$\theta_{13}^{\text{PMNS}} = \frac{\theta_C}{\sqrt{2}}$ from Grand Unified Theories

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 $\theta_{13}^{\text{PMNS}} = \theta_C / \sqrt{2} \text{ from GUTs}$

The Reactor Angle $\theta_{13}^{\text{PMNS}}$ and Model Building

until 2011 One of Flavour Model Builders' Favourites: Tri-Bimaximal Lepton Mixing

$$\sin^2\theta_{23}^{\text{PMNS}} = \frac{1}{2}, \ \sin^2\theta_{12}^{\text{PMNS}} = \frac{1}{3}, \ \theta_{13}^{\text{PMNS}} \approx 0^{\circ}$$

June 2011 Hints on $\theta_{13}^{MNS} \sim \theta_C$ from T2K

March 2012 Daya Bay: $\theta_{13}^{MNS} = 8.8^{\circ} \pm 1.0^{\circ}$ Striking resemblence with

$$\frac{\theta_{C}}{\sqrt{2}} = 9.2^{\circ}$$

with the Cabibbo angle θ_C

Question for Model Building

What can we do with the circumstance

$$heta_{13}^{\mathsf{PMNS}} = 8.8^{\circ} pprox rac{ heta_{\mathcal{C}}}{\sqrt{2}} = 9.2^{\circ}?$$

Do we

(a) have to fit this relation or(b) get this naturally out as some model prediction?

Remarks: Related Works

Cabibbo Haze

Use of Wolfenstein parameter λ_C as expansion parameter also in $U_{\rm PMNS}$ leads to

$$\theta_{13}^{\text{PMNS}} = \mathcal{O}(1)\theta_C.$$

[Datta, Everett, Ramond '05]

Quark-Lepton Complementarity

Assume $U_{CKM}U_{PMNS} = U_{BM}$, so e.g.

$$\theta_{12}^{\mathsf{PMNS}} + \theta_{\mathcal{C}} = 45^{\circ}$$

can also have a similar effect on $\theta_{13}^{\text{PMNS}}$.

[Minakata, Smirnov '04]

This talk: Simple conditions for GUT flavour models that lead to

$$\theta_{13}^{\text{PMNS}} \approx \frac{\theta_C}{\sqrt{2}}$$

as a direct consequence.

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Strategy to make $\theta_{13}^{\text{PMNS}} = \frac{\theta_c}{\sqrt{2}}$ rigid



Quick Rundown of Strategy:

- $\theta_{13}^{\text{PMNS}}$ dominantly coming from charged lepton contributions.
- θ_C mainly from mixing in Y_d .
- Y_e and Y_d are related by some minimal set of GUT field operators.

Condition I or "How to get $\sqrt{2}$?"



- U_{PMNS} product of two unitary matrices
- Leading order in small angles [e.g. Antusch, King '05]

$$\begin{split} s_{23}^{\text{PMNS}} e^{-i\delta_{23}^{\text{PMNS}}} &= s_{23}^{\nu} e^{-i\delta_{23}^{\nu}} - \theta_{23}^{e} c_{23}^{\nu} e^{-i\delta_{23}^{e}} \\ \theta_{13}^{\text{PMNS}} e^{-i\delta_{13}^{\text{PMNS}}} &= -\theta_{12}^{e} e^{-i\delta_{12}^{e}} (s_{23}^{\nu} e^{-i\delta_{23}^{\nu}} - \theta_{23}^{e} c_{23}^{\nu} e^{-i\delta_{23}^{e}}) \\ &+ \theta_{13}^{\nu} e^{-i\delta_{13}^{\nu}} - \theta_{13}^{e} c_{23}^{\nu} e^{-i\delta_{13}^{e}} \\ \theta_{13}^{e} &\approx 0 \Rightarrow \\ \hline \theta_{13}^{\text{PMNS}} &= \theta_{12}^{e} s_{23}^{\text{PMNS}} \end{split}$$

 θ_{13}^{ν}

Condition I or "How to get $\sqrt{2}$?"



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Condition II or "From where θ_C ?"



GUT relations between Y_e and $Y_d \Rightarrow \theta_C$ in Y_d

• Simplest possibility: Require

$$\theta_{12}^d = \theta_C = 13^\circ$$
, e.g. by requiring $\theta_{ij}^u \ll \theta_{12}^d$.

• Quark mixing sum rule (with $\theta_{13}^u, \theta_{13}^d \approx 0$):

$$heta_{12}^{d} pprox \left| heta_{12}^{\mathsf{CKM}} - rac{ heta_{13}^{\mathsf{CKM}}}{ heta_{23}^{\mathsf{CKM}}} e^{-i\delta^{\mathsf{CKM}}}
ight| pprox 12.0^{\circ}$$

[Antusch, King, Malinsky, Spinrath '09]

Condition II or "From where θ_C ?"

$$\begin{array}{c} \theta_{12}^{d} \approx \theta_{C} \end{array} Y_{u} \qquad \underbrace{U_{CKM} = U_{u}^{\dagger} U_{d}}_{\downarrow} Y_{d} \\ & & & \downarrow \\ & & & \downarrow \\ m_{\nu} \qquad \underbrace{U_{PMNS} = U_{e}^{\dagger} U_{\nu}}_{\downarrow} Y_{e} \end{array}$$

GUT relations between Y_e and $Y_d \Rightarrow \theta_C$ in Y_d

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Relations between Yukawa Couplings and Condition III

Example: SU(5) with $\bar{\mathbf{5}} = (d_R^{\dagger}, L)$, $\mathbf{10} = (Q, u_R^{\dagger}, e_R^{\dagger})$, $\bar{\mathbf{5}}_H = (H_d, ...)$

renormalizable:

 $\mathcal{L}_{\text{Yukawa}} \supset \quad y \quad \bar{\mathbf{5}}_2 \mathbf{10}_2 \bar{\mathbf{5}}_H \implies y_\mu = y_s$

non-renormalizable:

$$\mathcal{L}_{\text{Yukawa}} \supset \quad \frac{\mathcal{H}_{\text{GUT}}}{M} \quad \bar{\mathbf{5}}_2 \mathbf{10}_2 \bar{\mathbf{5}}_H \quad \Longrightarrow \ y_\mu = \left\{\mathbf{1}, \frac{3}{2}, 6, \dots\right\} y_s$$

[Georgi, Glashow '74] [Antusch, Spinrath '09]

 \Rightarrow To stay predictive:

Only one such operator per matrix element of Y_e/Y_d

Condition IV: GUT Scale Structure of Yukawa Matrices

$$\begin{array}{c} \theta_{12}^{d} \approx \theta_{C} \\ \theta_{12}^{o} \approx \theta_{C} \end{array} \qquad \begin{array}{c} Y_{u} \\ Y_{u} \\ & \downarrow_{CKM} = U_{u}^{\dagger} U_{d} \\ & \downarrow_{d} \\ & \downarrow_$$

- Condition II: $\theta_{12}^d = \theta_C$
- Condition III: Predictive relation between Y_e and Y_d

$$\Rightarrow$$
 Condition IV: $\theta_{12}^d = \theta_{12}^e$

Condition IV: GUT Scale Structure of Yukawa Matrices

- Condition II: $\theta_{12}^d = \theta_C$
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Condition IV in Pati-Salam

- Concentrate on 1-2 submatrices, neglect third generation mixing
- In Pati-Salam (PS):

$$Y_d = \begin{pmatrix} d & b \\ a & c \end{pmatrix} \Rightarrow \quad Y_e = \begin{pmatrix} c_d d & c_b b \\ c_a a & c_c c \end{pmatrix}$$

with group-theoretical Clebsch factors c_a, c_b, c_c, c_d .

• Mixing angles at leading order:

$$\theta_{12}^{d} = \left| \frac{b}{c} \right|, \ \theta_{12}^{e} = \left| \frac{c_b}{c_c} \right| \theta_{12}^{d}$$

[Antusch, Maurer '11]

$$\Rightarrow$$
 Condition IV in PS: $|c_b| = |c_c|$

Condition IV in SU(5)

• In SU(5):

$$Y_d = \begin{pmatrix} d & b \\ a & c \end{pmatrix} \Rightarrow \quad Y_e = \begin{pmatrix} c_d d & c_a a \\ c_b b & c_c c \end{pmatrix}$$

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$$\theta_{12}^d = \left| \frac{b}{c} \right|, \ \theta_{12}^e = \left| \frac{c_a}{c_c} \frac{a}{c} \right|$$

 \Rightarrow Have to relate *a* and *b* somehow

• Straightforward choice: Y_d, Y_e symmetric

$$\Rightarrow |a| = |b|, |c_a| = |c_b| = |c_c|$$

Blemish: Need $d \neq 0$ to get realistic mass ratios

• Second possibility is a bit less direct ...

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• Mixing angles and eigenvalues at leading order:

$$\theta_{12}^d = \left| \frac{b}{c} \right|, \ \theta_{12}^e = \left| \frac{c_a}{c_c} \frac{a}{c} \right|, \ y_s = |c|, \ y_d = \left| \frac{ab}{c} \right|$$

Yields

$$\theta_C^2 \approx \frac{1}{19 \pm 1} = \frac{m_d}{m_s} \approx \frac{y_d}{y_s} = \left| \frac{ab}{c^2} \right| = \left| \frac{a}{b} \right| \theta_C^2$$

[Gatto, Sartori, Tonin '68] [Leutwyler '96]

$$\Rightarrow |a| \approx |b|, |c_a| = |c_c|$$

• Only one consistent choice: $c_c = c_a = 6$, $c_b = \frac{1}{2}$.

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Corrections

Several corrections to
$$\theta_{13}^{\text{PMNS}} = \frac{\theta_C}{\sqrt{2}}$$
 apply:

- Deviation of $\theta_{23}^{\rm PMNS}$ from 45°: up to $\sim 0.8^{\circ}$
- Small angle approximation: $\sim 0.6^{\circ}$
- RG running of m_{ν} (normal hierarchy): $\leq 0.2^{\circ}$ for tan $\beta < 50$ (inverted hierarchy): ≈ 0

(for hierarchical neutrino masses)

• More model dependent:

 Y_{ν} effects in RG running, canonical normalisation

Daya Bay:
$$heta_{13}^{MNS}=8.8^\circ\pm1.0^\circ$$

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• Only a few conditions have to be met to obtain $\theta_{13}^{\text{PMNS}} = \frac{\theta_C}{\sqrt{2}}$



Corrections beyond leading order under control

Thank you for your attention!

Backup: Information on θ_{12}^{ν} from Lepton \mathcal{GP}

Using

$$\begin{split} \theta_{13}^{\text{PMNS}} &\approx s_{23}^{\text{PMNS}} \theta_{12}^{e} \\ \theta_{12}^{\text{PMNS}} &\approx \theta_{12}^{\nu} - s_{23}^{\text{PMNS}} \theta_{12}^{e} \cos \delta^{\text{PMNS}} \end{split}$$

one can get

$$\theta_{12}^{\nu} \approx \theta_{12}^{\rm PMNS} - \frac{\theta_{C}}{\sqrt{2}} \cos \delta^{\rm PMNS}$$



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Backup: Small Angle Approximation and $\theta_{12}^d \neq \theta_C$

In SU(5):

$$Y_d = \begin{pmatrix} 0 & b \\ a & c \end{pmatrix} \Rightarrow \quad Y_e = \begin{pmatrix} 0 & 6 a \\ \frac{1}{2} b & 6 c \end{pmatrix}$$

Fitting y_e/y_μ and θ_C to exact diagonalisation:

$$\theta_{13}^{\text{PMNS}}=9.8^\circ$$
 instead of 9.2°

In PS:

$$Y_d = \begin{pmatrix} d & b \\ b & c \end{pmatrix} \Rightarrow \quad Y_e = \begin{pmatrix} 9d & -3b \\ -3b & -3c \end{pmatrix}$$

with Quark Mixing Sum Rule $\theta_{12}^d = \left| \theta_{12}^{CKM} - \frac{\theta_{13}^{CKM}}{\theta_{23}^{CKM}} e^{-i\delta^{CKM}} \right| = 12.0^\circ$: $\left| \theta_{13}^{PMNS} = 8.6^\circ \right|$

Backup: Deviations of $\theta_{23}^{\text{PMNS}}$ from 45°

Actual formula:

$$\theta_{13}^{\mathsf{PMNS}} \approx \theta_{C} s_{23}^{\mathsf{PMNS}}$$

Current global fit data:

$$heta_{23}^{\mathsf{PMNS}} = 46.1^\circ \pm 3.4^\circ$$

[Schwetz, Tortola, Valle '11]

 \Rightarrow Introduces \sim 6% error in relation for $\theta_{13}^{\text{PMNS}}$

Other global fits find even larger uncertainties

$$\theta_{23}^{\text{PMNS}} = 40.4^{\circ}_{-1.6^{\circ}}^{+4.6^{\circ}}$$

[Fogli, Lisi, Marrone, Palazzo, Rotunno '11]

$$\Rightarrow$$
 Shift for $\theta_{13}^{\text{PMNS}}$ of -0.8°

• Reminder: $\theta_{13}^{MNS} = 8.8^{\circ} \pm 1.0^{\circ}$

Backup: Renormalisation Group Running

Starting from

$$heta_{13}^{\mathsf{PMNS}} pprox heta_C s_{23}^{\mathsf{PMNS}}$$

at the GUT scale, the correction to the relations takes the form

$$\begin{split} \theta_{13}^{\mathsf{PMNS}}|_{\mathcal{M}_{Z}} &= \theta_{C} \boldsymbol{s}_{23}^{\mathsf{PMNS}}|_{\mathcal{M}_{\mathsf{GUT}}} + \Delta \theta_{13}^{\mathsf{PMNS}} \\ &= \theta_{C} (\boldsymbol{s}_{23}^{\mathsf{PMNS}}|_{\mathcal{M}_{Z}} - \Delta \boldsymbol{s}_{23}^{\mathsf{PMNS}}) + \Delta \theta_{13}^{\mathsf{PMNS}} \end{split}$$

with $\Delta x = x|_{M_Z} - x|_{M_{GUT}}$.

• Considering strong u-mass hierarchy $m_{
u,\text{lightest}} = 0$, one finds

Normal: $\Delta \theta_{13}^{\text{PMNS}} \approx \theta_C \Delta s_{23}^{\text{PMNS}} + 0.2^{\circ} \cos(\delta^{\text{PMNS}} - \phi_2^{\text{PMNS}}) \left(\frac{\tan \beta}{50}\right)^2$ Inverted: $\Delta \theta_{13}^{\text{PMNS}} \approx \theta_C \Delta s_{23}^{\text{PMNS}}$

[Antusch, Kersten, Lindner, Ratz '03]

- More model dependent:
 - Y_{ν} effects in RG running and canonical normalisation effects

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