Little Higgs models in light of the 126 GeV discovery at the LHC

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Motivation

How to constrain a generic model in HEP?

- direct searches of resonances
- electroweak precision tests
- flavour constraints
- nowadays: Higgs sector

Higgs sector is the key to understand EW-scale physics (and beyond?)

Quicksand of the Standard Model

a new resonance at 126 GeV has just been discovered, but... radiative corrections to SM Higgs mass are quadratically divergent



Big Hierarchy Problem

highly unnatural for $\Lambda \gg {\rm TeV}$

Quicksand of the Standard Model

need new physics to prevent fine-tuning! in particular (for 10% fine-tuning), new physics should appear at

 $\Lambda_{\rm top} \lesssim 2 \ {\rm TeV}, \quad \Lambda_{\rm gauge} \lesssim 5 \ {\rm TeV}, \quad \Lambda_{\rm quartic} \lesssim 10 \ {\rm TeV}$

but... severe constraints on the low-energy effects of new-physics:

Little Hierarchy Problem

no fine-tuning implies $\Lambda \lesssim 1\,{\rm TeV}$ vs. no dangerous corrections implies $\Lambda \gtrsim 10\,{\rm TeV}$

Quicksand of the Standard Model

need new physics to prevent *fine-tuning*! in particular (for 10% fine-tuning), new physics should appear at

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but... severe constraints on the low-energy effects of new-physics:

broken symmetry	operators	scale Λ
B, L	$(QQQL)/\Lambda^2$	10^{13} TeV
flavor (1,2 nd family), CP	$(\bar{d}s\bar{d}s)/\Lambda^2$	$1000 { m TeV}$
flavor $(2,3^{rd} \text{ family})$	$m_b (\bar{s} \sigma_{\mu\nu} F^{\mu\nu} b) / \Lambda^2$	$50 { m TeV}$
custodial $SU(2)$	$(h^{\dagger}D_{\mu}h)^2/\Lambda^2$	$5 { m TeV}$
none (S-parameter)	$(D^2 h^\dagger D^2 h)/\Lambda^2$	$5 { m TeV}$

Little Hierarchy Problem

no fine-tuning implies $\Lambda \lesssim 1\,{\rm TeV}_{\rm VS.}$ no dangerous corrections implies $\Lambda \gtrsim 10\,{\rm TeV}$

Two paradigms for EWSB

hierarchy problem as guideline to answer the following question:

what is the dynamical origin of EWSB?

- weakly coupled answer \rightarrow Supersymmetry
- strongly coupled answer \rightarrow Composite Higgs, Little Higgs...

Strongly coupled answer





Composite/Little Higgs Ansatz

Higgs as pNGB of a new (approximate) global symmetry which is spontaneously broken at a scale $\Lambda \sim 4\pi f$

Strongly coupled answer





Composite/Little Higgs Ansatz

Higgs as pNGB of a new (approximate) global symmetry which is spontaneously broken at a scale $\Lambda \sim 4\pi f$

Summary

The Little Higgs paradigm:

- Higgs as a pNGB of a global SSB at $\Lambda \sim 4\pi f$ (like pions!)
- $\bullet\,$ new fermionic/vector states with masses $\sim f$ besides SM-ones
- *EWSB* is triggered *naturally* (Collective Symmetry Breaking), i.e. $v \sim \mathcal{O}(100 \text{ GeV})$ for $f \sim 1 \text{ TeV}$ with only logarithmic sensitivity to the cut-off $\Lambda \Rightarrow$ Big Hierarchy solved!
- $\Lambda \sim 10$ TeV for $f \sim 1$ TeV \Rightarrow Little Hierarchy solved!
- it is an effective theory valid up to $\Lambda:$ no UV-completion of the strongly coupled regime $E>\Lambda$

Statistical analysis

We considered the three most popular Little Higgs models:

- Simplest Little Higgs (SLH) [Schmaltz]
- Littlest Higgs (L^2H) [Arkani-Hamed et al.]
- Littlest Higgs with T-parity (LHT) [Low et al.]

and realized a χ^2 analysis on their parameter spaces, taking into account the whole set of 7+8 TeV Higgs searches by *ATLAS* and *CMS*, and by fitting 21 different *EW* Precision Observables:

$$\chi^2 = \sum_i \frac{(\mathcal{O}_i - \mathcal{O}_i^{\mathsf{exp}})^2}{\sigma_i^2}$$

where \mathcal{O}_i depends on the free parameters of the model considered.

Data used: Higgs sector

the Higgs results are expressed in terms of a signal strength modifier

$$\mu_{i} = \frac{\sum_{p} \epsilon^{p}_{i} \sigma_{p}}{\sum_{p} \epsilon^{p}_{i} \sigma_{p}^{SM}} \cdot \frac{BR(h \to X_{i}X_{i})}{BR(h \to X_{i}X_{i})_{SM}}$$

we included in our χ^2 analysis the best-fit values of μ_i reported by the Collaborations for all the different 7+8 TeV channels *i*:



Data used: EWPD

every extension of SM has to satisfy at least the precision constraints of the electroweak sector:

low-energy observables

e.g. $\nu\text{-scattering, parity violation observables}\ldots$

 $\bullet \ Z{\text -}{\sf pole \ observables}$

e.g. m_Z , Γ_Z , Z-pole asymmetries...



LH Smoking guns

Where do the LH corrections to the SM quantities come from?

- new decay channels of the Higgs, e.g. $h \rightarrow A_H A_H$ in LHT
- modified Higgs couplings with SM fermions and vector bosons

e.g.
$$2 \frac{m_W^2}{v} y_W h W^+ W^-$$
, $y_W = \begin{cases} 1 & SM \\ 1 + \mathcal{O}\left(v^2/f^2\right) & LH \end{cases}$

• interaction terms of Higgs with new fermions/vector bosons

e.g.
$$\frac{m_T}{v} y_T h \bar{T} T$$
 $m_T \sim f, y_T \sim \mathcal{O} \left(v^2 / f^2 \right)$

modified neutral- and charged-currents

e.g.
$$\frac{g}{c_W} \sum_f \bar{f} \gamma^{\mu} \Big((g_L^{SM} + \delta g_L) P_L + (g_R^{SM} + \delta g_R) P_R \Big) f Z_{\mu}$$

LHT results



- free parameters: *f SSB* scale, *R* ratio of Yukawa couplings in top sector
- $f_{\min}^{99\%} = 487.0$ GeV, translates into lower bounds on new states' masses, e.g.

$m_{W'}$	\gtrsim	323.5 GeV
m_T	\gtrsim	743.1 GeV

beyond the current limits from direct searches

 $\bullet\,$ min. required fine tuning: $\,\sim\,10\%$, defined as

$$\Delta = \frac{|\delta \mu^2|}{\mu_{\rm obs}^2}$$

• results mainly driven by EWPD (see next slide)

Higgs data vs. EWPD



- the shape of the combined result is driven by the *EW* constraints (much smaller uncertainties)
- $\bullet\,$ Higgs data only: for $v/f\gtrsim 0.6$ decay $h\to A_HA_H$ open and dominant
- $\bullet\,$ Higgs data only: subdominant dependence on R w.r.t. f is a consequence of the Collective Symmetry Breaking mechanism
- *EWPD* only: in the central region around $R \sim 1.5$ there is a cancellation among the different contributions to the *T*-parameter (best fit: $T = 0.05 \pm 0.12$)

Conclusions

- *Little Higgs* models are an appealing solution to the hierarchy problem, alternative to weakly coupled solutions like *SUSY*
- most of the parameter space of three popular Little Higgs models is still compatible at $\sim 99\%~CL$ with the early results of the 7+8 TeV Higgs searches
- electroweak precision data represent still the most severe constraints
- fine-tuning as a guideline to understand the naturalness of a model: Little Higgs models require a minimum level of $\sim 10\%$ of fine tuning
- new data on the Higgs sector with increasing luminosity will reduce the uncertainties and thus give more precise information

Thank you for your attention!

SLH results



$$\chi^2_{\min}$$
/d.o.f. = 1.02
 χ^2_{SM} /d.o.f. = 0.97

- free parameters: f~SSB scale, t_β ratio of vevs of scalar fields $\phi_{1,2}$
- $f_{\min}^{99\%} = 2.72$ TeV, translates into lower bounds on new states' masses, e.g.

$m_{W'}$	\gtrsim	1.23 TeV
m_T	\gtrsim	2.68 TeV

beyond the current limits from direct searches

 $\bullet\,$ min. required fine tuning: $\sim 1\%$, defined as

$$\Delta = \frac{|\delta\mu^2|}{\mu_{\rm obs}^2}$$

• results mainly driven by EWPD

L^2H results



- free parameters: f SSB scale, c mixing angle in gauge sector
- $f_{\min}^{99\%} = 3.12$ TeV, translates into lower bounds on new states' masses, e.g.

$m_{W'}$	\gtrsim	2.07 TeV
m_T	\gtrsim	4.38 TeV

beyond the current limits from direct searches

 $\bullet\,$ min. required fine tuning: $\sim 0.1\%$, defined as

$$\Delta = \frac{|\delta\mu^2|}{\mu_{\rm obs}^2}$$

• results mainly driven by EWPD

Partial decay widths in LH

• 1-loop decays

$$\begin{split} \Gamma(h \to gg)_{LH} &\sim \quad \frac{\alpha_s^2 m_h^3}{32\pi^3 v^2} \Big| \sum_{\rm f,col} -\frac{1}{2} F_{\frac{1}{2}}(x_f) \, y_f \Big|^2 \\ \Gamma(h \to \gamma\gamma)_{LH} &\sim \quad \frac{\alpha^2 m_h^2}{256\pi^3 v^2} \Big| \sum_{\rm f,ch} \frac{4}{2} F_{\frac{1}{2}}(x_f) \, y_f + \sum_{\rm v,ch} F_1(x_v) \, y_v + \sum_{\rm s,ch} F_0(x_s) \, y_s \Big|^2 \end{split}$$

where $x_i = \frac{4m_i^2}{m_h^2}$, $F_i(x_i)$ are loop functions, y_i the modified Yuk. couplings

$$\Rightarrow \quad \text{narrow-width approximation:} \ \frac{\sigma_{LH}}{\sigma_{SM}}(gg \to h) = \frac{\Gamma(h \to gg)_{LH}}{\Gamma(h \to gg)_{SM}}$$

• tree-level decays

$$\Gamma(h \to VV)_{LH} \sim \Gamma(h \to VV)_{SM} \left(\frac{g_{hVV}}{g_{hVV}^{SM}}\right)^2$$

$$\Gamma(h \to f\bar{f})_{LH} \sim \Gamma(h \to f\bar{f})_{SM} \left(\frac{g_{hff}}{g_{hff}^{SM}}\right)^2$$

where $g_{hVV} = \frac{m_V^2}{v} y_V$ and $g_{hff} = \frac{m_f}{v} y_f$