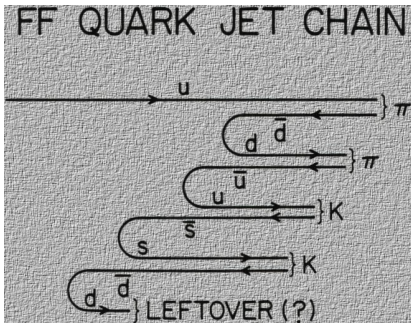
Independent fragmentation



Feynman–Field fragmentation ('78).

- ▶ $q\bar{q}$ pairs created from vacuum to dress bare quarks.
- ▶ Fragmentation function $f_{q \rightarrow h}(z) =$ density of momentum fraction z carried away by hadron h from quark q .
- ▶ Gaussian p_{\perp} distribution.

Independent fragmentation



Feynman–Field fragmentation ('78).

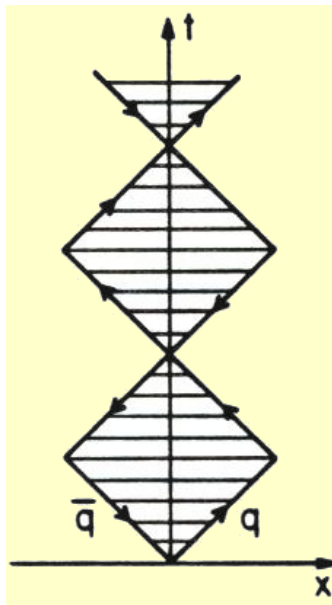
- ▶ $q\bar{q}$ pairs created from vacuum to dress bare quarks.
- ▶ Fragmentation function $f_{q \rightarrow h}(z)$ = density of momentum fraction z carried away by hadron h from quark q .
- ▶ Gaussian p_{\perp} distribution.
- ▶ Problems:
 - ▶ “last quark”.
 - ▶ not Lorentz invariant.
 - ▶ infrared safety.
 - ▶ ...
- ▶ Good at that time.
- ▶ Still useful for inclusive descriptions.

Lund string model

String model of mesons.

$L = 0$ mesons move in yoyo modes.

Area law: $m^2 \sim \text{area}$.



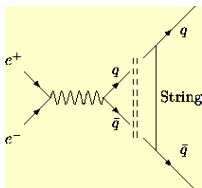
Lund string model

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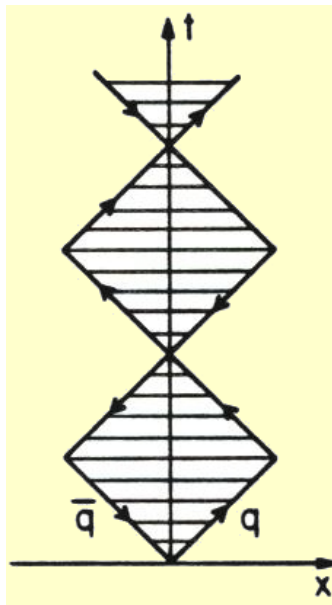
$L = 0$ mesons move in yoyo modes.

Area law: $m^2 \sim \text{area}$.

Simple model for particle production
in e^+e^- annihilation:



$q\bar{q}$ pair as pointlike source of string.

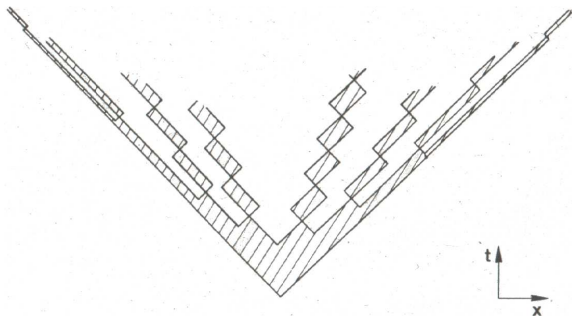


Lund string model

String energy \sim intense chromomagnetic field.

\rightarrow Additional $q\bar{q}$ pairs created by QM tunneling.

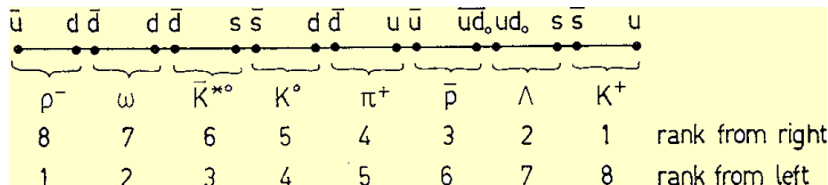
$$\frac{d\text{Prob}}{dxdt} \sim \exp\left(-\pi m_q^2 / \kappa\right) \quad \kappa \sim 1 \text{ GeV}.$$



String breaking expected long before yoyo point.

Lund string model

Ajacent breaks form hadrons.



Works in both directions (symmetry).

Lund symmetric fragmentation function

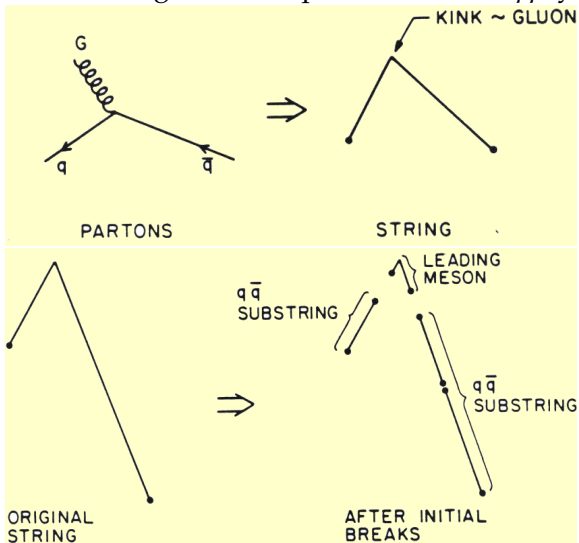
$$f(z, p_{\perp}) \sim \frac{1}{z} (1-z)^a \exp\left(-\frac{b(m_h^2 + p_{\perp}^2)}{z}\right)$$

a, b, m_h^2 main adjustable parameters.

Note: diquarks \rightarrow baryons.

Lund string model

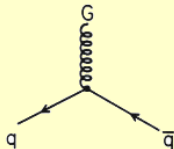
gluon = kink on string = motion pushed into the $q\bar{q}$ system.



Lund string model

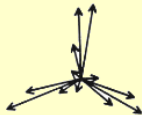
gluon = kink on string = motion pushed into the $q\bar{q}$ system.

SYMMETRIC PARTON CONFIGURATION



HADRONIZATION

INDEPENDENT
FRAGMENTATION

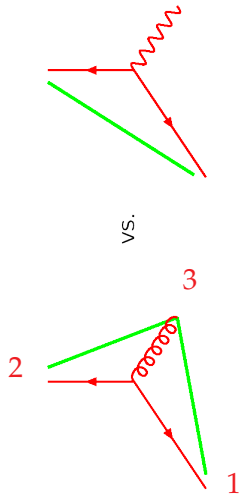


LUND
PICTURE

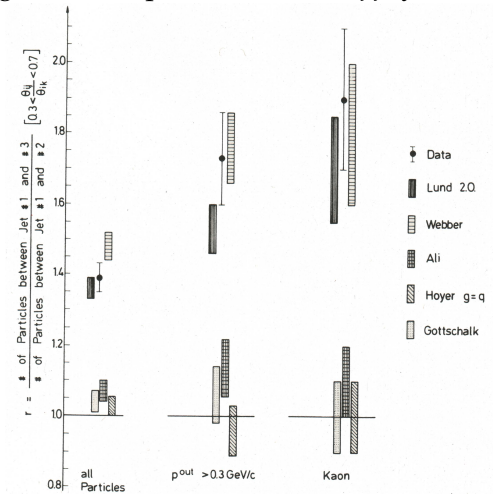


Lund string model

gluon = kink on string = motion pushed into the $q\bar{q}$ system.



“String effect”



Some remarks:

- ▶ Originally invented without parton showers in mind.

Some remarks:

- ▶ Originally invented without parton showers in mind.
- ▶ Strong physical motivation.
- ▶ Very successful description of data.
- ▶ Universal description of data
(fit at e^+e^- , transfer to hadron-hadron).
- ▶ Many parameters, ~ 1 per hadron.
- ▶ Too easy to hide errors in perturbative description?

Some remarks:

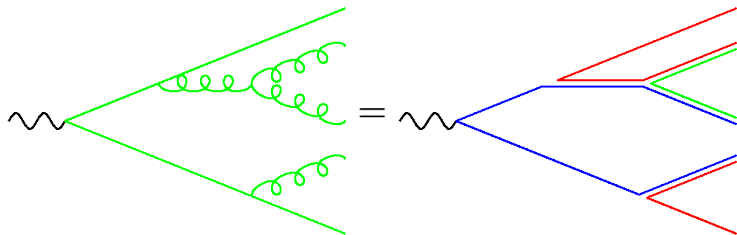
- ▶ Originally invented without parton showers in mind.
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(fit at e^+e^- , transfer to hadron-hadron).
- ▶ Many parameters, ~ 1 per hadron.
- ▶ Too easy to hide errors in perturbative description?

→ try to use more QCD information/intuition.

Colour preconfinement

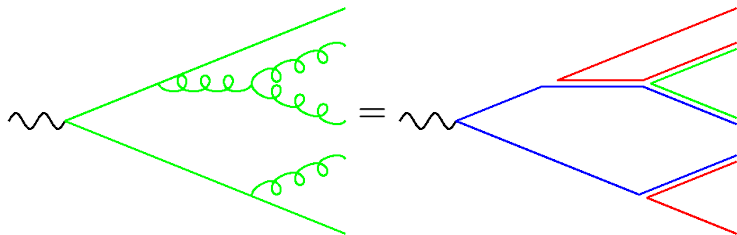
Large N_C limit \rightarrow planar graphs dominate.

Gluon = colour — anticolourpair



Colour preconfinement

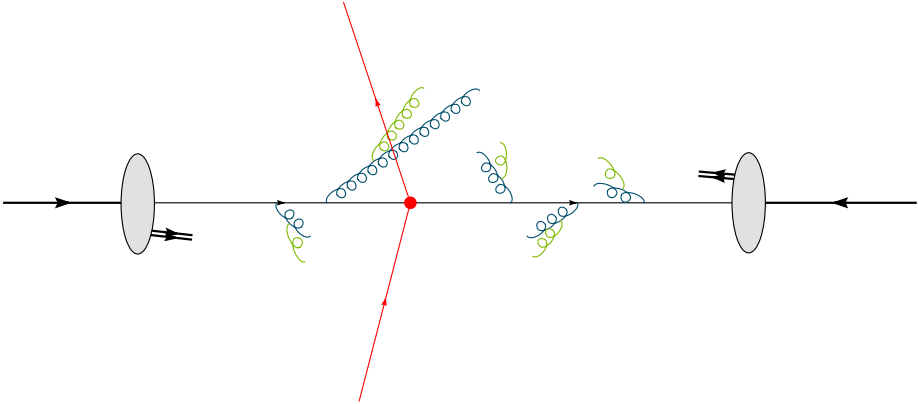
Large N_C limit \rightarrow planar graphs dominate.
Gluon = colour — anticoulourpair



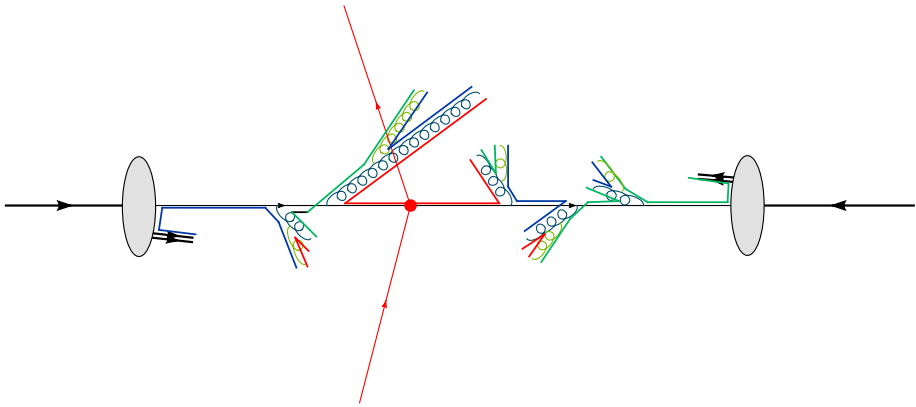
Parton shower organises partons in colour space. Colour partners (=colour singlet pairs) end up close in phase space.

\rightarrow Cluster hadronization model

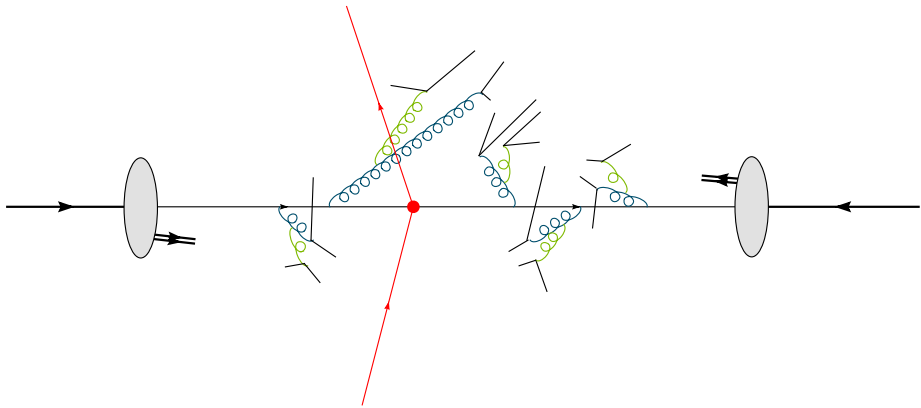
Cluster hadronization



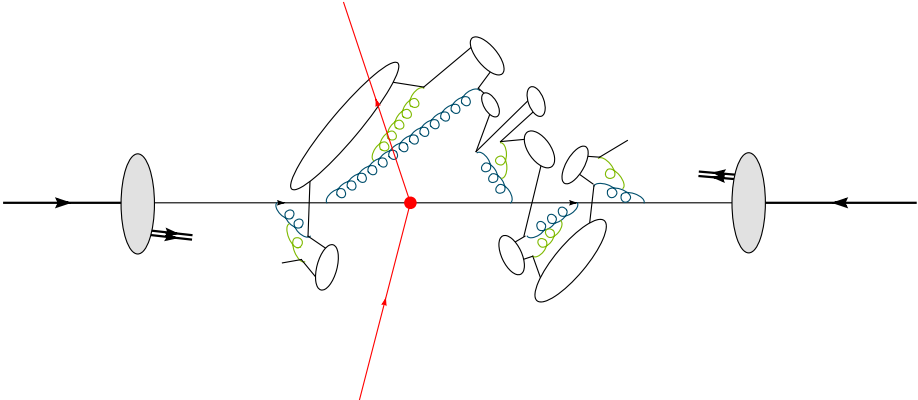
Cluster hadronization



Cluster hadronization



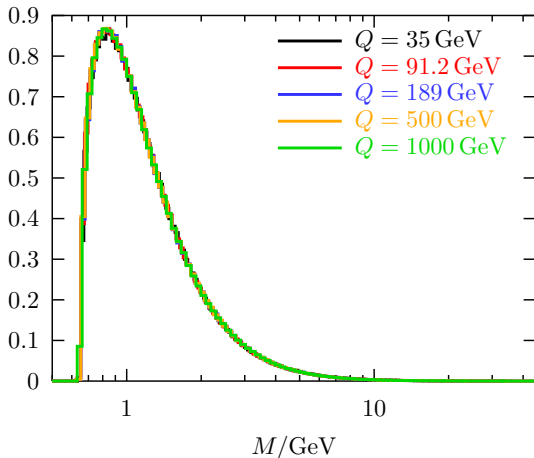
Cluster hadronization



Cluster hadronization

Primary cluster mass spectrum independent of production mechanism. Peaked at some low mass.

Primary Light Clusters



Cluster hadronization

Primary cluster mass spectrum independent of production mechanism. Peaked at some low mass.

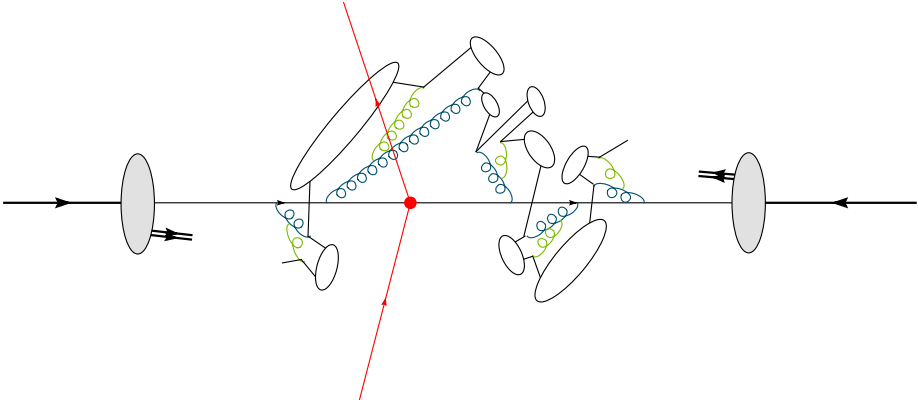
Cluster = continuum of high mass resonances.
Decay into well-known lighter mass resonances
= discrete spectrum of hadrons.

No spin information carried over, i.e. only phase space.

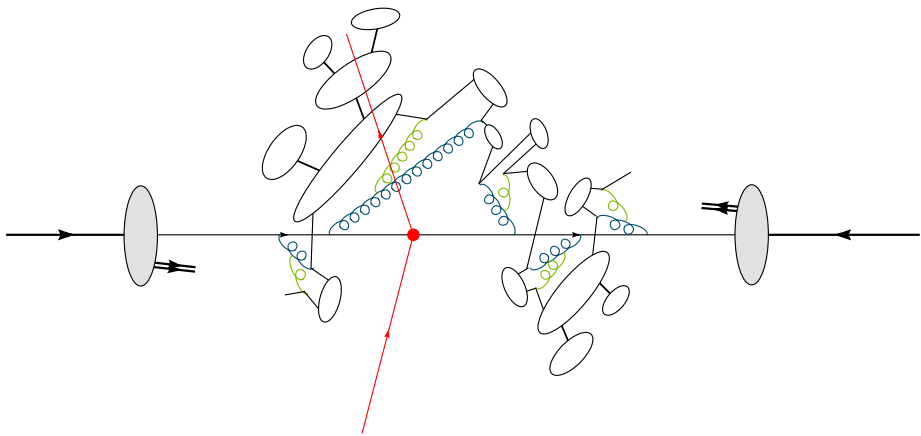
Suppression of heavier particles
(particularly baryons, can be problematic).

Cluster spectrum determined entirely by parton shower,
i.e. perturbation theory. Hence, t_0 crucial parameter.

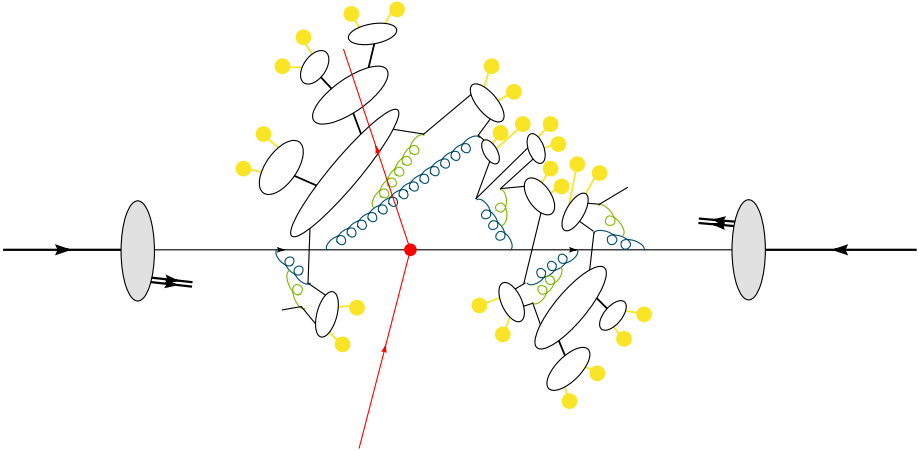
Cluster hadronization



Cluster hadronization



Cluster hadronization



Cluster hadronization in a nutshell

- ▶ **Nonperturbative $g \rightarrow q\bar{q}$ splitting** ($q = uds$) isotropically. Here, $m_g \approx 750 \text{ MeV} > 2m_q$.
- ▶ **Cluster formation**, universal spectrum (see below)
- ▶ **Cluster fission**, until

$$M^p < M_{\text{max}}^p + (m_1 + m_2)^p$$

where masses are chosen from

$$M_i = \left[\left(M^p - (m_i + m_3)^p \right) r_i + (m_i + m_3)^p \right]^{1/p},$$

with additional phase space constraints. Constituents keep moving in their original direction.

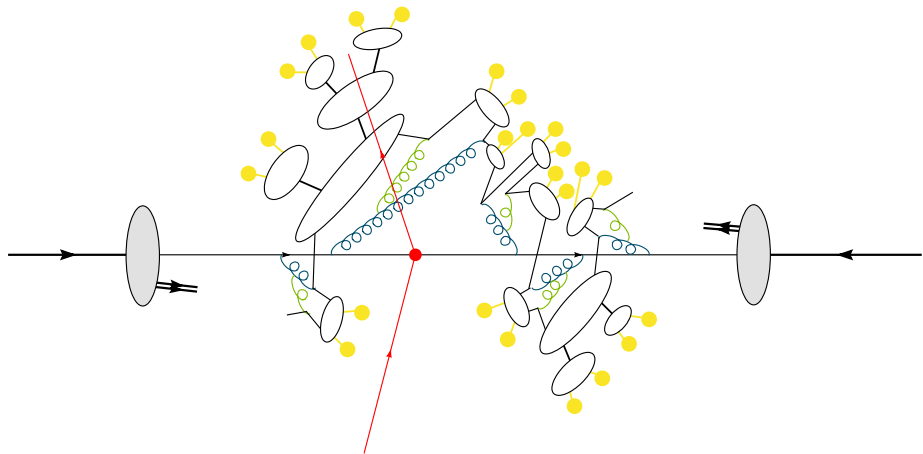
- ▶ **Cluster Decay**

$$P(a_{i,q}, b_{q,j} | i, j) = \frac{W(a_{i,q}, b_{q,j} | i, j)}{\sum_{M/B} W(c_{i,q'}, d_{q',j} | i, j)}.$$

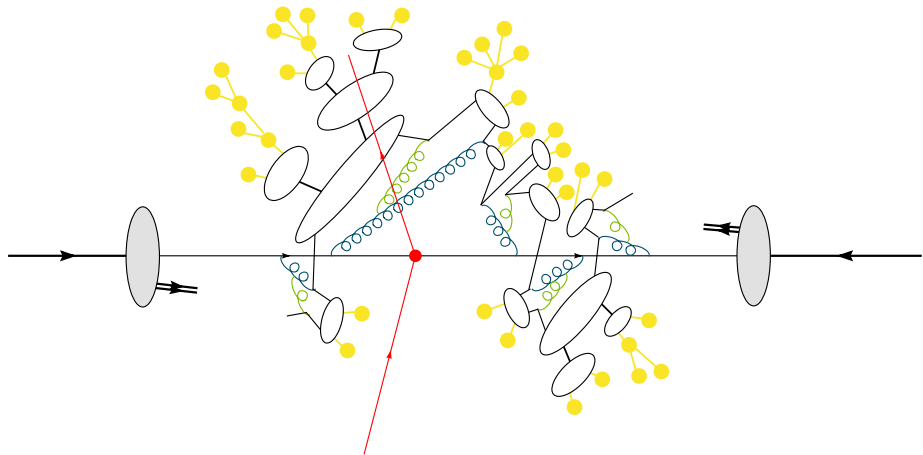
- ▶ Only string and cluster models used in recent MC programs.
Independent fragmentation only for inclusive observables.
- ▶ Strings started non-perturbatively, improved by parton shower.
- ▶ Cluster model started mostly on perturbative side, improved by string like cluster fission.

Hadronic Decays

Hadronic decays



Hadronic decays



Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

Many aspects:

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$$\hookrightarrow e^+ e^- \gamma$$

EM decay.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

Weak mixing.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

Weak decay.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

Strong decay.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

Weak decay, ρ^+ mass smeared.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

$$\hookrightarrow \pi^+ \pi^0$$

$$\hookrightarrow e^+ e^- \gamma$$

ρ^+ polarized, angular correlations.

Many aspects:

$$B^{*0} \rightarrow \gamma B^0$$

$$\hookrightarrow \bar{B}^0$$

$$\hookrightarrow e^- \bar{\nu}_e D^{*+}$$

$$\hookrightarrow \pi^+ D^0$$

$$\hookrightarrow K^- \rho^+$$

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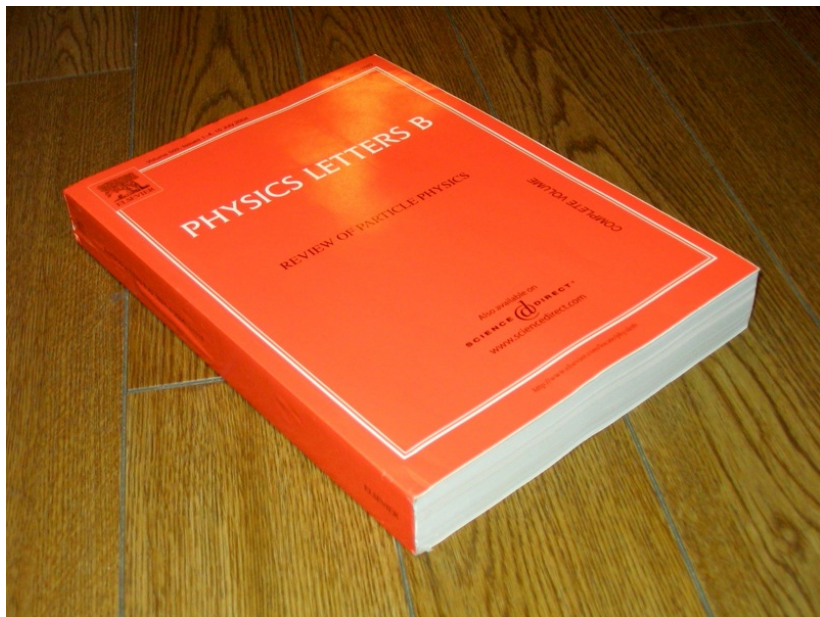
$$\hookrightarrow e^+ e^- \gamma$$

Dalitz decay, m_{ee} peaked.

Tedious.

100s of different particles, 1000s of decay modes,
phenomenological matrix elements with parametrized form
factors...

Hadronic decays



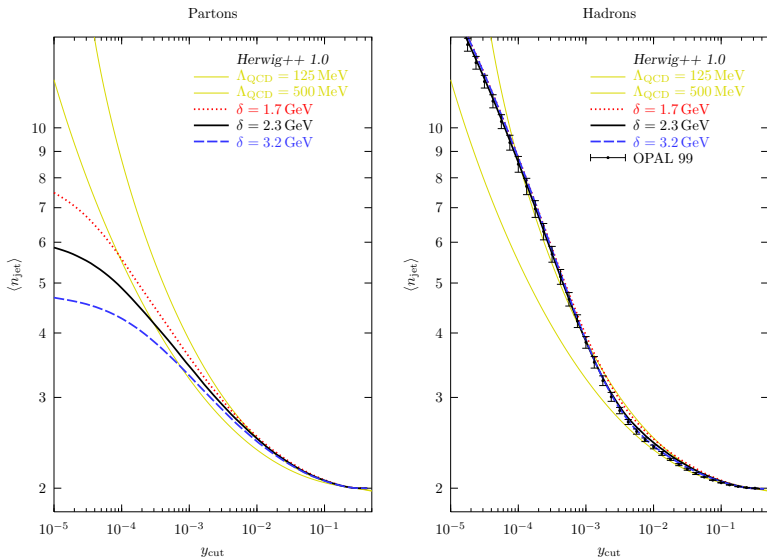
A few plots

How well does it work?

- ▶ $e^+e^- \rightarrow$ hadrons, mostly at LEP.
- ▶ Jet shapes, jet rates, event shapes, identified particles...
- ▶ 'Tuning' of parameters.
- ▶ Want to get *everything* right with *one* parameter set.
- ▶ Compare to literally 100s of plots.

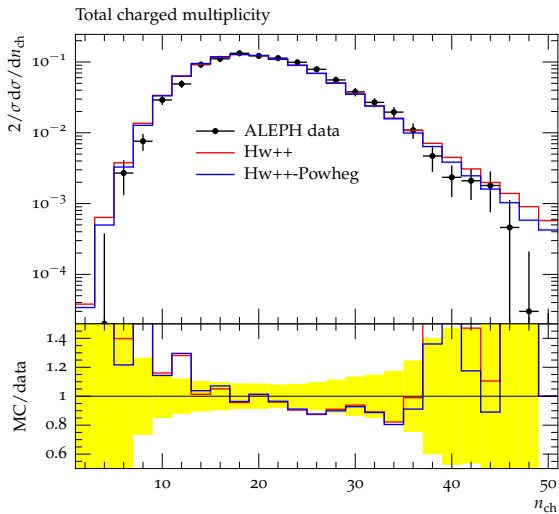
How well does it work?

Smooth interplay between shower and hadronization.



How well does it work?

N_{ch} at LEP. Crucial for t_0 (Herwig++ 2.5.2)



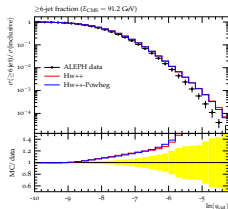
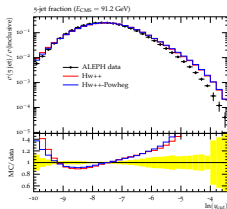
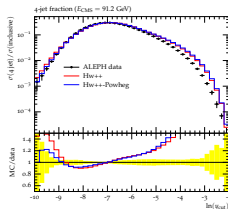
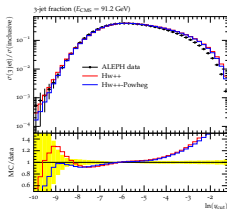
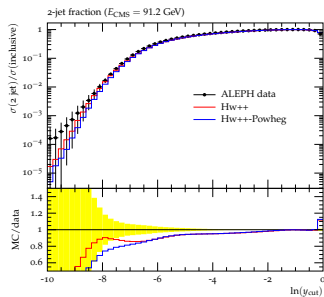
How well does it work?

Jet rates at LEP.

$$R_n = \sigma(n\text{-jets})/\sigma(\text{jets})$$

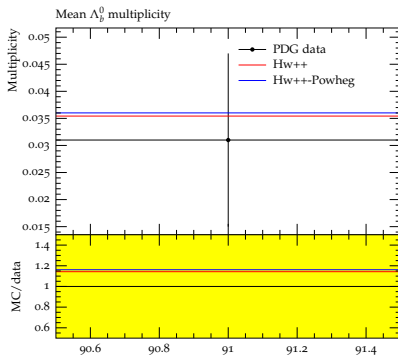
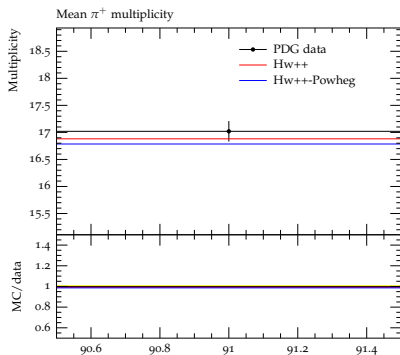
$$R_6 = \sigma(> 5\text{-jets})/\sigma(\text{jets})$$

(Herwig++ 2.5.2)



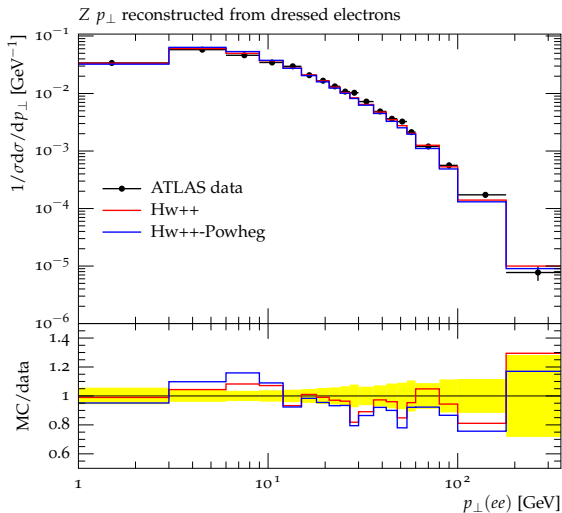
How well does it work?

Hadron Multiplicities at LEP (e.g. π^+ , Λ_b^0).

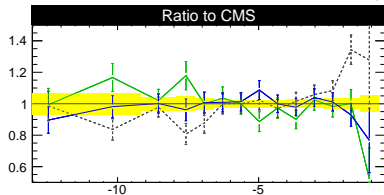
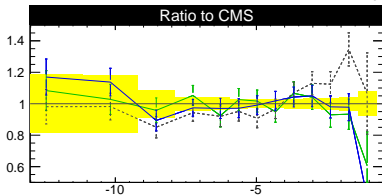
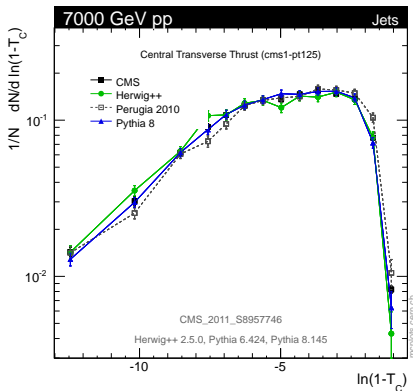
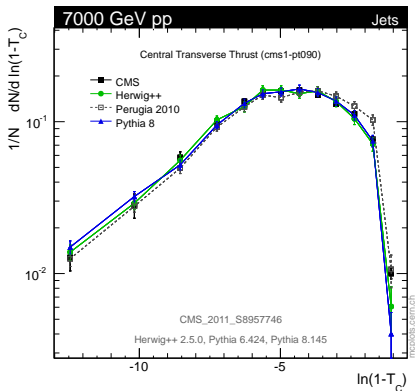


How well does it work?

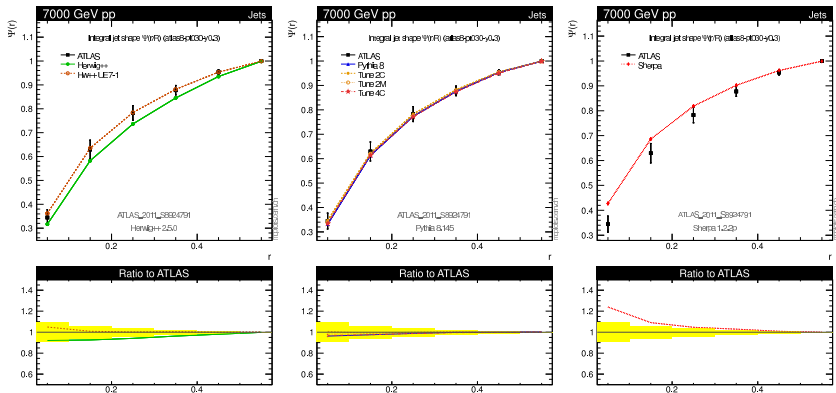
$p_{\perp}(Z^0) \rightarrow$ intrinsic k_{\perp} (LHC 7 TeV).



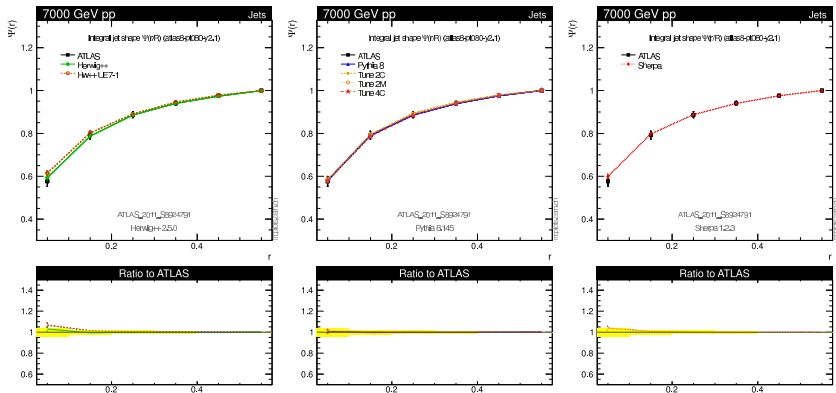
Transverse thrust



not too hard, central ($30 < p_T/\text{GeV} < 40; 0 < |y| < 0.3$)

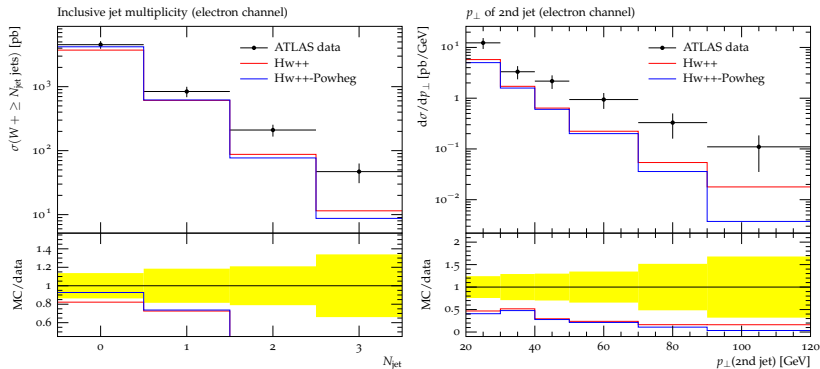


harder, more forward ($80 < p_T/\text{GeV} < 110; 1.2 < |y| < 2.1$)



Limits of parton shower

W + jets, LHC 7 TeV.



Higher jets not covered by parton shower only \rightarrow matching.

- ▶ Hadronization
 - ▶ Confinement
 - ▶ Flux tube model
 - ▶ Independent fragmentation
 - ▶ Lund string model
 - ▶ Cluster hadronization model
- ▶ Hadronic decays
- ▶ Comparison to data

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