



Maximal R-symmetry violating amplitudes in type IIB superstring theory

(based on [arXiv:1204.4208](https://arxiv.org/abs/1204.4208))

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...string scattering amplitudes in a flat background..



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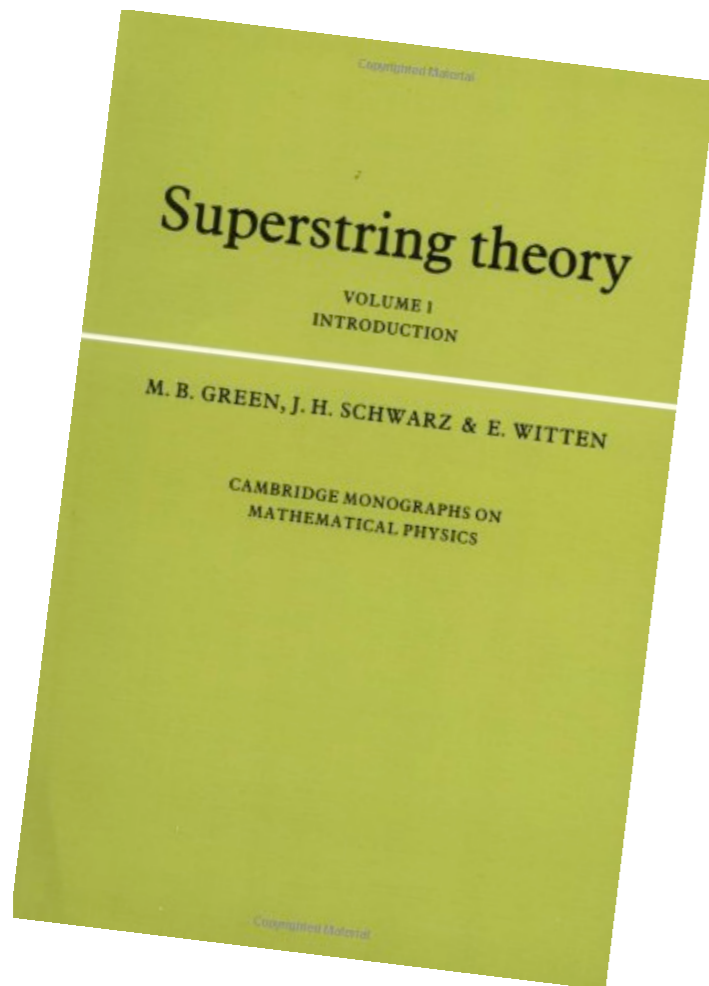
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- calculate correlation functions
- integrate over moduli



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(but go see Schlotterer’s talk!)



Goal: simple answers

- symmetry vs simplicity \rightarrow most (manifestly) symmetric answers are the simplest
- golden standard of simple scattering amplitudes:
four dimensional MHV amplitudes in tree level Yang-Mills



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- similar insight into perturbative string theory?
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D=4 vs D>4?



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- every leg has (k, η_a) coordinates \rightarrow on-shell superspace
- chirality constraints easy, reality conditions hard



Superamplitudes

- promote each leg of an amplitude $A(k_i) \rightarrow A(\{k_i, \eta_i\})$
- component amplitudes by fermionic integration



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- simple formulation of the on-shell susy Ward identities

$$\left. \begin{aligned} Q &= \sum_i Q_i \\ \bar{Q} &= \sum_i \bar{Q}_i \end{aligned} \right\} \boxed{QA = \bar{Q}A = 0}$$

**exact,
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- 3 massless particle exception: $\frac{3}{4} \mathcal{D}$ fermionic coordinates



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- delta function only: **simplest solutions** $\left\{ \begin{array}{l} \text{four massless legs} \\ \text{one massive, two massless} \end{array} \right.$
- immediate four point tree amplitudes /w massless matter:

$$A_{D=8, \text{YM}} \sim \frac{\delta^8(Q)}{st} \quad A_{D=10, \text{Grav.}} \sim \frac{\delta^{16}(Q)}{stu}$$



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- also three points, five points, on-shell recursion in paper



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• fermionic weight of amplitudes is related to $U(1)_R$ charge

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• $U(1)_R$ in $D=10 \rightarrow$ part of $SL(2,R)/U(1)$ of IIB not conserved

\rightarrow simple superamplitudes?



Massless on-shell superspace in D=10, type IIB

D=10: 256 states in the fundamental multiplet

$$\phi(\eta) = \phi_0 + \phi^a \eta_a + \dots + \bar{\phi}_0 (\eta)^8$$

field content:

	bosonic		fermionic
0	$\underline{1}_4$		
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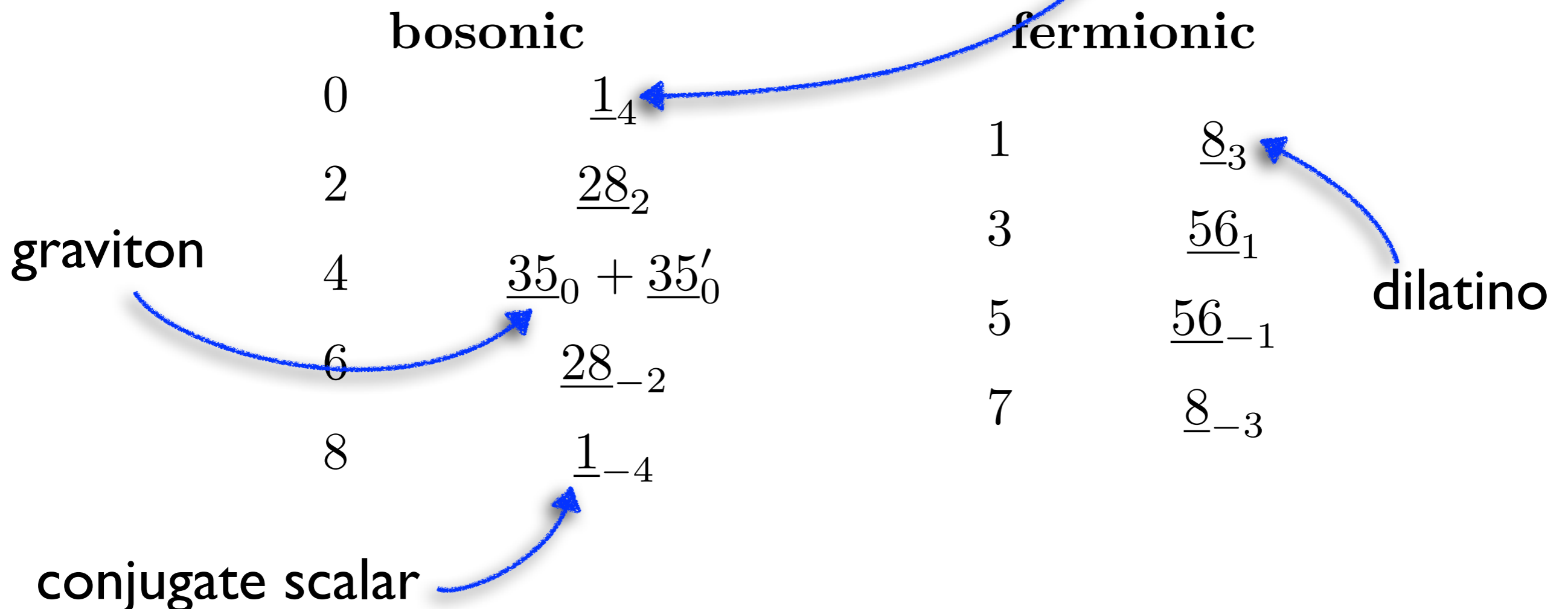
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 - only **massive** particle poles ($n>4$)
 - → **no** poles in field theory limit

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$$\tilde{A}_n^{\text{MRV}} = (g\alpha'^2)^{n-2} (\alpha'^3 c_0 + \alpha'^4 c_1 + \alpha'^5 c_2 + \mathcal{O}(\alpha'^6))$$

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$c_1 \rightarrow \left(\sum_i k_i\right)^2 = 0 \rightarrow [s_{12}^2]_n$

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- using results from: [Kawai-Lewellen-Tye, 86], [Stieberger-Taylor, 06], [Huber-Maitre, 07] + equivalence theorem



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- only odd zeta's: [Stieberger, 09], [Schlotterer-Stieberger, 12]



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“Soft dilaton theorem” [Ademollo et.al., 75], [Shapiro, 75]



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- differential operator annihilates gravitational coupling
- \rightarrow relates c_i for various multiplicities, up to degeneracy



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include more **stringy**
symmetries?



MRV amplitudes in field theory limit $A_n = \tilde{A}_n \delta^{16}(Q)$

more stringy symmetry in IIB: $SL(2, \mathbb{Z})$

[Green-Gutperle, 97], [Green et.al., 97-12]

$$\tau_b = a_b + i e^{-\phi_b} \quad f_{\beta}^k(\tau_b, \bar{\tau}_b) = \sum_{(l,m) \neq (0,0)} (l + m\tau_b)^{k-\beta} (l + m\bar{\tau}_b)^{-k-\beta}$$



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$$A_n^{\text{MRV}} \propto \delta^{16}(Q) (\alpha'^2 g)^{n-2} \left(\alpha'^3 f_{\frac{3}{2}}^{n-4} + \alpha'^5 f_{\frac{5}{2}}^{n-4} ([s_{12}^2]_n) + \mathcal{O}(\alpha'^6) \right)$$



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- analytic part of amplitude: the “no logs”-part
- guess for next order exists
- much work: relation to effective action, better normalization...



Summary, outlook

remarkably simple class of amplitudes in type IIB superstrings



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more examples of applications / more explicit amplitudes?

- how deep does analogy to MHV go?
- worldsheet picture? (pure spinor?)
- IIA? D=11? open strings? → `constrained superspaces`?



Your Question
Here?