Nonequilibrium QFT approach to leptogenesis

Tibor Frossard

MPIK-Heidelberg



INTERNATIONAL MAX PLANEX RESEARCH SCHOOL



In collaboration with M. Garny, A. Hohenegger, A. Kartavtsev, and D. Mitrouskas

Baryogenesis via Leptogenesis

 Matter-antimatter asymmetry in the Universe

$$\eta = \frac{\mathsf{CMB} + \mathsf{BBN}}{n_{\gamma}} \Rightarrow (6.2 \pm 0.2) \, 10^{-10}$$



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 - Baryon number violation
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Leptogenesis

SM + 3 heavy Majorana neutrinos:

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2}\bar{N}_i \left(i\partial \!\!\!/ - M_i\right)N_i - h_{\alpha i}\bar{\ell}_{\alpha}\bar{\phi}P_R N_i - h_{i\alpha}^{\dagger}\bar{N}_i\bar{\phi}^{\dagger}P_L\ell_{\alpha}$$

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- Expanding Universe + N_i decay

Generation of an asymmetry:

	production	washout
N_i (inverse)decay	$\mathcal{O}\left(h^{4} ight)$	$\mathcal{O}\left(h^{2} ight)$
$\Delta L = 2$ scattering	$\mathcal{O}\left(h^{6} ight)$	$\mathcal{O}\left(h^{4}\right)$
top scattering	$\mathcal{O}\left(\lambda_t^2 h^4\right)$	$\mathcal{O}\left(\lambda_t^2 h^2\right)$
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 \Rightarrow into the Boltzmann equation for the lepton number!

This approach suffers from several problems:

Double counting problem



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 - \rightarrow thermal masses?
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Nonequilibrium Quantum Field Theory approach is needed!!

Nonequilibrium QFT approach

Basics of nonequilibrium QFT

Right-handed neutrinos propagator on CTP:

$$S_{ij}^{\alpha\beta}(x,y) = < T_{\mathcal{C}} N_i^{\alpha}(x) \bar{N}_j^{\beta}(y) >$$



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Schwinger-Dyson equation

$$\hat{S}^{-1}(x,y) = \hat{S}_{0}^{-1}(x,y) - \hat{\Sigma}(x,y) , \qquad \Sigma_{ij}^{\alpha\beta}(x,y) = -\frac{\delta i \Gamma_{2}}{\delta S_{ji}^{\beta\alpha}(y,x)}$$

$$\Gamma_{2} = \frac{2\mathsf{PI} \text{ effective}}{\operatorname{action}} = \underbrace{\left(\underbrace{s}_{ji} \right)^{s} + \underbrace{\left(\underbrace{s}_{ji} \right)^{s}}_{s} + \underbrace{\left(\underbrace{s}_{ji} \right)^{s}}_{s} \right)}$$

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 $\rightarrow \text{doubling of the d.o.f.: } \hat{S}(x,y) = \underbrace{\hat{S}_F(x,y)}_{\text{statistical}} - \frac{i}{2} \text{sign}_{\mathcal{C}}(x^0 - y^0) \underbrace{\hat{S}_{\rho}(x,y)}_{\text{spectral}}$

Quantum Boltzmann equation

- Wigner transform ↔ Fourrier transform wrt relative coordinates
- Gradient expansion ↔ Expansion in slow relative to fast time-scales (H/T)
- Quasiparticle approximation ↔ Narrow width approximation
- Kadanoff-Baym ansatz \leftrightarrow One-particle distribution function

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Quantum Boltzmann equation

$$\frac{d}{dt}(n_{\ell}(t) - n_{\bar{\ell}}(t)) = g_w \int \frac{d^4p}{(2\pi)^4} \mathrm{tr}\Big[\Big(\underbrace{\Sigma_{\ell <}(t,p)(1-f_{\ell}^p)}_{\text{gain term}} + \underbrace{\Sigma_{\ell >}(t,p)f_{\ell}^p}_{\text{loss term}} \Big) S_{\ell\rho}(t,p) \Big]$$

$$\Sigma_{\ell} = \ell - \underbrace{ \begin{pmatrix} & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & & \\$$

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The amplitudes incorporate thermal masses and thermal widths but,

- \rightarrow Self-energy contribution to the decay rate is missing !
- → $s \times s$ and $t \times t$ channels of the scattering processes are missing! ⇒ need to go **beyond** the free propagator approximation!

Free (thermal-) propagators into the lepton self-energy:

 $S_{\ell F} = (\frac{1}{2} - f_{\ell}^{p}) P_{L} p \operatorname{sign}(p^{0})(2\pi) \delta(p^{2} - m_{\ell}^{2}) \quad \Delta_{\phi F} = (\frac{1}{2} + f_{\phi}^{k}) \operatorname{sgn}(k^{0})(2\pi) \delta(k^{2} - m_{\phi}^{2})$ $S_{NF}^{ij} = \delta^{ij}(\frac{1}{2} - f_{N_{i}}^{q})(q + M_{i}) \operatorname{sign}(q^{0})(2\pi) \delta(q^{2} - M_{i}^{2})$



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In **equilibrium** the *full* statistical propagator \hat{S}_F can be expressed in terms of the flavor diagonal propagators \hat{S} and flavor off-diagonal self-energy $\hat{\Sigma}^{od}$:

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 \rightarrow assume to be valid out of equilibrium

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 \rightarrow insert the resummed Majorana propagator into the one-loop lepton self-energy:



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The internal Majorana propagators are automatically RIS-subtracted \Rightarrow don't need to neglect the quantum statistical factors!

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Numerical comparison with canonical approach



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