

# Nonlinear processes in atoms

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# What seeding might do for us ...

- Category A: smaller bandwidth → “long”, spectrally narrow pulses
  - “Resonant by design”  
→ x-ray driven Rabi oscillations
  - “Resonant by accident”  
→ nonsequential (direct) two-photon absorption  
→ enhanced multiple ionization
- Category B: shorter pulses → transform-limited sub-femtosecond pulses
  - Suppression of electronic damage in coherent diffractive imaging
  - Electronic coherence and correlation in the time domain
  - Imaging of electronic quantum motion

# Acknowledgments I

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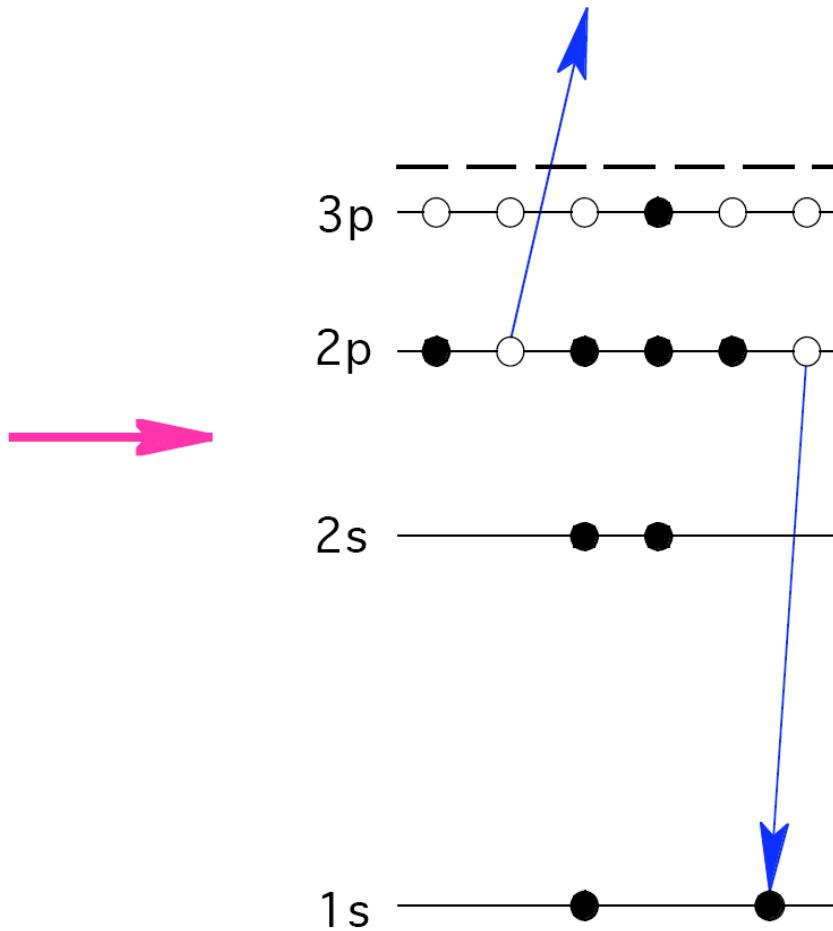
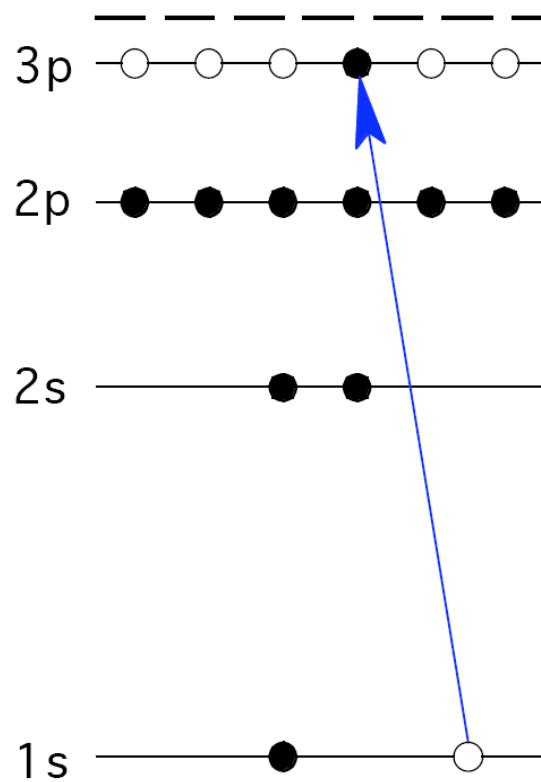
<sup>17</sup> Tohoku University, Sendai 980-8577, Japan



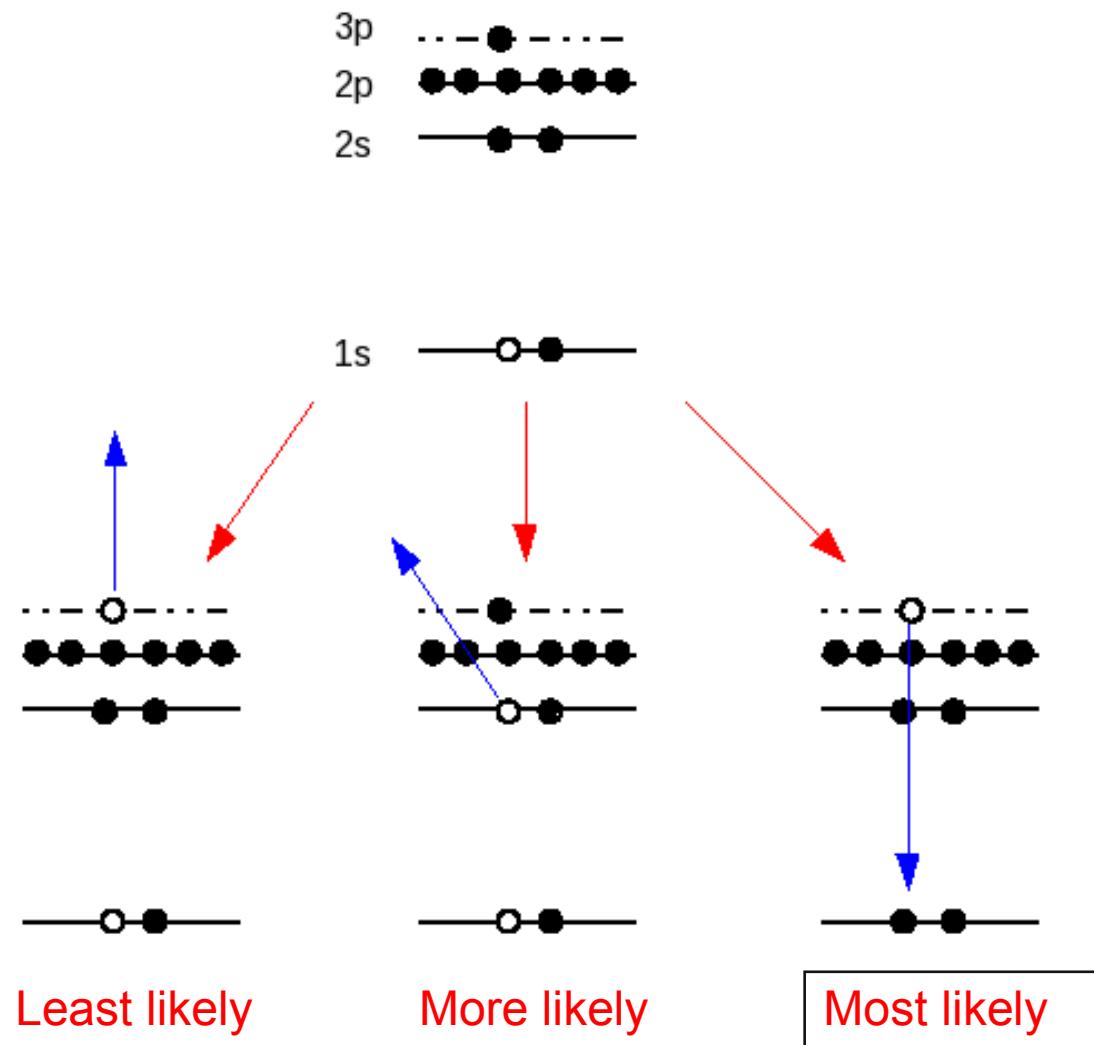
## “Resonant by design”

- X-ray driven Rabi oscillations involving inner-shell electrons
- Modification of bound electronic dynamics leads to modification of Auger electron line profile
- Step towards the development of nonlinear spectroscopy in the x-ray domain

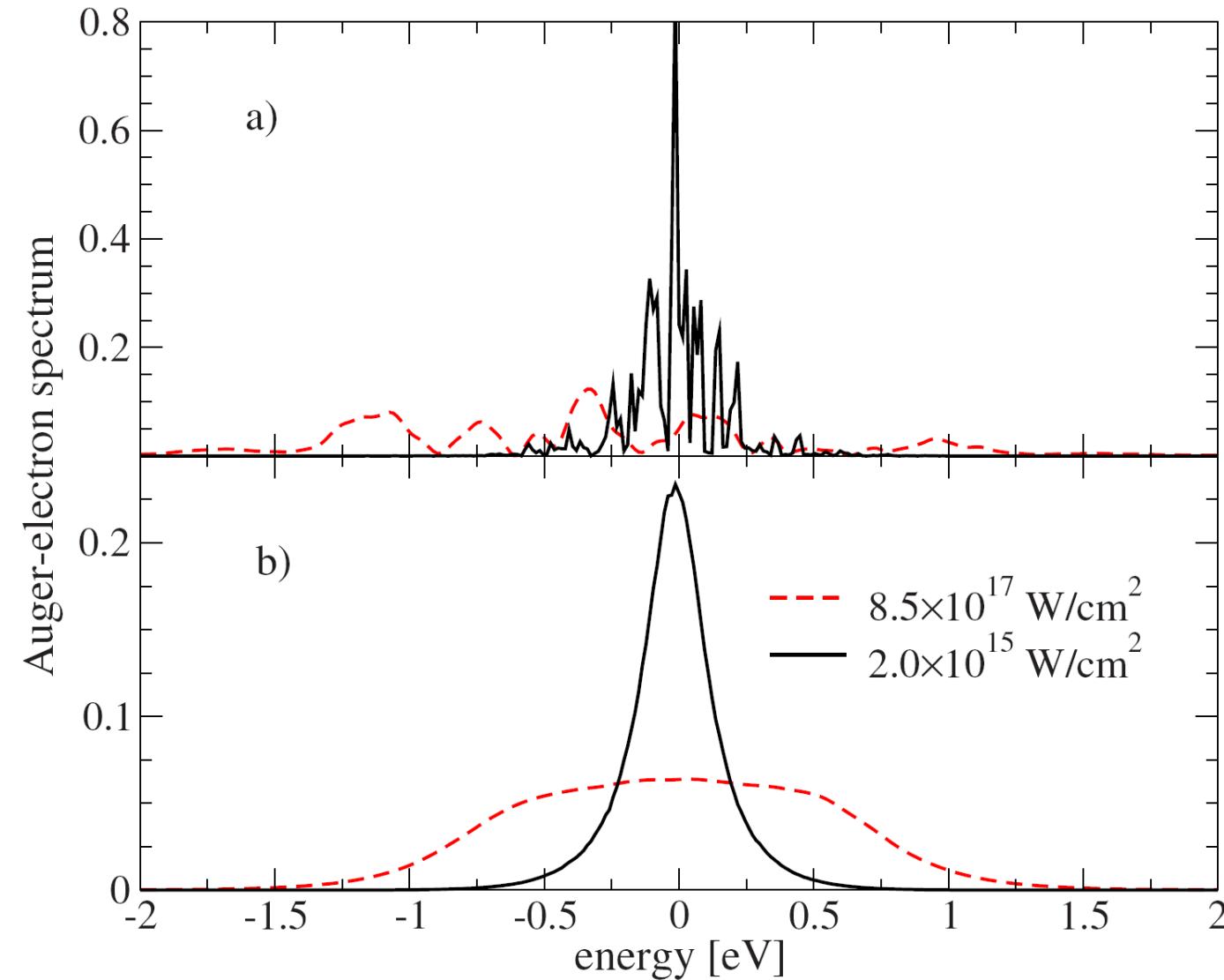
# Resonant Auger effect



# X-ray driven processes competing with resonant Auger decay



# Resonant Auger line profile (SASE, 230 fs)

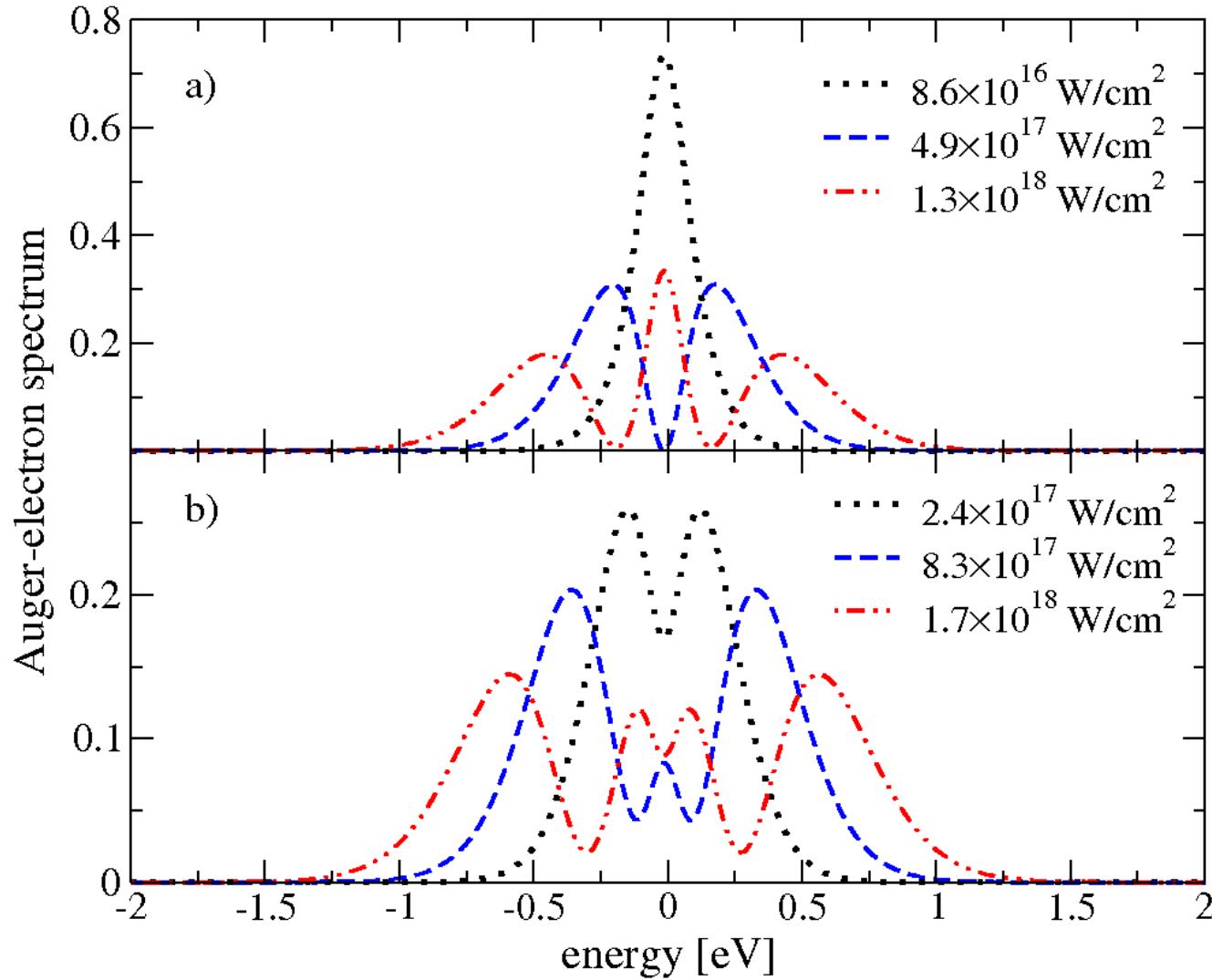


a) Single shot

b) Ensemble average

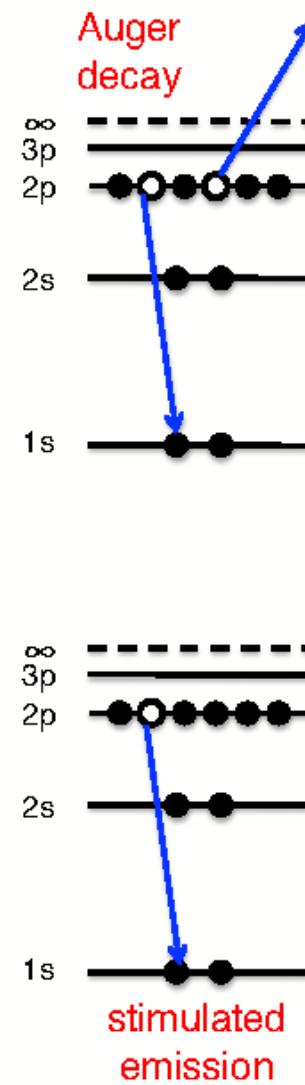
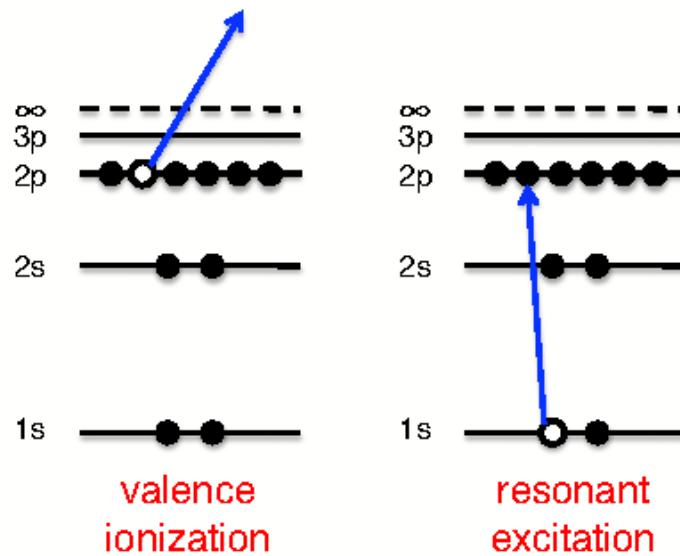
Rohringer, Santra,  
Phys. Rev. A **77**,  
053404 (2008)

# Resonant Auger line profile (Gaussian pulse, 2 fs)



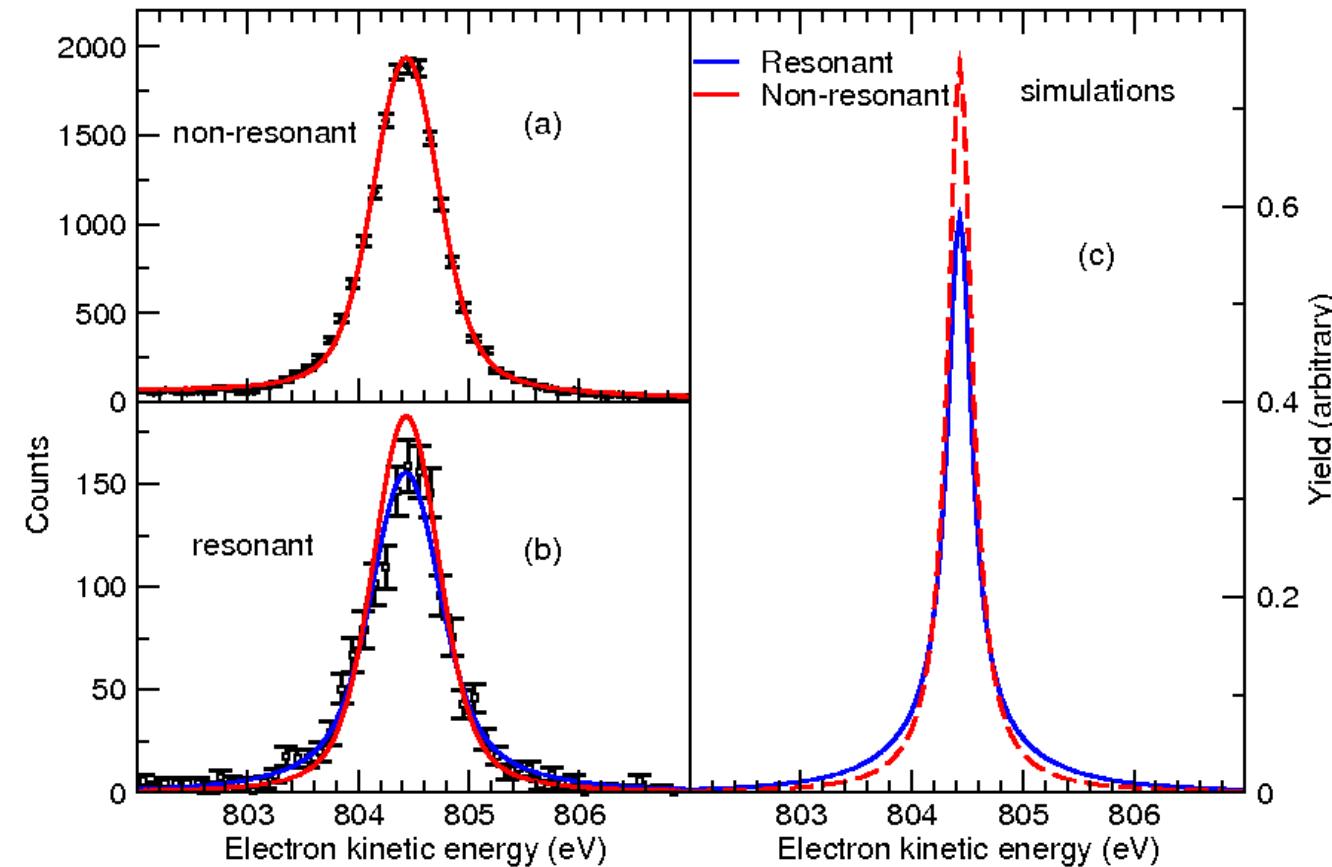
Rohringer, Santra,  
Phys. Rev. A **77**,  
053404 (2008)

# Neon interacting with a high-intensity x-ray pulse at 848 eV



E. P. Kanter *et al.*,  
Phys. Rev. Lett.  
**107**, 233001 (2011)

# Modification of Auger line profile by high-intensity x rays

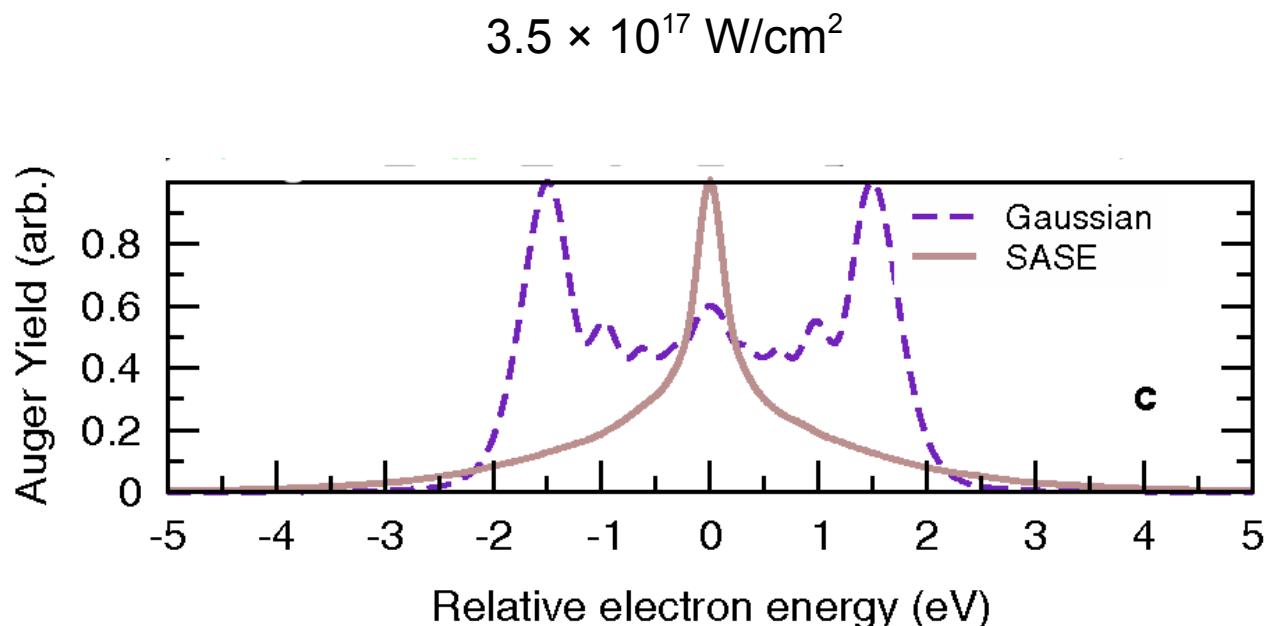


Resonant: 848 eV  
Non-resonant: 930 eV

Resonant Rabi period  
approaching 1 fs

E. P. Kanter *et al.*,  
Phys. Rev. Lett.  
**107**, 233001 (2011)

# SASE vs. Gaussian (theory)

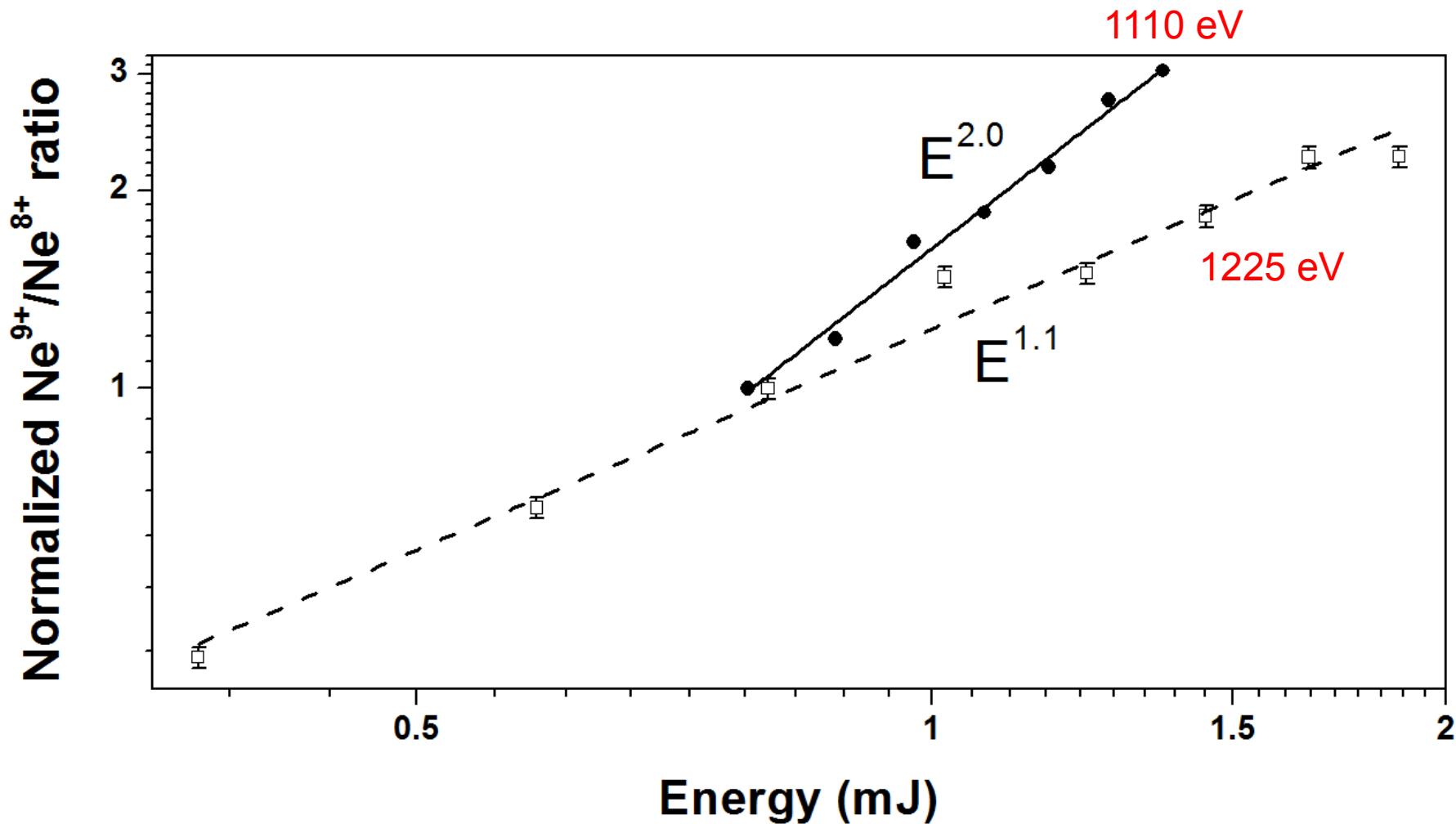


E. P. Kanter *et al.*,  
Phys. Rev. Lett.  
**107**, 233001 (2011)

The resonant Auger line shape generated by an ensemble of SASE pulses (averaged Gaussian temporal profile of 8.5 fs FWHM, 6 eV bandwidth) and a longitudinally coherent Gaussian pulse (8.5 fs FWHM, transform-limited).

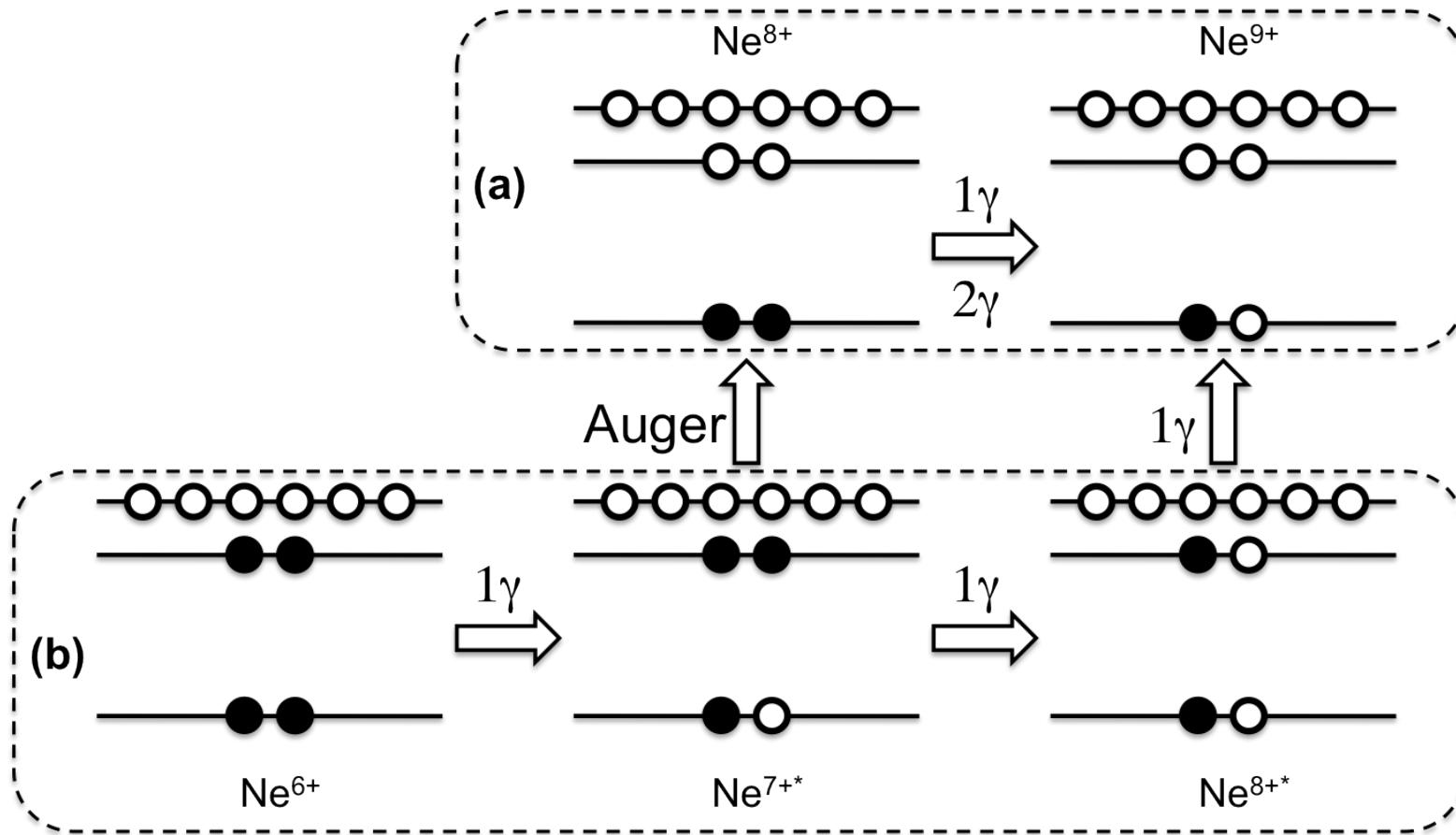
“Resonant by accident”  
→ nonsequential (direct) two-photon absorption

# Nonlinear production of Ne<sup>9+</sup>: observation



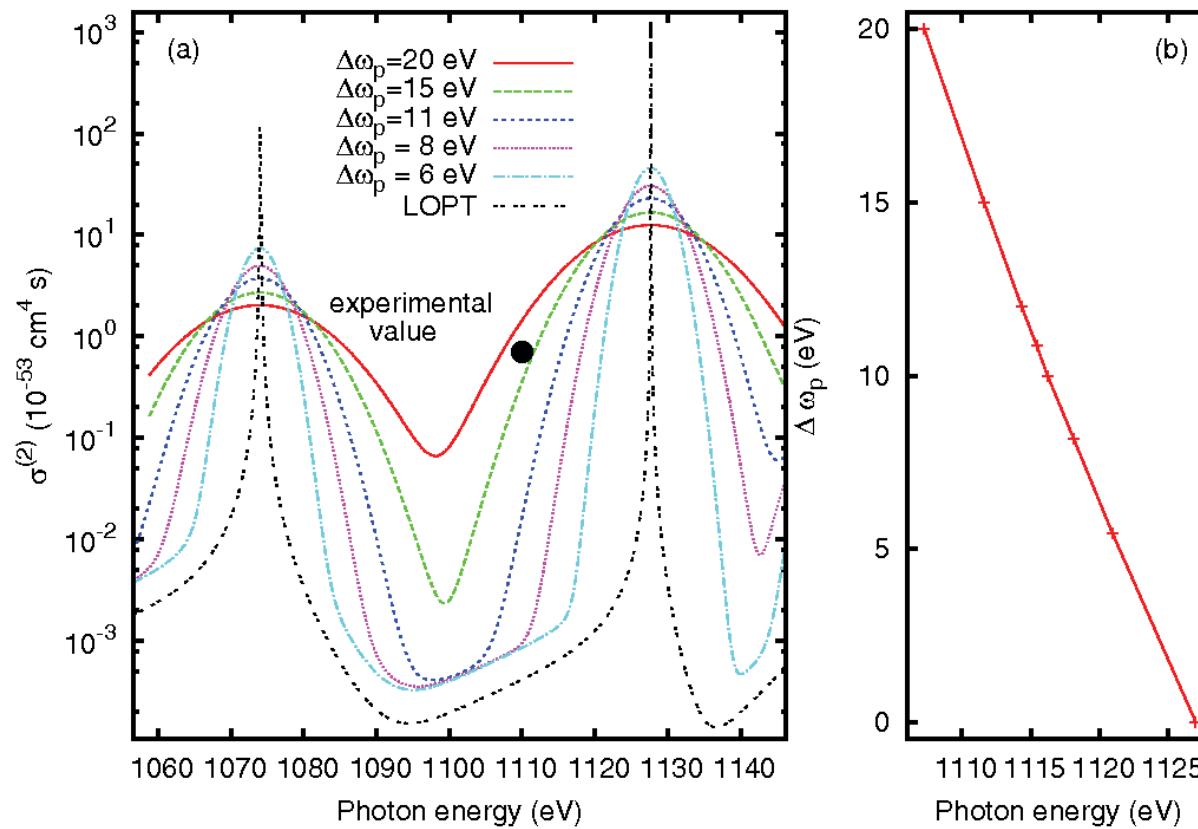
G. Doumy *et al.*, Phys. Rev. Lett. **106**, 083002 (2011)

# Nonlinear production of $\text{Ne}^{9+}$ : mechanisms



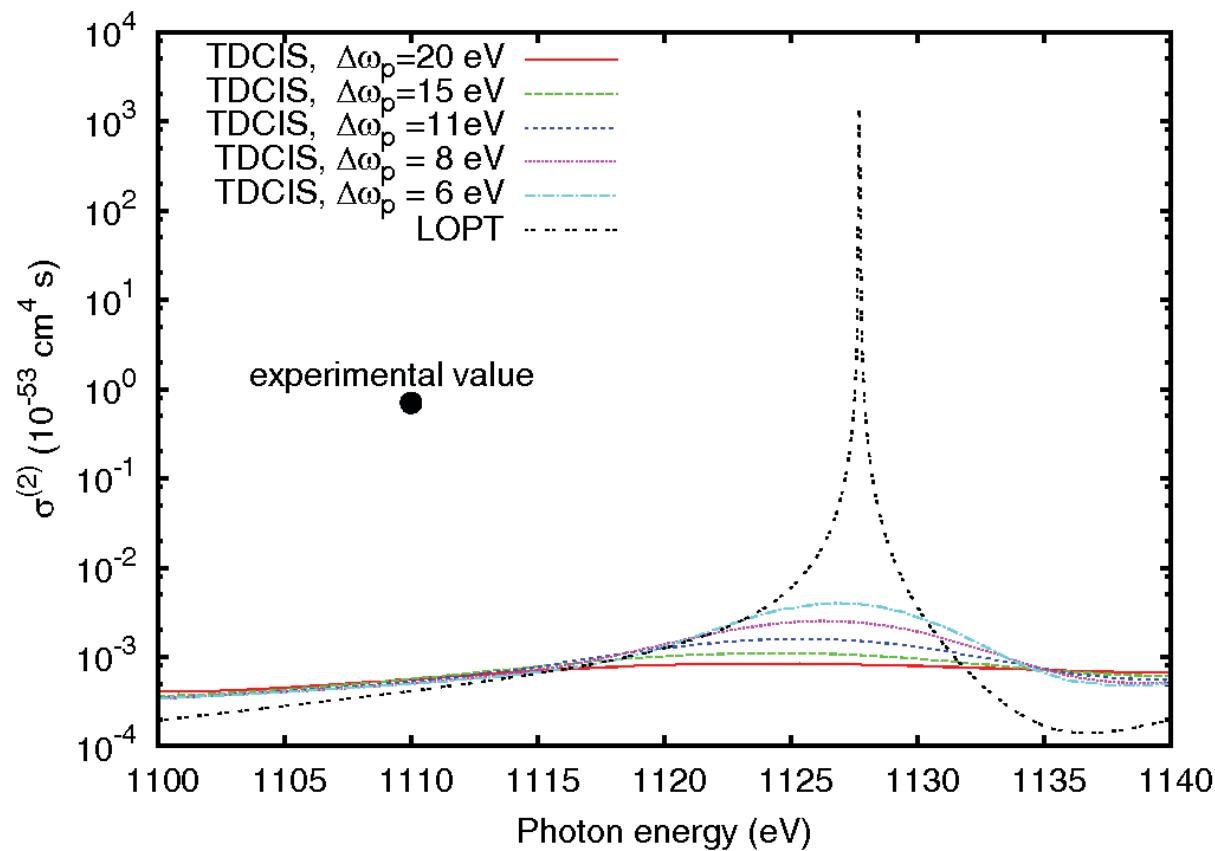
G. Doumy *et al.*, Phys. Rev. Lett. **106**, 083002 (2011)

# Effective two-photon cross section of $\text{Ne}^{8+}$ : SASE



A. Sytcheva *et al.*, Phys. Rev. A **85**, 023414 (2012)

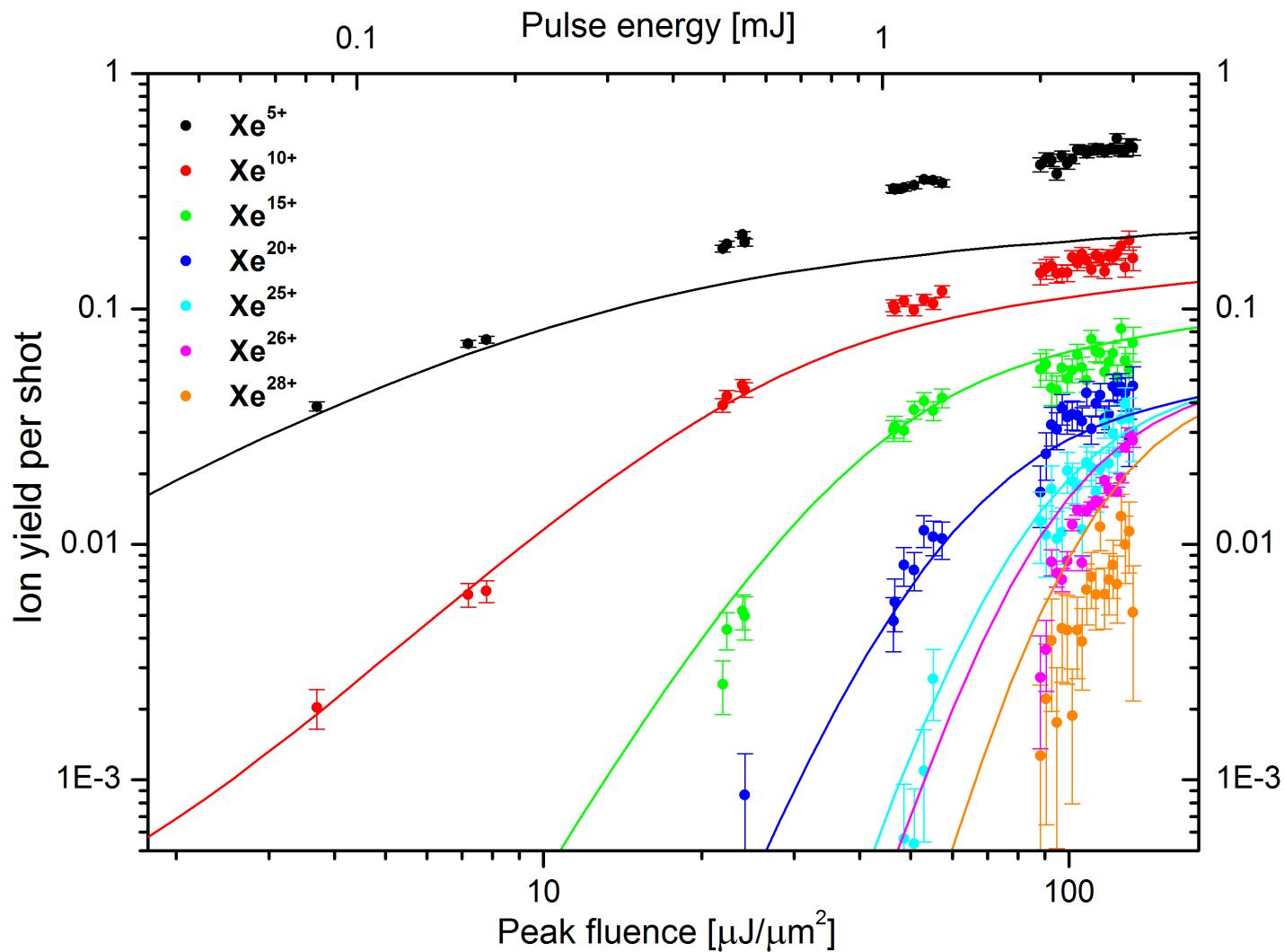
# Effective two-photon cross section of $\text{Ne}^{8+}$ : coherent case



A. Sytcheva *et al.*, Phys. Rev. A **85**, 023414 (2012)

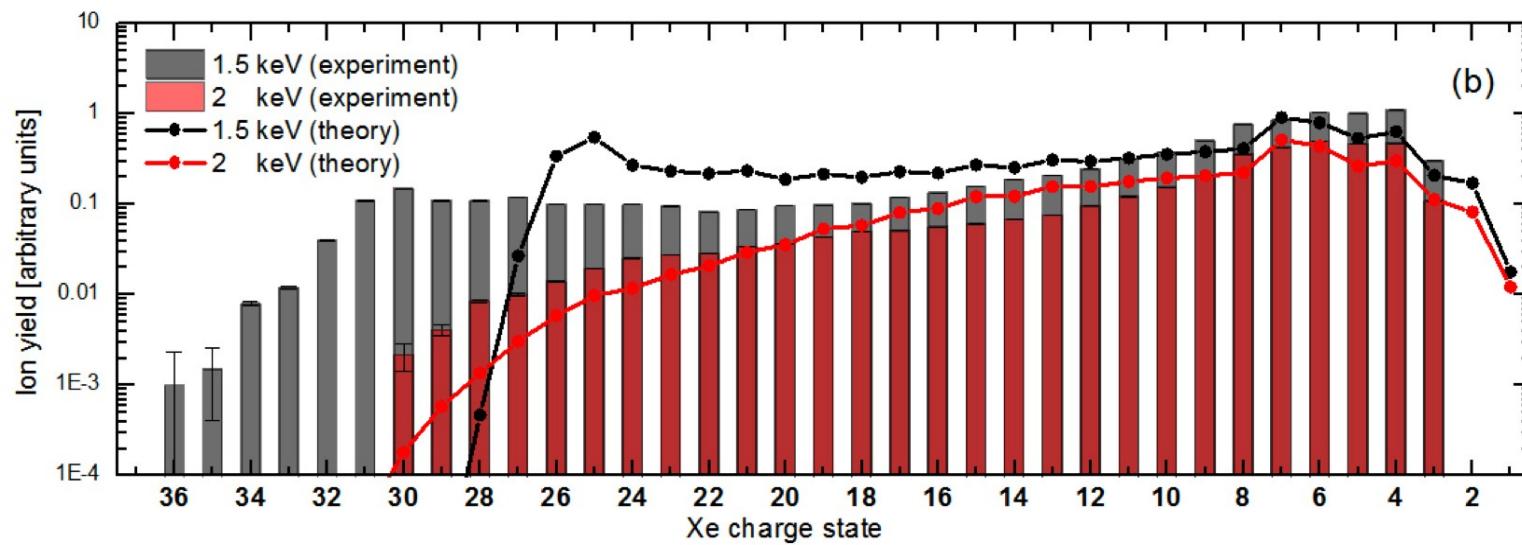
“Resonant by accident”  
→ enhanced multiple ionization

# X-ray multiphoton ionization of xenon at 2 keV



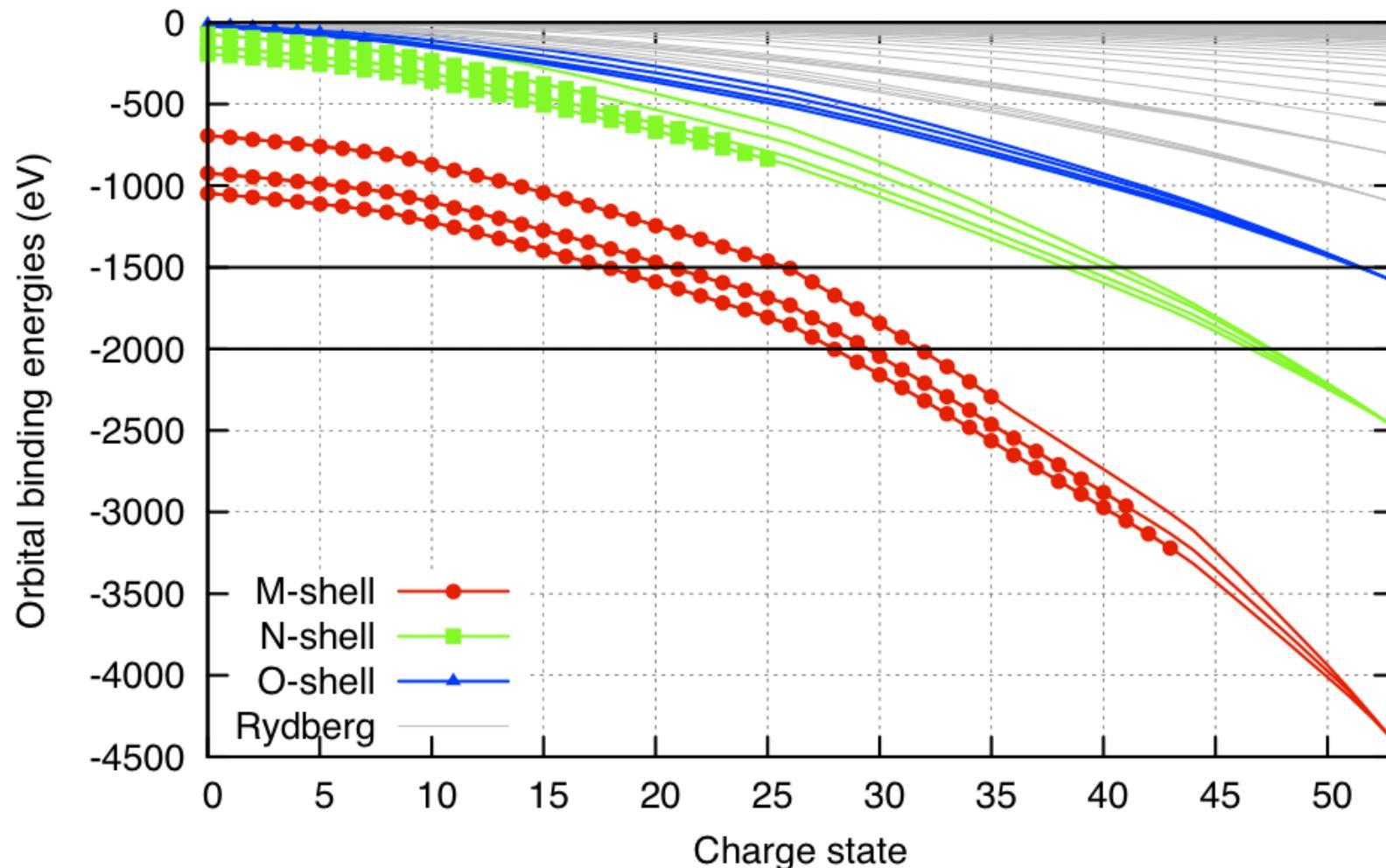
Comparison  
between  
experiment and  
theory

# Enhanced multiple ionization at 1.5 keV



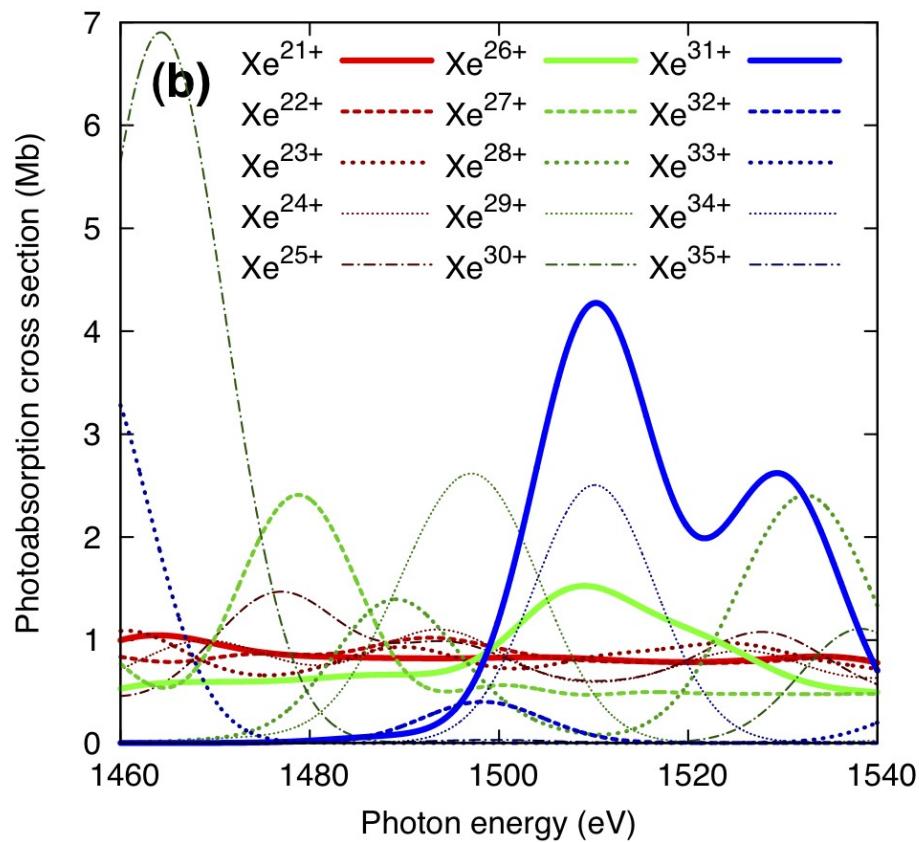
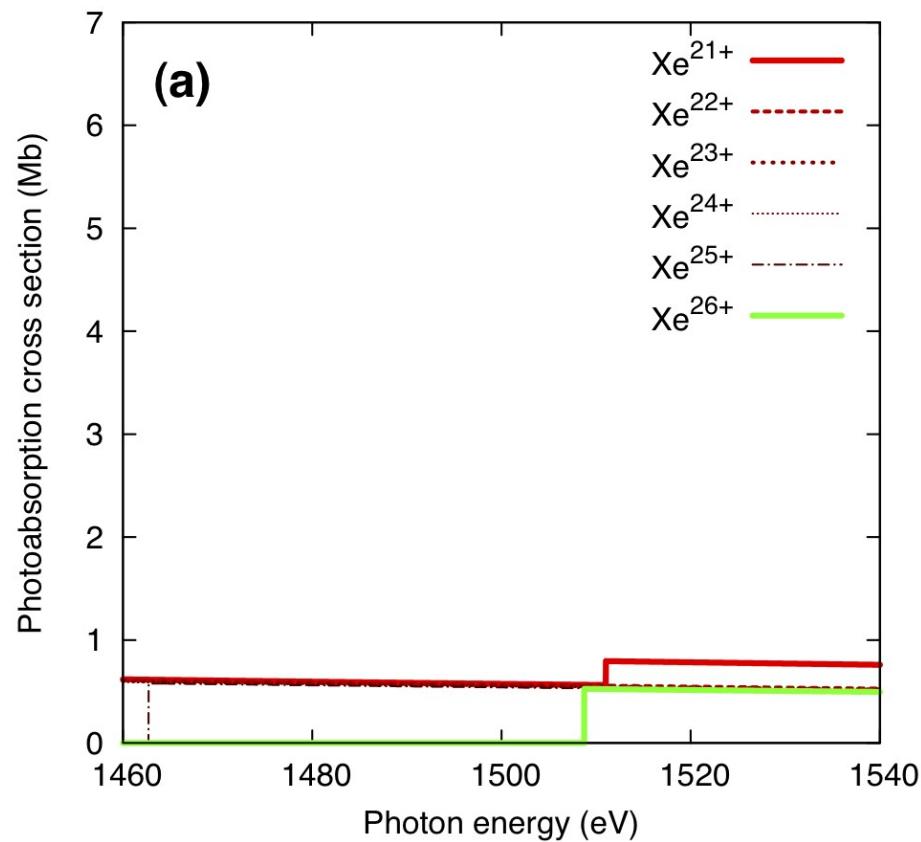
B. Rudek, S.-K. Son *et al.*, submitted

# Orbital binding energies of the ground configuration of $\text{Xe}^{q+}$



S.-K. Son and R. Santra, Phys. Rev. A 85, 063415 (2012)

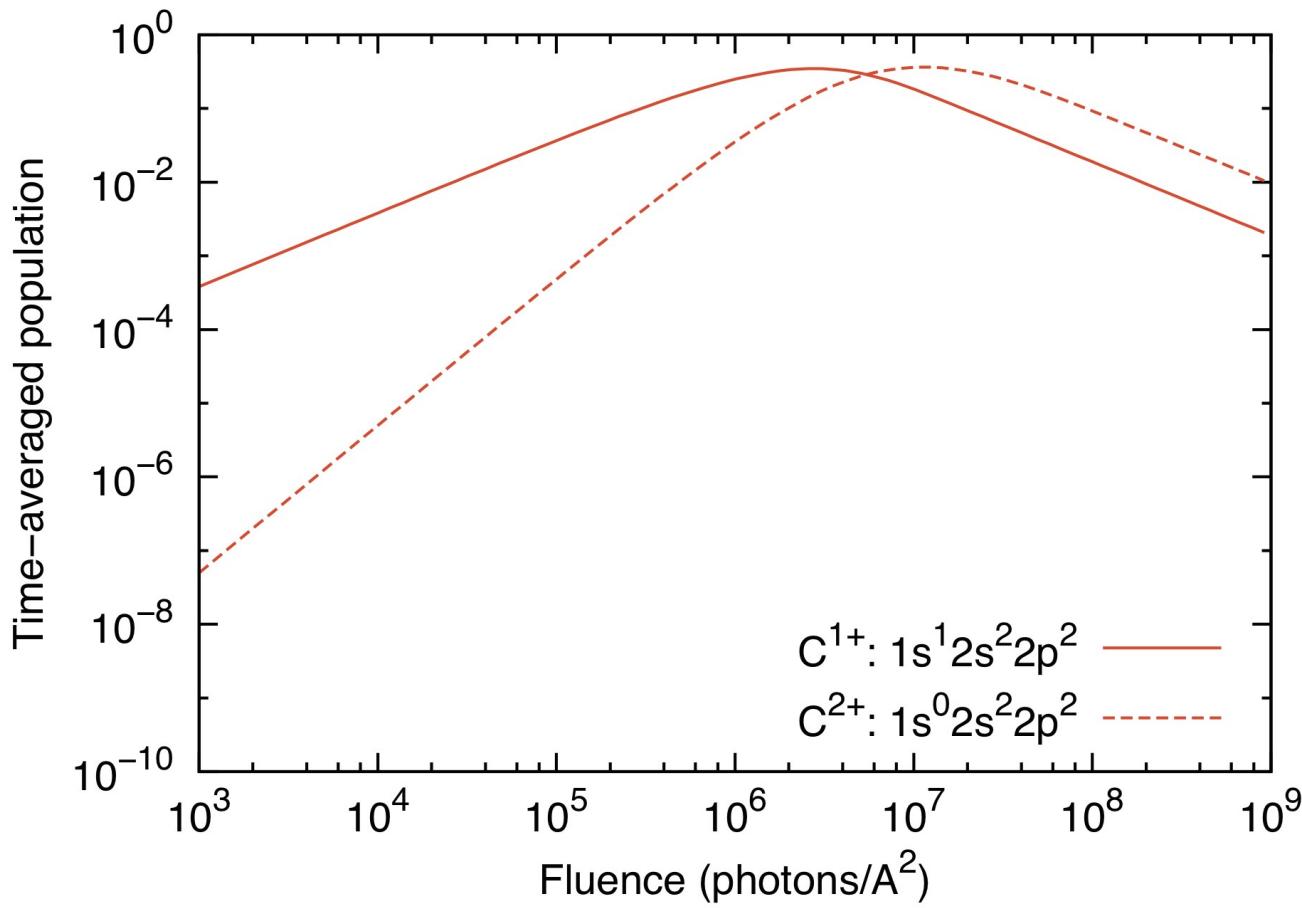
# Resonantly enhanced multiple ionization caused by large SASE bandwidth



B. Rudek, S.-K. Son *et al.*, submitted

# Suppression of electronic damage in coherent diffractive imaging

# Time-averaged population of core-hole configurations

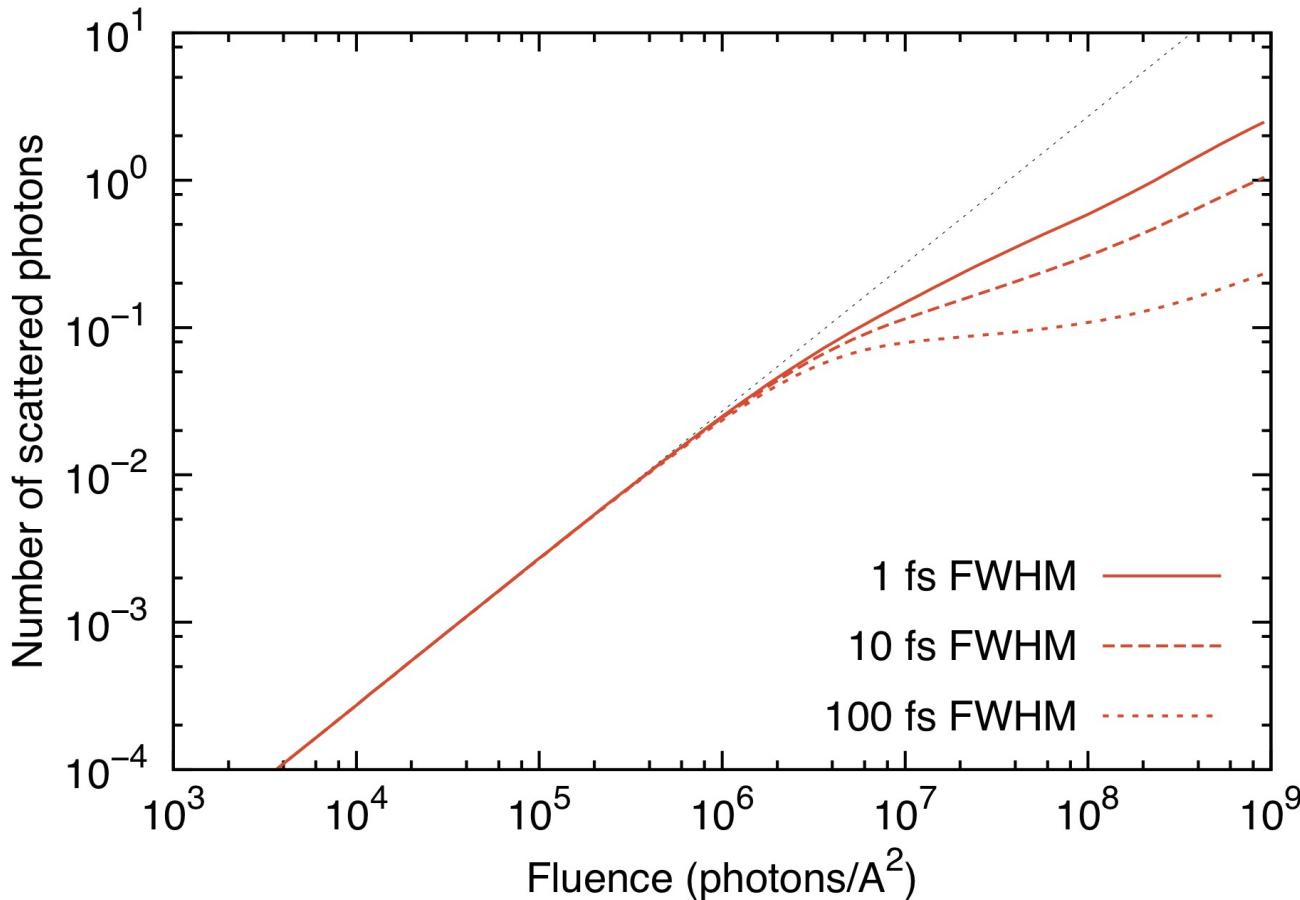


8 keV and 1 fs  
FWHM

$10^6 \text{ photons}/\text{\AA}^2$   
=  
 $10^{12} \text{ photons}/(100 \text{ nm})^2$

S.-K. Son, L. Young,  
and R. Santra,  
Phys. Rev. A **83**,  
033402 (2011)

# Number of photons scattered by a single carbon atom

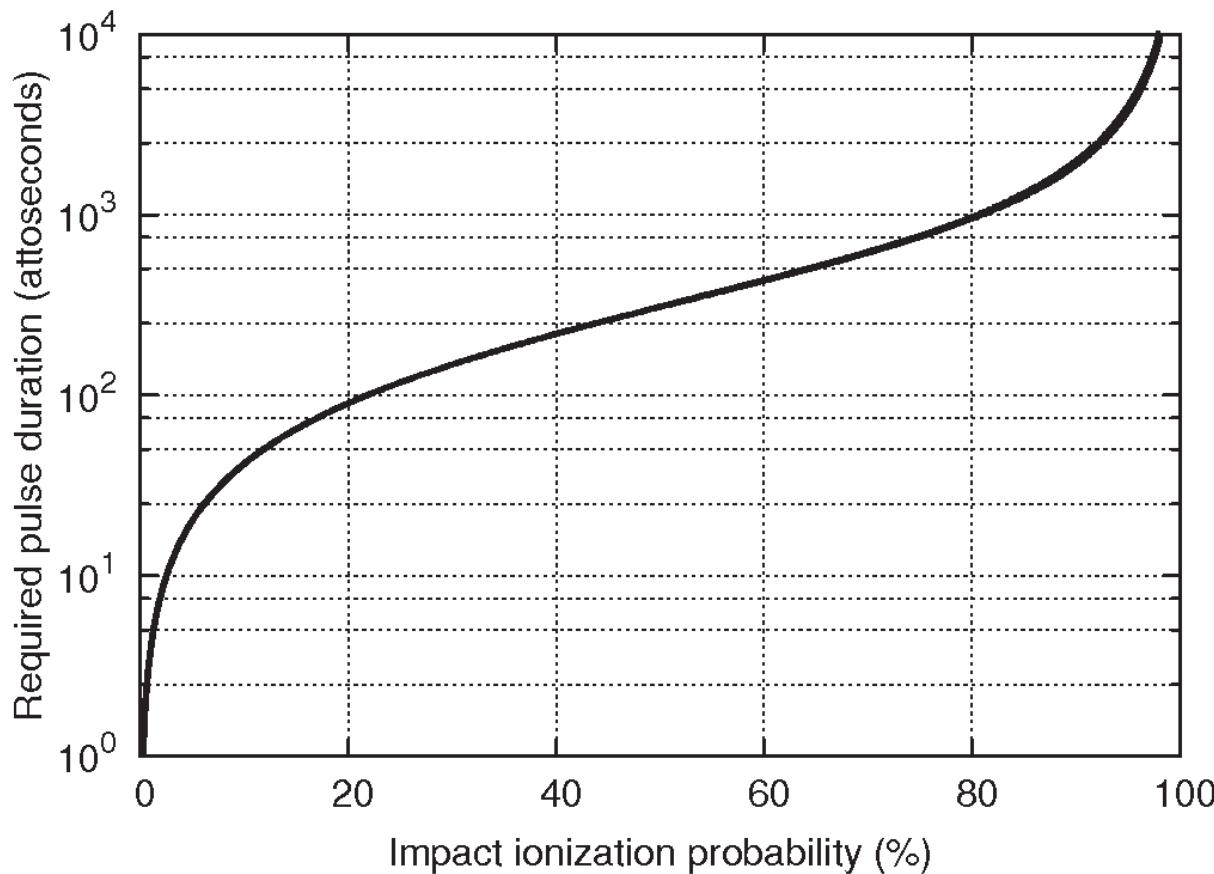


8 keV with spatial  
resolution of 1.7 Å

S.-K. Son, L. Young,  
and R. Santra,  
Phys. Rev. A **83**,  
033402 (2011)

“x-ray transparency” [Nature **466**, 56 (2010)]  
“frustrated absorption” [Phys. Rev. Lett. **104**, 253002 (2010)]

# Suppression of electron impact ionization

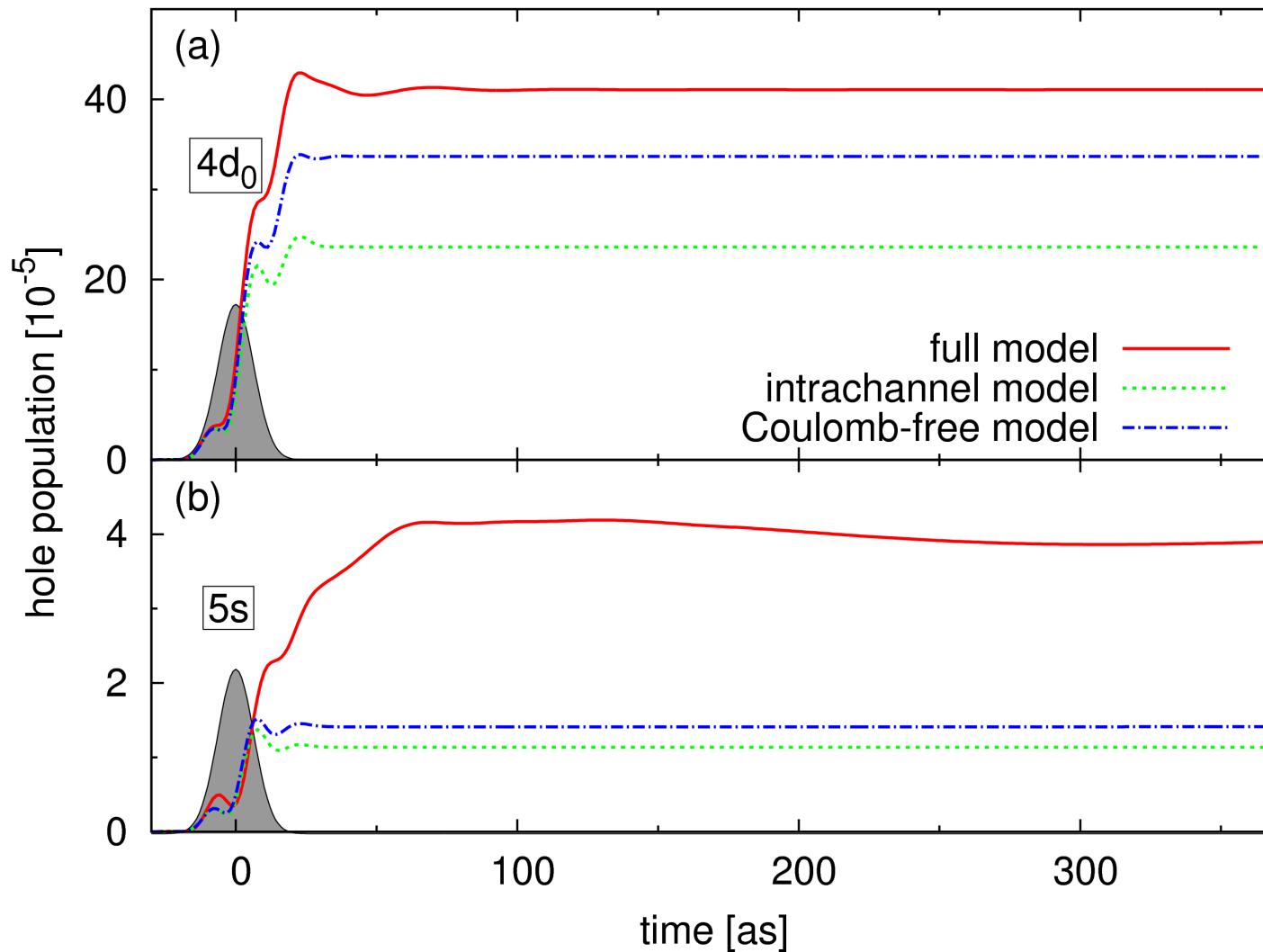


S.-K. Son, L. Young,  
and R. Santra,  
Phys. Rev. A **83**,  
033402 (2011)

Pulse duration required for a given impact ionization probability, for a photoelectron with a kinetic energy of 12 keV in a carbon-based medium with a mean free path of 13 nm

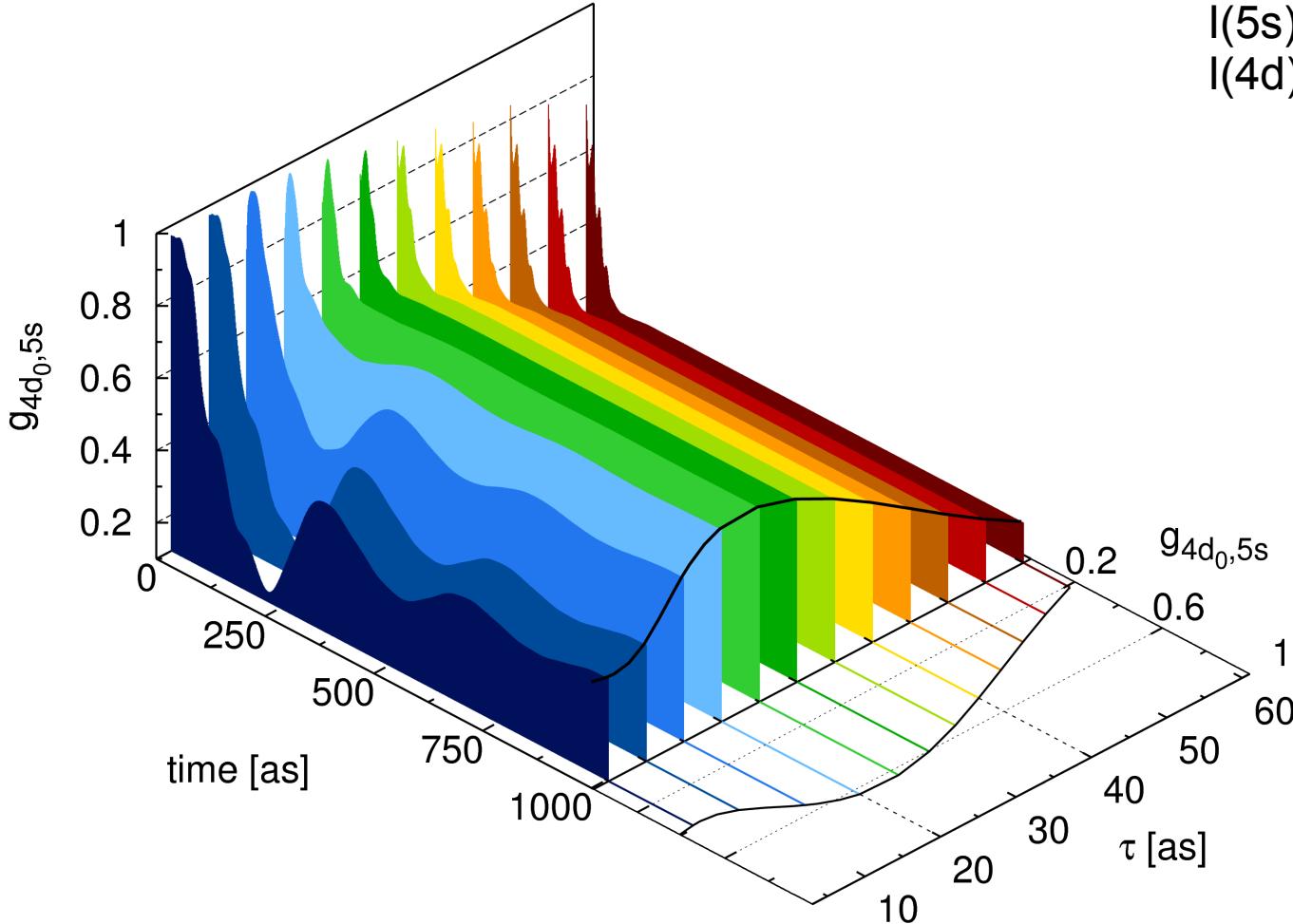
# Electronic coherence and correlation in the time domain

# Attosecond photoionization of xenon: hole populations



S. Pabst *et al.*,  
Phys. Rev. Lett.  
**106**, 053003  
(2011)

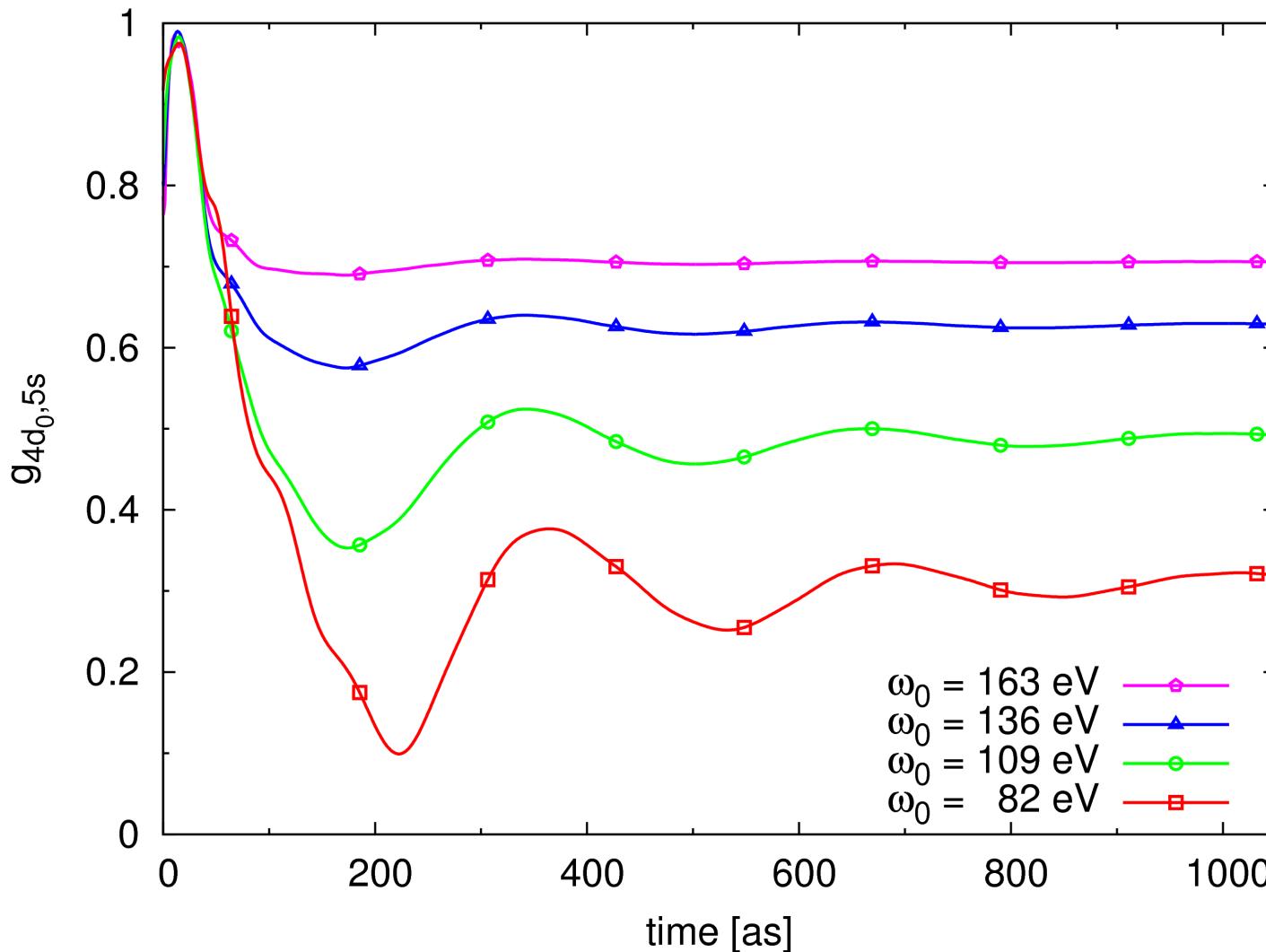
# Time evolution of the degree of coherence (full theory)



Photon energy 136 eV  
 $I(5s) \approx 20$  eV  
 $I(4d) \approx 70$  eV

S. Pabst *et al.*,  
Phys. Rev. Lett.  
**106**, 053003  
(2011)

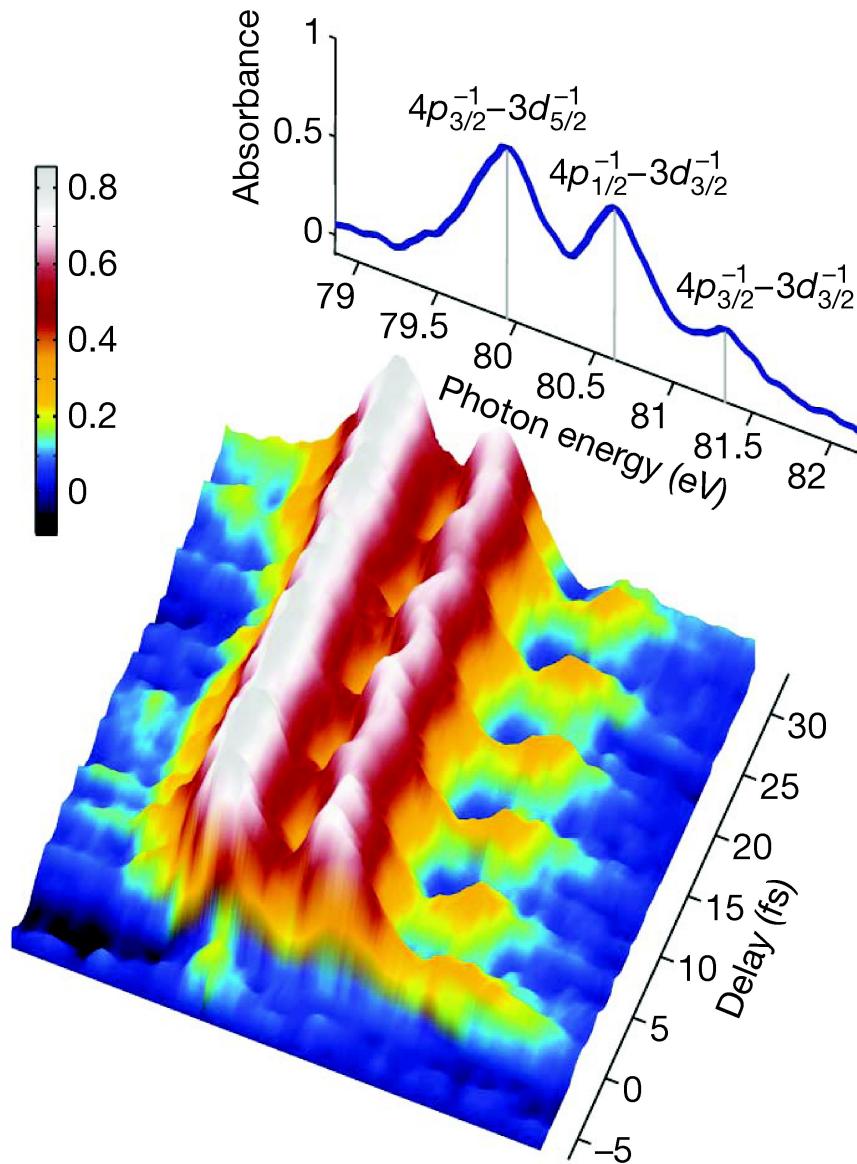
# Time evolution of the degree of coherence (full theory)



S. Pabst *et al.*,  
Phys. Rev. Lett.  
**106**, 053003  
(2011)

# Attosecond transient absorption spectroscopy

Experiment  
on krypton



E. Goulielmakis *et al.*,  
Nature **466**, 739 (2010)

# Imaging of electronic quantum motion

# Time-resolved x-ray scattering theory

If x rays could be treated as a classical field, then one would expect

$$\begin{aligned}\frac{dP}{d\Omega} &= \frac{dP_e}{d\Omega} \left| \int d^3x \rho(\mathbf{x}, t) e^{i\mathbf{Q}\cdot\mathbf{x}} \right|^2 \\ &= \frac{dP_e}{d\Omega} \int d^3x \int d^3x' \langle \Psi, t | \hat{n}(\mathbf{x}') | \Psi, t \rangle \langle \Psi, t | \hat{n}(\mathbf{x}) | \Psi, t \rangle e^{i\mathbf{Q}\cdot(\mathbf{x}-\mathbf{x}')}\end{aligned}$$

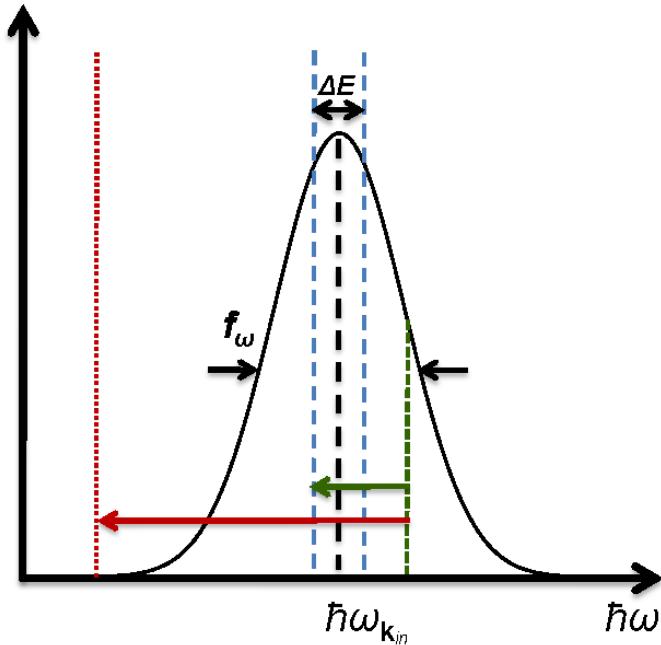
Quantum electrodynamics, on the other hand, predicts

$$\frac{dP}{d\Omega} = \frac{dP_e}{d\Omega} \int d^3x \int d^3x' \langle \Psi, t | \hat{n}(\mathbf{x}') C(\hat{H}) \hat{n}(\mathbf{x}) | \Psi, t \rangle e^{i\mathbf{Q}\cdot(\mathbf{x}-\mathbf{x}')}$$

$$C(\hat{H}) = \frac{\tau \Delta E}{\hbar \sqrt{\pi} 8 \ln 2} \exp\left(-\frac{\tau^2}{8 \ln 2 \hbar^2} (\hat{H} - \langle \hat{H} \rangle_t)^2\right)$$

G. Dixit, O. Vendrell, and R. Santra, PNAS, in press

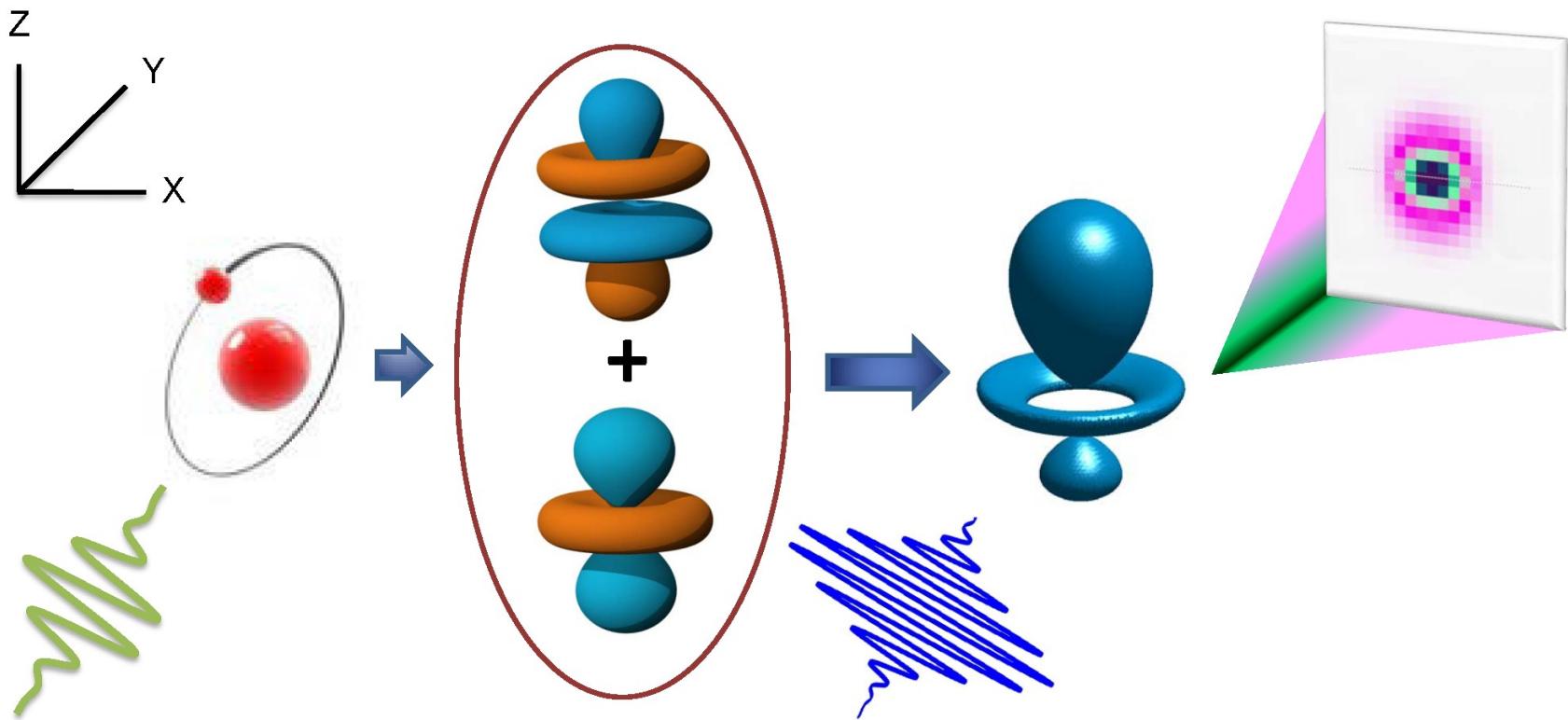
# Key idea



- Probing dynamics requires a finite spectral bandwidth.
- As a consequence of this bandwidth, it becomes impossible to decide whether the detected photon was scattered elastically or inelastically.
- This problem persists even if the scattering detector has perfect energy resolution.

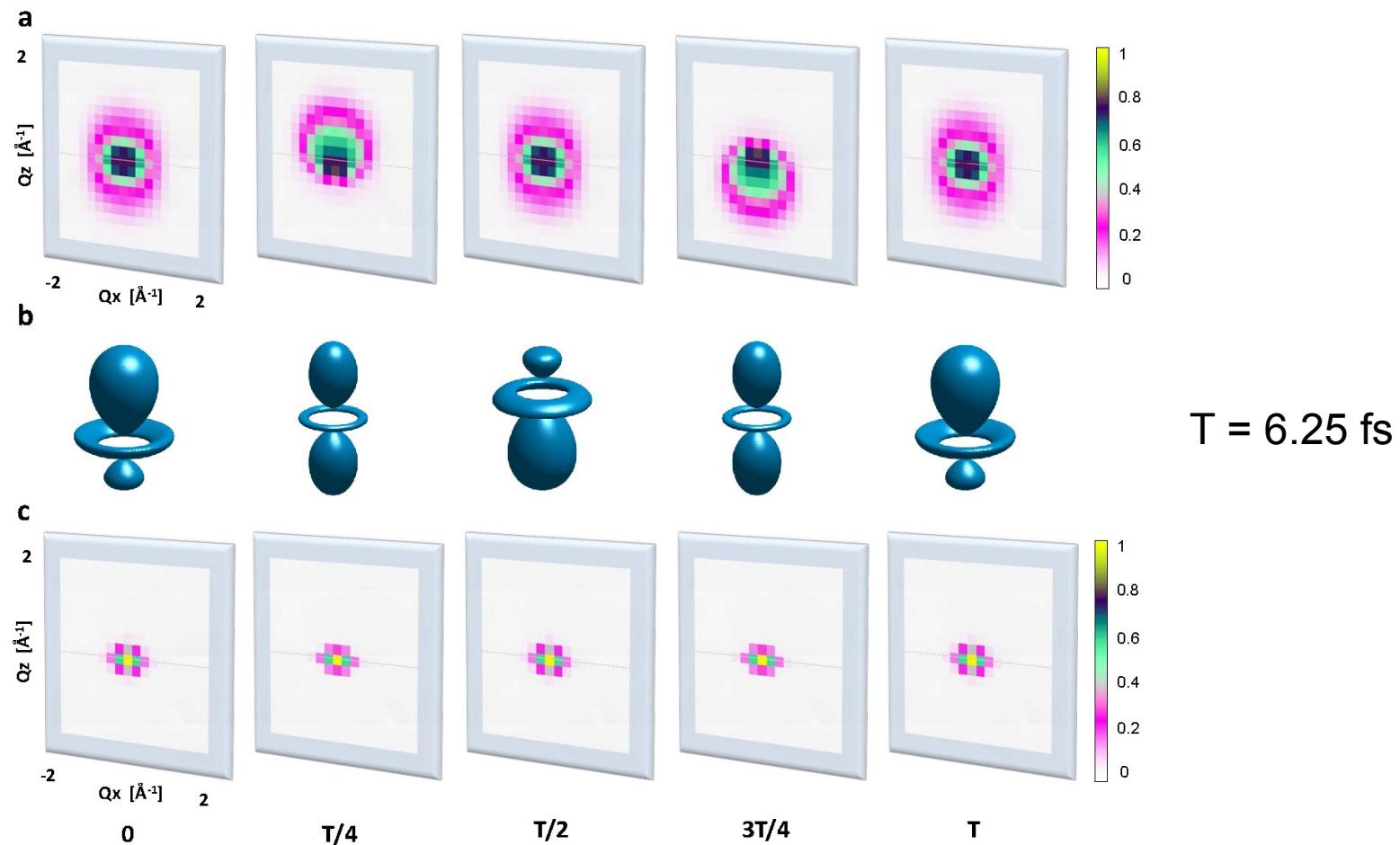
G. Dixit, O. Vendrell, and R. Santra, PNAS, in press

# Example



G. Dixit, O. Vendrell, and R. Santra, PNAS, in press

# QED vs. semiclassical theory



G. Dixit, O. Vendrell, and R. Santra, PNAS, in press

# Conclusions

X-ray free-electron lasers are great.

If they produced shorter, more coherent pulses, they would be even greater!