P.Buividovich, T.K., M.I.Polikarpov, ArXiv: 1111.6733 (to appear in PRD)T.K., "Chiral superfluidity of the quark-gluon plasma", ArXiv: 1208.0012T.K., "Experimental predictions of the chiral superfluidity" (in preparation)

Chiral Superfluidity for the Heavy-Ion Collisions



Tigran Kalaydzhyan

DESY THEORY WORKSHOP 25-28 September 2012, Lessons from the first phase of the LHC DESY Hamburg, Germany

QCD vacuum

 $G^{a\,\mu
u} ilde{G}^a_{\mu
u}$



 $\rho_R \neq \rho_L$



Positive topological charge density

Negative topological charge density



Right-handed



 Spins parallel to B



- Spins parallel to B
- Momenta antiparallel



- Spins parallel to B
- Momenta antiparallel

• If
$$\rho_5 \equiv \rho_L - \rho_R \neq 0$$

then we have a
net electric
current parallel
to B (CME)

Chiral Magnetic Effect



Kharzeev, McLerran, Warringa (2007)

For a local strong parity violation see e.g. 0909.1717 (STAR) and 1207.0900 (ALICE)

Electromagnetic fields



Huge electromagnetic fields, never observed before!

Black curves are from W.-T. Deng and X.-G. PRC 85, 044907



T_c ~ 170 MeV











Build an effective model for QCD at T_c < T < 2 T_c (did we lose anything?)



- Build an effective model for QCD at T_c < T < 2 T_c (did we lose anything?)
- Find hydrodynamic equations corresponding to the effective Lagrangian



- Build an effective model for QCD at T_c < T < 2 T_c (did we lose anything?)
- Find hydrodynamic equations corresponding to the effective Lagrangian
- Solve the hydro equations in gradient expansion, satisfying thermodynamic laws



- Build an effective model for QCD at T_c < T < 2 T_c (did we lose anything?)
- Find hydrodynamic equations corresponding to the effective Lagrangian
- Solve the hydro equations in gradient expansion, satisfying thermodynamic laws
- Extract phenomenological output for the heavy-ion collisions

Insight from the lattice

Spectrum of the Dirac operator

$$\hat{D}\psi_{\lambda} = \lambda\psi_{\lambda}$$



Chiral properties are described by near-zero modes

Insight from the lattice

Spectrum of the Dirac operator

$$\hat{D}\psi_{\lambda} = \lambda\psi_{\lambda}$$



- Chiral properties are described by near-zero modes
- There are two separated parts of the spectrum at intermediate temperatures!

- Euclidean functional integral for $\,{\rm SU}(N_c) \times U_{\rm em}(1)\,$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(D\!\!\!/ - im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

$$D = -i(\partial + A + gG + \gamma_5 A_5).$$

- Euclidean functional integral for $\,{\rm SU}(N_c) \times U_{\rm em}(1)\,$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(D\!\!\!/ - im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

where we define the Dirac operator as

$$D = -i(\partial + A + gG + \gamma_5 A_5).$$

integrate out quarks below a cut-off Dirac eigenvalue Λ.

- Euclidean functional integral for ${\rm\,SU}(N_{\rm c}) imes U_{\rm em}(1)$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(\not{\!\!D}-im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

$$D = -i(\partial + A + gG + \gamma_5 A_5).$$

- integrate out quarks below a cut-off Dirac eigenvalue Λ.
- add gauge-invariant terms to the Lagrangian to match the chiral anomaly

- Euclidean functional integral for ${\rm\,SU}(N_{\rm c}) imes U_{\rm em}(1)$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(\not{\!\!D}-im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

$$D = -i(\partial + A + gG + \gamma_5 A_5).$$

- integrate out quarks below a cut-off Dirac eigenvalue Λ.
- add gauge-invariant terms to the Lagrangian to match the chiral anomaly
- consider a pure gauge $A_{5\mu} = \partial_{\mu}\theta$ for the auxiliary axial field

- Euclidean functional integral for ${\rm\,SU}(N_{\rm c}) imes U_{\rm em}(1)$ is given by

$$\int D\bar{\psi}D\psi \exp\left\{-\int_{V} d^{4}x \,\bar{\psi}(\not{\!\!D}-im)\psi + \frac{1}{4}G^{a\mu\nu}G^{a}_{\mu\nu} + \frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right\},\,$$

$$D = -i(\partial + A + gG + \gamma_5 A_5).$$

- integrate out quarks below a cut-off Dirac eigenvalue Λ.
- add gauge-invariant terms to the Lagrangian to match the chiral anomaly
- consider a pure gauge $A_{5\mu} = \partial_{\mu}\theta$ for the auxiliary axial field
- and the chiral limit $m \rightarrow 0$

The total effective Euclidean Lagrangian reads as

$$\mathcal{L}_{E}^{(4)} = \frac{1}{4} G^{a\mu\nu} G^{a}_{\mu\nu} + \frac{1}{4} F^{\mu\nu} F_{\mu\nu} - j^{\mu} A_{\mu}$$
$$+ \frac{\Lambda^{2} N_{c}}{4\pi^{2}} \partial^{\mu} \theta \partial_{\mu} \theta + \frac{g^{2}}{16\pi^{2}} \theta G^{a\mu\nu} \widetilde{G}^{a}_{\mu\nu} + \frac{N_{c}}{8\pi^{2}} \theta F^{\mu\nu} \widetilde{F}_{\mu\nu}$$
$$+ \frac{N_{c}}{24\pi^{2}} \theta \Box^{2} \theta - \frac{N_{c}}{12\pi^{2}} \left(\partial^{\mu} \theta \partial_{\mu} \theta \right)^{2}$$

So we get an axion-like field with decay constant $f = \frac{2\pi}{\pi} \sqrt{N_c}$ and a negligible mass $m_{\theta}^2 = \lim_{V \to \infty} \frac{\langle Q^2 \rangle}{f^2 V} \equiv \chi(T)/f^2$.

Dynamical axion-like internal degree of freedom in QCD!

• From the quartic Lagrangian at $N_c = N_f = 1$ we get

$$\rho_5 = \frac{1}{2} \left(\frac{\Lambda}{\pi}\right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

• From the quartic Lagrangian at $N_c = N_f = 1$ we get

$$\rho_5 = \frac{1}{2} \left(\frac{\Lambda}{\pi}\right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

• Free quarks (see 0808.3382):

 $\Lambda = \pi \sqrt{\frac{2}{3}} \sqrt{T^2 + \frac{\mu^2}{\pi^2}} \propto \rho_{sphaleron}^{-1}$

• From the quartic Lagrangian at $N_c = N_f = 1$ we get

$$\rho_5 = \frac{1}{2} \left(\frac{\Lambda}{\pi}\right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

• Free quarks (see 0808.3382):

$$\Lambda = \pi \sqrt{\frac{2}{3}} \sqrt{T^2 + \frac{\mu^2}{\pi^2}} \propto \rho_{sphaleron}^{-1}$$

• Free quarks and a strong B-field: $\Lambda = 2\sqrt{|eB|}$

• From the quartic Lagrangian at $N_c = N_f = 1$ we get

$$\rho_5 = \frac{1}{2} \left(\frac{\Lambda}{\pi}\right)^2 \mu_5 + \frac{1}{3\pi^2} \mu_5^3$$

Free quarks (see 0808.3382):

$$\Lambda = \pi \sqrt{\frac{2}{3}} \sqrt{T^2 + \frac{\mu^2}{\pi^2}} \propto \rho_{sphaleron}^{-1}$$

- Free quarks and a strong B-field: $\Lambda = 2\sqrt{|eB|}$
- Dynamical fermions (1105.0385): $\Lambda \simeq 3 \,\mathrm{GeV} \gg \Lambda_{QCD}$

A "hidden" non-perturbative scale!

Considering EOM for the Minkowski effective Lagrangian and only the color-singlet states, we obtain:

$$\partial_{\mu}T^{\mu\nu} = F^{\nu\lambda}J_{\lambda} \,,$$

$$\partial_{\mu}J^{\mu} = 0 \,,$$

$$\partial_{\mu}J_5^{\mu} = CE^{\mu}B_{\mu} \,,$$

$$u^{\mu}\partial_{\mu}\theta + \mu_5 = 0\,,$$

Considering EOM for the Minkowski effective Lagrangian and only the color-singlet states, we obtain:



Similar to the superfluid dynamics!

 Solving the hydrodynamic equations in the gradient expansion, we obtain the constitutive relations:

$$T^{\mu\nu} = (\epsilon + P) u^{\mu}u^{\nu} + Pg^{\mu\nu} + f^2 \partial^{\mu}\theta \partial^{\nu}\theta + \tau^{\mu\nu} ,$$

$$J^{\mu} = \rho u^{\mu} + C \widetilde{F}^{\mu\kappa} \partial_{\kappa} \theta + \nu^{\mu} ,$$

$$J_5^{\mu} = f^2 \partial^{\mu} \theta + \nu_5^{\mu} \,.$$

 Solving the hydrodynamic equations in the gradient expansion, we obtain the constitutive relations:



Notice the additional current

An additional electric current induced by the θ -field:

$$j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial\theta \cdot B)$$



An additional electric current induced by the θ -field:

$$j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial\theta \cdot B)$$

• Chiral Magnetic Effect (electric current along B-field)



An additional electric current induced by the θ -field:

 $j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial\theta \cdot B)$

- Chiral Magnetic Effect (electric current along B-field)
- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)



An additional electric current induced by the θ -field:

$$j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial\theta \cdot B)$$

- Chiral Magnetic Effect (electric current along B-field)
- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)
- Chiral Dipole Wave (dipole moment induced by B-field)



An additional electric current induced by the θ -field:

$$j_{\lambda} = -C\mu_5 B_{\lambda} + C\epsilon_{\lambda\alpha\kappa\beta} u^{\alpha} \partial^{\kappa} \theta E^{\beta} - Cu_{\lambda} (\partial\theta \cdot B)$$

- Chiral Magnetic Effect (electric current along B-field)
- Chiral Electric Effect (electric current transverse to E-field and to both normal and superfluid velocities)
- Chiral Dipole Wave (dipole moment induced by B-field)
- The field θ(x) itself: Chiral Magnetic Wave (propagating imbalance between the number of left- and righthanded quarks)



Thank you for the attention!

and

Have a good time!

All comments on the papers are welcome! Also feel free to ask questions about the experimental observables during the coffee-breaks.