

Measurement of the inclusive $t\bar{t}$ cross section and search for additional scalars in $t\bar{t}$ final states at the CMS experiment

PhD Disputation at Hamburg University
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13.10.2025



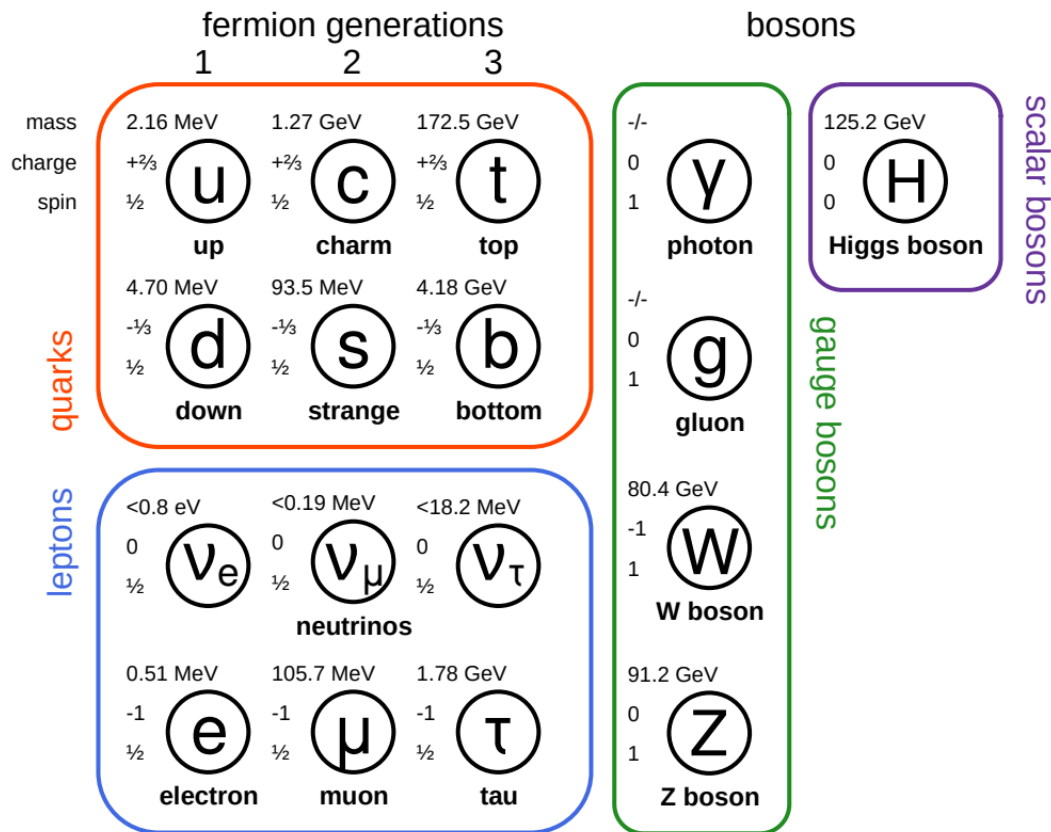
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→ we investigate Nature on the smallest known scales

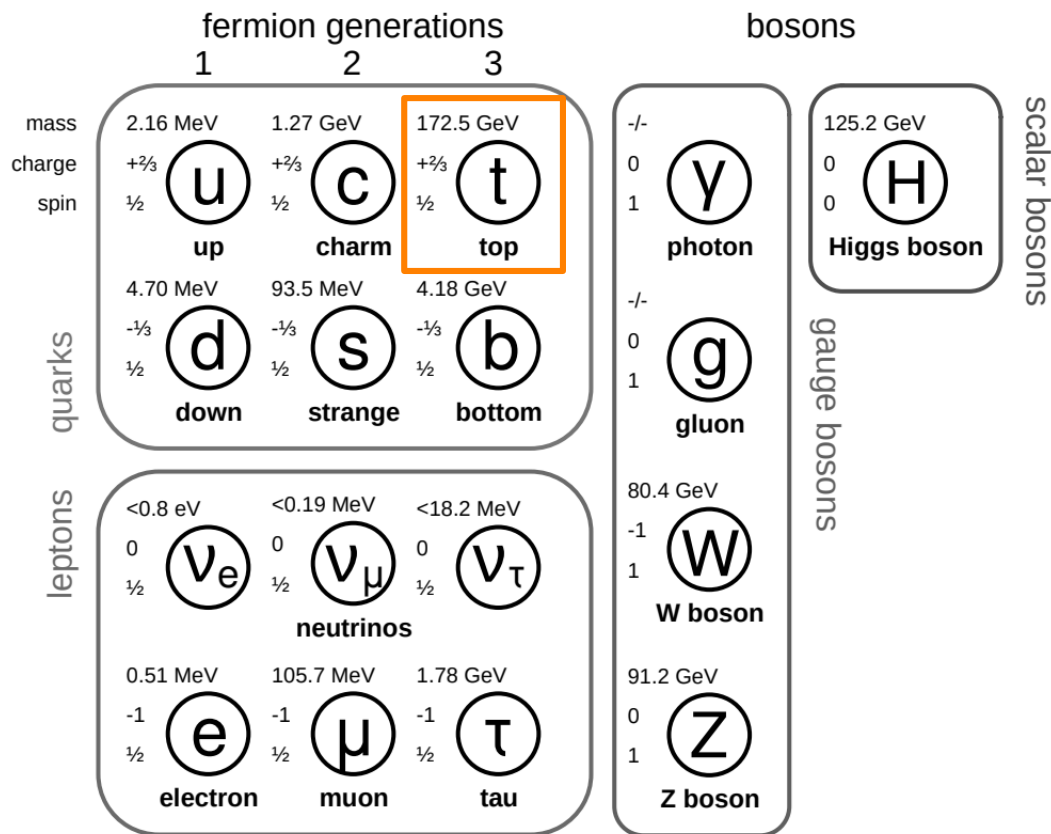
The Standard Model

- Relativistic Quantum Field Theory
- Particles = fundamental excitations of fields
- Most successful theory of our universe at small scales
- But: known to be incomplete!
 - neutrino masses, dark matter, ...



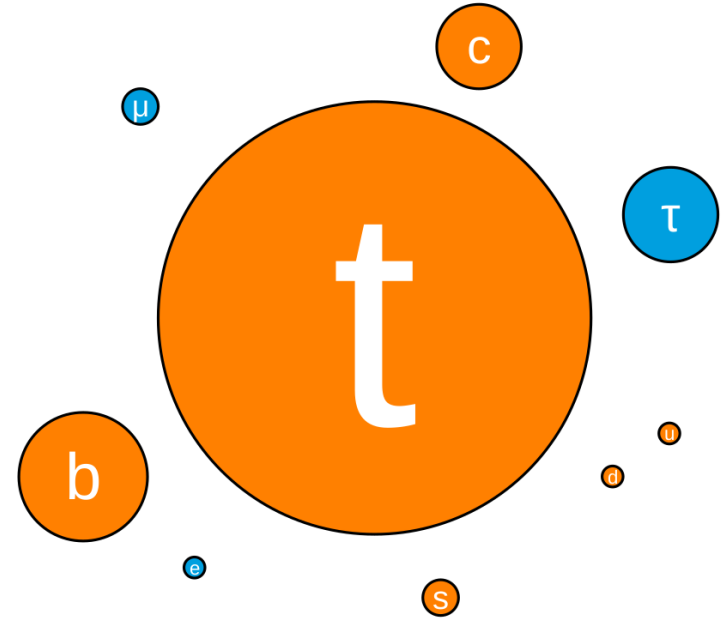
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The Top Quark

- The top quark is special:
 - Heaviest elementary particle
 - Largest coupling to Higgs
 - ➔ gateway to new physics:
e.g. additional Higgs sectors
 - Short lifetime $\sim 10^{-25}$ s
 - No hadronization - bare quark
 - ➔ access to spin properties
- Understanding the top is crucial for SM and BSM physics!



What do we do in High Energy Physics?

- we investigate Nature on the smallest known scales
 - **test** our theory: the Standard Model
 - **search** for physics beyond the SM

How do we do this?

- collider experiments

The Large Hadron Collider

- The largest particle accelerator built (so far)
 - 27 km circumference
 - 8 T magnetic fields from superconducting magnets
 - 3×10^{14} protons in the machine simultaneously
- Four large experiments: ATLAS, CMS, ALICE, LHCb

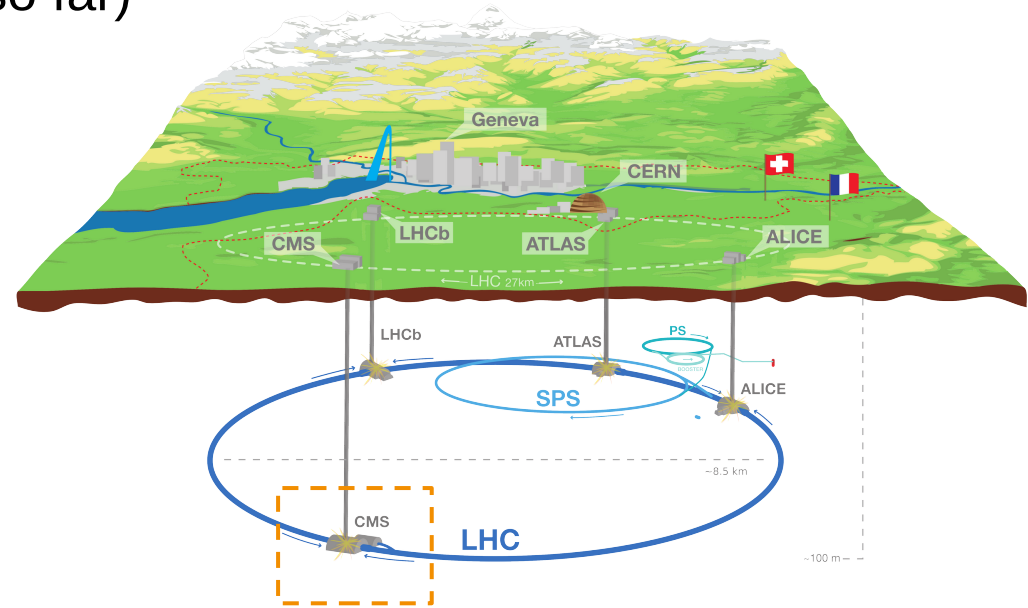
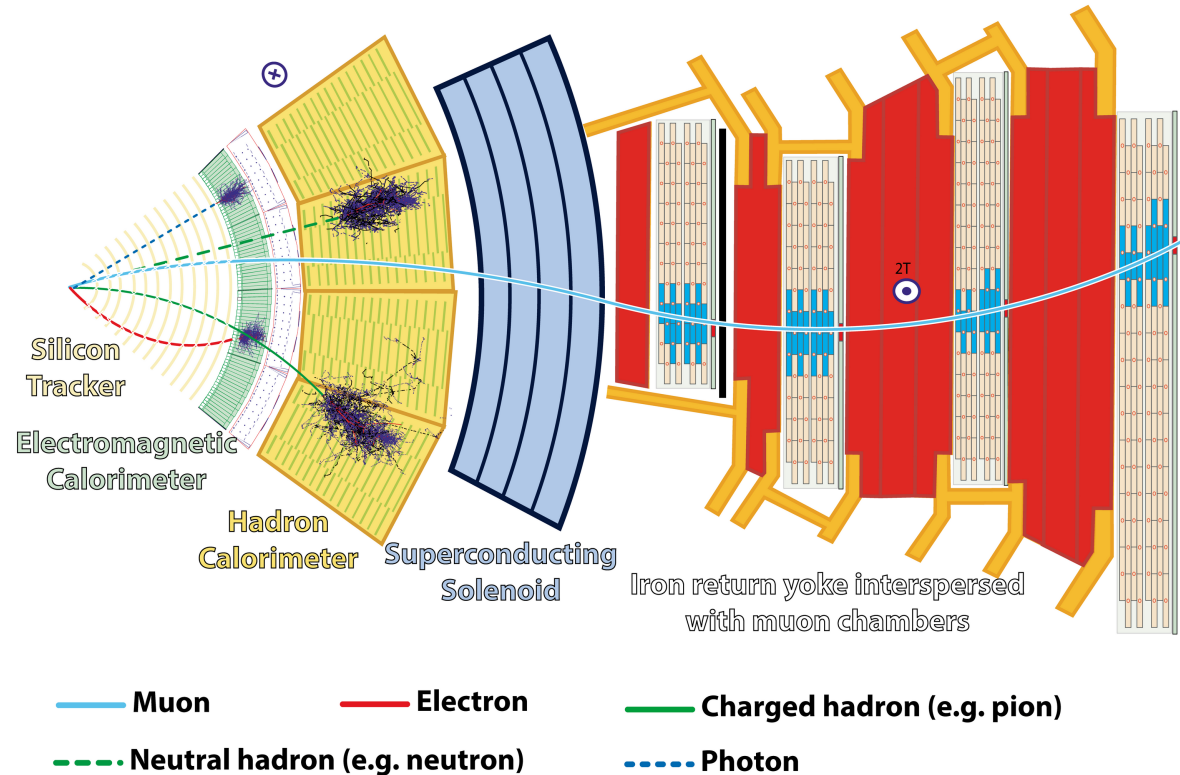


Image source: <https://cds.cern.ch/record/1708849>

The CMS Experiment

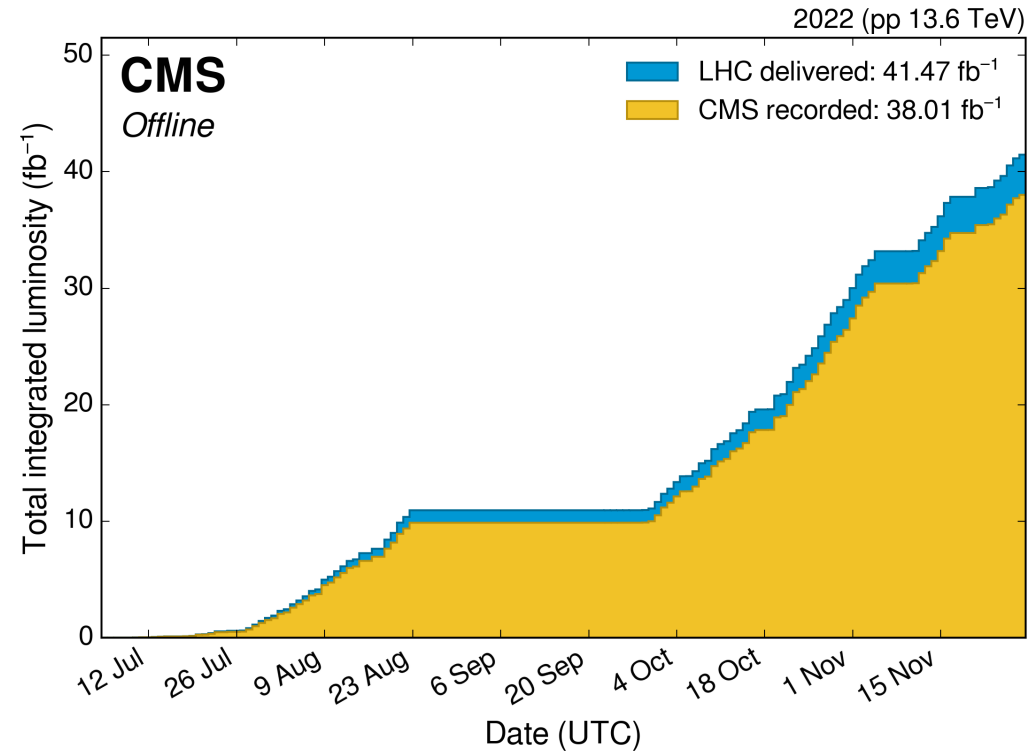
Image source: <https://cds.cern.ch/record/2120661>

- General-purpose hermetic detector with a superconducting solenoid magnet
- Subdetectors for different particles and functions
- So far, three data-taking runs

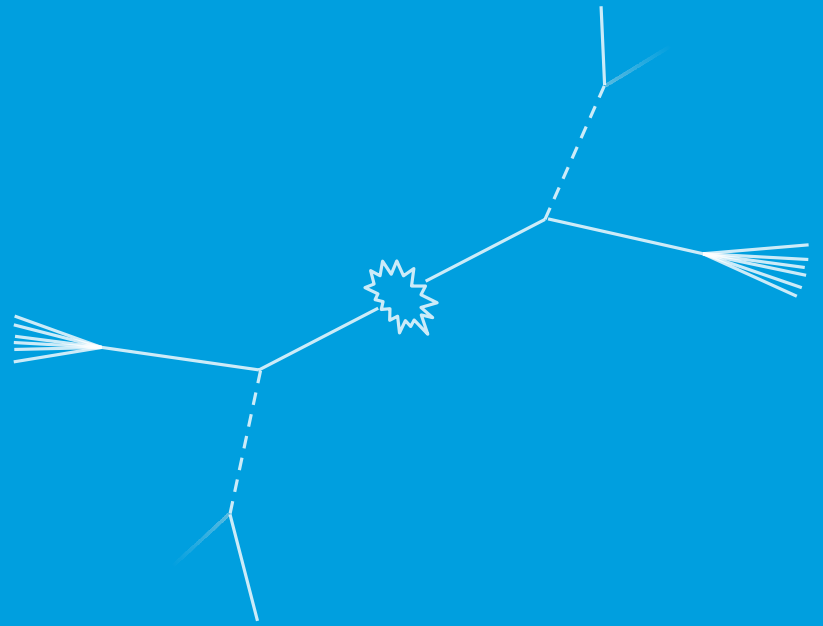


The Start of LHC Run 3

- July 2022: **start of LHC Run 3**
 - After three years of shutdown
 - Higher c.o.m. energy: 13.6 TeV
 - Upgraded detector components
 - Rapidly growing integrated luminosity
- Old calibrations no longer valid after long shutdown
- **Need for early measurement to demonstrate good data quality**



Part I



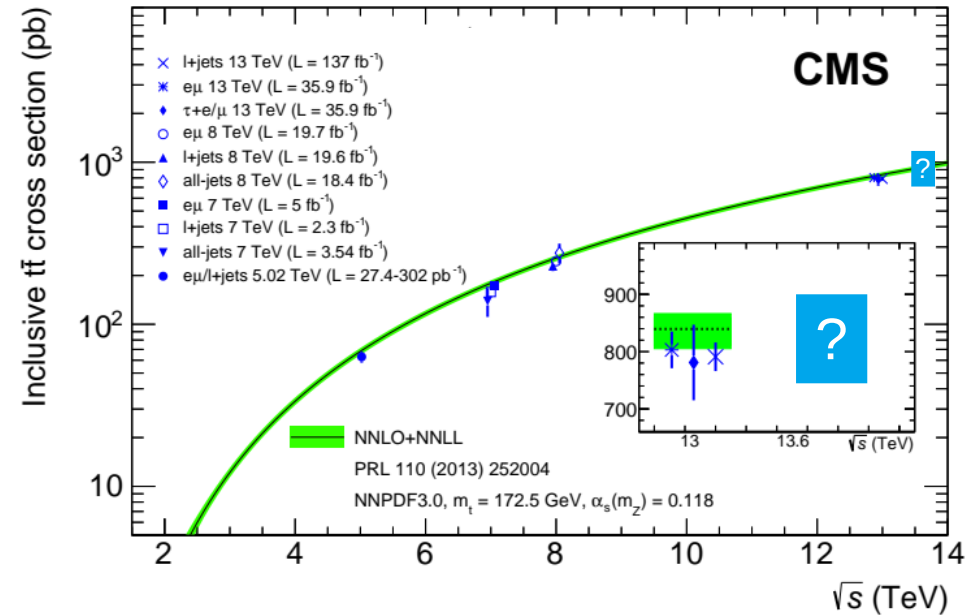
First measurement of the inclusive $t\bar{t}$
cross section in Run 3

Inclusive $t\bar{t}$ cross section

- top covers important objects at CMS: electrons, muons, (b-)jets
- Easiest top-related observable:
inclusive $t\bar{t}$ production cross section
 - rises by $\sim 10\%$ from 13 \rightarrow 13.6 TeV

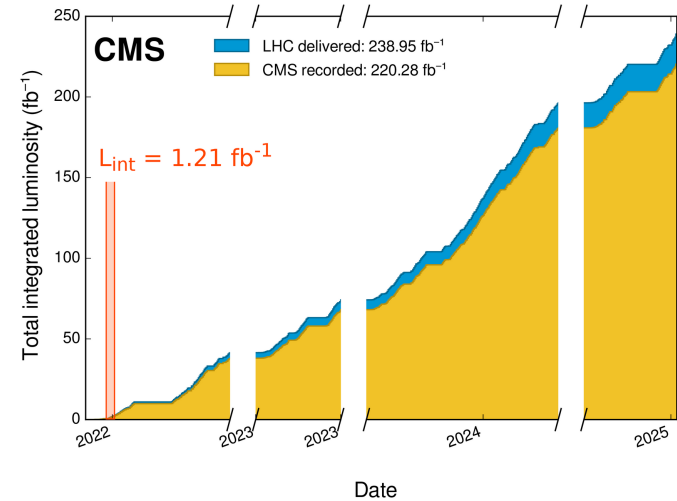
➔ Measure inclusive $t\bar{t}$ cross section in early Run 3

- Extremely short timeline:
 - Setup analysis & test on Run 2 data before Run 3 start
 - Measurement itself performed in only 2 months



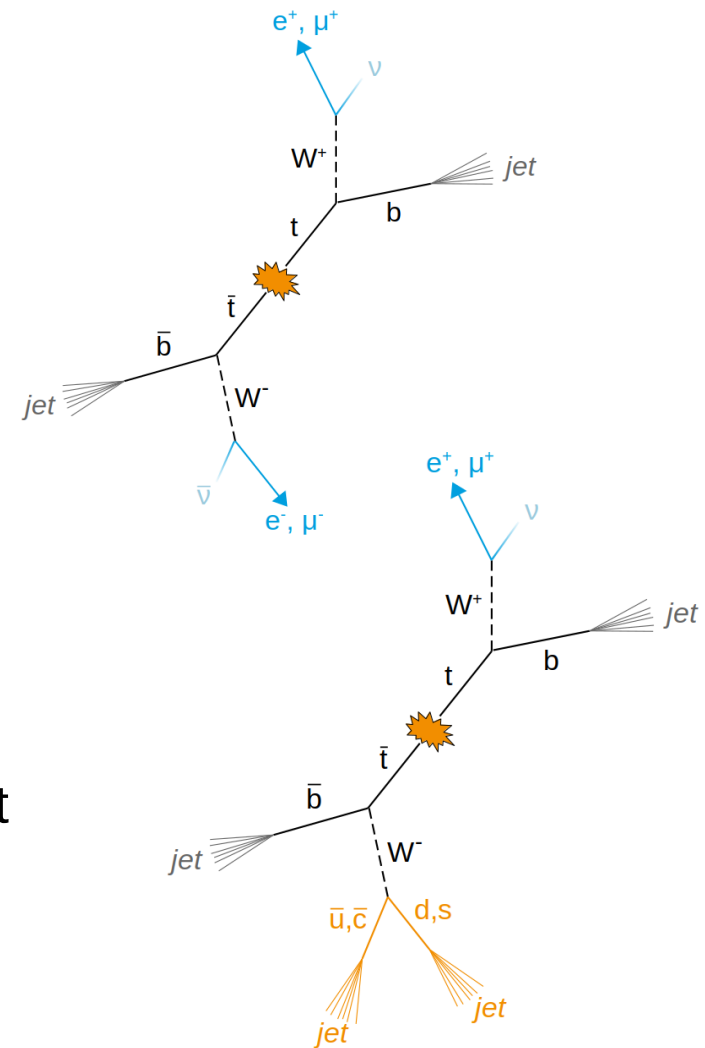
Analysis strategy

- Strategy designed from scratch, targeting early analysis
- Use ~1 week of data from Jul-Aug 2022: 1.21 fb^{-1}
- Central calibrations & corrections not yet available
- need to estimate them or constrain them from data



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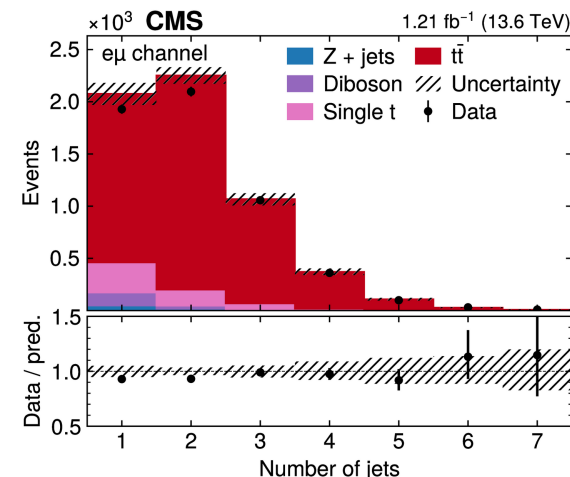
- Strategy designed from scratch, targeting early analysis
- Use ~1 week of data from Jul-Aug 2022: **1.21 fb⁻¹**
- Central calibrations & corrections not yet available
 - need to estimate them or constrain them from data
- For the first time: combine **dilepton** (ee, eμ, μμ) and **lepton+jets** (e, μ) decay channels in one likelihood fit
 - Dilepton: high purity, handle on b tagging
 - lepton+jets: high statistics
 - together: constrain lepton efficiencies



Channel definition

Dilepton channel: 2 opposite-sign leptons

- At least 1 jet
- For ee and $\mu\mu$: **large Z+jets BG**
 - require at least 1 b jet
 - cut away Z peak: $|m_{\ell\ell} - m_Z| > 15 \text{ GeV}$



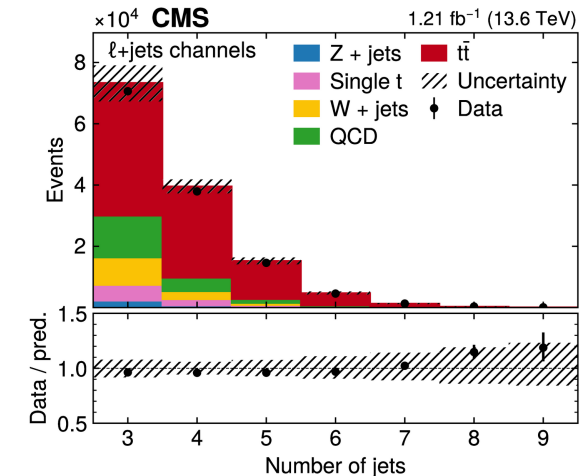
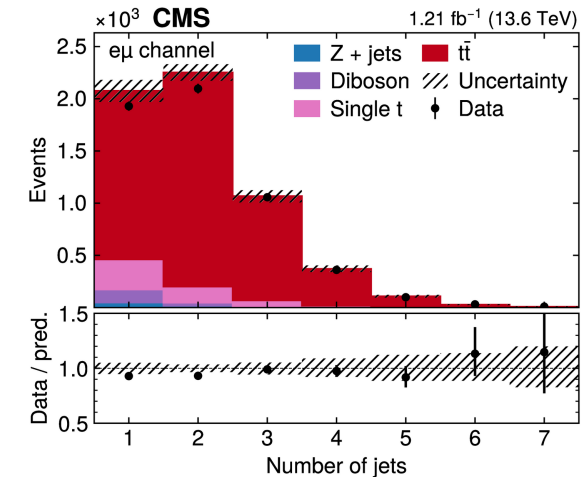
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Lepton+jets channel: exactly 1 lepton

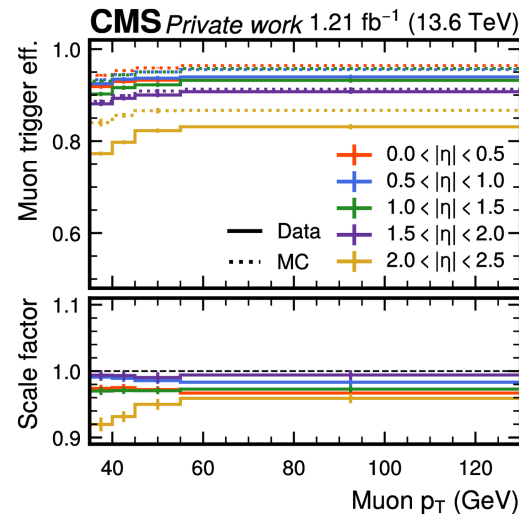
- At least 3 jets
- At least 1 b tagged jet



Corrections

Most corrections derived as part of this work:

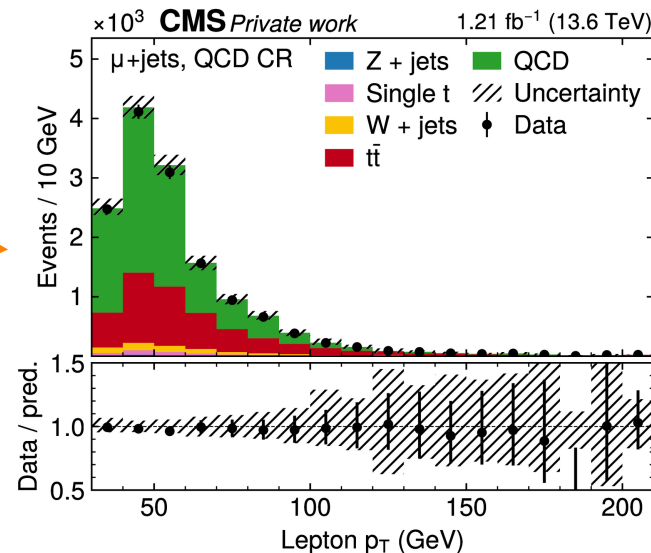
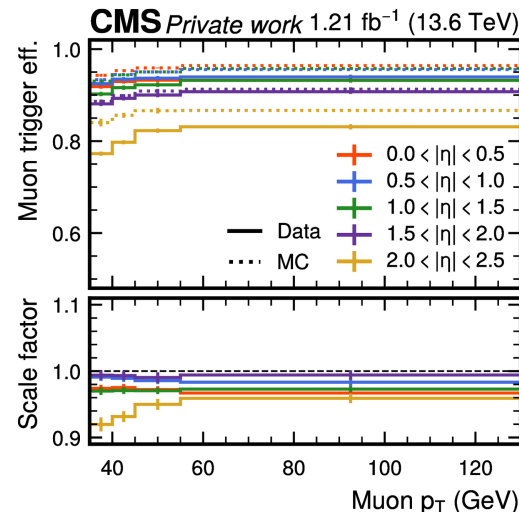
- **b tagging efficiencies** – in situ from data in the likelihood fit
- **Trigger efficiencies** – derived with tag&probe method
- **Pileup corrections** – data-driven reweighting to data



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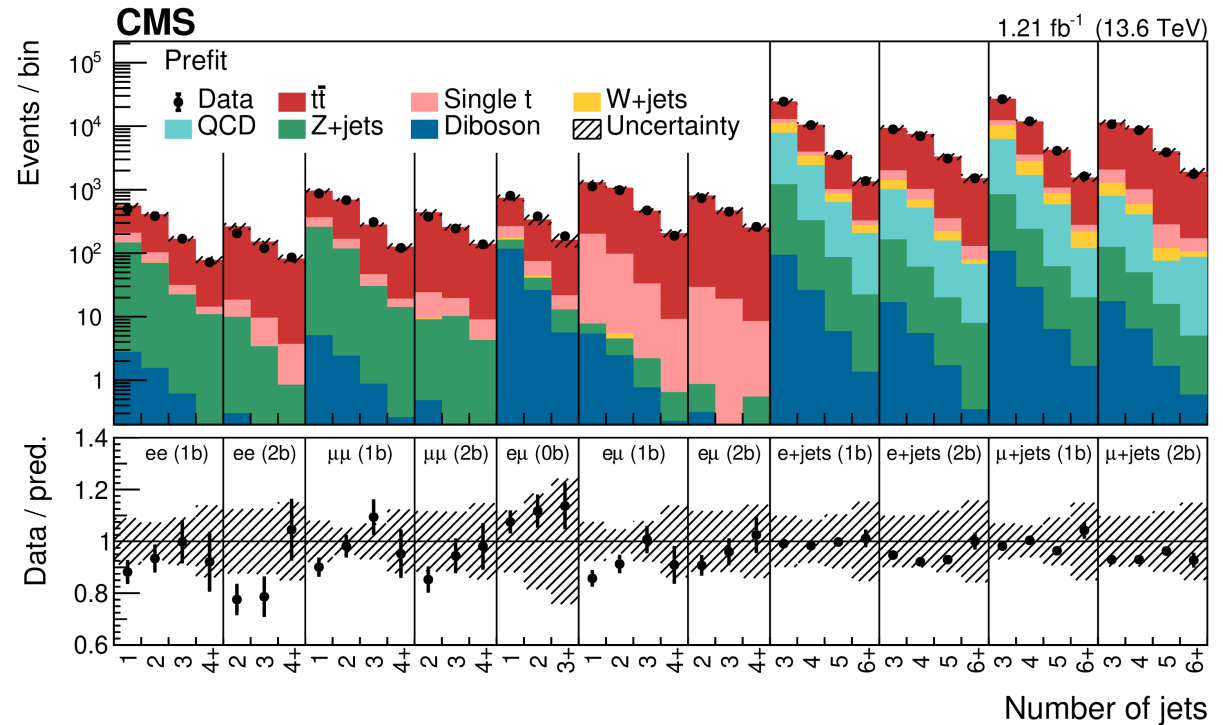
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- **Pileup corrections** – data-driven reweighting to data
- **QCD multijet BG with fake/non-prompt lepton** – data-driven using ABCD method in lepton isolation sideband
- **Normalization of Z+jets background** – from data close to Z peak



Results: prefit

- Channels defined by
 - Lepton content
- Further separated by
 - b jet content:
0, 1, or 2 b jets
- Coarsely binned in
 - Number of jets

Note: no b jet SF applied, no b tagging uncertainties

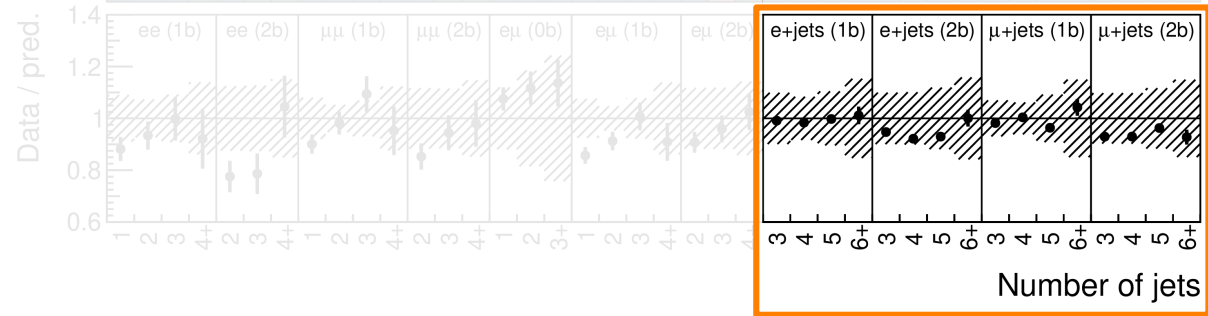


Results: prefit

- Difference in b tagging efficiency ϵ_b in data and MC
- Sensitivity through categorization in number of b tags:
 - Events with 2 b tags $\sim \epsilon_b^2$
 - Events with 1 b tag $\sim \epsilon_b (1 - \epsilon_b)$
- Simultaneous determination of efficiency and cross section!

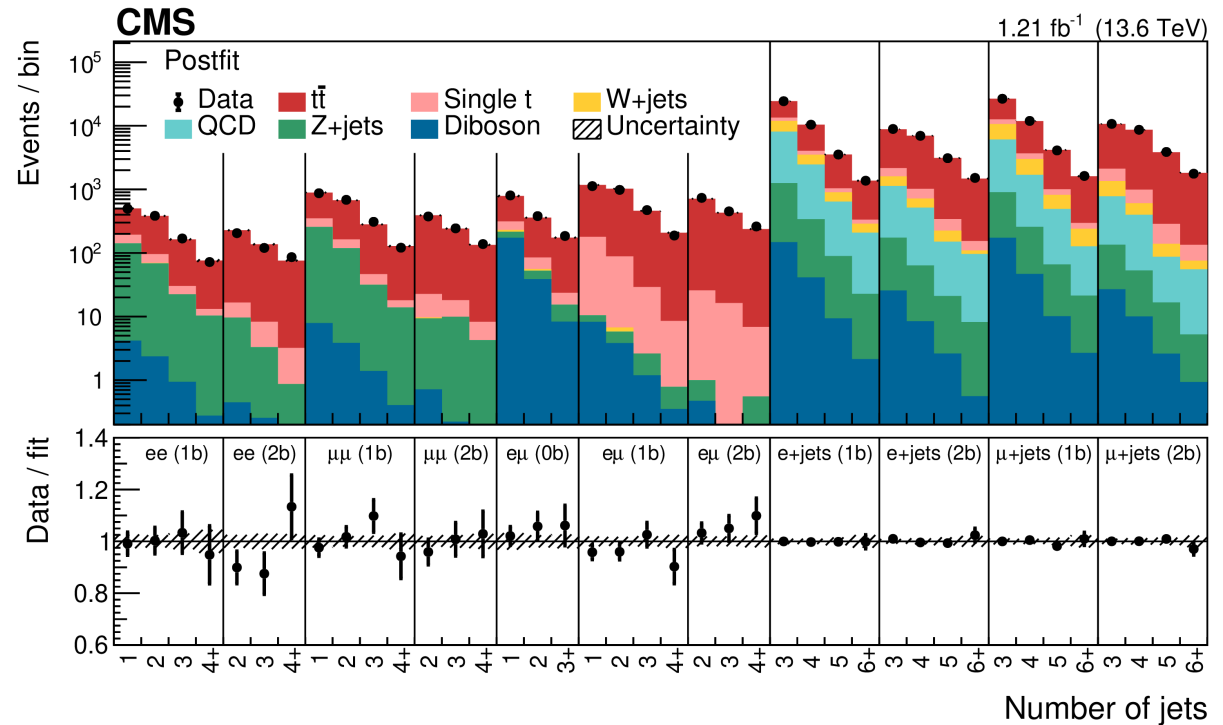
Number of jets

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Results: postfit

- **Profile likelihood fit** with nuisance parameters
- b tag efficiency freely floating
- Postfit agreement is improved significantly!



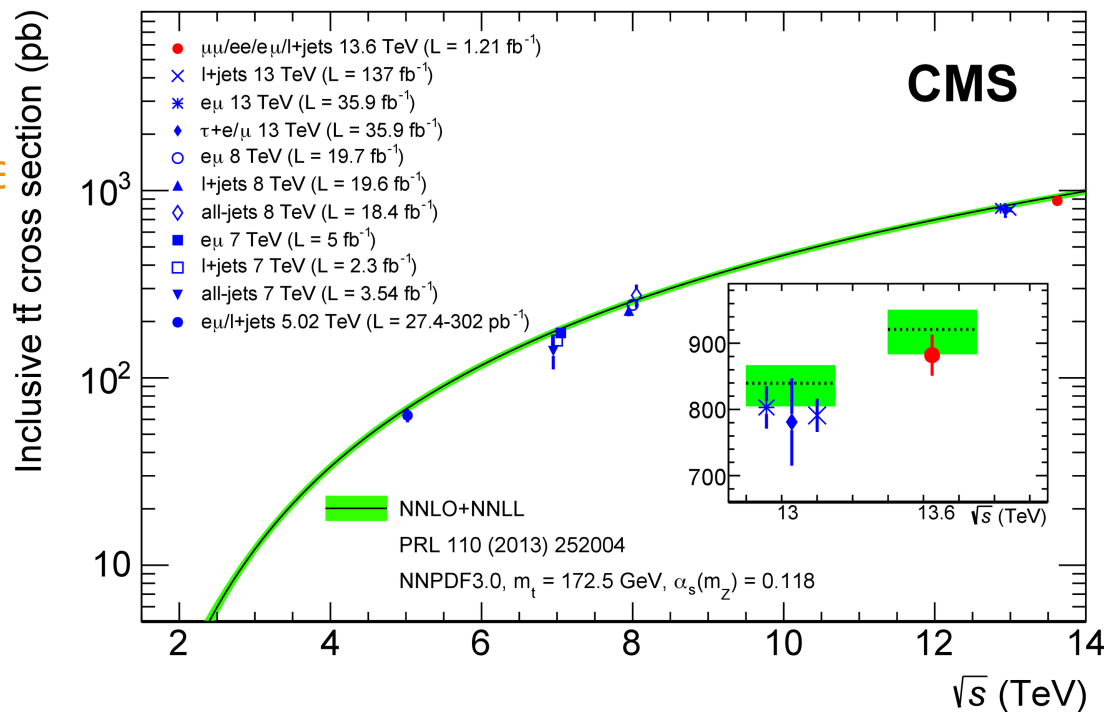
Results: cross section

JHEP 08 (2023) 204

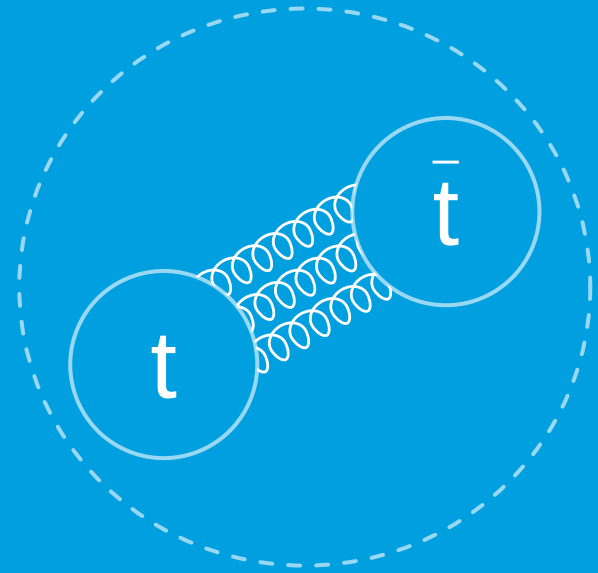
$$\sigma_{t\bar{t}} = 882 \pm 23(\text{stat} + \text{syst}) \pm 20(\text{lumi})\text{pb}$$

- ~ 3.5% total uncertainty
- Competetive with previous measurements despite early timeline
- Dominant uncertainty sources:
 - Luminosity
 - Lepton ID efficiency
 - b tag efficiency
- In agreement with theory:

$$\sigma_{t\bar{t}}^{\text{pred}} = 921_{-37}^{+29}\text{pb}$$



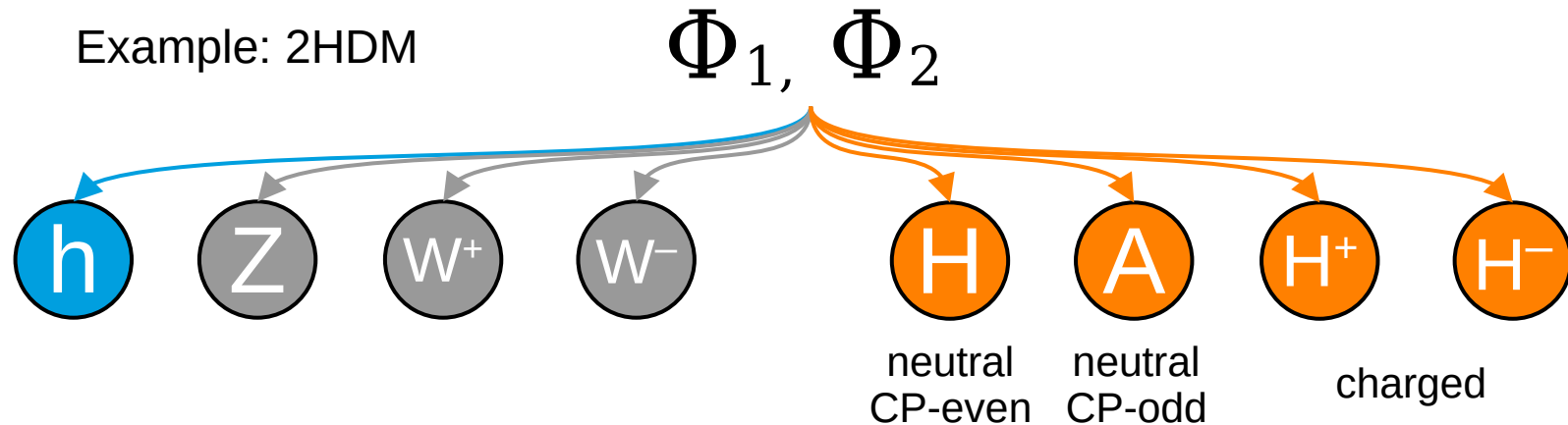
PART II



Search for (pseudo)scalar bosons in $t\bar{t}$ events and $t\bar{t}$ bound states

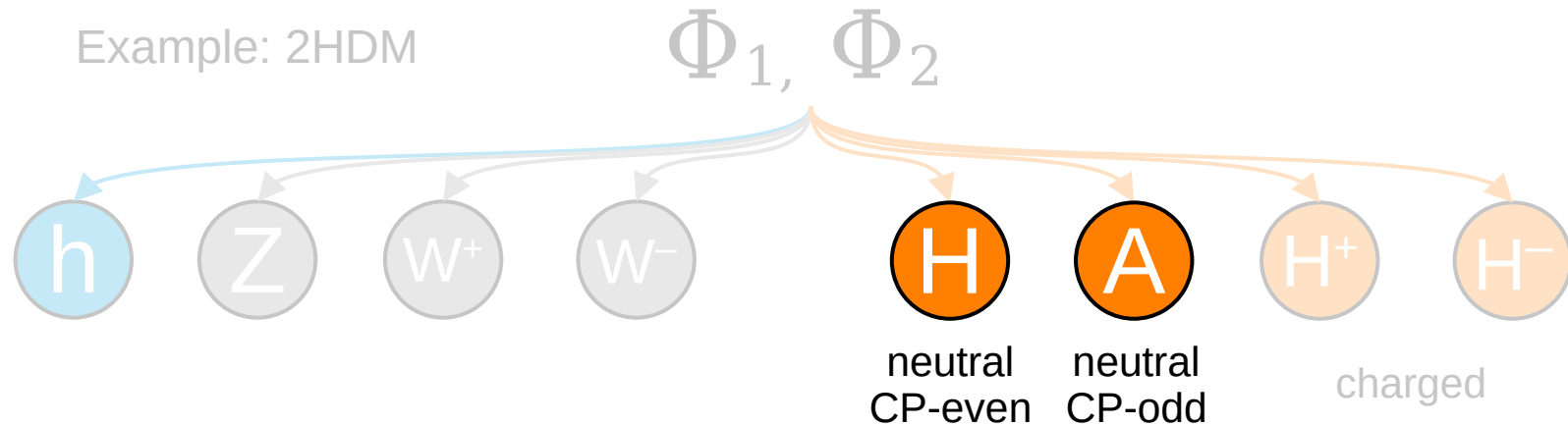
Additional Higgs bosons

- SM Higgs: complex SU(2) doublet $\phi \rightarrow$ one real scalar after symmetry breaking
- Many BSM models predict additional Higgs bosons
e.g. Two-Higgs Doublet Model (2HDM), Supersymmetry, ...



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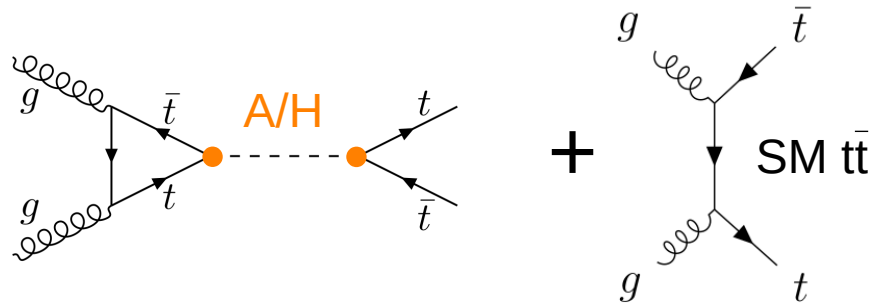
- If $m_{A/H} > 2m_t$: decay to $t\bar{t} \rightarrow$ search in $t\bar{t}$ final states

Generic parameterization

- Generic heavy **pseudoscalar (A)** or **scalar (H)** coupling solely to top quarks
- Production in gluon fusion via top quark loop

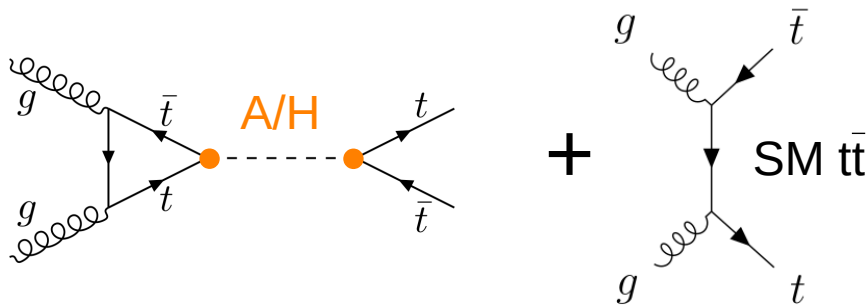
$$\mathcal{L}_A^{\text{int}} = ig_{A t \bar{t}} \frac{m_t}{v} \bar{t} \gamma_5 t A$$

$$\mathcal{L}_H^{\text{int}} = -g_{H t \bar{t}} \frac{m_t}{v} \bar{t} t H$$



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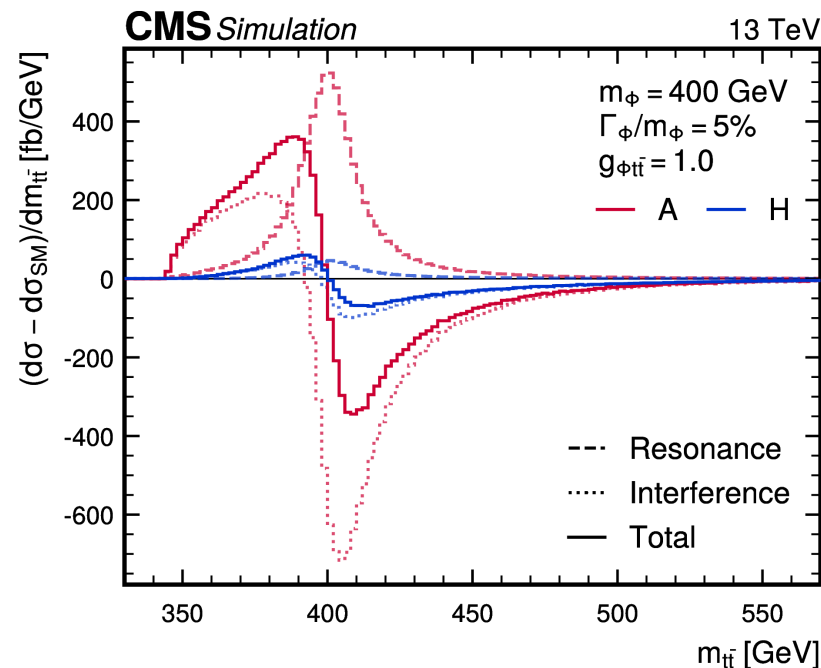
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- Same final state as SM $t\bar{t} \rightarrow$ interference
 \rightarrow **peak-dip structure** in invariant mass $m_{t\bar{t}}$

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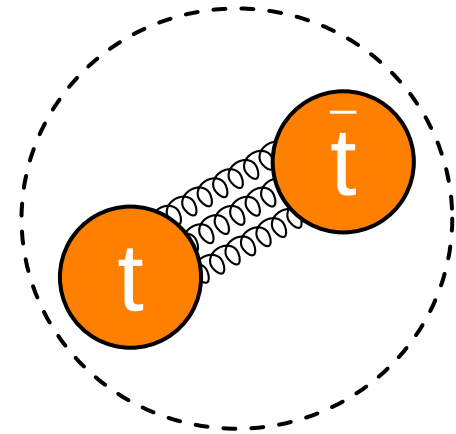


$t\bar{t}$ bound states?

- All quarks other than the top are known to form bound states (quarkonia)
- Simple estimate by analogy to hydrogen atom with QCD potential:

binding energy $E_b = -\frac{1}{2} \frac{m_q}{2} \left(C^{[1]} \alpha_S \right)^2$ with $C^{[1]} = \frac{4}{3}$ for color-singlet

for $t\bar{t}$: $E_b \approx -2 \text{ GeV}$ system should be bound!



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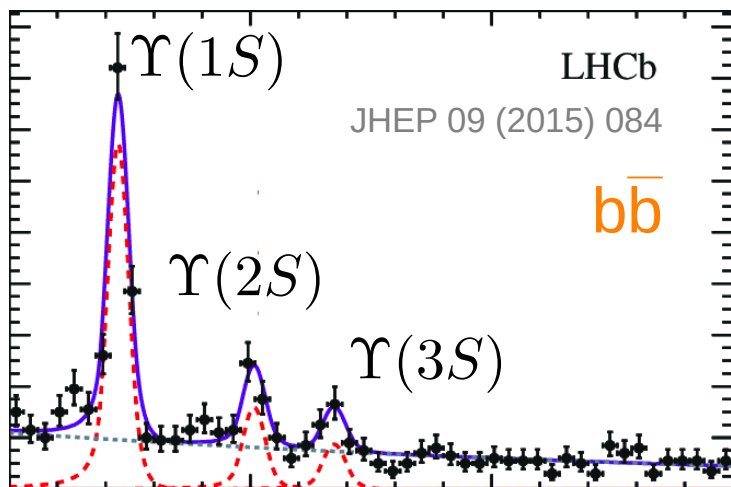
- However: **lifetime of the bound state shorter than classical revolution time**

$$\tau_{t\bar{t}} = \frac{1}{2} \tau_t = 2.4 \times 10^{-25} \text{ s} < \tau_{\text{rev}} = 1 \times 10^{-24} \text{ s}$$

- Fraction of $t\bar{t}$ systems that live long enough: $\approx 1\%$ small effect!

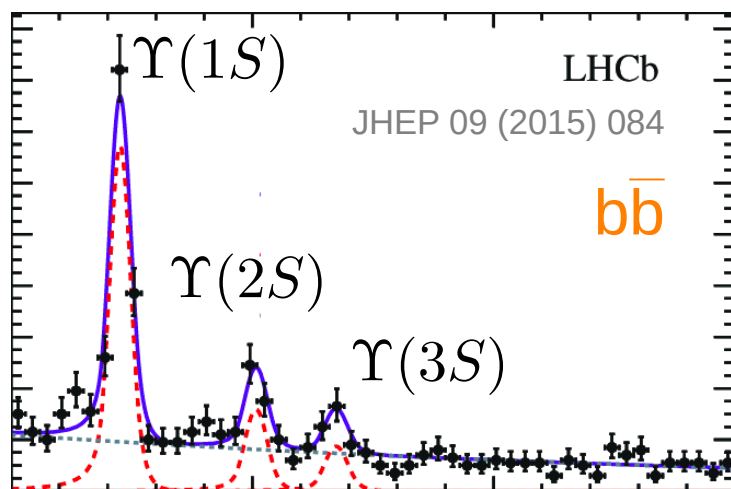
$t\bar{t}$ bound states: NRQCD

- more quantitative: non-relativistic QCD (NRQCD)
 - factorize hard scattering and long-distance effects (from Schrödinger equation)

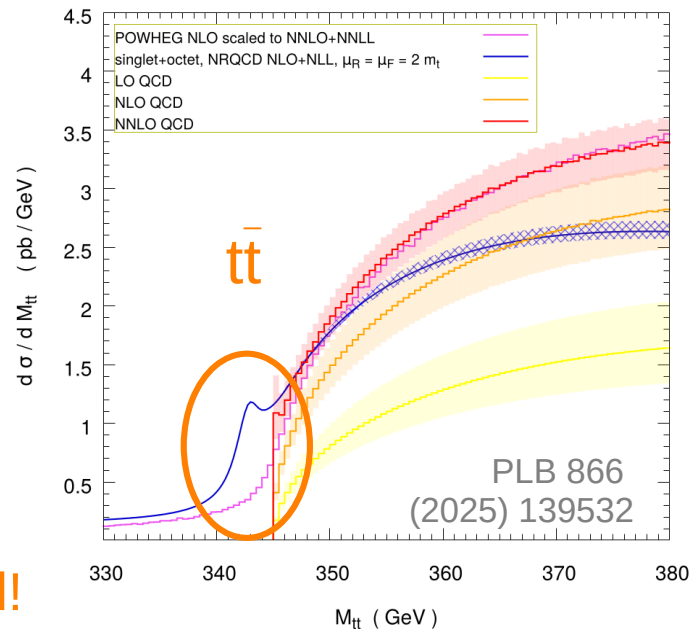


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finite width
→

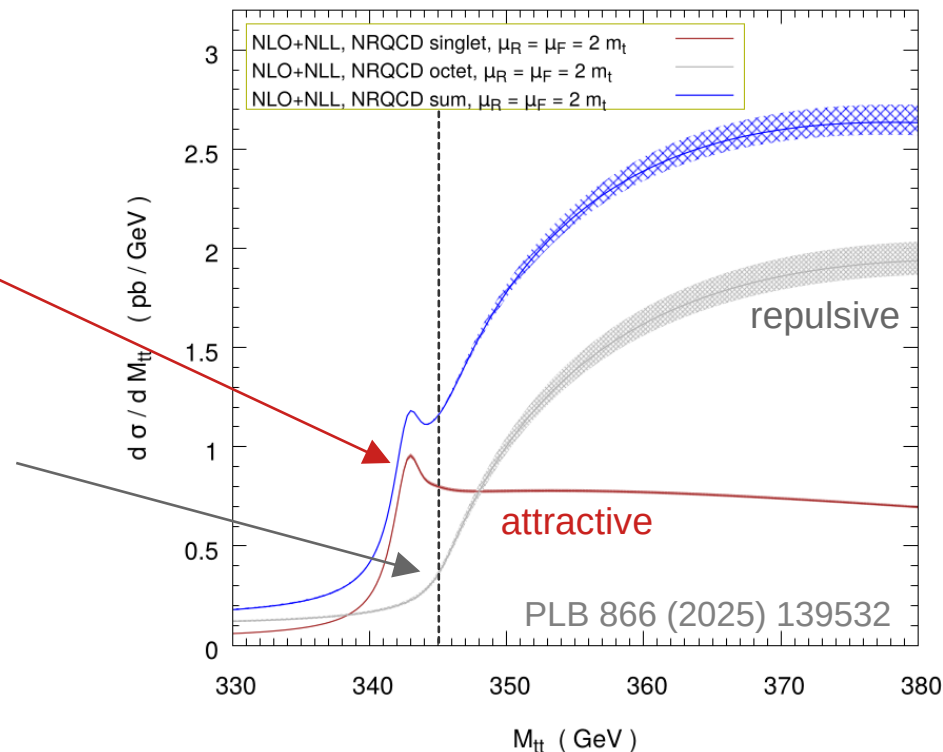


one broad peak around the $t\bar{t}$ threshold!
"quasi-bound state"

$t\bar{t}$ bound states: NRQCD

See e.g.
PRD 110 (2024) 5, 054032
JHEP 03 (2024) 099
PRD 104 (2021) 3, 034023
etc.

- Results of NRQCD calculations:
 - Color-singlet ($^1S_0^{[1]}$) - **attractive**
→ Peak below the $t\bar{t}$ threshold
CP-odd / pseudoscalar spin state!
 - Color-octet ($^1S_0^{[8]}$ or $^3S_1^{[8]}$) - **repulsive**
→ Suppressed below the $t\bar{t}$ threshold



$t\bar{t}$ bound states: modeling in MC

- Cannot use NRQCD to produce events for analysis...
- Use **simplified model for MC simulation: η_t** (PRD 104 (2021) 034023, JHEP 03 (2024) 099)

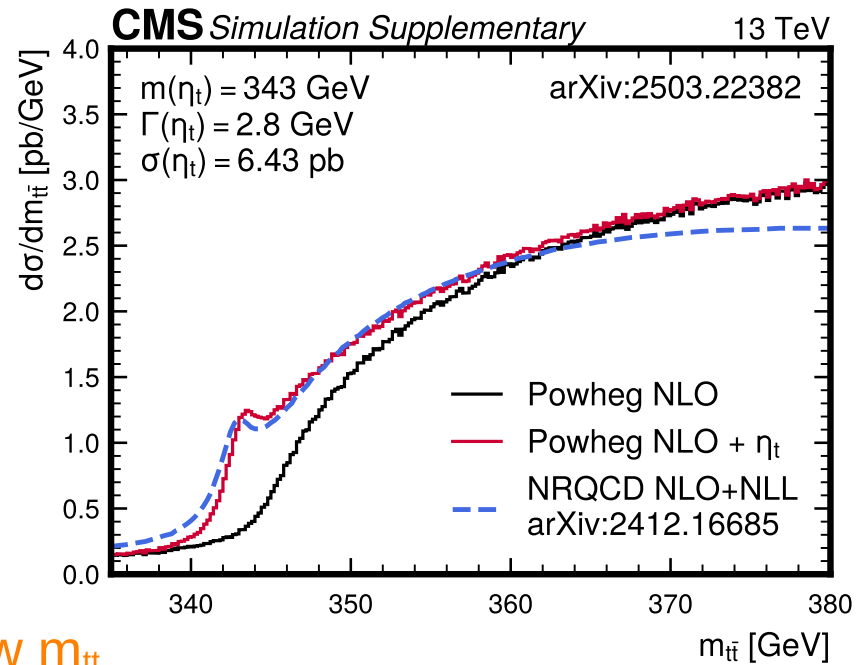
- Generic spin-0, color-singlet state η_t
- Couplings to gluons and tops (pseudoscalar)
- Fit mass and width from NRQCD:

$$m_{\eta_t} - 2m_t = -2 \text{ GeV} \quad \Rightarrow \quad m_{\eta_t} = 343 \text{ GeV}$$

$$\Gamma_{\eta_t} \approx 2\Gamma_t = 2.8 \text{ GeV}$$

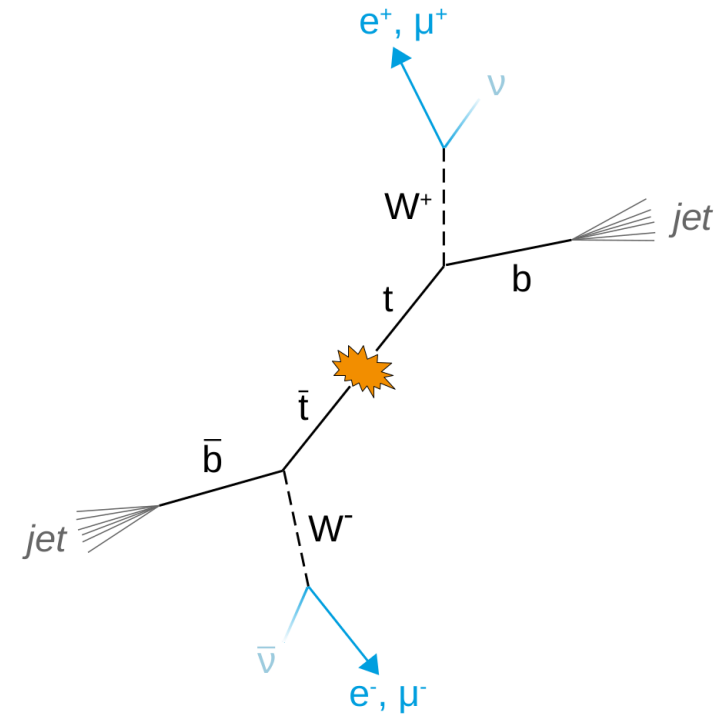
- Stack on top of ordinary pQCD $t\bar{t}$ simulation!
Fits well to NRQCD prediction

- Result: **pseudoscalar-like enhancement at low $m_{t\bar{t}}$**



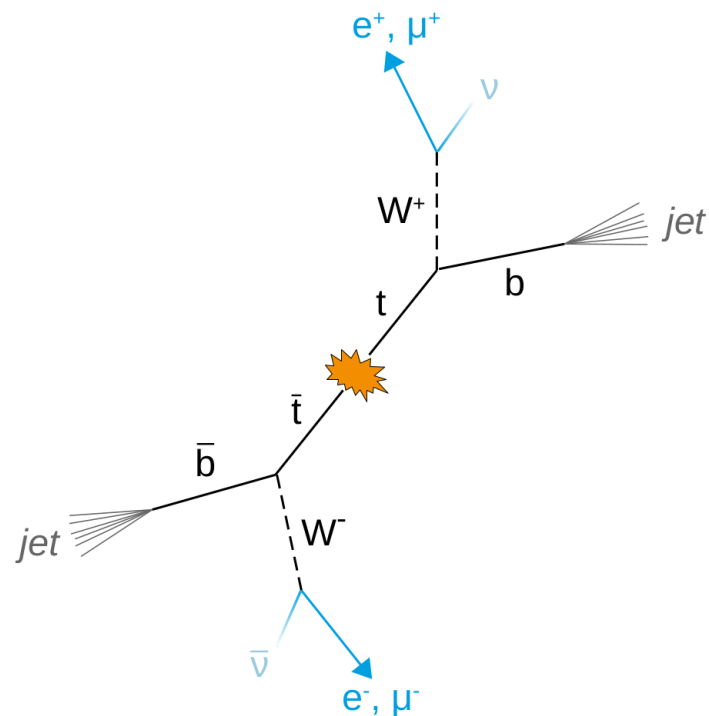
Analysis strategy

- Consider **dilepton final state of $t\bar{t}$**
 - **Full Run 2 dataset:** 138 fb^{-1}
 - Selection: 2 leptons, ≥ 2 jets, ≥ 1 b tag
in $ee/\mu\mu$: reject Z peak, require $p_{\text{T}}^{\text{miss}} > 40 \text{ GeV}$



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 - Selection: 2 leptons, ≥ 2 jets, ≥ 1 b tag
in $ee/\mu\mu$: reject Z peak, require $p_{\text{T}}^{\text{miss}} > 40 \text{ GeV}$
- **Analytic reconstruction of $t\bar{t}$ system:**
 - Assign b jets using likelihood based on $m_{\ell b}$
 - Assumptions: all $p_{\text{T}}^{\text{miss}}$ from $\nu\nu$, tops and Ws on-shell
 - Solve fourth-order polynomial equations
 - Finite detector resolution: repeat reconstruction 100 times with randomly smeared inputs, take weighted average
- Four-momenta of top and antitop
→ invariant $t\bar{t}$ mass, and...



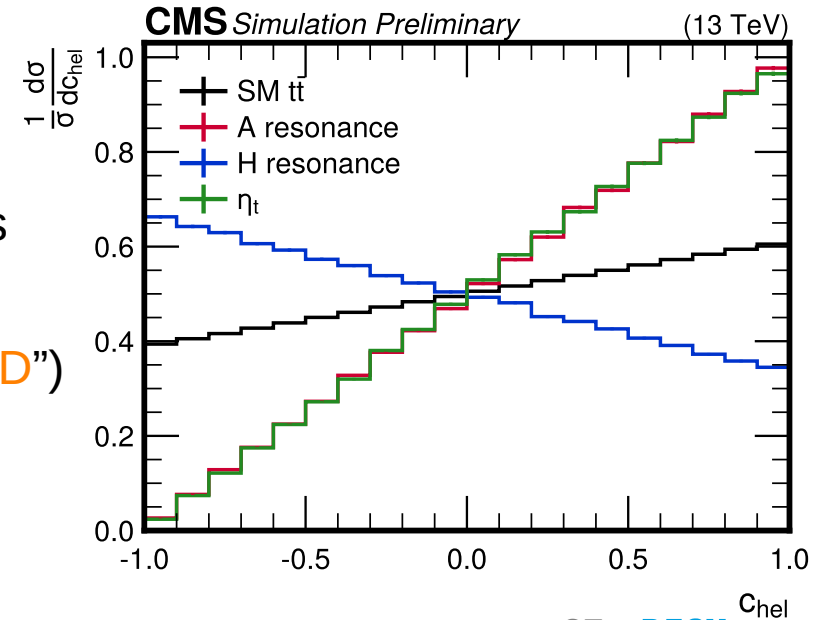
Spin correlation observables

- Both A/H and η_t predict $t\bar{t}$ production in a **pure $t\bar{t}$ spin state**:
 1S_0 or 3P_0 (from A / η_t resp. H)
- Top decays before hadronization \rightarrow transfer spin information to decay products
- Construct **spin correlation observables** from tops & leptons

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- Variable #1: **C_{hel}**
 - Boost leptons into rest frames of their parent tops \rightarrow Scalar product between directions of flight
 - Straight line with slope sensitive to $t\bar{t}$ spin state ("**D**")
 - Maximal for 1S_0 (from A / η_t) – **separates from SM**

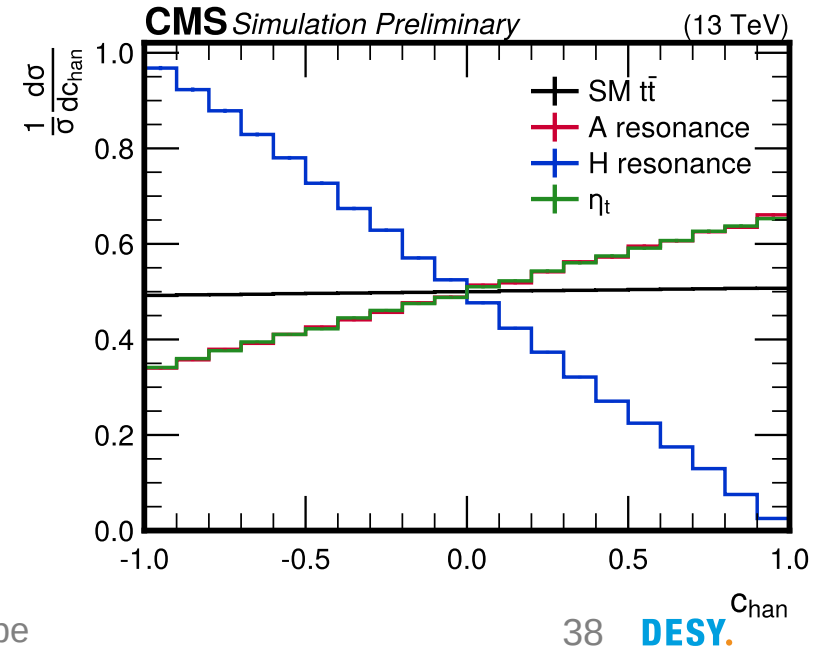


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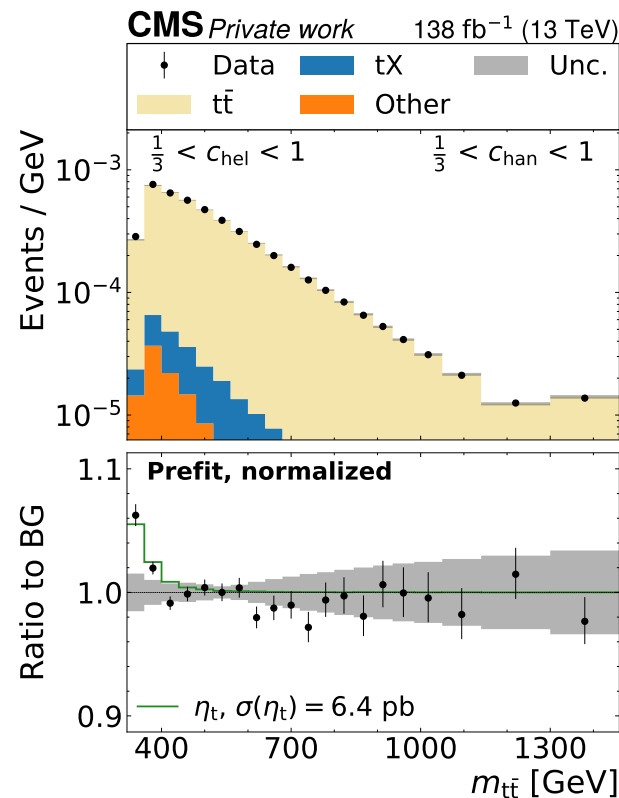
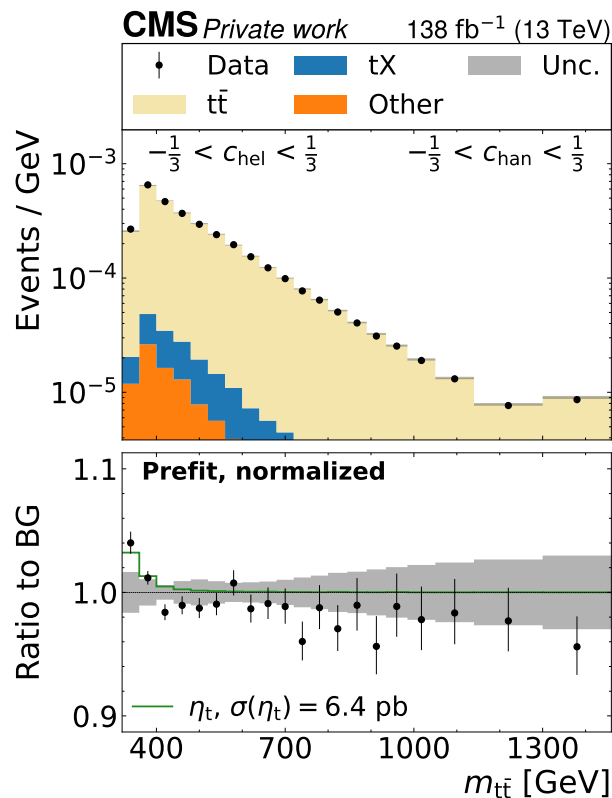
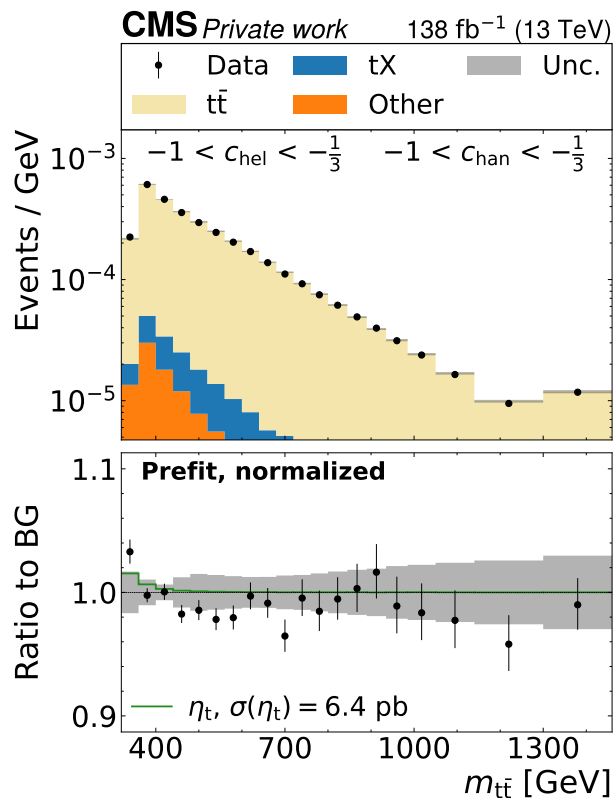
- Variable #2: **C_{chan}**
 - Similar as c_{hel} , separating scalars from SM
 - Maximally negative slope for 3P_0 state
 - Construct similarly from lepton momenta, with sign flip for component parallel to top momentum

\rightarrow 3 search variables: $m_{t\bar{t}} \times C_{\text{hel}} \times C_{\text{chan}}$



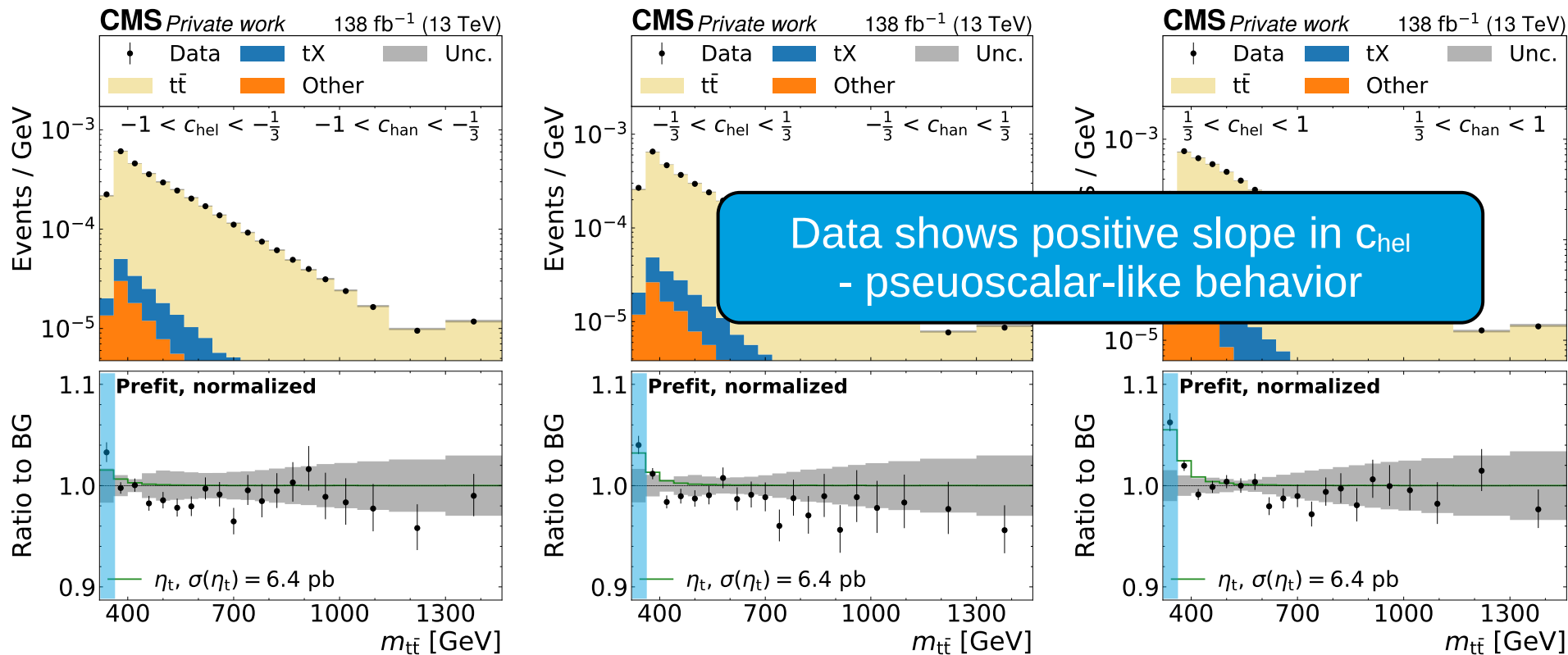
Prefit distributions

Differences between data and prediction observed in low $m_{t\bar{t}}$ bins!



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Cross section measurement

- Extract cross section using the simplified η_t color-singlet model
- “cross section” = difference to perturbative prediction

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$$\sigma(\eta_t) = 8.7 \pm 0.5(\text{stat}) \pm 1.0(\text{syst}) \text{ pb} = 8.7 \pm 1.1 \text{ pb}$$

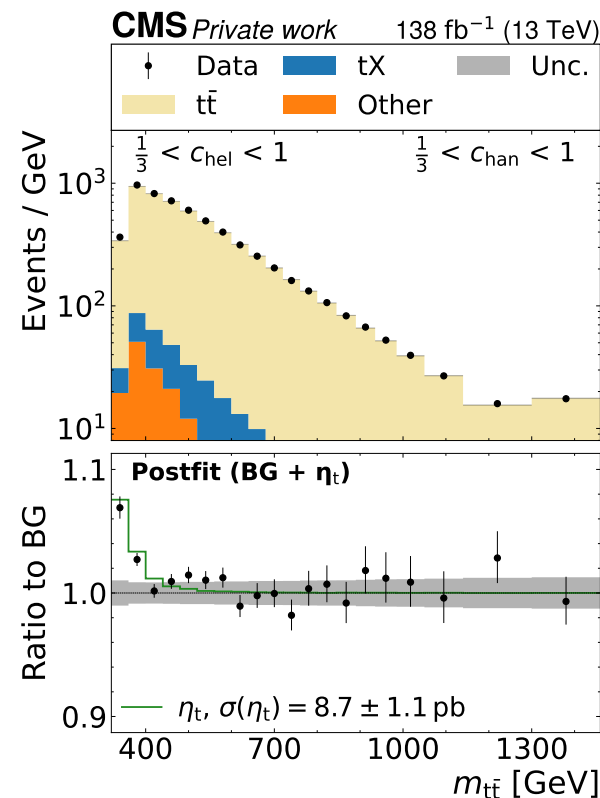
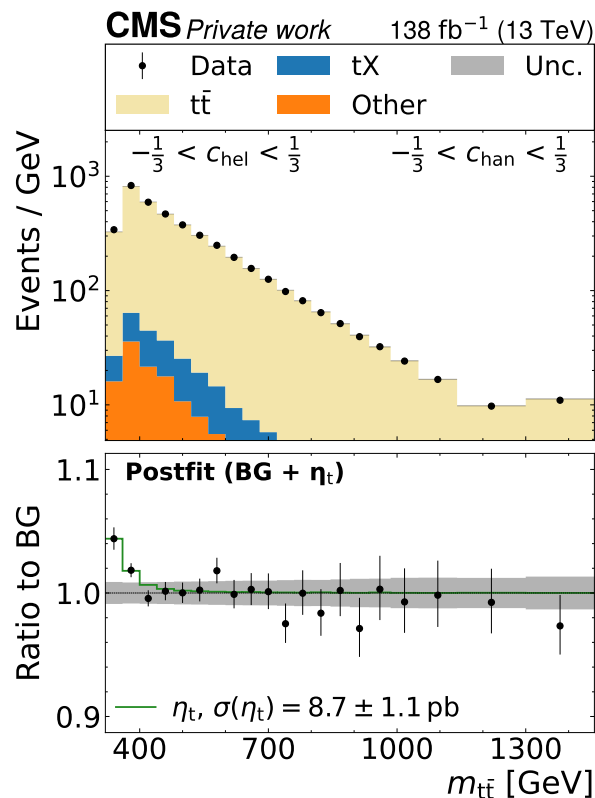
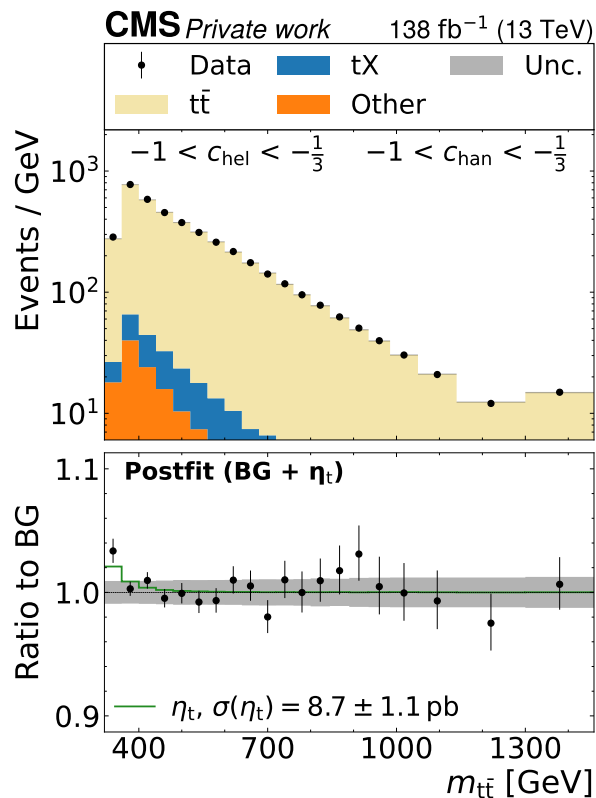
> 5σ significance!

- Main uncertainties: $t\bar{t}$ background modeling (PRD 104 (2021) 034023)
- Same order of magnitude as NRQCD estimate: $\sigma(\eta_t)^{\text{pred}} = 6.4 \text{ pb}$
- Confirmed by ATLAS in preliminary result (ATLAS-CONF-2025-008)

published as RoPP 88 087801 (2025)

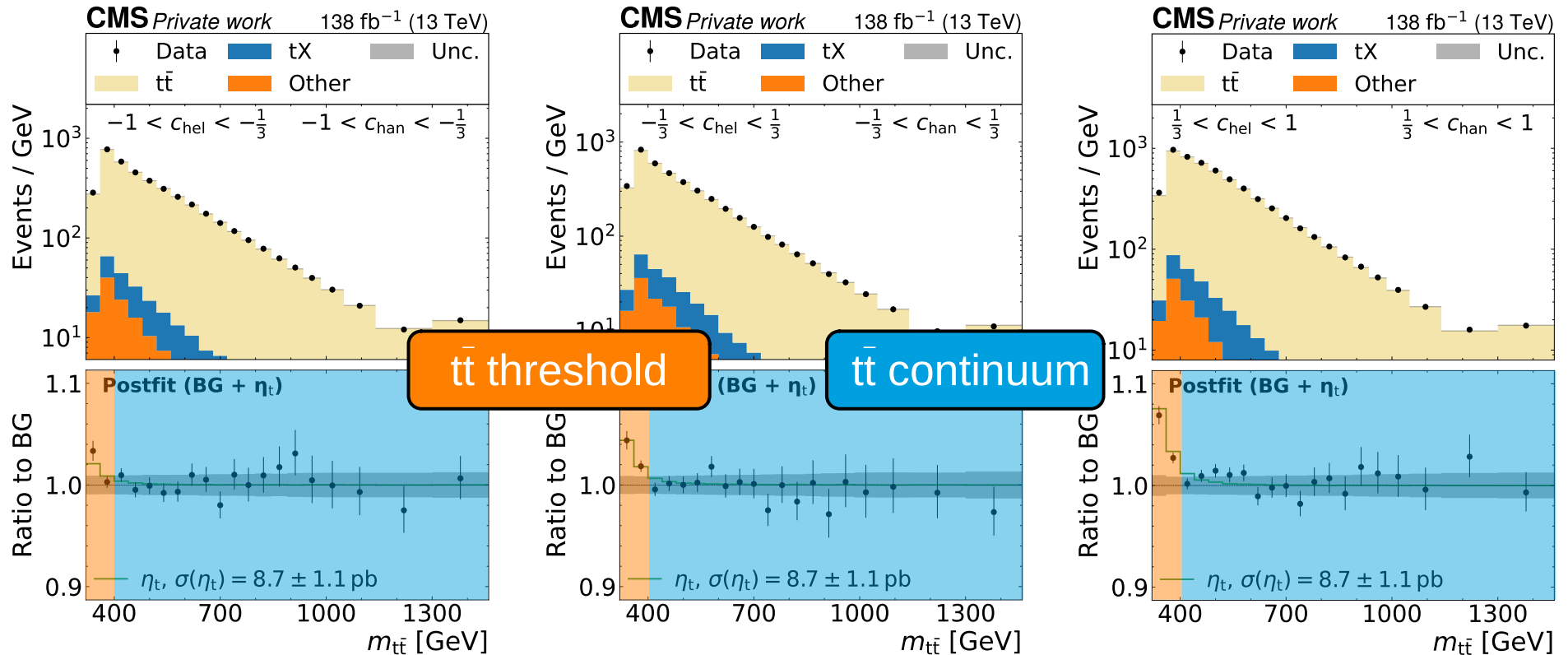
Postfit distributions: η_t

η_t model describes the data well after the fit

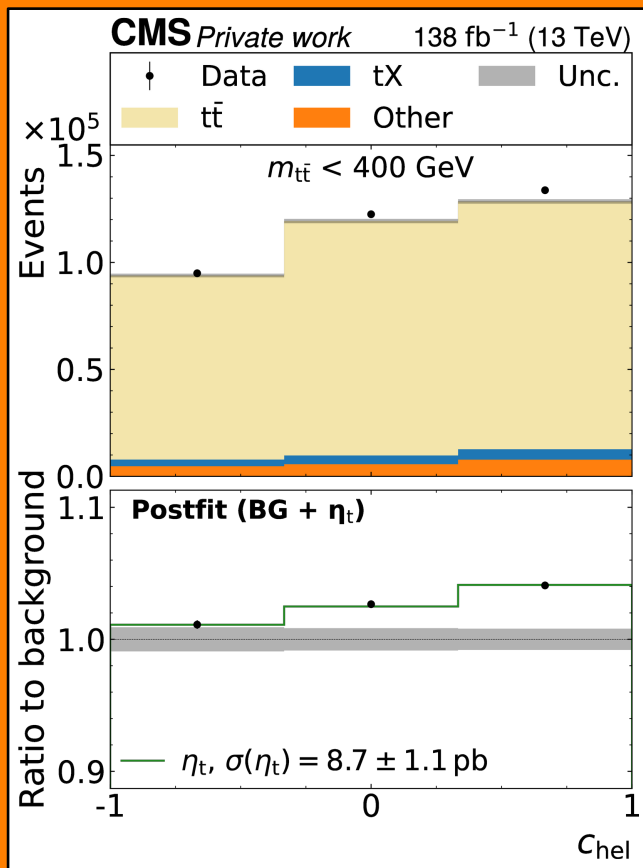


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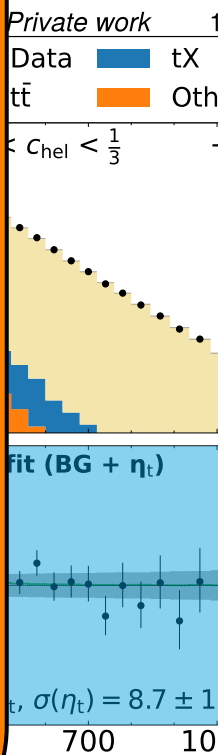
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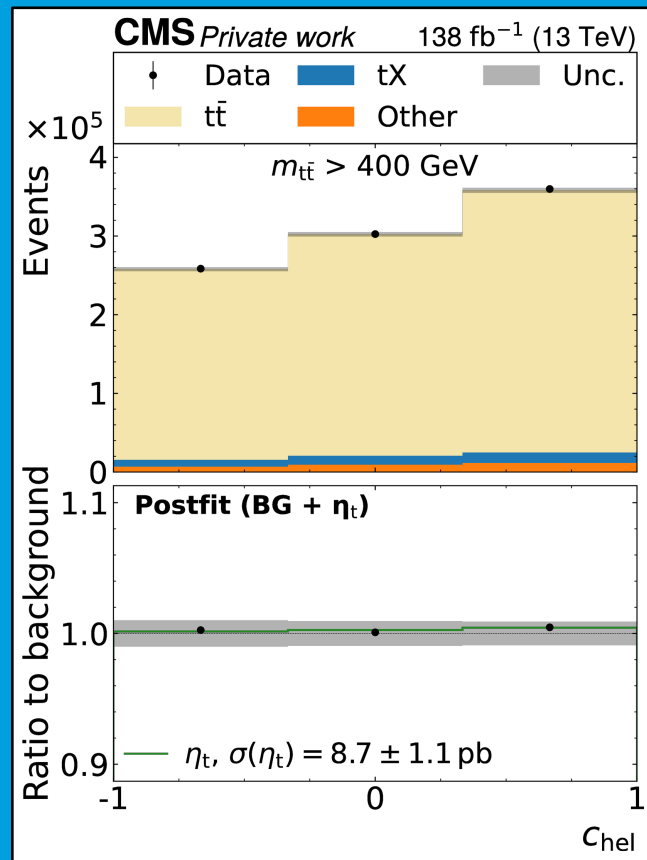
$t\bar{t}$ threshold:
Data shows slope in C_{hel}



ns: $m_{t\bar{t}}$
is the data

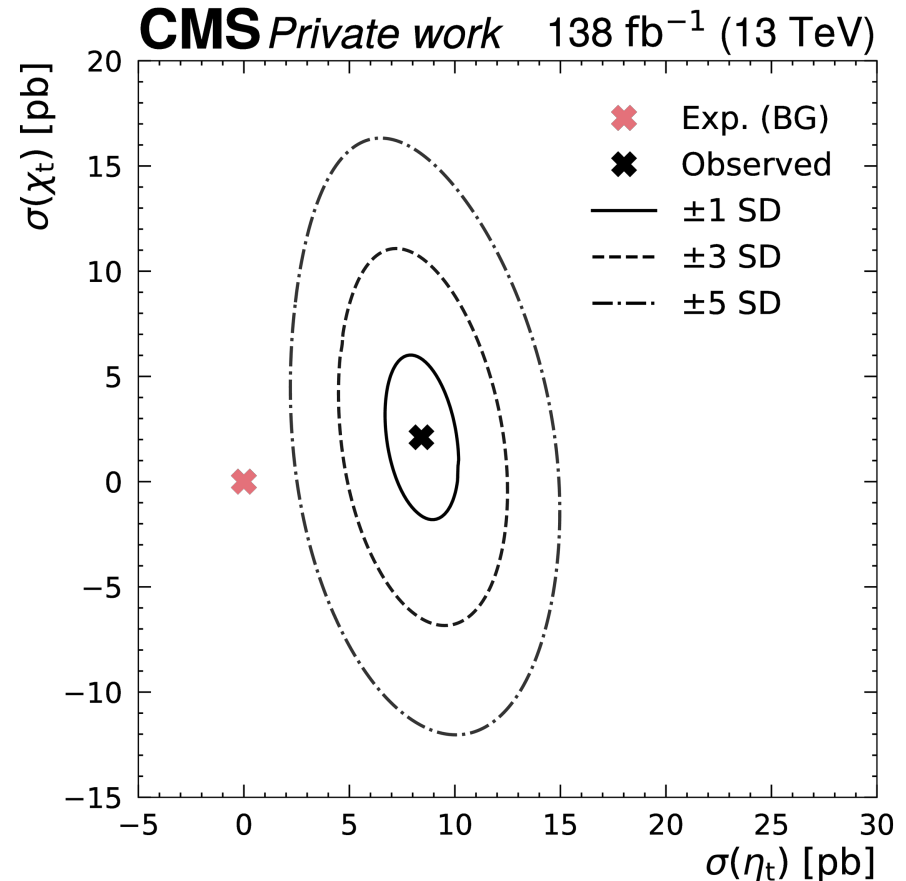


$t\bar{t}$ continuum:
No slope in data



Scalar vs. pseudoscalar

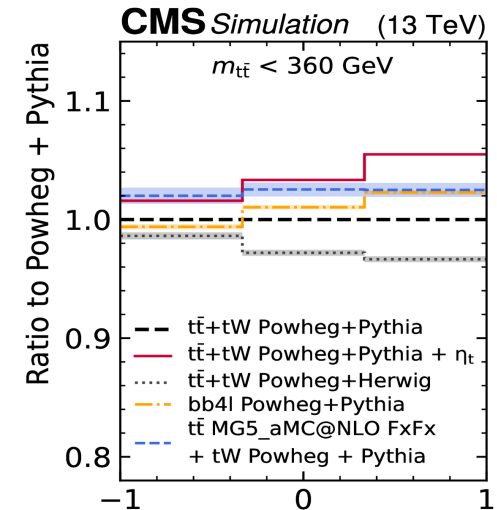
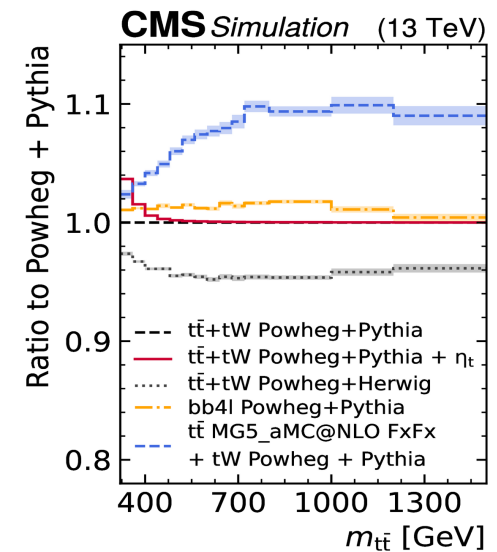
- Can we quantify whether the excess is scalar or pseudoscalar?
- introduce χ_t : scalar $t\bar{t}$ bound state (3P_0 $t\bar{t}$ spin state)
 - similar simplified model as η_t
- Perform 2D fit with η_t and χ_t as signals
- non-zero η_t by > 5 standard deviations
- χ_t compatible with zero by 1 sigma



Checks of the result

How do we make sure the excess is real?

Check **alternative generators** for continuum $t\bar{t} + tW$
nominal: Powheg + Pythia with NNLO QCD/NLO EW reweighting



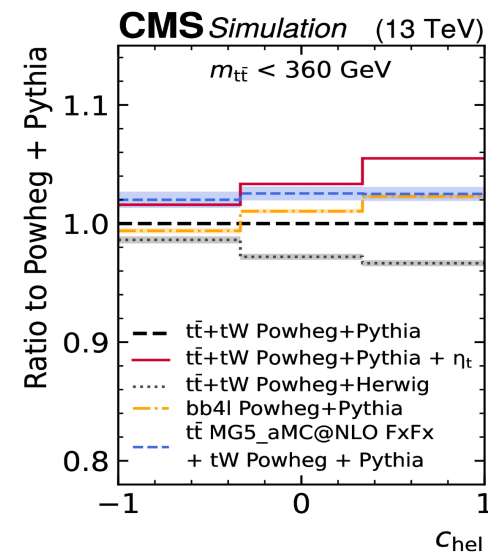
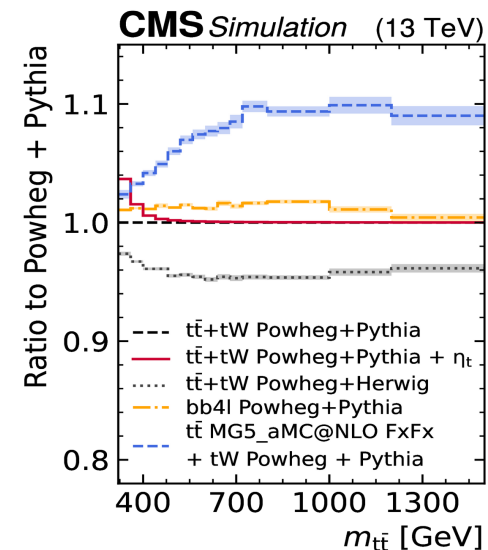
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- **bb4l**: full off-shell $pp \rightarrow b\bar{b}l\bar{l}\nu\nu$ at NLO in QCD, including $t\bar{t}/tW$ interference
 - **validated & implemented myself in CMS** (service work)
 - Enhanced slope in c_{hel} w.r.t nominal Powheg $t\bar{t} + tW$ - similar to η_t
 - Reduces extracted η_t cross section:

$$\begin{array}{ccc}
 \text{bb4l} & & \text{nominal Powheg } t\bar{t} + tW \\
 \sigma(\eta_t) = 6.6 \pm 1.4 \text{ pb} & \longleftrightarrow & \sigma(\eta_t) = 8.7 \pm 1.1 \text{ pb}
 \end{array}$$

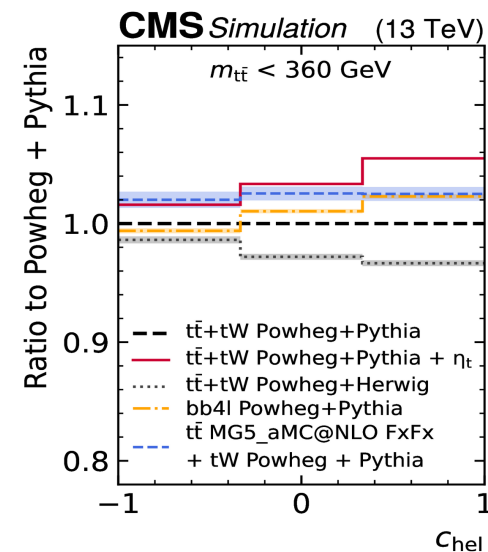
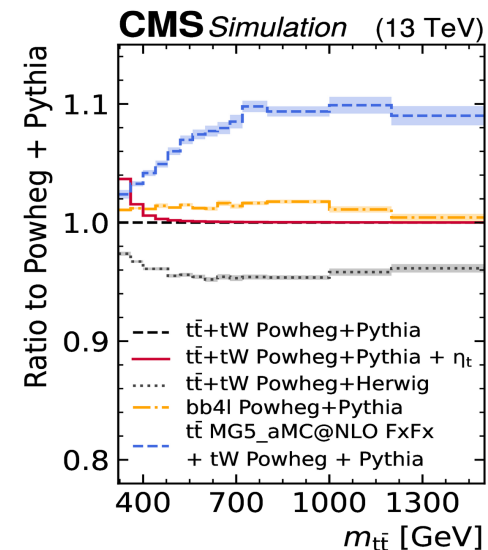


Checks of the result

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Powheg+Herwig instead of Powheg+Pythia
 - Herwig predicts more events at $t\bar{t}$ threshold but less spin correlation \rightarrow distinguishable from η_t
 - no large change in cross section

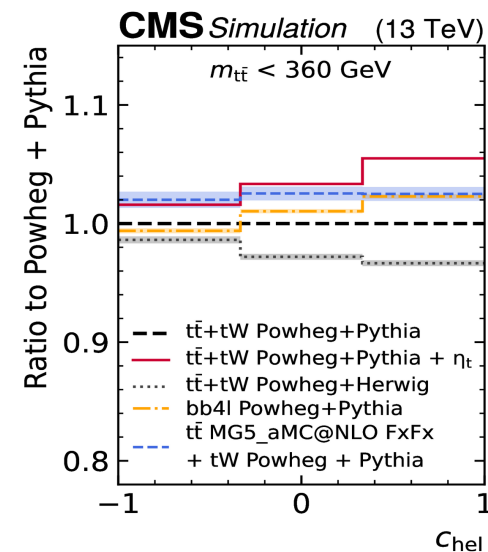
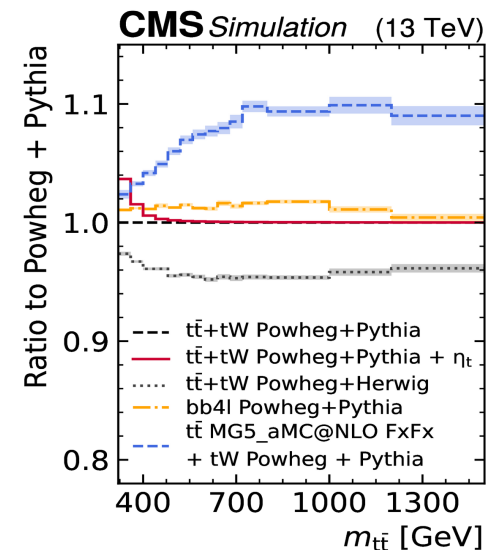


Checks of the result

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- Parton shower for $t\bar{t}$:
Powheg+Herwig instead of Powheg+Pythia
 - Herwig predicts more events at $t\bar{t}$ threshold but less spin correlation \rightarrow distinguishable from η_t
 - no large change in cross section
- Include bb4l and Herwig as additional systematics:
 $\sigma(\eta_t) = 8.8_{-1.4}^{+1.2} \text{ pb}$ slight increase in uncertainty



$t\bar{t}$ bound state or BSM?

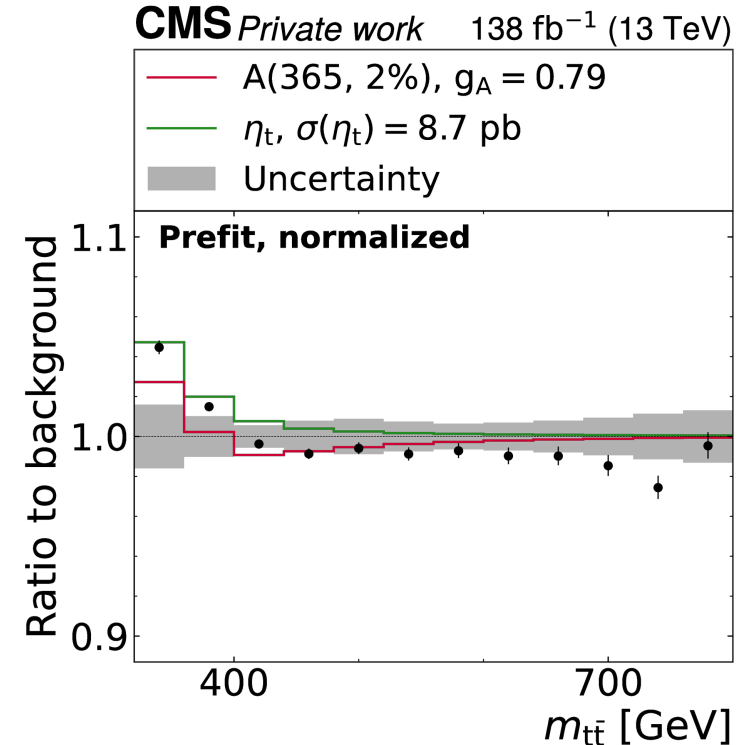
- same excess can be interpreted as (BSM) pseudoscalar A:

best fit

$$m_A = 365 \text{ GeV}, \quad \Gamma_A/m_A = 2\%, \quad g_{A t\bar{t}} = 0.79$$



lowest mass point probed in the analysis!



$t\bar{t}$ bound state or BSM?

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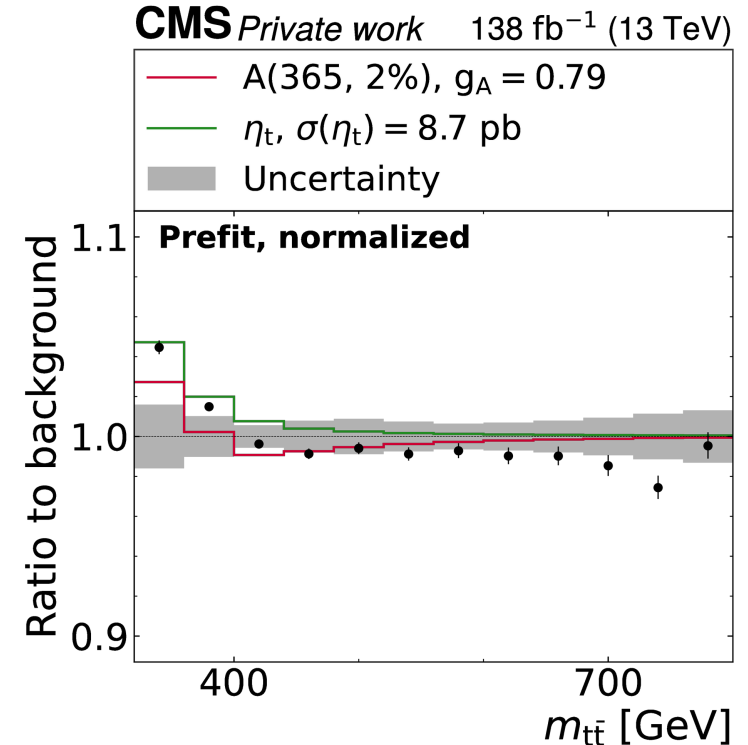
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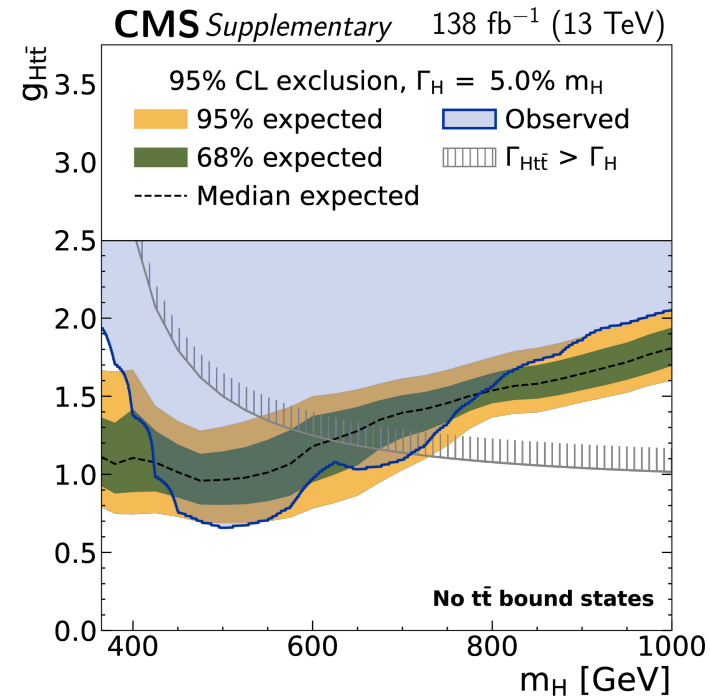
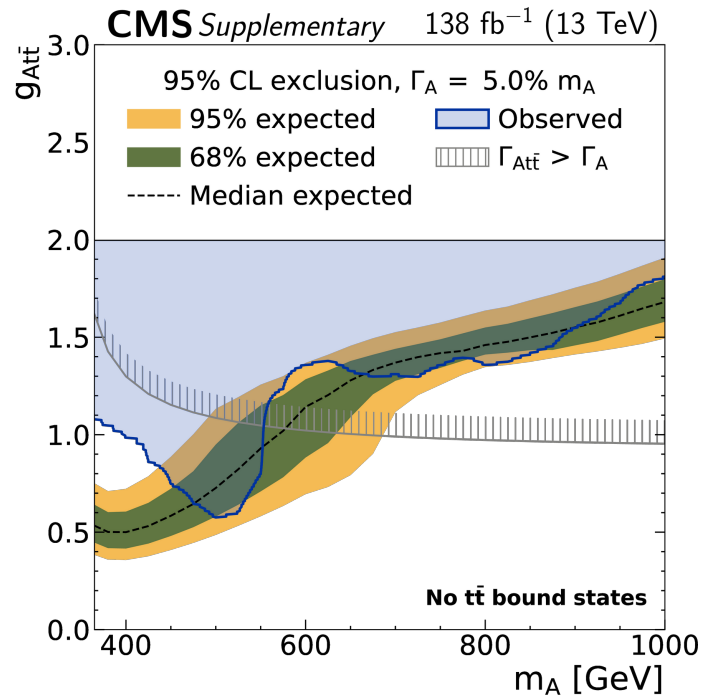
lowest mass point probed in the analysis!

- in general: could be any combination from bound state effects and BSM
- fit slightly prefers η_t bound state, but only by 1σ



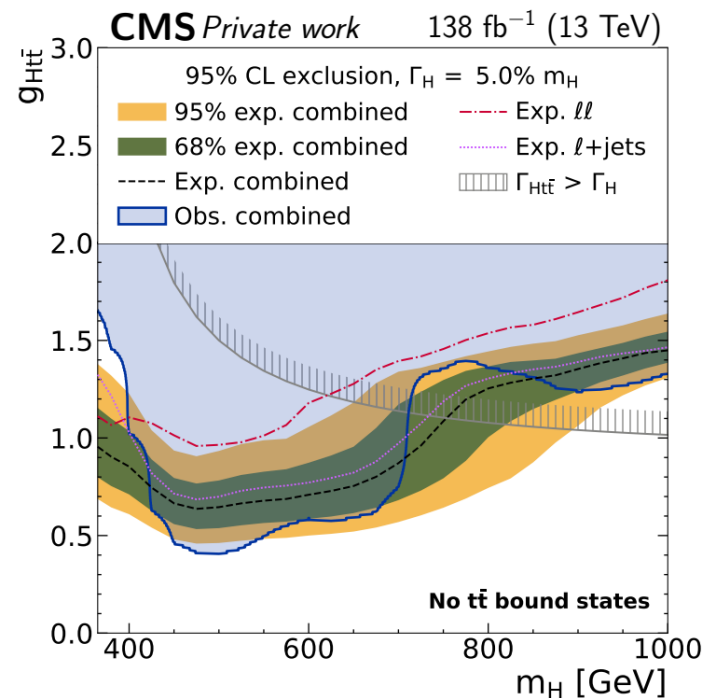
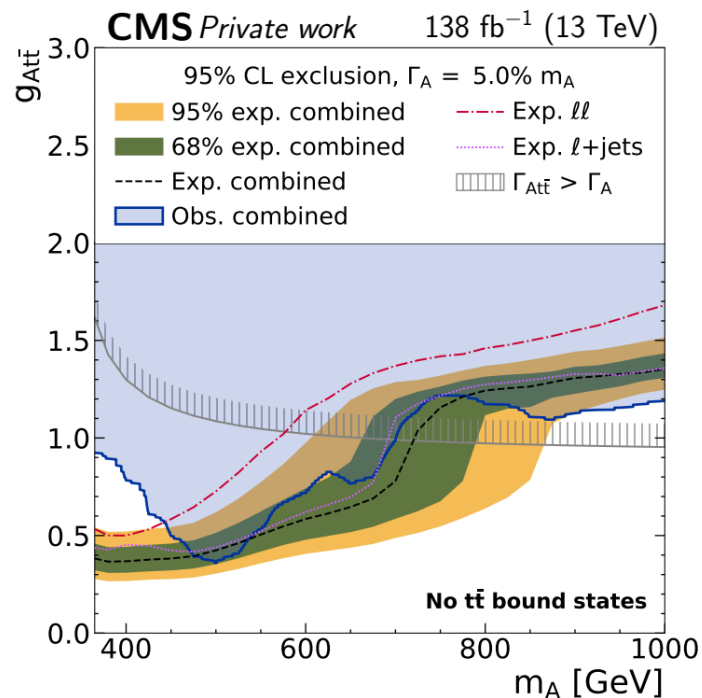
Limits on A/H

- Set limits on A/H -top coupling over large mass range
- No bound state effects considered \rightarrow excess shown in limits



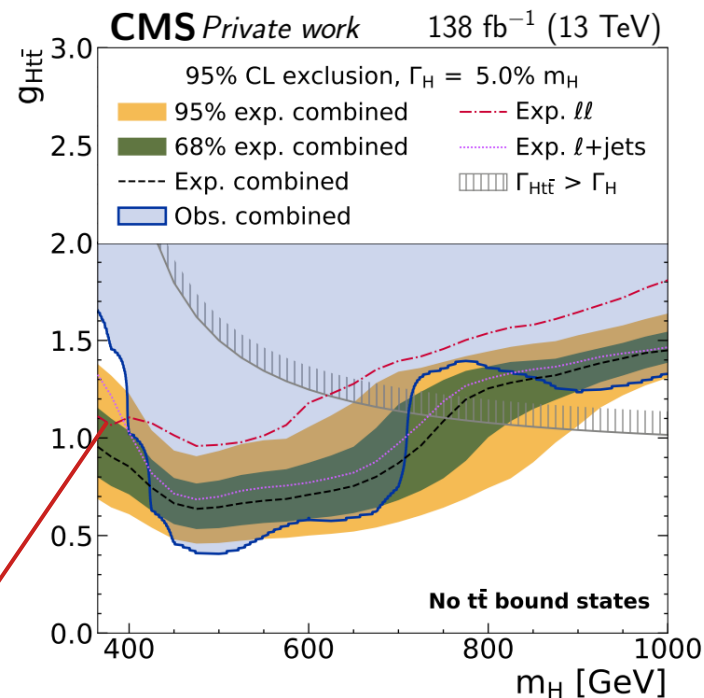
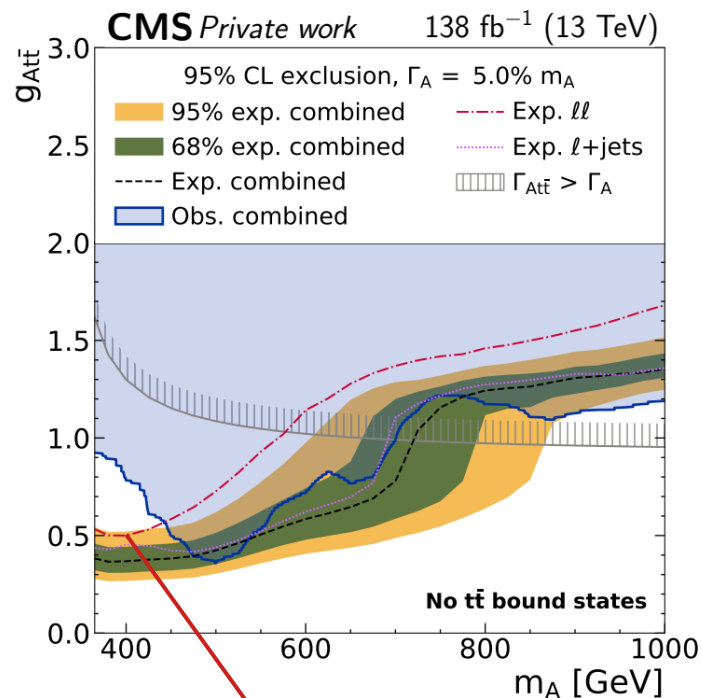
Limits on A/H : combination

- Statistical combination with **lepton+jets decay channel**



Limits on A/H : combination

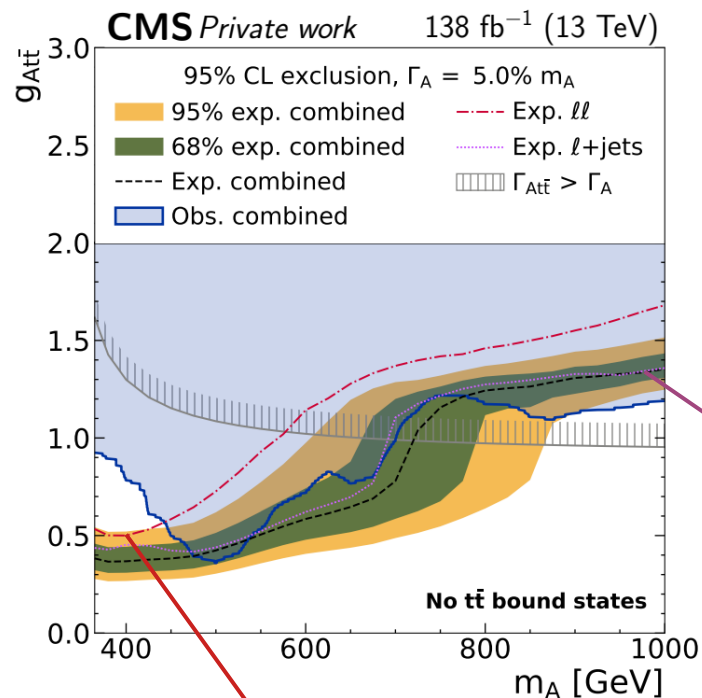
- Statistical combination with **lepton+jets decay channel**



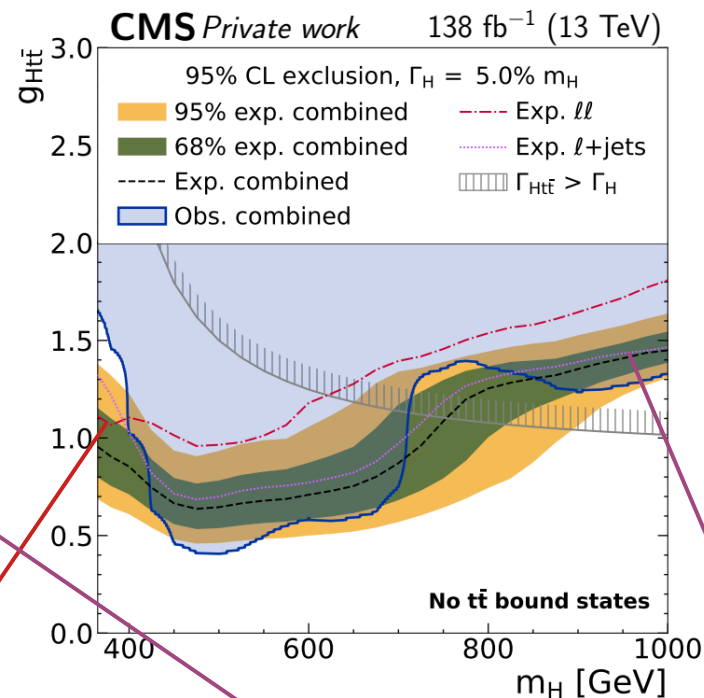
at low masses: dilepton channel comparable
or better than ℓ +jets

Limits on A/H : combination

- Combine A/H limits with **lepton+jets decay channel**



at low masses: dilepton channel comparable
or better than ℓ +jets



at high masses:
 ℓ +jets channel dominates

Axion-Like Particles

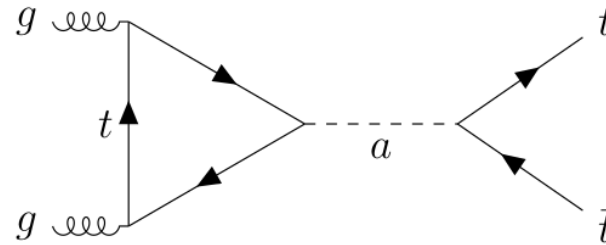
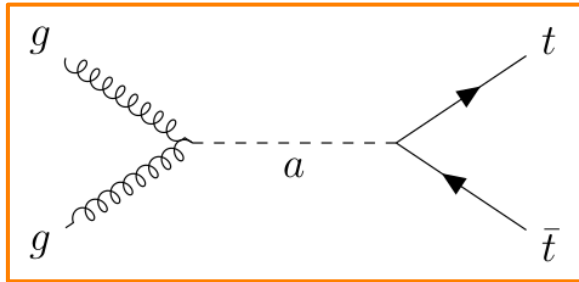
- Further interpretation: heavy **Axion-Like Particles** (ALPs)
 - could solve strong CP problem
 - if $m_{\text{ALP}} > 2 m_t$: decay to $t\bar{t}$ just like pseudoscalar A
 - possible **additional ALP-gluon coupling** in Lagrangian

$$\mathcal{L}^{\text{ALP}} = \frac{1}{2}(\partial_\mu a)(\partial^\mu a) + \frac{m_a^2}{2}a^2 - \boxed{c_{\tilde{G}} \frac{a}{f_a} G_{\mu\nu}^a \tilde{G}^{a\mu\nu}} + ic_t m_t \frac{a}{f_a} \bar{t} \gamma^5 t$$

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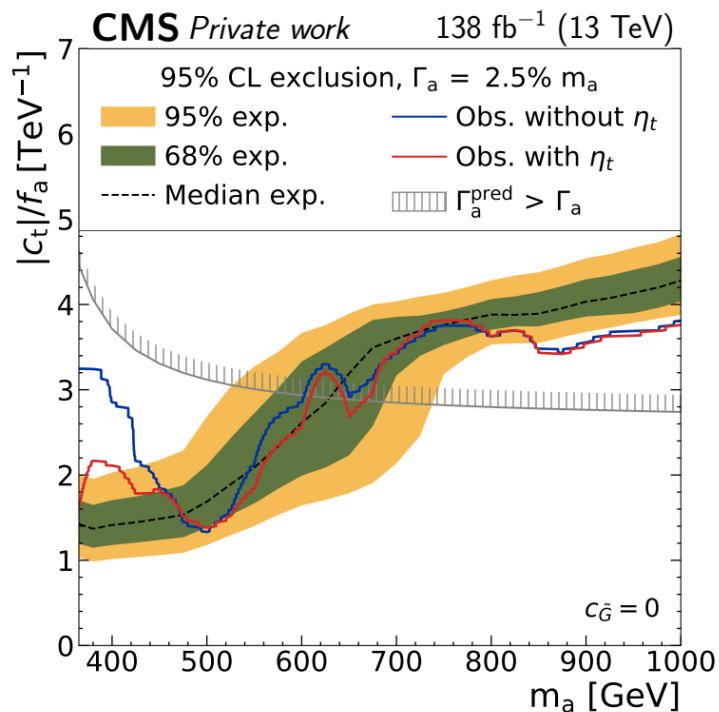
additional production diagram

→ different $m_{t\bar{t}}$ spectrum

Axion-Like Particles

JHEP 12 (2024) 197

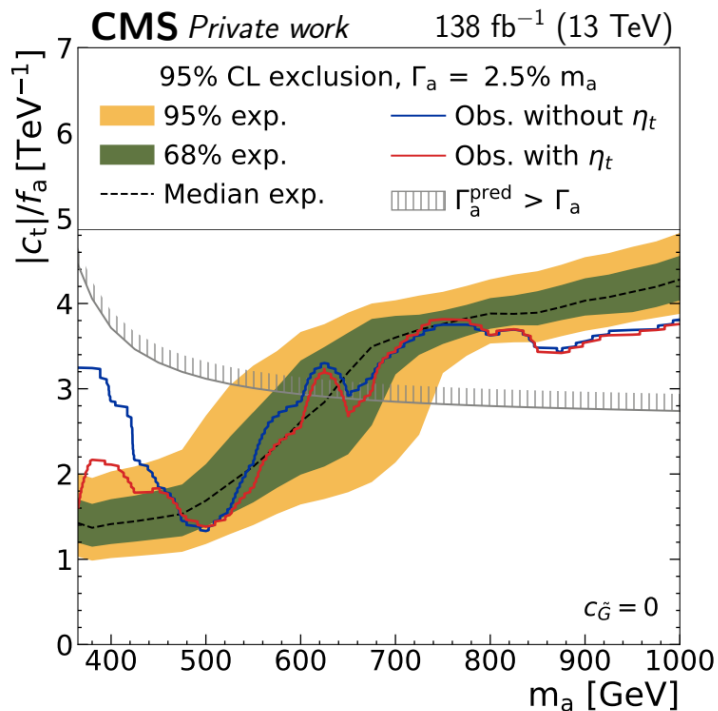
no ALP-gluon coupling $c_{\tilde{G}}$:
set experimental limits



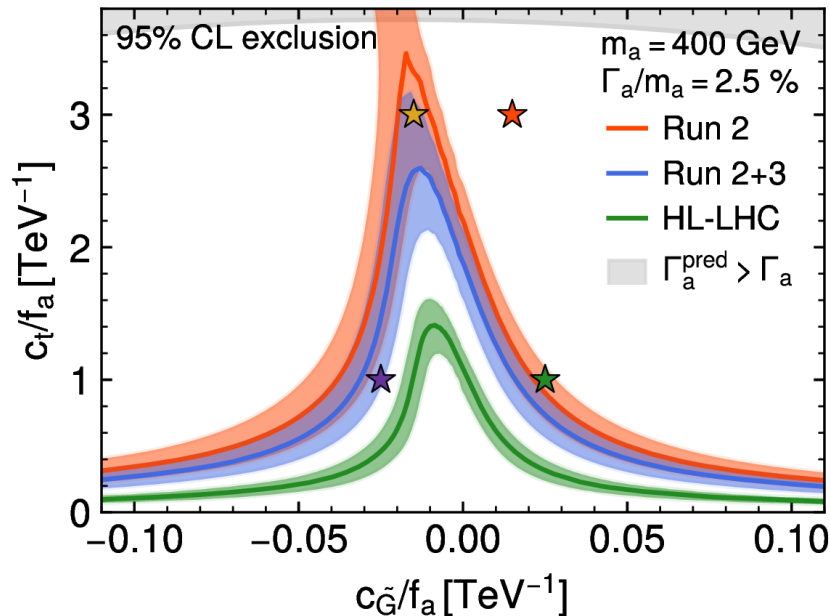
Axion-Like Particles

JHEP 12 (2024) 197

no ALP-gluon coupling $c_{\tilde{G}}$:
set experimental limits



with ALP-gluon coupling $c_{\tilde{G}}$:
projected limits for different
eras of LHC



Summary

- Early measurement of the inclusive $t\bar{t}$ cross section in Run 3
 - First public physics result of Run 3, only ~2 months after start of datataking
 - Required estimation of many needed corrections
 - Comparable precision to Run 2 measurements *JHEP* 08 (2023) 204
- Search for new spin-0 states in $t\bar{t} \rightarrow$ dilepton events
 - Low- $m_{t\bar{t}}$ excess observed & interpreted as $t\bar{t}$ bound state $\eta_t (> 5\sigma!)$
 - Further interpretations as generic (pseudo)scalars in combination with lepton+jets channel or Axion-Like Particles *RoPP* 88 087801 (2025)
arXiv:2507.05119
(submitted to *RoPP*)
JHEP 12 (2024) 197

Backup

$t\bar{t}$ bound states: NRQCD

See e.g.

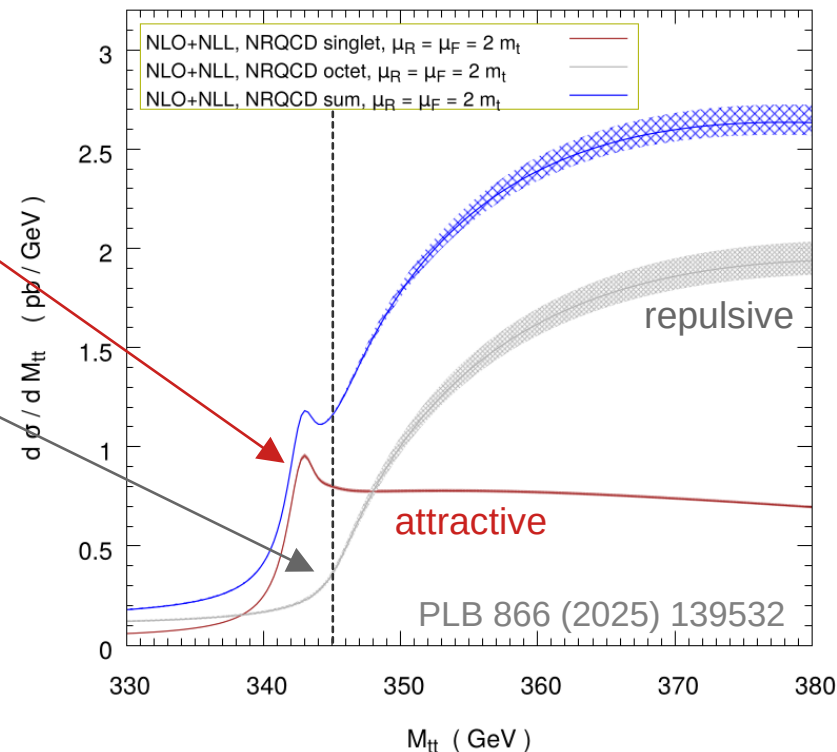
PRD 110 (2024) 5, 054032

JHEP 03 (2024) 099

PRD 104 (2021) 3, 034023

etc.

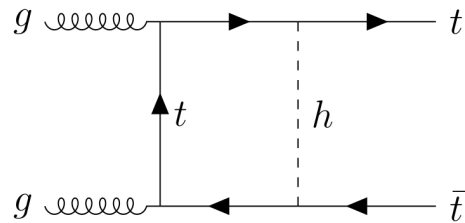
- Results of NRQCD calculations:
 - Color-singlet ($^1S_0^{[1]}$) - **attractive**
→ **Peak below the $t\bar{t}$ threshold**
CP-odd / pseudoscalar spin state!
 - Color-octet ($^1S_0^{[8]}$ or $^3S_1^{[8]}$) - **repulsive**
→ **Suppressed below the $t\bar{t}$ threshold**
- Difficulties: matching to relativistic calculation
soft gluon emissions, ...
- Exact lineshape and width below
experimental resolution



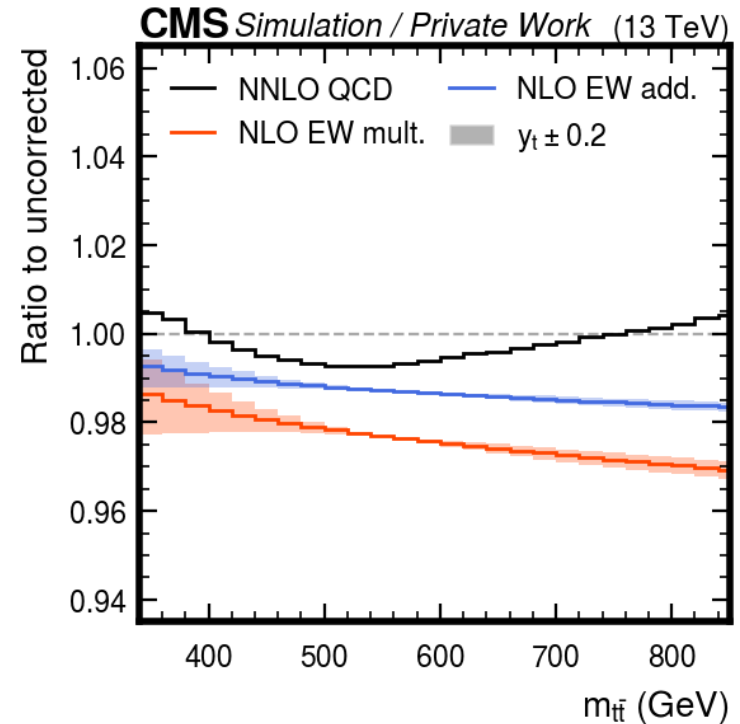
Background modeling

- Major irreducible background: **SM $t\bar{t}$**
 - Model from **NLO MC** (Powheg+Pythia)
 - Correct to **NNLO QCD** and **NLO EW** from fixed-order predictions by reweighting in 2D bins of $m_{t\bar{t}}$ and top scattering angle

(EPJC 78 (2018) 537, EPJC 51 (2007) 37)



- Normalize to NNLO+NNLL cross section
(CPC 185 (2014) 2930)



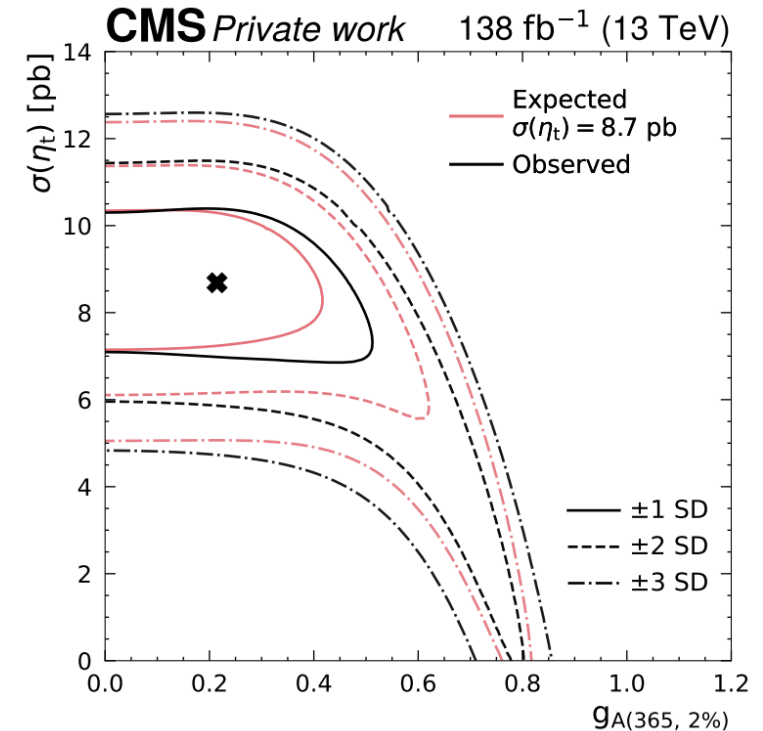
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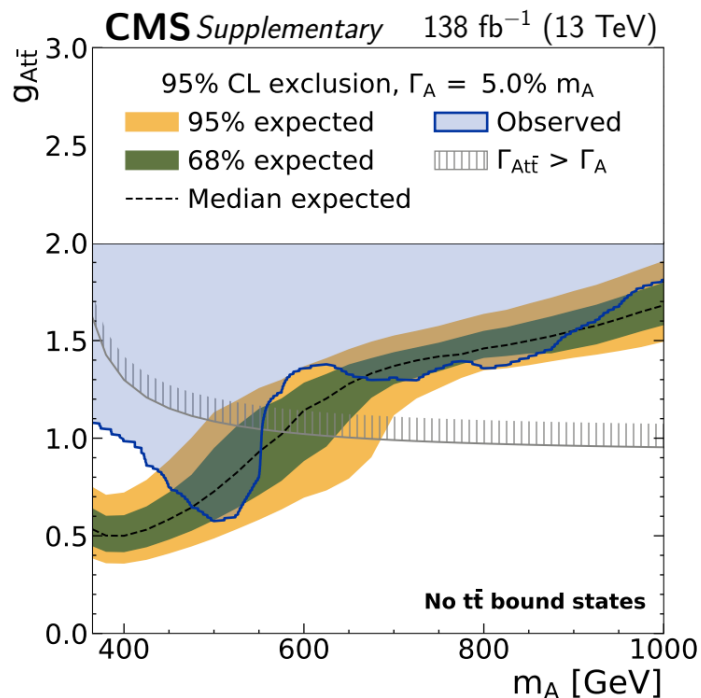
lowest mass point probed in the analysis!

- in general: could be any combination from bound state effects and BSM
- fit prefers η_t bound state, but only by 1σ



Limits on A/H

- Set limits on couplings of A/H to top over large mass range
- Two scenarios on how to treat η_t :



no η_t :

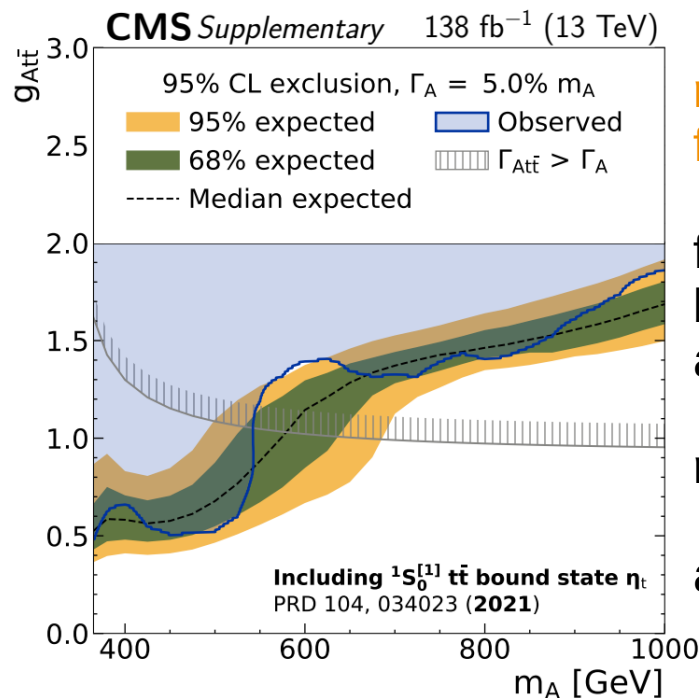
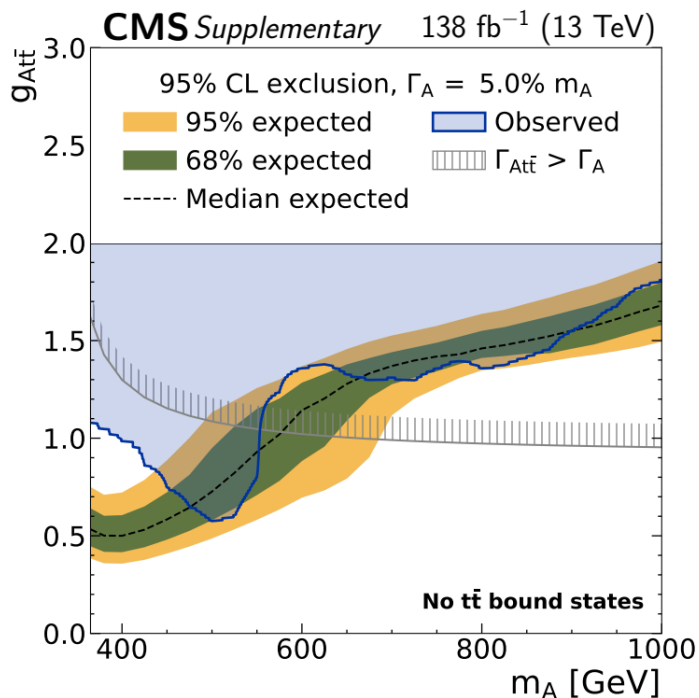
excess visible
in limits

Limits on A/H

- Set limits on couplings of A/H to top over large mass range
- Two scenarios on how to treat η_t :

no η_t :

excess visible
in limits



η_t as additional
free-floating BG:

fit can choose
between A/H
and η_t

η_t is preferred
→ excess is
absorbed

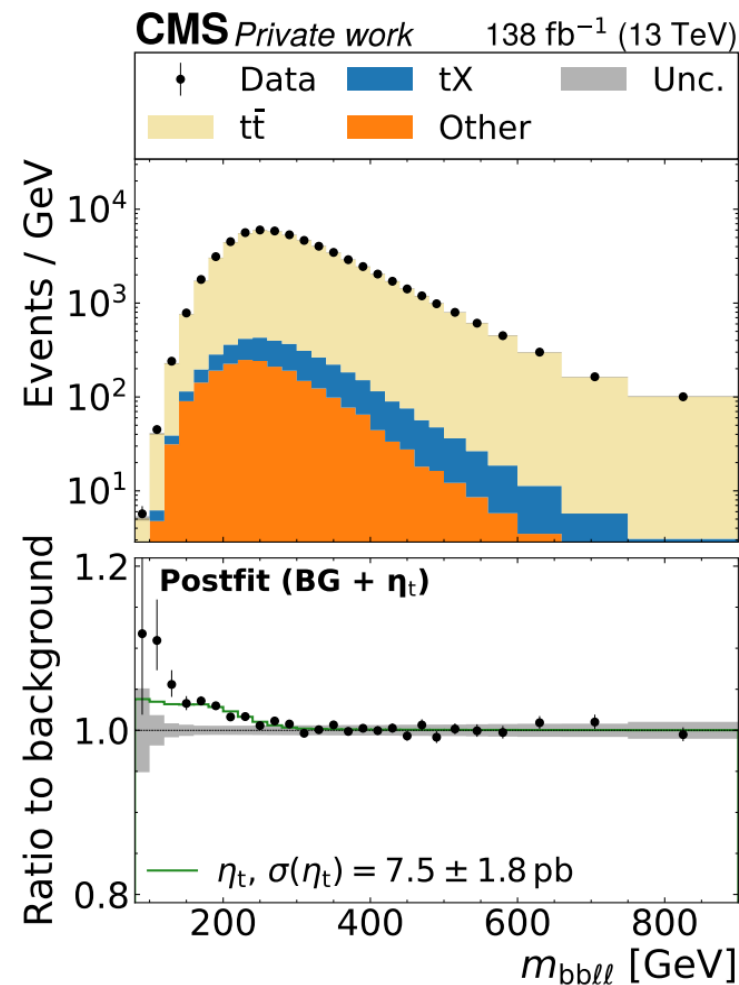
Checks of the result

How do we make sure the excess is real?

Check **alternative variables**: $m_{bb\ell\ell}$ instead of $m_{t\bar{t}}$

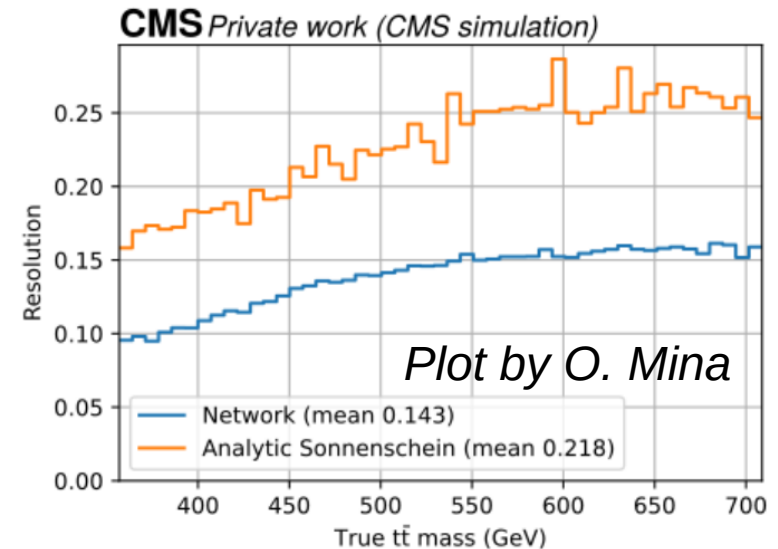
- Sidestep possible biases from $t\bar{t}$ reconstruction
- Repeat full fit: consistent with main result

... and many, many more



Future: ML for $t\bar{t}$ reconstruction

- Analytic dilepton $t\bar{t}$ reconstruction has several problems
 - Assumes top mass constraints – questionable in off-shell regions
 - b jet assignment done in simplistic way, etc...
- We should be able to do better!
- Use modern machine learning techniques:
Lorentz-invariant transformers (L-GATr)
[arXiv:2405.14806]
- Idea: also estimate confidence in the result?
Could reject poorly-reconstructed events...
- Supervision of two summer student projects working on ML for $t\bar{t}$ reco.
+ one combined summer student+B.Sc. project working on the classic algorithm



How it was done: columnar analysis

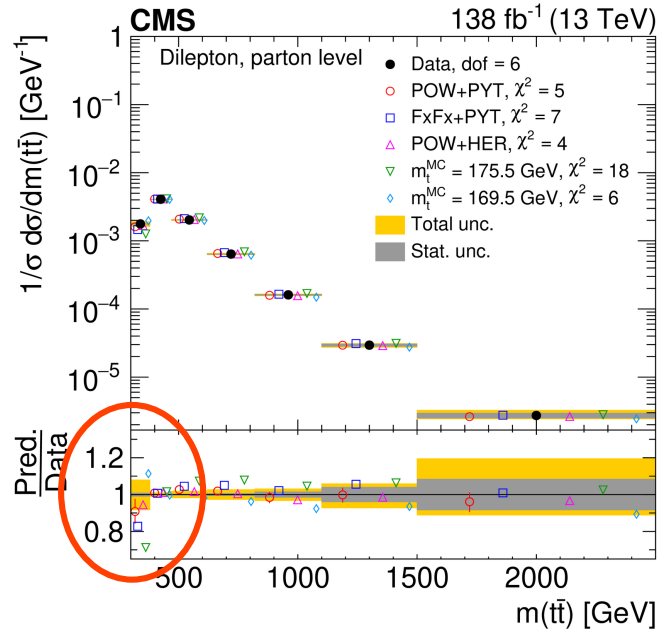
- Measurement was performed with pepper:
fully **python coffea-based columnar analysis framework**
- Fully generic, starting from central nanoAOD
- Developed at DESY, by now also used at other CMS institutes
- **I am one of two maintainers** – implemented many new features & improvements
 - e.g. changes for Run 3, improved HTCondor submission, ...
- Supervised **combined summer student & Bachelor project**:
profile & improve performance of histogramming → up to 150% speedup



SM: $t\bar{t}$ differential measurements

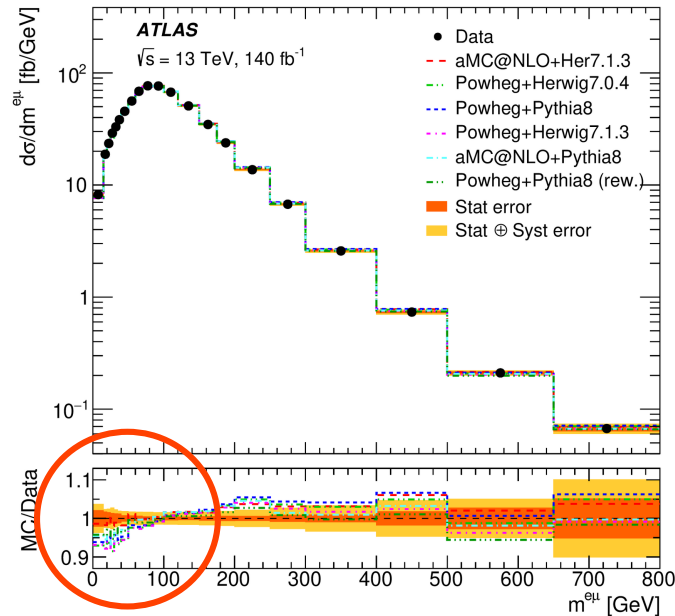
CMS $\ell\ell - m_{t\bar{t}}$

CMS-PAS-TOP-20-006



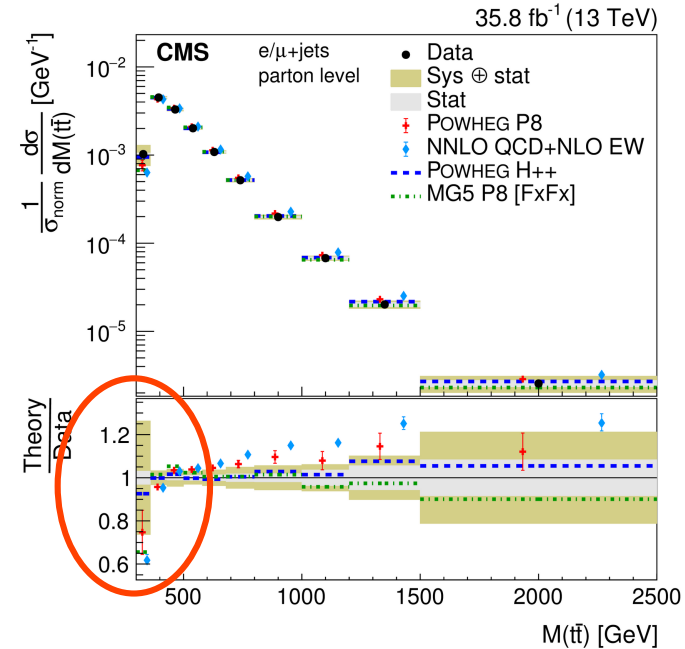
ATLAS $e\mu - m_{e\mu}$

JHEP 07 (2023) 141



CMS ℓ +jets - $m_{t\bar{t}}$

PRD 97 (2018) 112003



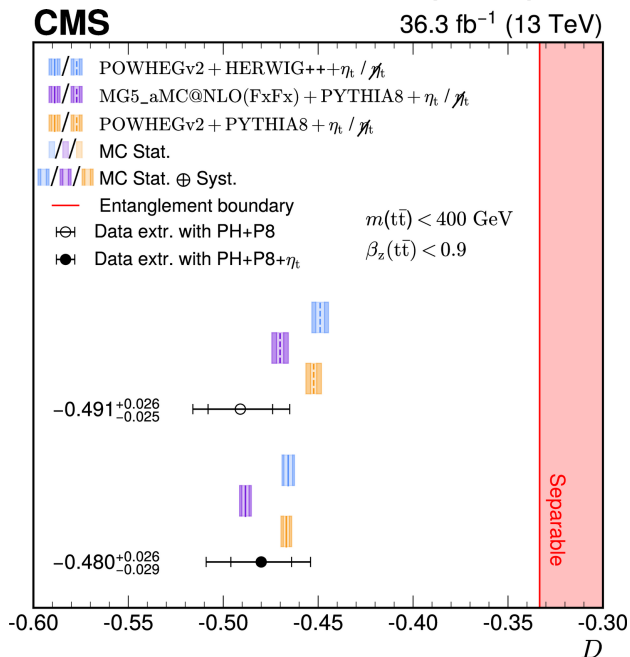
→ good description by theory except for **excess in data in threshold region**

SM: $t\bar{t}$ spin entanglement

- Measured quantity: D “ \approx strength of $t\bar{t}$ spin correlation”

somewhat oversimplified,
more later...

ROPP 87 117801 (2024)

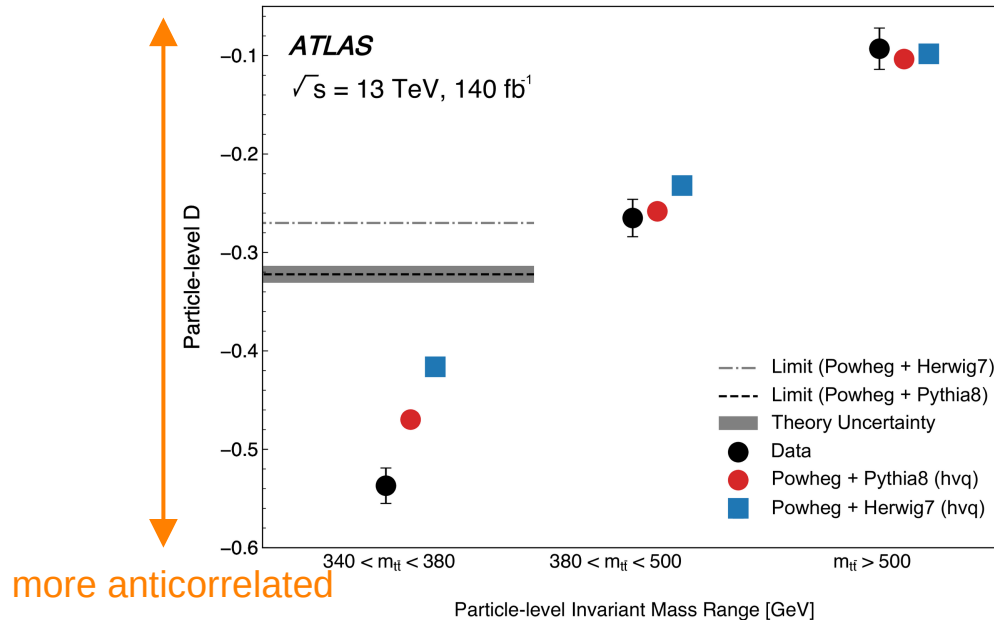


more anticorrelated

less correlated

less correlated

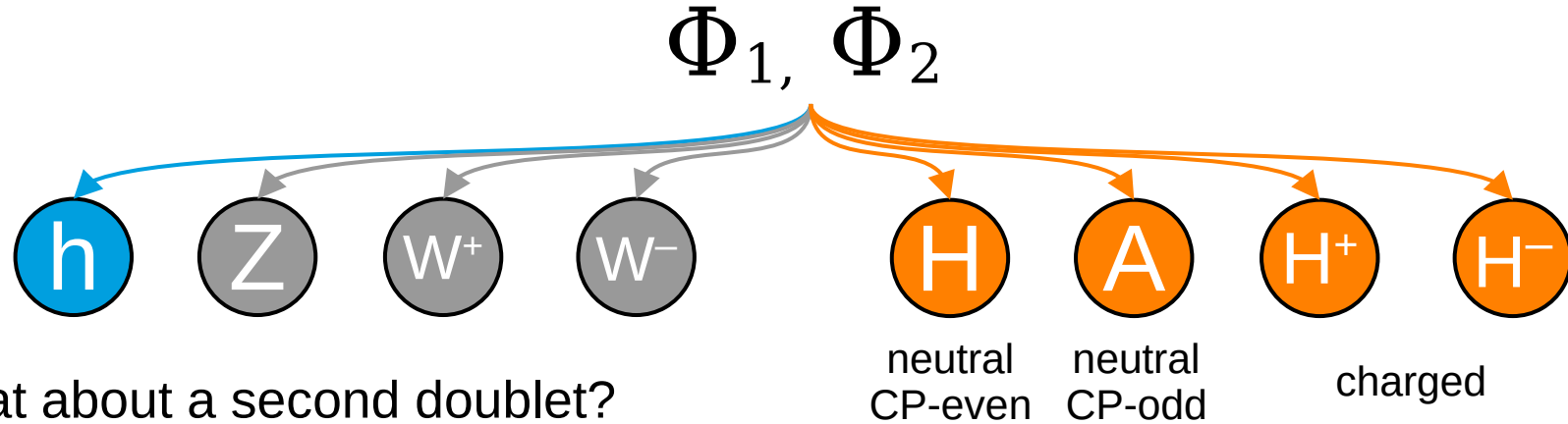
Nature 633 (2024) 542



more anticorrelated

BSM: Two-Higgs Doublet Model

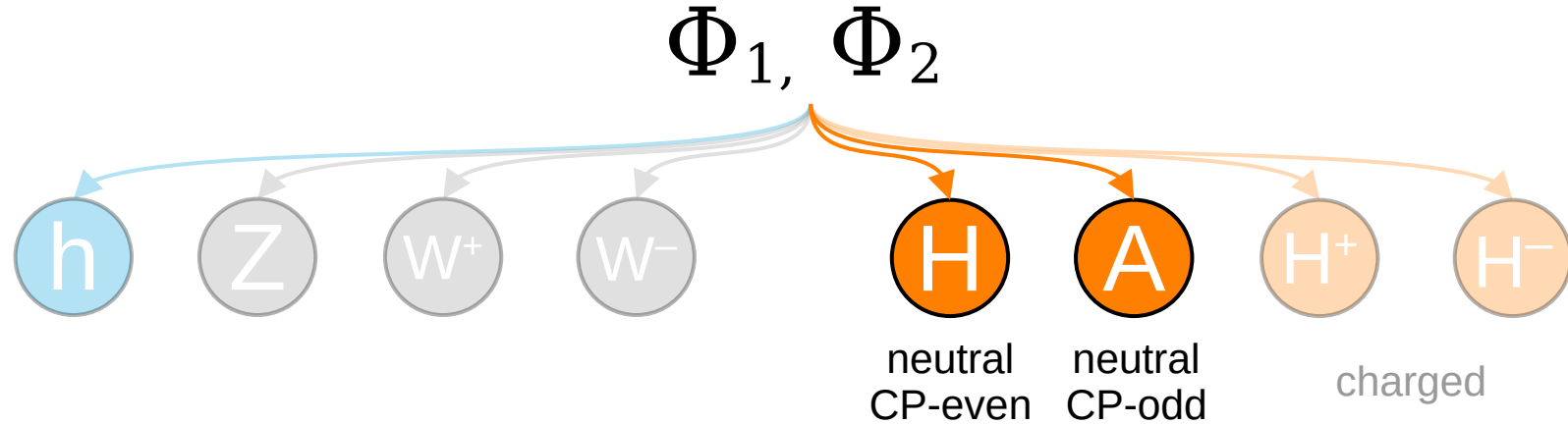
- SM Higgs: complex SU(2) doublet ϕ
→ becomes real scalar after symmetry breaking



- What about a second doublet?
→ Two-Higgs Doublet Model (2HDM)
 - Simplest UV-complete extension of the Higgs sector
 - 4 additional degrees of freedom: H , A , H^+ , H^-
 - Only a starting point – e.g. included in SUSY

BSM: Two-Higgs Doublet Model

- SM Higgs: complex $SU(2)$ doublet ϕ
→ becomes real scalar after symmetry breaking

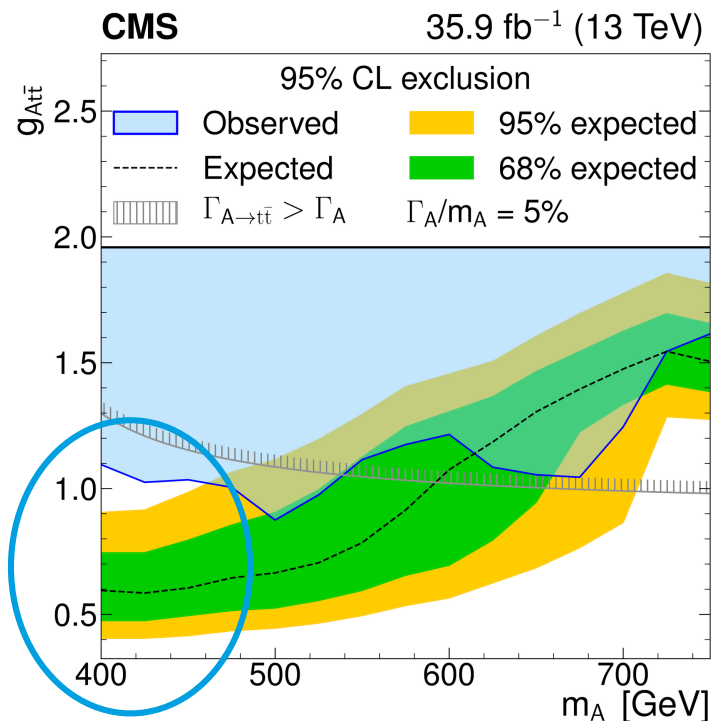
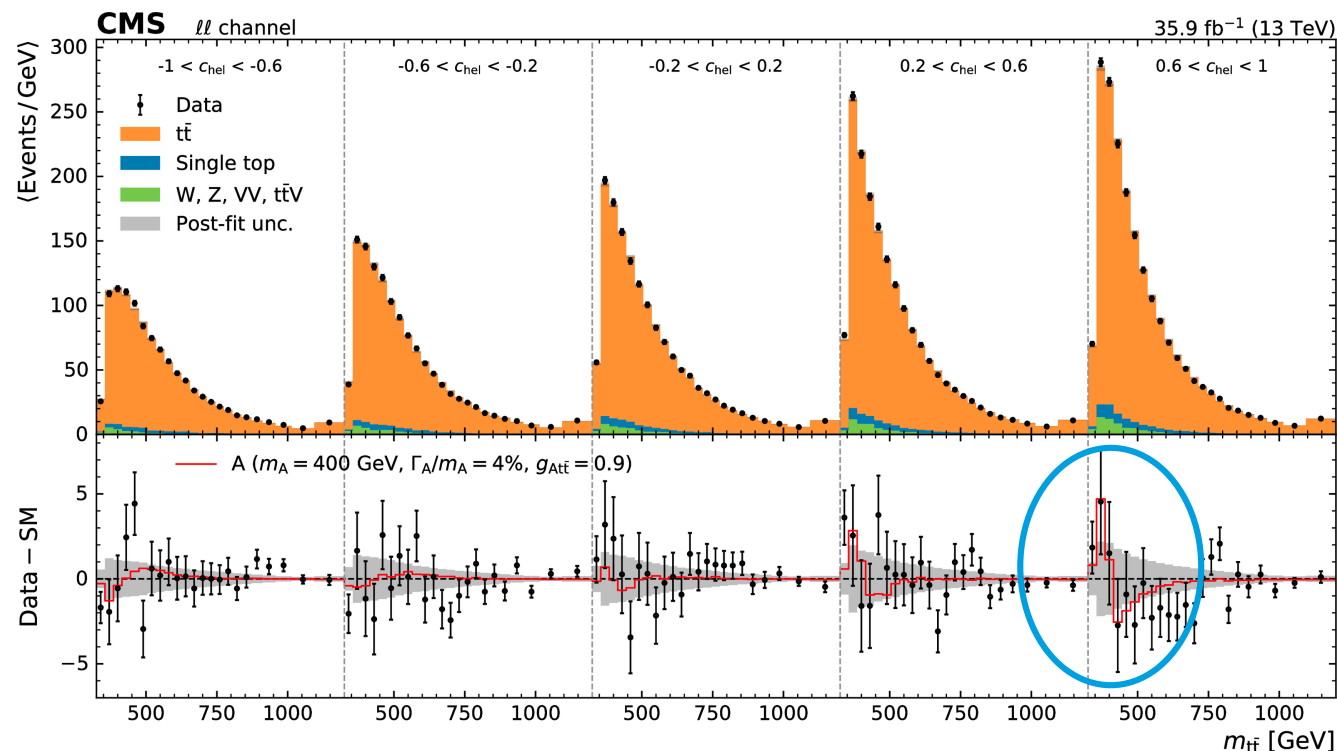


- Most models: couplings are Yukawa-like → largest coupling to top quark!
- For neutral states A and H : **decay to $t\bar{t}$ dominates** if $m_{A/H} > 2m_t$
→ **search in $t\bar{t}$ final states!**

BSM: Previous work

Search for (pseudo)scalars in $t\bar{t}$ events with 2016 data

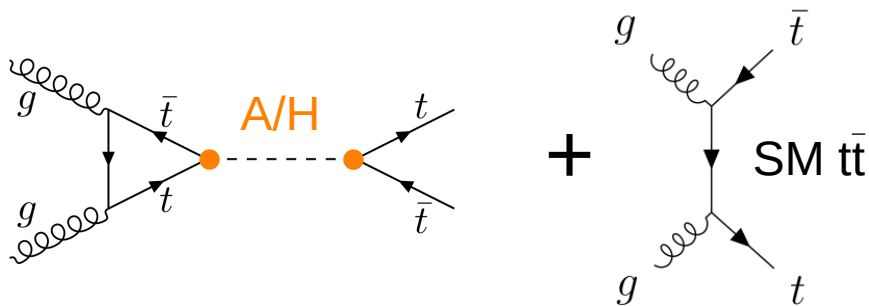
JHEP 04 (2020) 171



Excess for low pseudoscalar masses ($\sim 3\sigma$ local)

Signal modeling: A/H

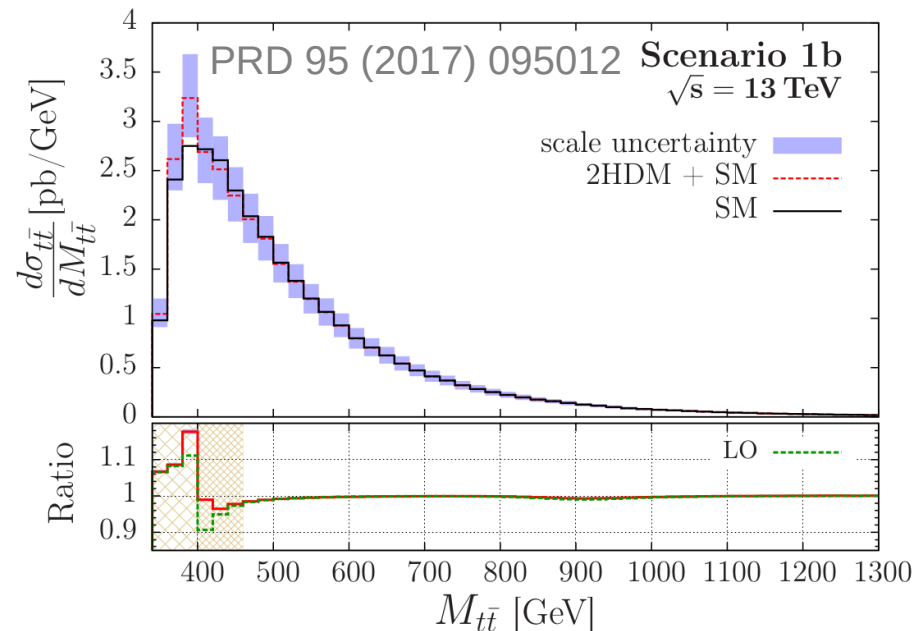
- Generic heavy **pseudoscalar (A)** or **scalar (H)** coupling solely to top quarks
- Production in gluon fusion via top quark loop



- Same final state as SM $t\bar{t}$ → interference
→ **peak-dip structure** in $m_{t\bar{t}}$
- Free parameters: masses, widths, coupling modifiers g_A / g_H

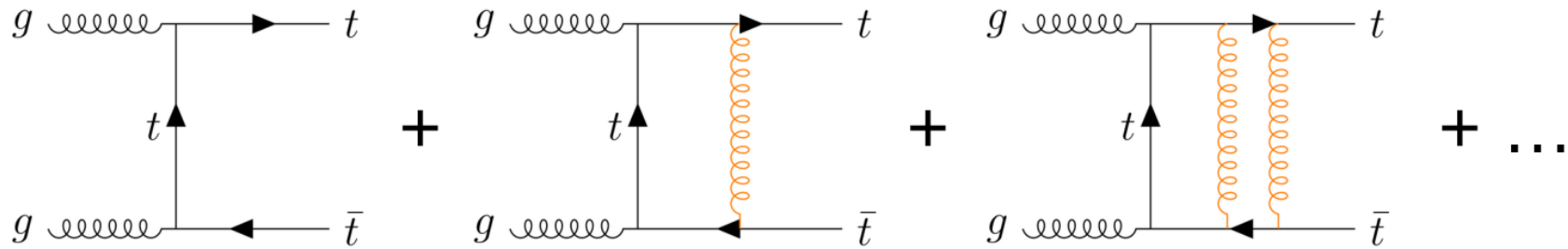
$$\mathcal{L}_A^{\text{int}} = ig_{A t \bar{t}} \frac{m_t}{v} \bar{t} \gamma_5 t A$$

$$\mathcal{L}_H^{\text{int}} = -g_{H t \bar{t}} \frac{m_t}{v} \bar{t} t H$$



Modeling: $t\bar{t}$ bound states

- How to calculate bound states of heavy quarks?
 → **resummation of soft gluons** (Coulomb singularities) in α_S/β ↖ top velocity

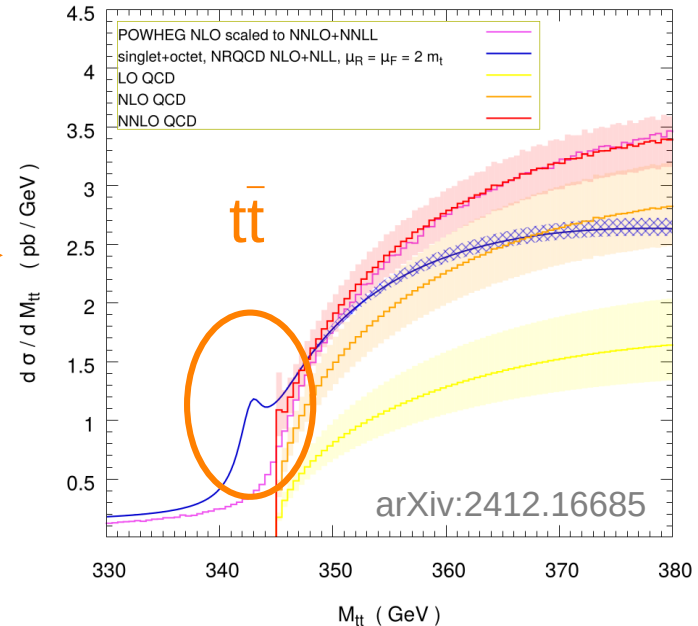
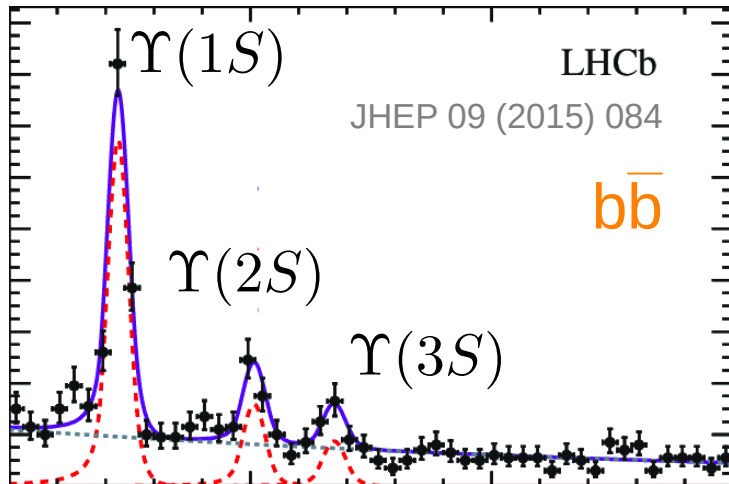


- equivalent to solving **non-relativistic Schrödinger equation** with QCD potential
 momentum space
 like e.g. hydrogen atom!

$$V_{\text{QCD}}(|\vec{q}|) = -\frac{4\pi C^{[1,8]} \alpha_S(|\vec{q}|)}{\vec{q}^2} + \mathcal{O}(\alpha_S^2)$$

Modeling: $t\bar{t}$ bound states

- Need to consider finite top width
→ **one broad peak** instead of multiple narrow bound states



- “**virtual bound state**” or “**quasi-bound state**”
- Short top lifetime → bound state disassociates by one top decaying to Wb

Spin density matrix

- Both A/H and η_t predict $t\bar{t}$ production in a **pure $t\bar{t}$ spin state**: 1S_0 or 3P_0 (from A / η_t resp. H)
- Encoded in **spin density matrix**:

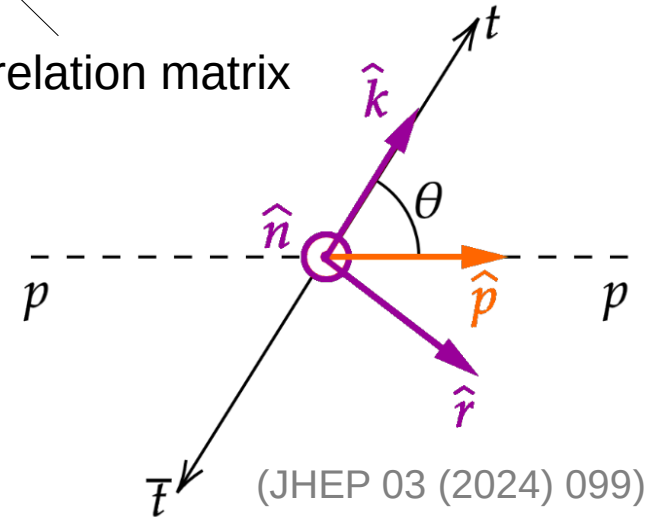
$$\mathbf{R} = A + B_i \sigma_i + \bar{B}_i \bar{\sigma}_i + \sigma_i C_{ij} \bar{\sigma}_j$$

cross section

polarization vectors

correlation matrix

- Choose **helicity basis** $\{\hat{k}, \hat{r}, \hat{n}\}$:
 - \hat{k} : direction of flight of the top quark
 - \hat{r} and \hat{n} : orthogonal to \hat{k}



Comparison with ATLAS

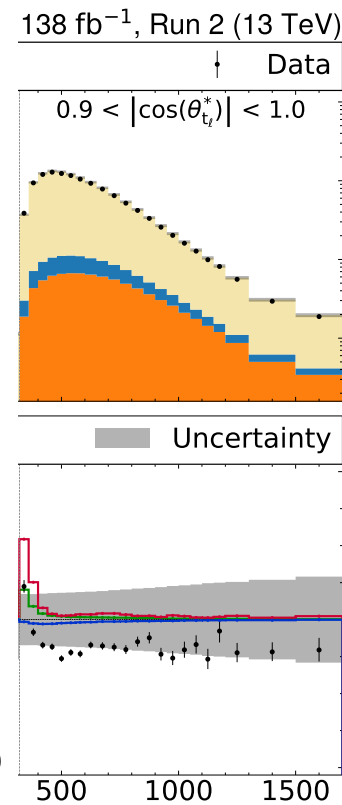
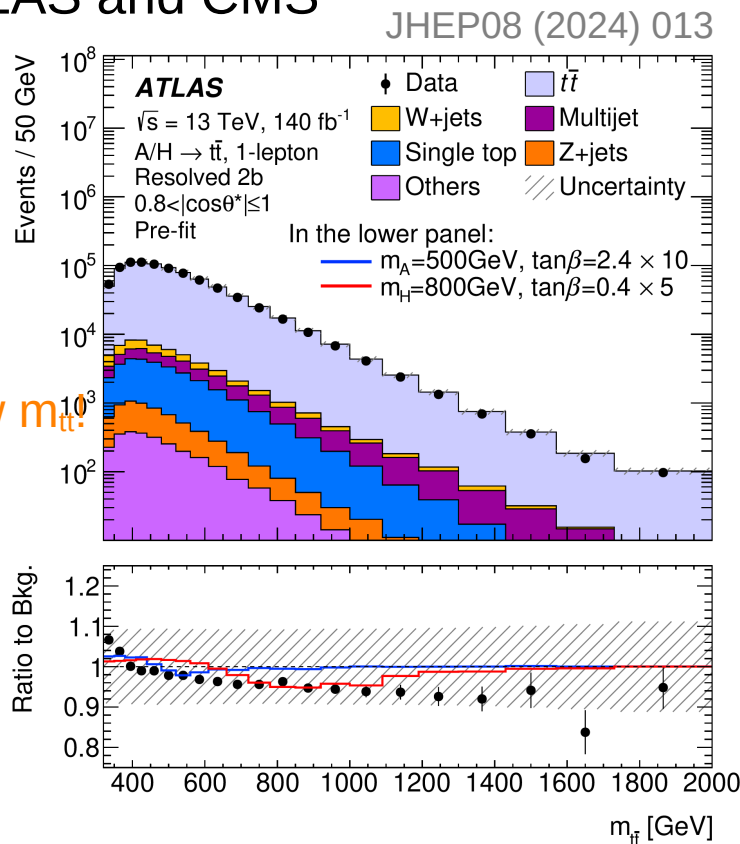
- Different channel definitions in ATLAS and CMS
- ℓ +jets resolved:
 - ATLAS: $1\ell, 1b, \geq 4$ jets
 - CMS: $1\ell, 2b, 3$ jets
 - both: $1\ell, 2b, \geq 4$ jets
→ compare pre-fit distributions!

Comparison with ATLAS

- Different channel definitions in ATLAS and CMS

- ℓ +jets resolved:

- ATLAS: $1\ell, 1b, \geq 4$ jets
- CMS: $1\ell, 2b, 3$ jets
- both: $1\ell, 2b, \geq 4$ jets
- compare pre-fit distributions!
- Similar prefit excess in data at low $m_{t\bar{t}}$!



Comparison with ATLAS

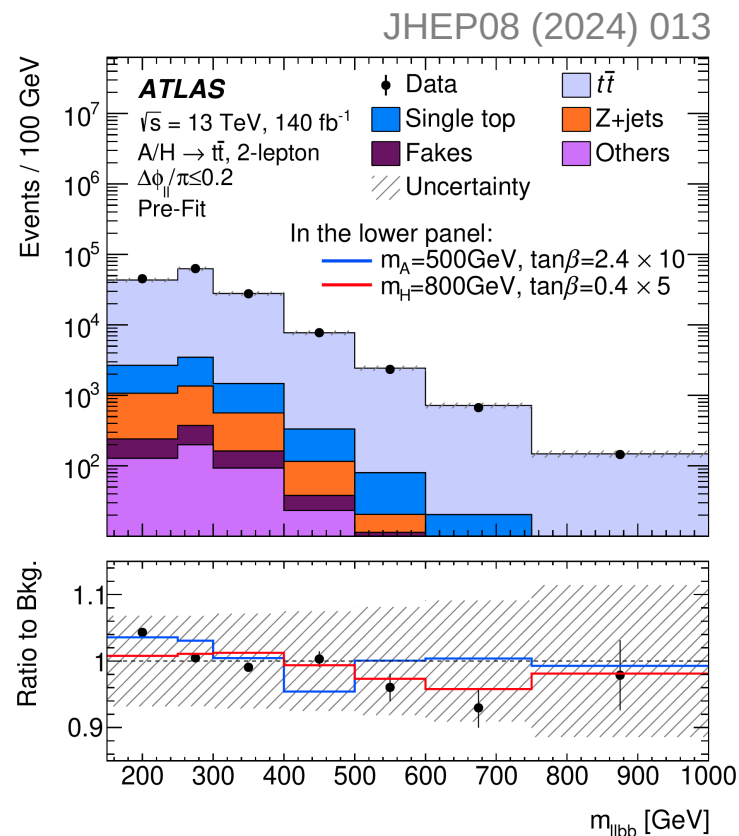
- Different channel definitions in ATLAS and CMS

- ℓ +jets resolved:

- ATLAS: $1\ell, 1b, \geq 4$ jets
- CMS: $1\ell, 2b, 3$ jets
- both: $1\ell, 2b, \geq 4$ jets
→ compare pre-fit distributions!
- Similar prefit excess in data at low $m_{t\bar{t}}$!

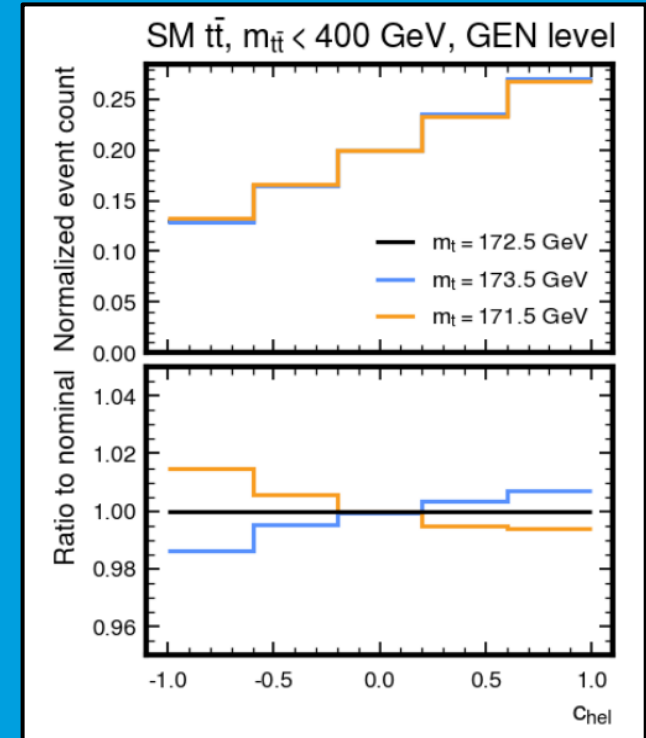
- dilepton: **very different strategy**

- CMS: reconstruct $m_{t\bar{t}} \times C_{\text{hel}} \times C_{\text{chan}}$
drives sensitivity at the $t\bar{t}$ threshold!
- ATLAS: no top quark reconstruction
instead: $m_{\ell\bar{\ell}bb} \times \Delta\phi_{\ell\bar{\ell}}$



Comparison with ATLAS

- From internal studies:
inclusion of **spin correlations** in CMS helps to disentangle signal and systematic uncs.
 - e.g. for **downwards shift of top mass**:
 - More events at threshold \rightarrow like signal :(
 - BUT less spin correlation \rightarrow unlike signal :)
 - Similar for many other uncertainties
e.g. Pythia vs. Herwig in the $t\bar{t}$ BG
- \rightarrow (one) reason why **CMS dilepton is more sensitive at the $t\bar{t}$ threshold**



Comparison with ATLAS

- Similar prefit excess in comparable lepton+jets channel
 - Very different strategy in dilepton channels
 - drives the CMS sensitivity at the threshold
 - subdominant for ATLAS (no top reconstruction)
 - Some different uncertainty treatment
-
- We are comparing in detail in the LHC Top Working Group!

List of systematic uncertainties

Experimental

- Jet energy corrections - split into 11 subsources
- Jet energy resolution
- Unclustered p_{miss}^T (uncorrelated between years)
- Luminosity – correlated and decorrelated parts between years
- Pileup
- Trigger efficiencies (separate for $\ell\ell$ / ℓj)
- Electron efficiencies (reco. & ID)
- Muon efficiencies – split into syst. and stat.
- B tagging and mistagging efficiencies
 - B tagging split into subsources
- L1 ECAL prefiring (where applicable)
- Data-driven EW+QCD BG ($\ell+jets$) : shape & rate (50%) uncorrelated between channels
- Data-driven Z+jets normalization ($\ell\ell$)

Theory

- Factorization & renormalization scales:
 - $t\bar{t}$, tW, tq, Z+jets; η_t (BG or signal), A/H signal
 - Uncorrelated between processes
 - $t\bar{t}$: including cross section variation
- Same for initial & final state radiation PS scales
- MC top mass: $\pm 1\text{GeV}$ (interpolated from $\pm 3\text{GeV}$)
 - Also including cross section variations
- ME-PS matching (h_{damp})
- Underlying event tune
- Color reconnection: 3 different samples
- PDF: PCA performed on final templates from 100 replicas → only leading component considered
- PDF α_s
- Electroweak corrections:
 - SM Higgs-Top Yukawa coupling (1 +0.11 -0.12)
 - EW correction scheme (additive v. multiplicative)
- Minor BG cross sections: 15% for tW and tq; 30% for Diboson and $t\bar{t}+X$

List of MC generators

Process	QCD order	ME Generator
$t\bar{t}$	NLO	POWHEG v2 (h _v q)
tW	NLO	POWHEG v2 (ST_wtch)
Z+jets	NNLO	POWHEG v2 (Zj MiNNLO)
t -channel single top	NLO	POWHEG v2 (ST_tch) + MADSPIN
s -channel single top	NLO	MG5_AMC@NLO
$t\bar{t}W$	NLO	MG5_AMC@NLO
$t\bar{t}Z$	NLO	MG5_AMC@NLO
WW, WZ & ZZ	LO	PYTHIA 8.2
A/H signal	LO	MG5_AMC@NLO
η_t signal	LO	MG5_AMC@NLO

Data-driven Z+jets normalization

- b jets in Z+jets are known to be badly modeled in MC – might lead to wrong normalization after requiring ≥ 1 btag
- Take normalization from Z peak sideband ($R_{in/out}$ method)
- Use weaker assumption than standard $R_{in/out}$ (“ratio of ratios”):
Get $R_{in/out}$ in 0 b tag sideband; take “ratio of ratios” for ≥ 1 and 0 btags from MC

$$\frac{(R_{in/out}^{\geq 1b})_{data}}{(R_{in/out}^{\geq 1b})_{MC}} = \frac{(R_{in/out}^{0b})_{data}}{(R_{in/out}^{0b})_{MC}} \quad \longrightarrow \quad SF = \frac{(N_{out}^{\geq 1b})_{data}}{(N_{out}^{\geq 1b})_{MC}} = \frac{(N_{in}^{\geq 1b})_{data}}{(N_{in}^{\geq 1b})_{MC}} \frac{(R_{in/out}^{0b})_{MC}}{(R_{in/out}^{0b})_{data}}.$$

$$\text{with } N_{data} = N_{data}^{\ell\ell} - 0.5 N_{data}^{e\mu} k_{\ell\ell}, \quad \text{where } k_{ee} = \frac{1}{k_{\mu\mu}} = \sqrt{\frac{N_{data}^{ee}}{N_{data}^{\mu\mu}}}.$$

EW corrections to $t\bar{t}$

- Our EW correction (Hathor) is NLO in EW but LO in QCD
- Ambiguity on how to apply EW corrections to (N)NLO simulation
- Nominal choice: multiplicative

$$\sigma^{\text{rew.}} = \overset{\text{Powheg}}{\sigma_{\text{NLO QCD}}^{\text{LO EW}}} \times \overset{\text{Hathor}}{\frac{\sigma_{\text{LO QCD}}^{\text{NLO EW}}}{\sigma_{\text{LO QCD}}^{\text{LO EW}}}}$$

MadGraph

- Alternate choice: additive

$$\sigma^{\text{rew.}} = \sigma_{\text{NLO QCD}}^{\text{LO EW}} + \sigma_{\text{LO QCD}}^{\text{NLO EW}} - \sigma_{\text{LO QCD}}^{\text{LO EW}}$$

- Difference treated as systematic uncertainty