

Multi-Component Dark Matter Systems and Their Observation Prospects

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DESY THEORY WORKSHOP 2012

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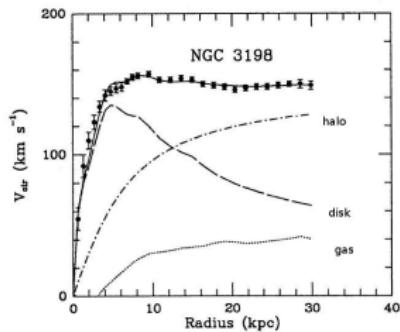
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 FOR PRECISION TESTS
OF FUNDAMENTAL
SYMMETRIES



MAX-PLANCK-GESELLSCHAFT

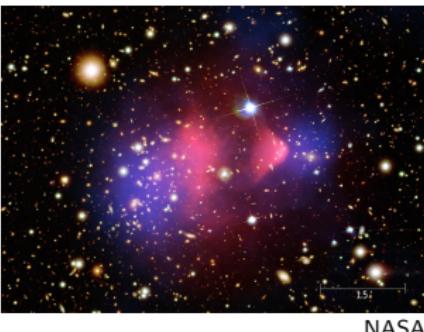
Hints for DM

Galaxy rotation curves

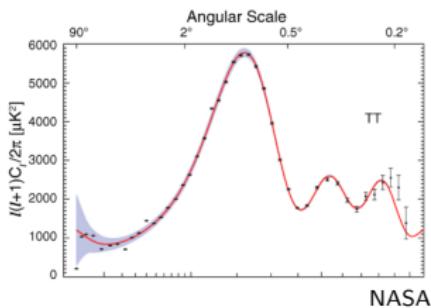


Begeman et al.,
Mon. Not. R. astr. Soc. (1991) **249**, 523

Bullet cluster

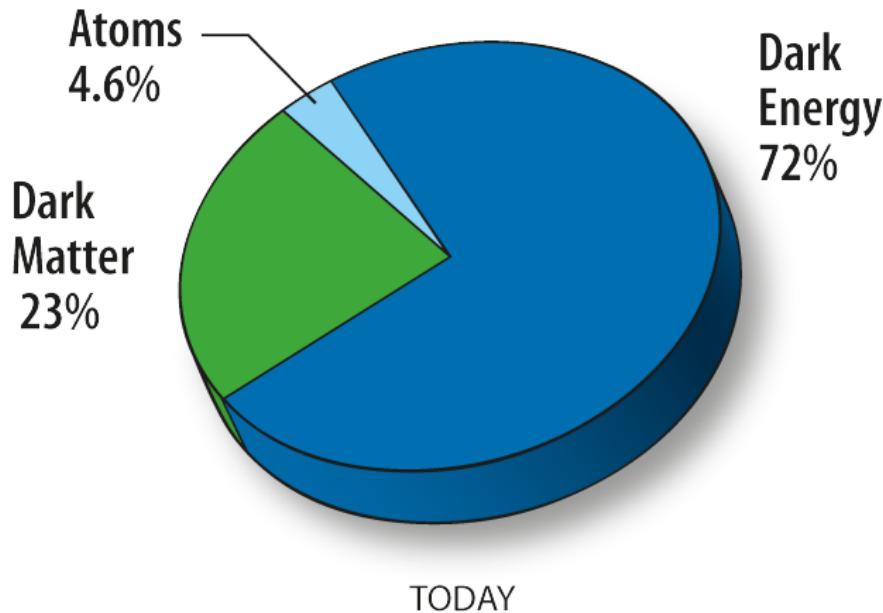


CMB



Consistent hints on all scales.

Content of the Universe



Wikipedia

Do you believe there exists only one kind of DM?

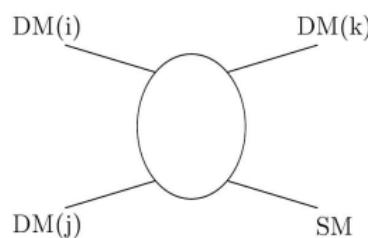
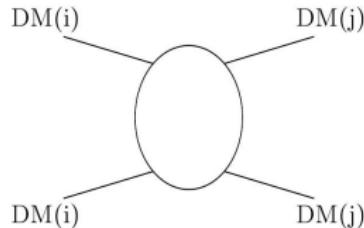
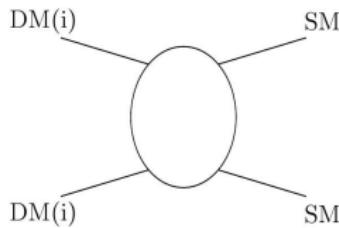
Multi-component DM systems

- How can a multi-DM system arise?
 - $Z_N (N \geq 4)$
 - Product of two or more Z_2 :
If $(Z_2)^l$ is unbroken, we have at least $K = l$ stable DM particles. In a kinematically fortunate situation, $2^l - 1$ stable DM particles may exist.
 - ...
- Non-standard annihilations

standard
annihilation

conversion

semi-
annihilation



Coupled Boltzmann equations

- standard annihilation

$$\frac{dY_i}{dx} = -0.264 g_*^{1/2} \left[\frac{\mu M_{\text{PL}}}{x^2} \right] \left\{ \langle \sigma(ii; X_i X'_i) v \rangle (Y_i Y_i - \bar{Y}_i \bar{Y}_i) \right.$$

$$+ \sum_{i>j} \langle \sigma(ii; jj) v \rangle \left(Y_i Y_i - \frac{Y_j Y_j}{\bar{Y}_j \bar{Y}_j} \bar{Y}_i \bar{Y}_i \right) - \sum_{j>i} \langle \sigma(jj; ii) v \rangle \left(Y_j Y_j - \frac{Y_i Y_i}{\bar{Y}_i \bar{Y}_i} \bar{Y}_j \bar{Y}_j \right)$$

$$+ \sum_{j,k} \langle \sigma(ij; kX_{ijk}) v \rangle \left(Y_i Y_j - \frac{Y_k Y_k}{\bar{Y}_k \bar{Y}_k} \bar{Y}_i \bar{Y}_j \right) - \sum_{j,k} \langle \sigma(jk; iX_{jki}) v \rangle \left(Y_j Y_k - \frac{Y_i Y_i}{\bar{Y}_i \bar{Y}_i} \bar{Y}_j \bar{Y}_k \right) \}$$

- DM conversion

- DM semi-annihilation

$$Y_i = n_i/s \quad \mu = \left(\sum_i m_i^{-1} \right)^{-1}$$

see also F. D'Eramo, J. Thaler, JHEP **06** (2010) 109, G. Belanger, K. Kannike, A. Pukhov, M. Raidal, JCAP **04** (2012) 010.

A fictive three-component DM system

stand. annihilation

$$\langle \sigma(ii; X_i X_i) v \rangle = \sigma_{0,i}$$

conversion

$$\langle \sigma(11; 22) v \rangle = \sigma_{0,12}$$

$$\langle \sigma(11; 33) v \rangle = \sigma_{0,13}$$

$$\langle \sigma(22; 33) v \rangle = \sigma_{0,23}$$

semi-annihilation

$$\langle \sigma(12; 3X_{123}) v \rangle = \sigma_{0,123}$$

$$\langle \sigma(23; 1X_{231}) v \rangle = \sigma_{0,231}$$

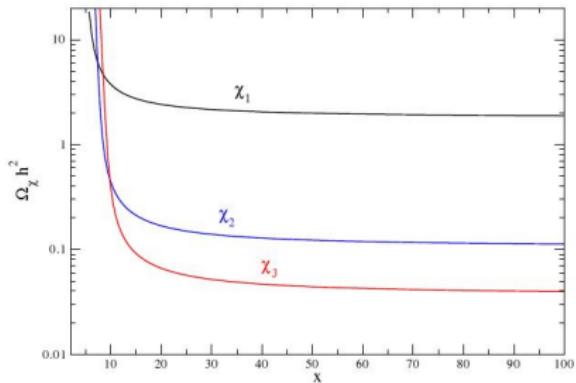
$$\langle \sigma(31; 2X_{312}) v \rangle = \sigma_{0,312}$$

$\times 10^{-9} \text{ GeV}^{-2}$

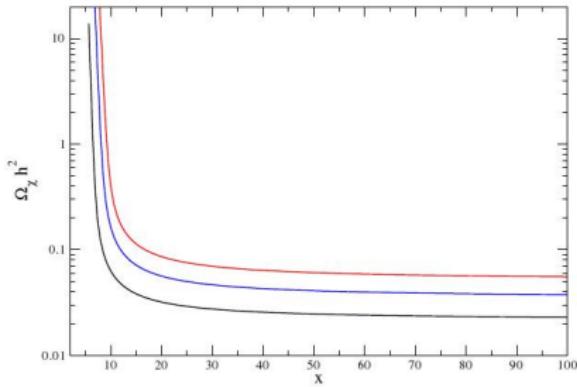
$$m_1 > m_2 > m_3 \quad \text{and} \quad m_2 + m_3 > m_1.$$

Temperature evolution

only standard



standard + conversion

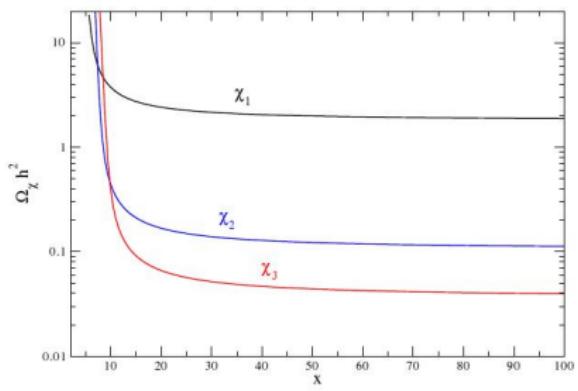


- standard: $\sigma_{0,1} = 0.1$, $\sigma_{0,2} = 2$, $\sigma_{0,3} = 6$
- conversion: $\sigma_{0,12} = \sigma_{0,13} = \sigma_{0,23} = 5.2$

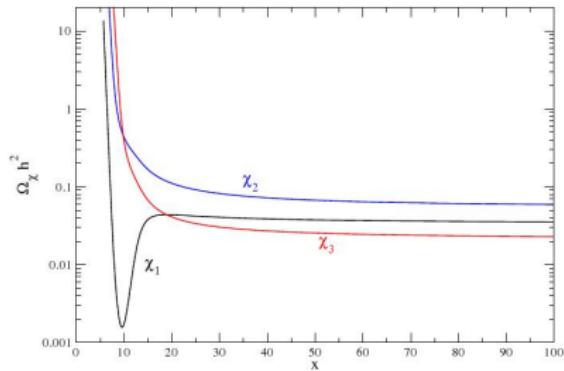
$$m_1 = 200 \text{ GeV}, m_2 = 160 \text{ GeV}, m_3 = 140 \text{ GeV}; x = \mu/T.$$

Temperature evolution

only standard



standard +
semi-annihilation



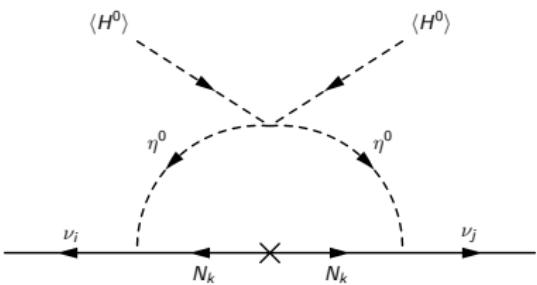
- standard: $\sigma_{0,1} = 0.1$, $\sigma_{0,2} = 2$, $\sigma_{0,3} = 6$
- conversion: $\sigma_{0,123} = \sigma_{0,312} = \sigma_{0,231} = 5.1$

$m_1 = 200 \text{ GeV}$, $m_2 = 160 \text{ GeV}$, $m_3 = 140 \text{ GeV}$; $x = \mu/T$.

The Ma model

Extend SM to $SU(2)_L \times U(1)_Y \times Z_2$,
 introducing $N_i \sim (1, 0; -)$ and
 $(\eta^+, \eta^0) \sim (2, 1/2; -)$.

Z_2 is exact $\rightarrow \langle \eta \rangle = 0$.



Higgs potential

$$\begin{aligned} V = m_1^2 H^\dagger H + m_2^2 \eta^\dagger \eta + \frac{1}{2} \lambda_1 (H^\dagger H)^2 + \frac{1}{2} \lambda_2 (\eta^\dagger \eta)^2 + \lambda_3 (H^\dagger H)(\eta^\dagger \eta) + \\ \lambda_4 (H^\dagger \eta)(\eta^\dagger H) + \frac{1}{2} \lambda_5 [(H^\dagger \eta)^2 + \text{H.c.}] \end{aligned}$$

Neutrino mass

$$(\mathcal{M}_\nu)_{ij} = \sum_k \frac{h_{ik} h_{jk} M_k}{16\pi^2} \left[\frac{m_R^2}{m_R^2 - M_k^2} \ln \frac{m_R^2}{M_k^2} - \frac{m_I^2}{m_I^2 - M_k^2} \ln \frac{m_I^2}{M_k^2} \right]$$

E. Ma, Phys. Rev. D73 (2006) 077301, arXiv:hep-ph/0601225

The Ma model – DM

N_R DM studied by

- Krauss, Nasri, Trodden, Phys. Rev. **D67** (2003) 085002
- Kubo, Ma, Suematsu, Phys. Lett. **B642** (2006) 18 ...

η DM studied by

- Barbieri, Hall, Rychkov, Phys. Rev. **D74** (2006) 015007
- Lopez Honorez, Nezri, Oliver, Tytgat, JCAP **02** (2007) 028
- Dolle, Su, Phys. Rev. **D80** (2009) 055012 ...

Promotion of Z_2 to $Z_2 \times Z'_2$

→ Promotion to a three-component DM system.

Extension of the Ma model

New particles

Add Majorana fermion χ and scalar ϕ with interaction $Y_k^X \chi N_k \phi$.

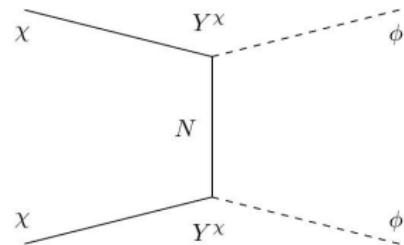
DM candidates

field	$SU(2)_L$	$U(1)_Y$	Z_2	Z'_2
N_i^c	1	0	—	+
$\eta = (\eta^+, \eta^0)$	2	1/2	—	+
χ	1	0	+	—
ϕ	1	0	—	—

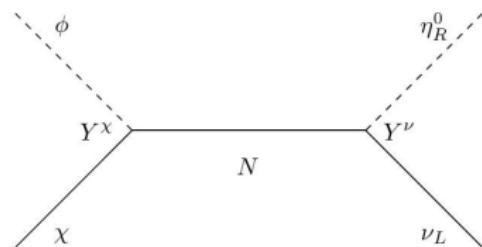
Our DM particles

η_R^0, χ, ϕ .

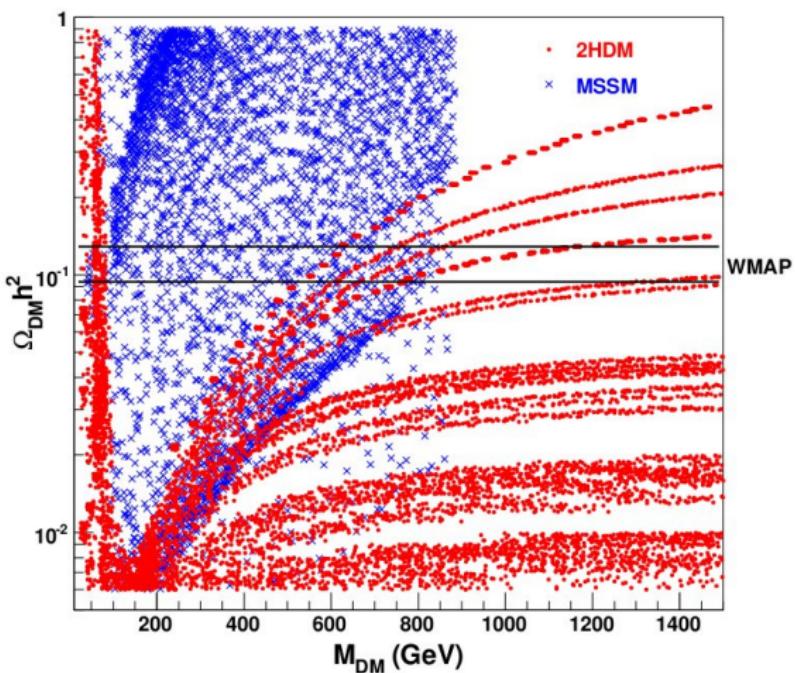
Conversion



Semi-annihilation

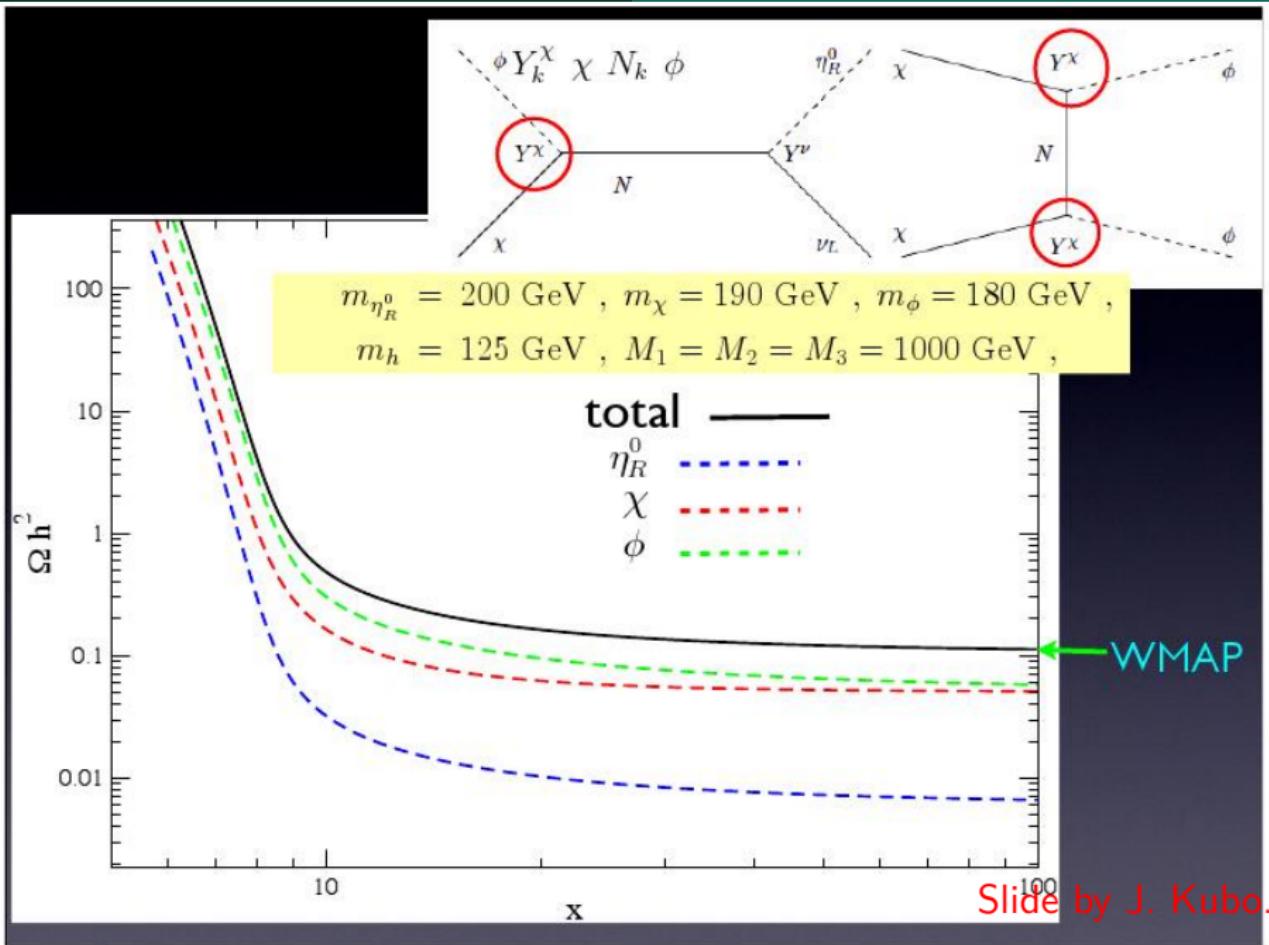


Inert Doublet Model



L. Lopez Honorez, E. Nezri, J. Oliver, M. Tytgat, JCAP 02 (2007) 028

Only low and high mass regimes are allowed.



Including all constraints

IDM (only η_R^0 is DM):

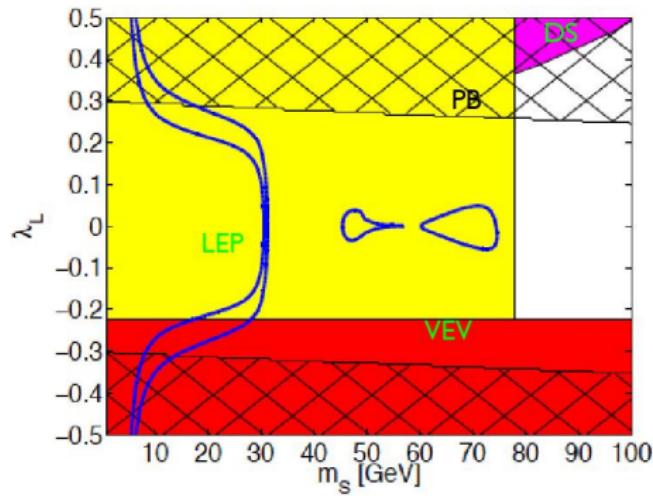
E. Dolle, S. Su, Phys. Rev. D80 (2009) 055012

$$60 \text{ GeV} \leq m_{\eta_R^0} \leq 80 \text{ GeV} \text{ or } m_{\eta_R^0} > 500 \text{ GeV}$$

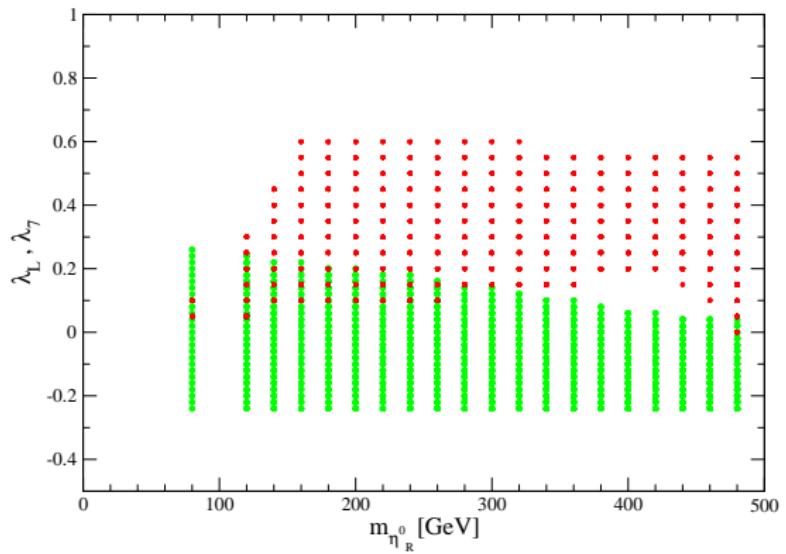
Parameters

$$\delta_1 = m_{\eta^\pm} - m_{\eta_R^0} = 10 \text{ GeV}$$

$$\delta_2 = m_{\eta_I^0} - m_{\eta_R^0} = 10 \text{ GeV}$$



With χ and ϕ



Parameters

green: λ_L and red: λ_7

$$\delta_1 = m_{\eta^\pm} - m_{\eta_R^0} = 10 \text{ GeV}$$

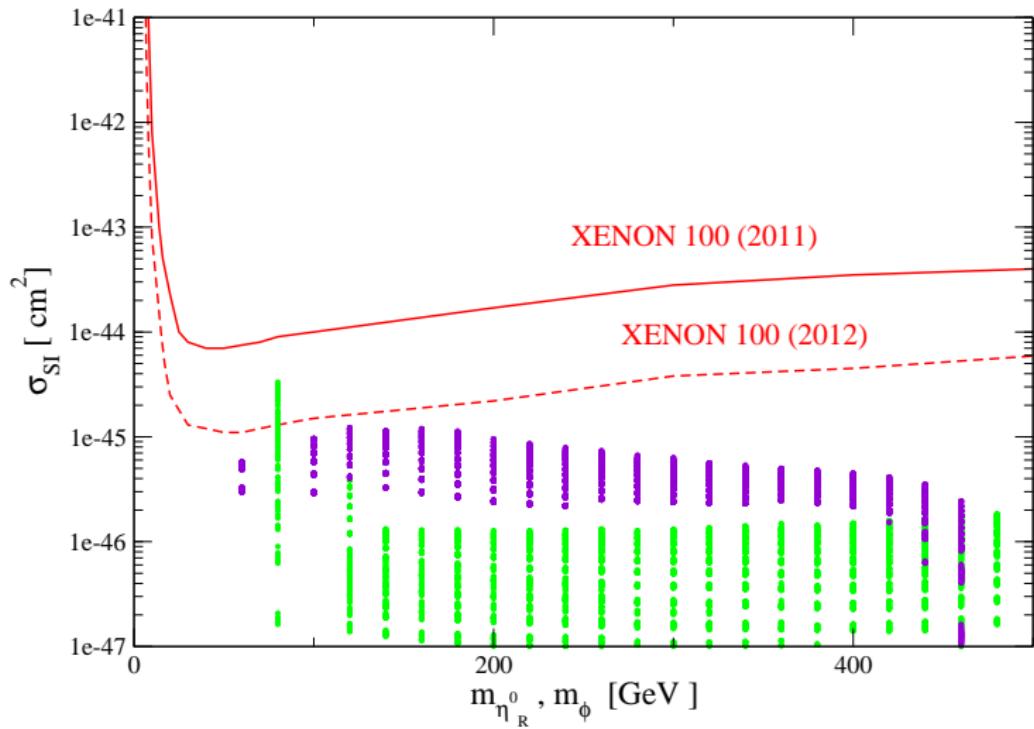
$$\delta_2 = m_{\eta_I^0} - m_{\eta_R^0} = 10 \text{ GeV}$$

$$m_\chi = m_{\eta_R^0} - 10 \text{ GeV}$$

$$m_\phi = m_{\eta_R^0} - 20 \text{ GeV}$$

$$M_k = 1000 \text{ GeV}$$

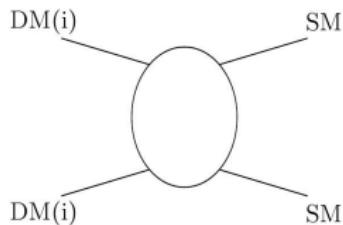
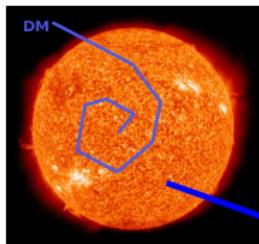
Direct detection



green: η_R^0 and violet: ϕ

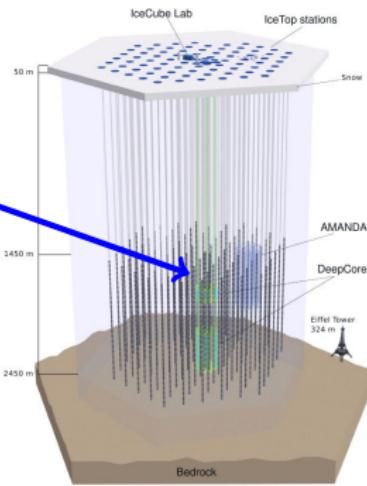
Indirect search at neutrino telescopes

DM can be captured and annihilated in the Sun, producing neutrinos that can escape from the Sun.



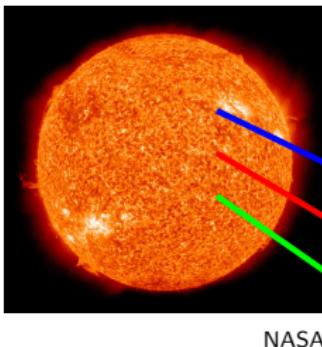
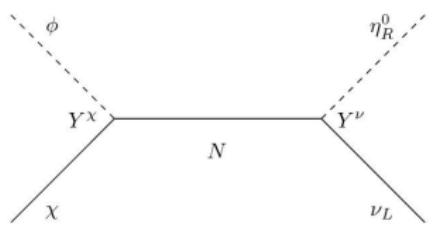
produces diffuse ν 's

IceCube detector



Monochromatic ν 's from the Sun

Semi-annihilations



Monochromatic neutrinos:

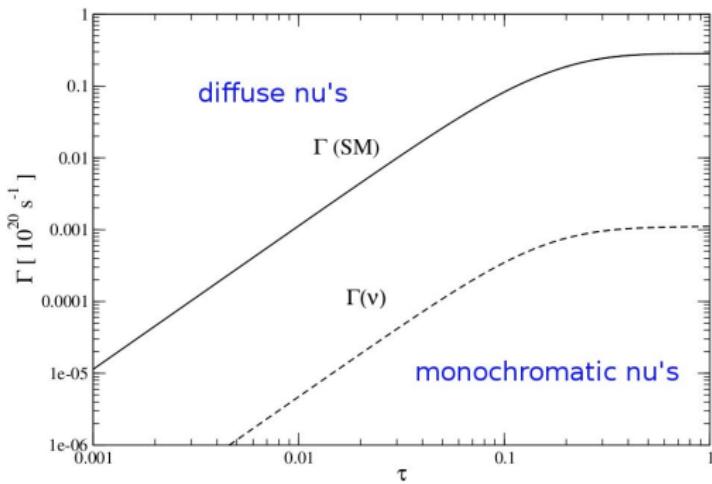
$$\begin{aligned} E_\nu &= m_\eta + m_\phi - m_\chi \\ E_\nu &= m_\eta + m_\chi - m_\phi \\ E_\nu &= m_\chi + m_\phi - m_\eta \end{aligned}$$

Time evolution of the numbers of DM in the Sun:

$$\begin{aligned} \dot{N}_\eta &= C_\eta - C_A(\eta\eta \leftrightarrow \text{SM})N_\eta^2 - C_A(\eta\eta \leftrightarrow \phi\phi)N_\eta^2 - C_A(\eta\chi \leftrightarrow \phi\nu_L)N_\eta N_\chi \\ &\quad - C_A(\eta\phi \leftrightarrow \chi\nu_L)N_\eta N_\phi + C_A(\phi\chi \leftrightarrow \eta\nu_L)N_\chi N_\phi, \\ \text{analog for } \dot{N}_\chi \text{ and } \dot{N}_\phi \end{aligned}$$

Time evolution of the annihilation rates

$$\Gamma(\text{SM}) = C_A(\eta\eta \leftrightarrow \text{SM}) N_\eta^2 / 2 + C_A(\phi\phi \leftrightarrow \text{SM}) N_\phi^2 / 2$$



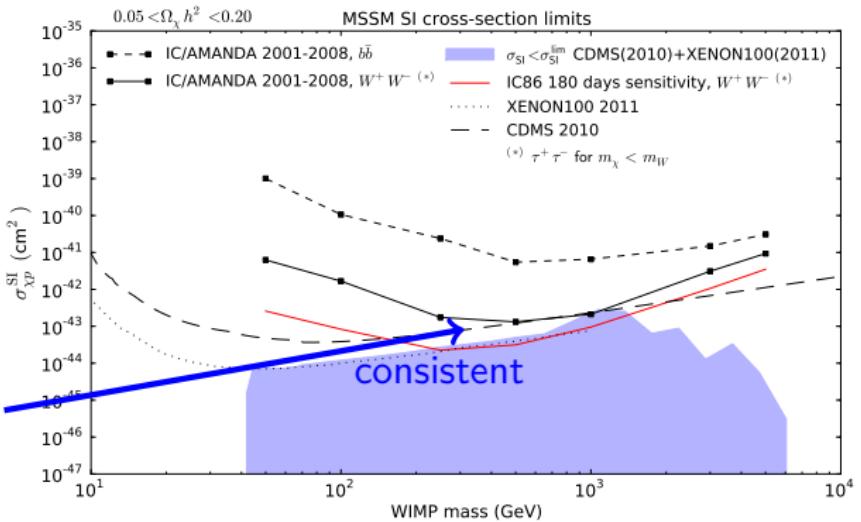
$$\Gamma(\nu) = C_A(\eta\phi \leftrightarrow \chi\nu) N_\eta N_\phi + C_A(\eta\chi \leftrightarrow \phi\nu) N_\eta N_\chi + C_A(\chi\phi \leftrightarrow \eta\nu) N_\chi N_\phi$$

Input parameters: $m_{\eta_R^0} = 200 \text{ GeV}$, $m_\chi = 190 \text{ GeV}$, $m_\phi = 180 \text{ GeV}$,
 $m_h = 125 \text{ GeV}$, $M_k = 1000 \text{ GeV} \rightarrow E_\nu \approx 200 \text{ GeV}$.

Limits from neutrino telescopes

- IceCube/Amanda
- ANTARES
- Super-K

$$\Gamma(\text{SM}) \approx 0.3 \times 10^{20} / \text{sec}$$



$\Gamma(\text{monochromatic } \nu) \approx 0.001 \times 10^{20} \text{ sec}$
 $\rightarrow 0.05 \text{ events per year at Ice Cube}$

Conclusions

- Non-standard annihilations of DM can play an important role for the relic abundance of DM and for indirect observation of DM.
- The detection of monochromatic neutrinos from the Sun may give a hint for multi-component DM in the Universe.

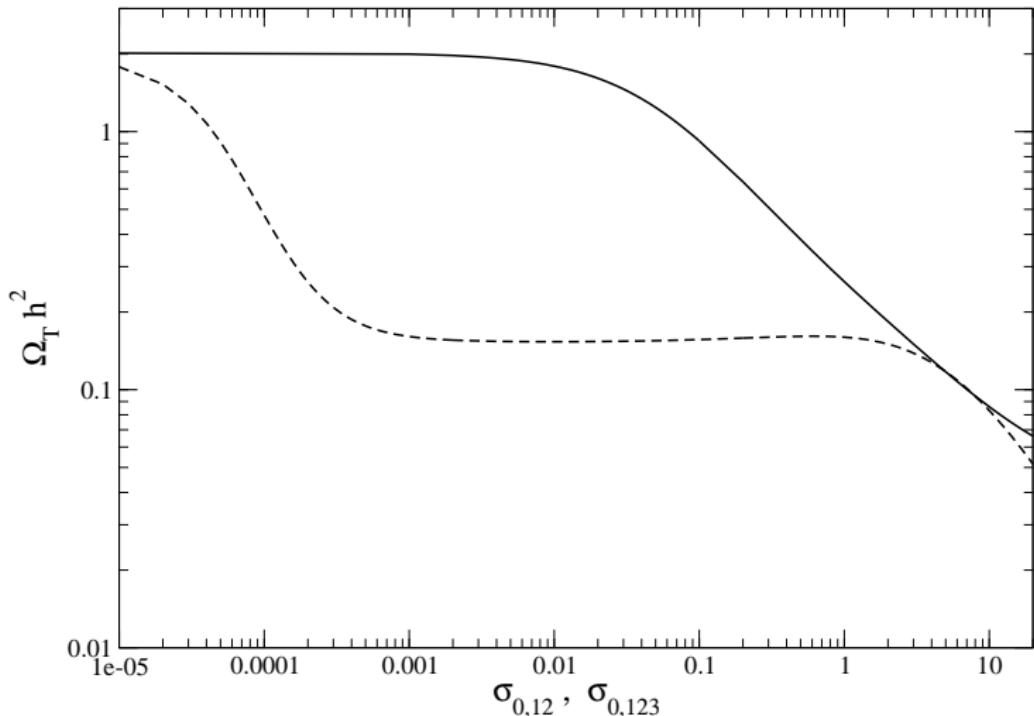
Conclusions

- Non-standard annihilations of DM can play an important role for the relic abundance of DM and for indirect observation of DM.
- The detection of monochromatic neutrinos from the Sun may give a hint for multi-component DM in the Universe.

Thanks for your attention!

Backup slides

Dependence on the non-standard annihilations



dashed = semi-annihilation, full = conversion.