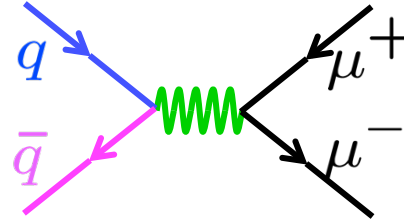


What can more Drell-Yan data tell us about QCD (or Drell-Yan Overview)?

Paul E. Reimer
Physics Division
Argonne National Laboratory



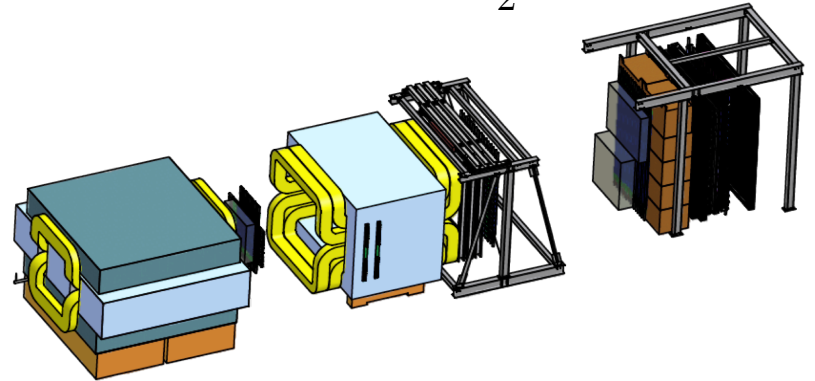
HERMES
U. Elschenbroich

1. The Drell-Yan Process: The Big Picture
2. Unpolarized Measurements

1. Longitudinal parton distributions $\frac{d\sigma^{pd}}{d\sigma^{pp}} = \frac{1}{2} \left(1 + \frac{\bar{d}}{\bar{u}} \right)$
2. Angular distributions $\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$

3. Polarized Measurements

$$f_{1T}^{\perp} |_{DIS} = - f_{1T}^{\perp} |_{DY}$$



This work is supported in part by the U.S. Department of Energy, Office of Nuclear Physics, under Contract No. DE-AC02-06CH11357.

The Big Picture



11/13/2011

Early Muon Pair Data—soon to be called Drell-Yan

Observation of Massive Muon Pairs in Hadron Collisions*

J. H. Christenson, G. S. Hicks, L. M. Lederman, P. J. Limon, and B. G. Pope

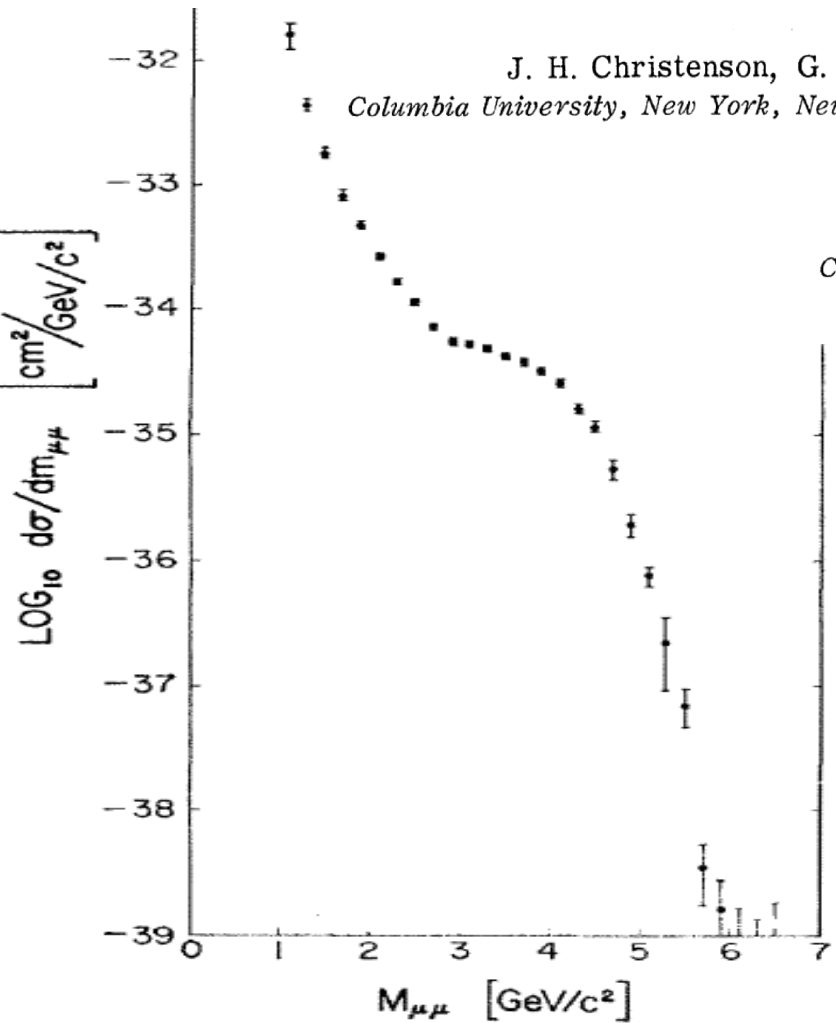
Columbia University, New York, New York 10027, and Brookhaven National Laboratory, Upton, New York 11973

and

E. Zavattini

CERN Laboratory, Geneva, Switzerland

(Received 8 September 1970)



Muon Pairs in the mass range $1 < m_{\mu\mu} < 6.7 \text{ GeV}/c^2$ have been observed in collisions of high-energy protons with uranium nuclei. At an incident energy of 29 GeV, **the cross section varies smoothly as $d\sigma/dm_{\mu\mu} \approx 10^{-32} / m_{\mu\mu}^5 \text{ cm}^2 (\text{GeV}/c)^{-2}$ and exhibits no resonant structure.** The total cross section increases by a factor of 5 as the proton energy rises from 22 to 29.5 GeV.

Early Muon Pair Data—soon to be called Drell-Yan

Observation of Massive Muon Pairs in Hadron Collisions*

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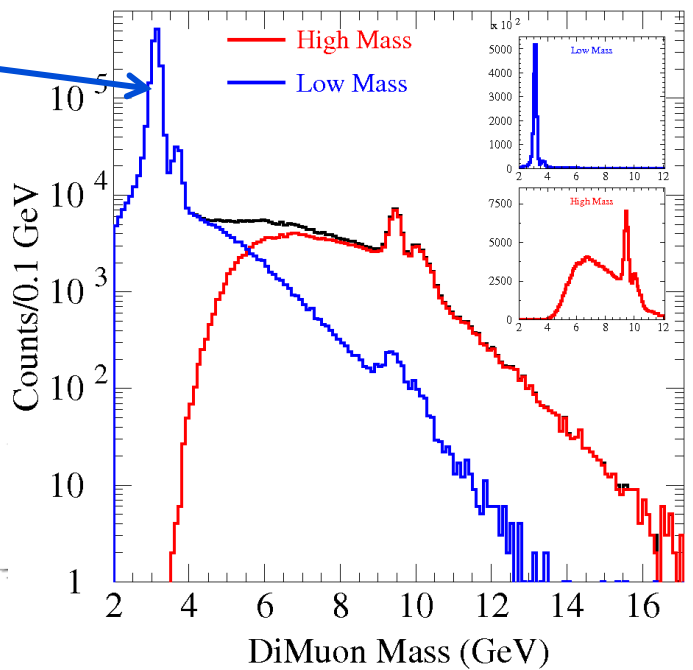
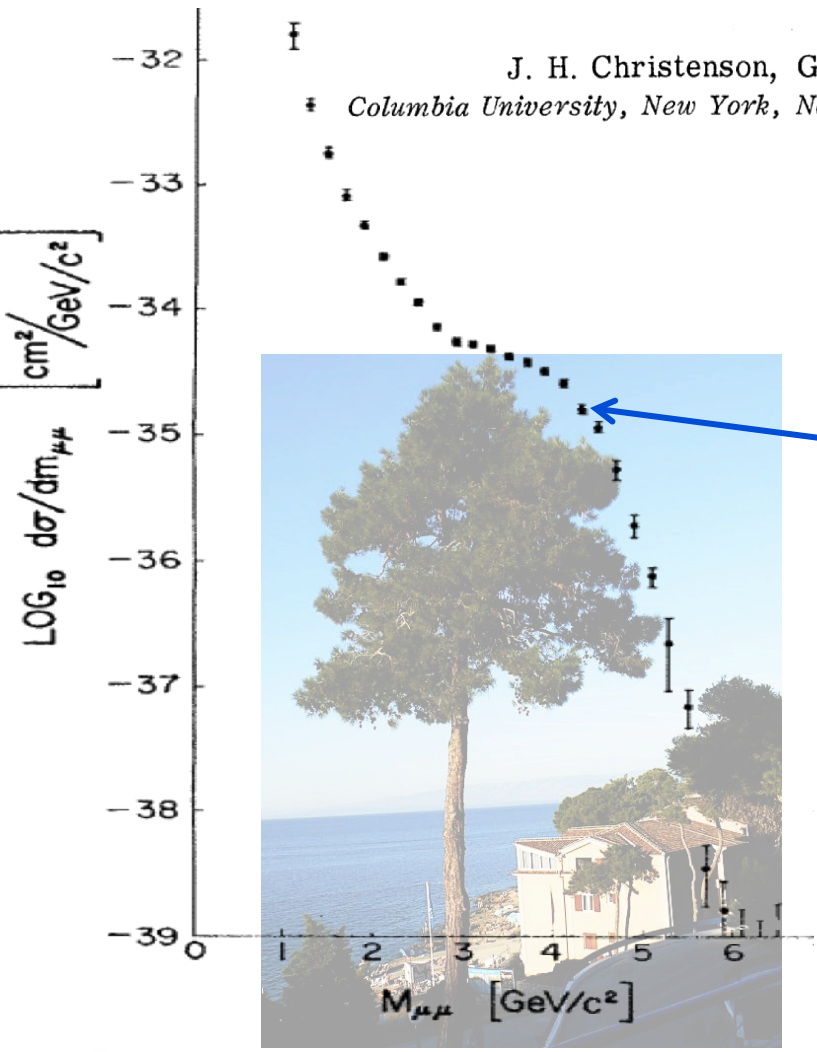
Columbia University, New York, New York 10027, and Brookhaven National Laboratory, Upton, New York 11973

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Don't miss the tree when looking at the beautiful scenery

Paul E. Reimer QCD-N12 Drell-Yan



Drell and Yan's explanation

VOLUME 25, NUMBER 5

PHYSICAL REVIEW LETTERS

3 AUGUST 1970

MASSIVE LEPTON-PAIR PRODUCTION IN HADRON-HADRON COLLISIONS AT HIGH ENERGIES*

Sidney D. Drell and Tung-Mow Yan

Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305

(Received 25 May 1970)

On the basis of a parton model studied earlier we consider the production process of large-mass lepton pairs from hadron-hadron inelastic collisions in the limiting region, $s \rightarrow \infty$, Q^2/s finite, Q^2 and s being the squared invariant masses of the lepton pair and the two initial hadrons, respectively. General scaling properties and connections with deep inelastic electron scattering are discussed. In particular, a rapidly decreasing cross section as $Q^2/s \rightarrow 1$ is predicted as a consequence of the observed rapid falloff of the inelastic scattering structure function νW_2 near threshold.

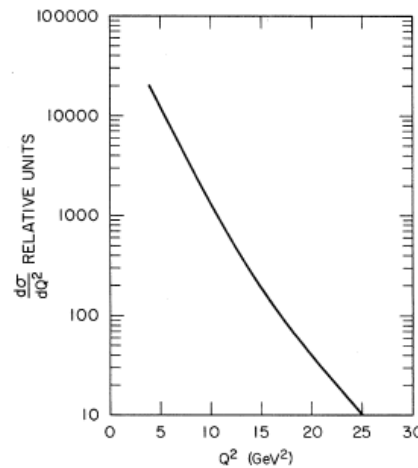
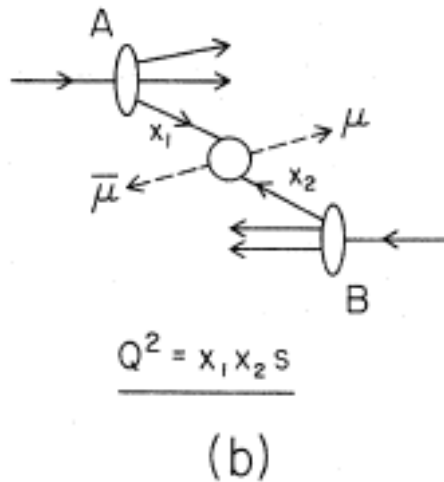
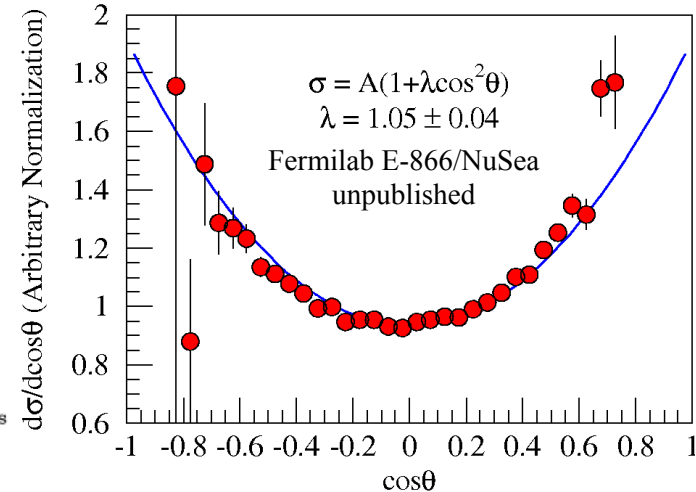


FIG. 2. $d\sigma/dQ^2$ computed from Eq. (10) assuming identical parton and antiparton momentum distributions and with relative normalization.

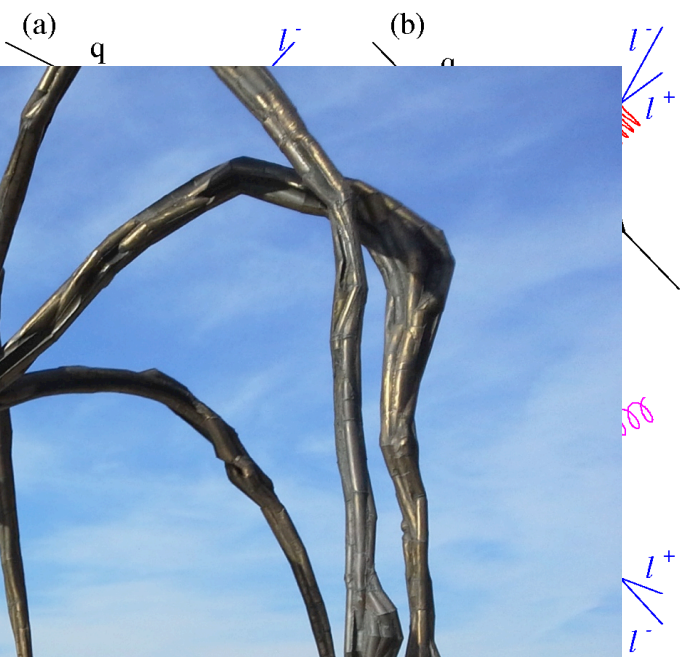
Also predicted
 $\lambda(1+\cos^2\theta)$
angular distributions.



Drell-Yan Cross Section

- These **of the**
- Parton**
- Intrinsic (altho
- Soft glu**

- $\frac{d\sigma}{d\Omega}$



Paul E. Reimer QCD-N12 Drell-Yan



Drell-Yan Cross Section

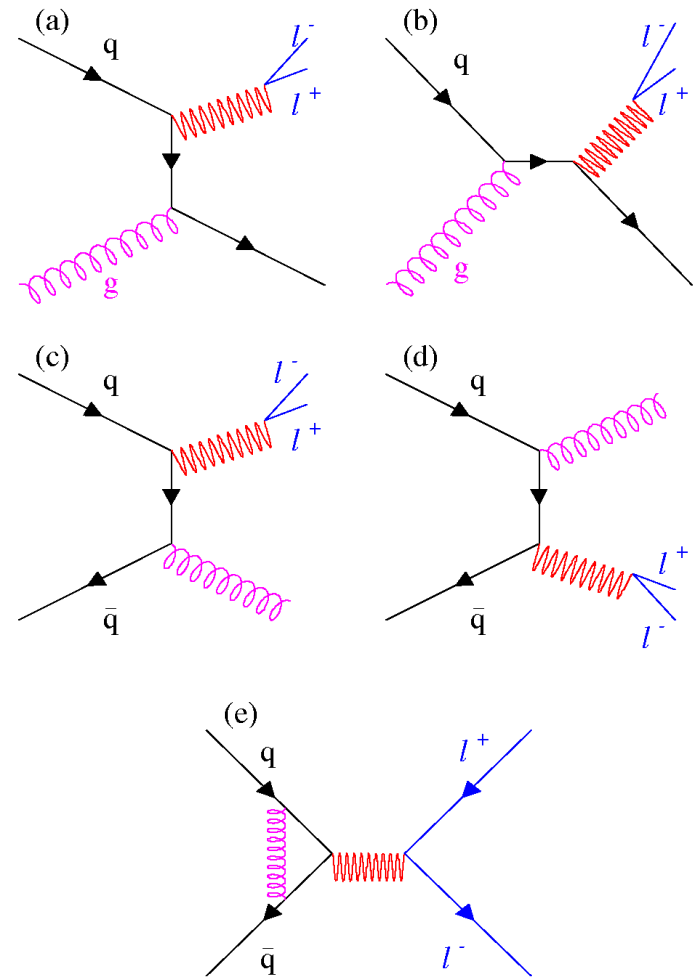
Next-to-Leading Order

- These diagrams are responsible for up to **50% of the measured cross section**
- Parton distributions are Universal!**
- Intrinsic transverse momentum of quarks (although a small effect, $\lambda > 0.8$)
- Soft gluon resummation at all orders**

Angular Distributions

- $$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

Higher Twist??

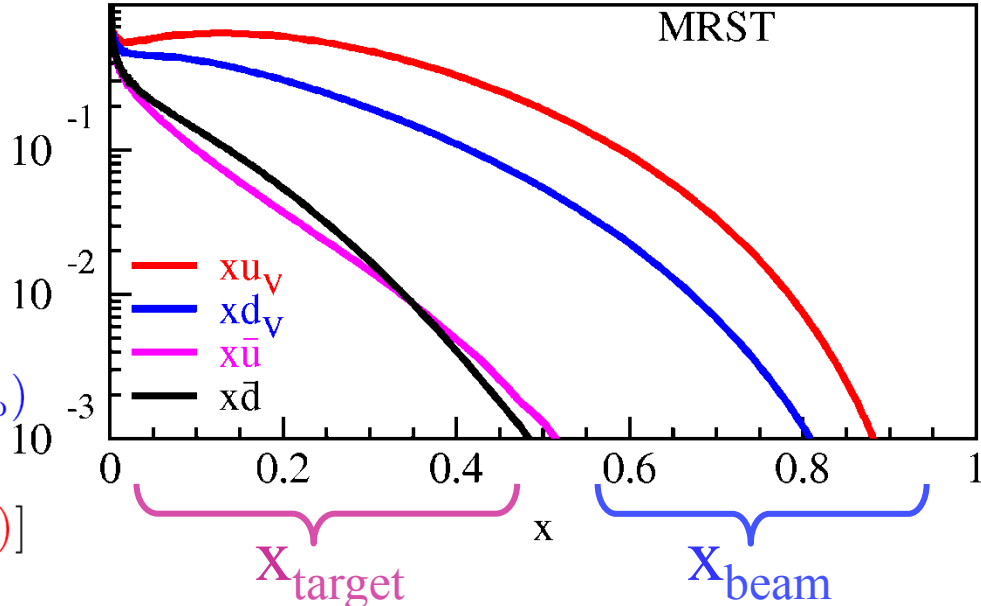


Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions
- Proton Beam**
 - Target antiquarks and beam

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b) + \bar{q}_b(x_b) q_t(x_t)]$$

u-quark dominance
 $(2/3)^2$ vs. $(1/3)^2$



Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions

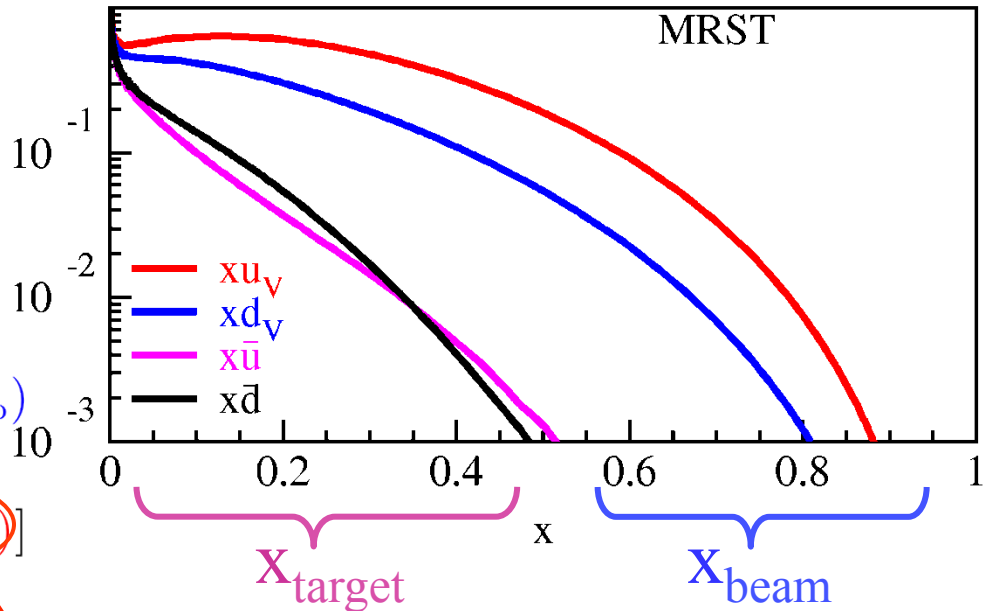
- Proton Beam**

- Target antiquarks and beam

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b) + \bar{q}_b(x_b) q_t(x_t)]$$

- u-quark dominance
 $(2/3)^2$ vs. $(1/3)^2$

$$+ \bar{q}_b(x_b) q_t(x_t)$$



Acceptance limited

Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions

- Proton Beam**

- Target antiquarks and beam

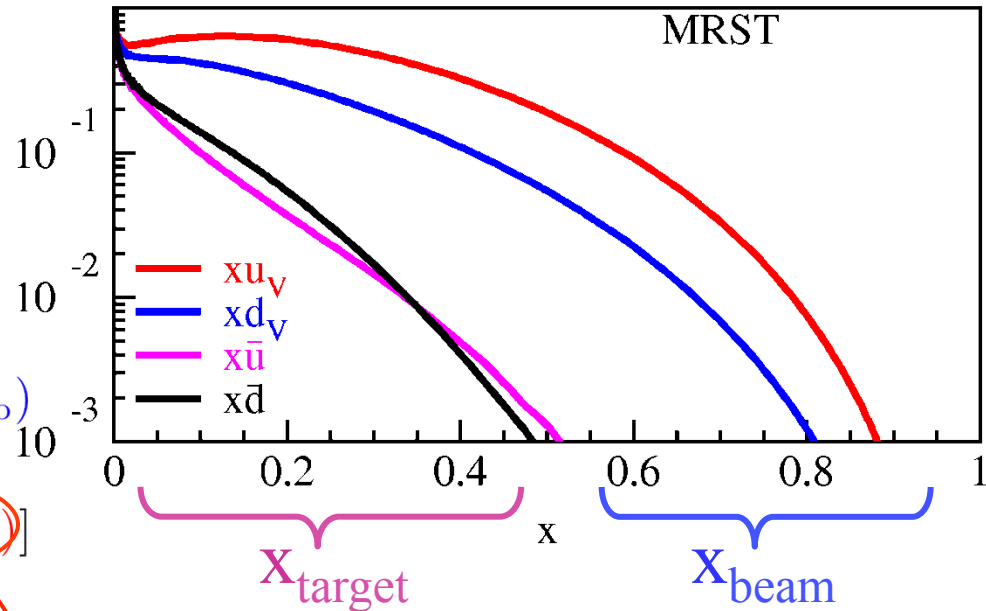
$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b)]$$

- u-quark dominance
(2/3)² vs. (1/3)²

- π beam**

- Valence beam anti-u quark and u target quark

$$\left. \frac{d^2\sigma}{dx_\pi dx_N} \right|_{\pi^- N} = \frac{4\pi\alpha^2}{x_\pi x_N s} \left[\frac{4}{9} \bar{u}_\pi(x_\pi) u_N(x_N) + \frac{1}{9} d_\pi(x_\pi) \bar{d}_N(x_N) + \frac{4}{9} u_\pi(x_\pi) \bar{u}_N(x_N) + \frac{1}{9} \bar{d}_\pi(x_\pi) d_N(x_N) \right]$$



Acceptance limited

Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions

- Proton Beam**

- Target antiquarks and beam

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b)]$$

- u-quark dominance
(2/3)² vs. (1/3)²

- π beam**

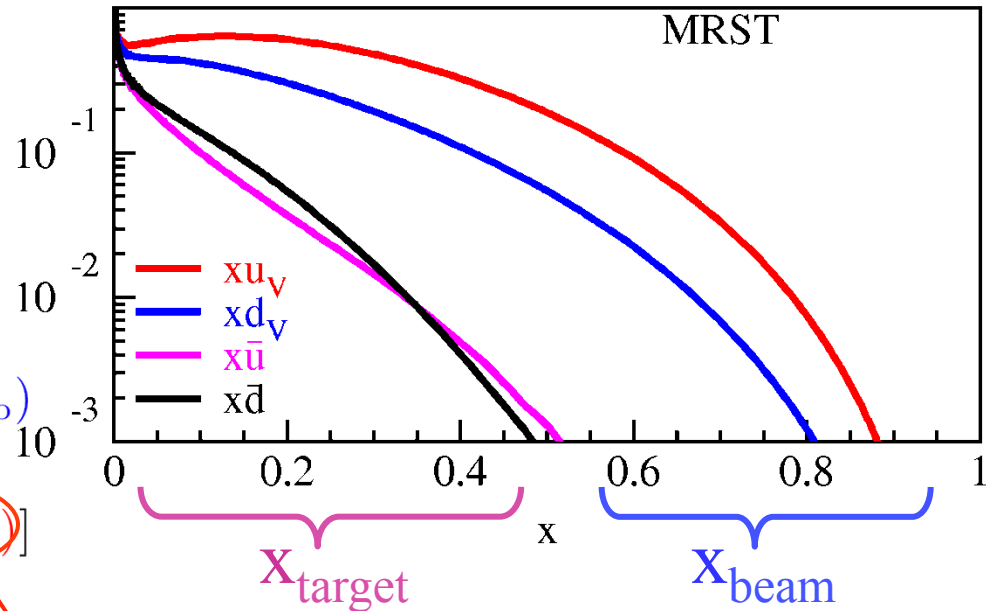
- Valence beam anti-u quark and u target quark

$$\left. \frac{d^2\sigma}{dx_\pi dx_N} \right|_{\pi^- N} = \frac{4\pi\alpha^2}{x_\pi x_N s} \left[\begin{aligned} & \frac{4}{9} \bar{u}_\pi(x_\pi) u_N(x_N) \\ & + \frac{1}{9} d_\pi(x_\pi) \bar{d}_N(x_N) \\ & + \frac{4}{9} u_\pi(x_\pi) \bar{u}_N(x_N) \\ & + \frac{1}{9} \bar{d}_\pi(x_\pi) d_N(x_N) \end{aligned} \right]$$

Valence × Valence → $\frac{4}{9} \bar{u}_\pi(x_\pi) u_N(x_N)$

Valence-sea × 1/4 → $\frac{1}{9} d_\pi(x_\pi) \bar{d}_N(x_N)$

Sea-Sea → $\frac{1}{9} \bar{d}_\pi(x_\pi) d_N(x_N)$



Acceptance limited

Beam	Target	Experiment
Hadron	Beam valence quarks Target antiquarks	Fermilab E-906, RHIC (forward acpt.) J-PARC
Anti-Hadron	Beam val. antiquarks Target valence quarks	GSF-FAIR Fermilab Collider
Meson	Beam val. antiquarks Target valence quarks	COMPASS

Unpolarized Drell-Yan—Boring?



Unpolarized Drell-Yan—No Exciting!!!!

- Measurements of longitudinal structure of target and beam
- Measurements of Boer-Mulders through angular distributions



Drell-Yan Cross Section

- Measured cross section is a convolution of beam and target parton distributions

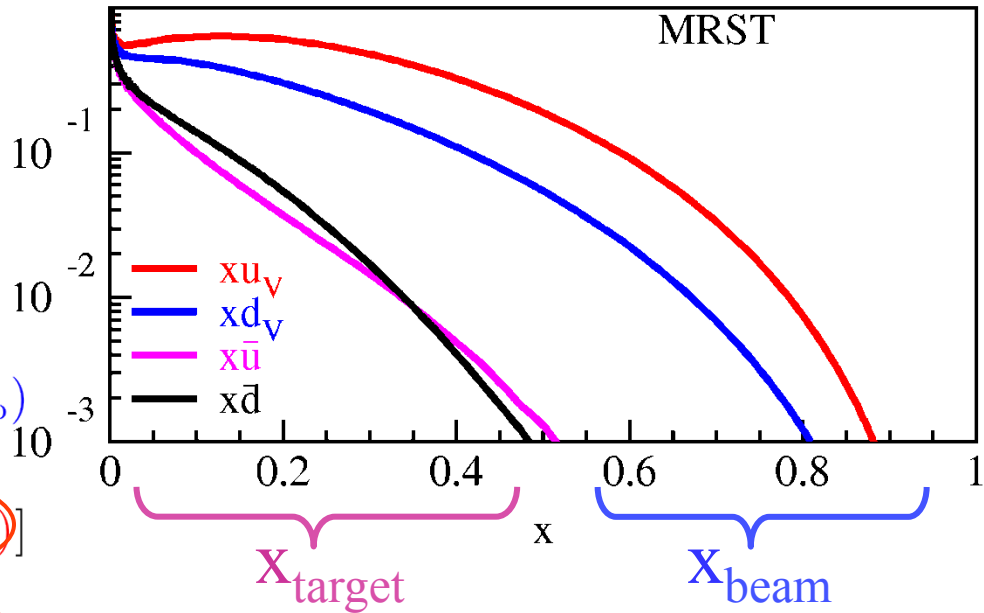
- Proton Beam**

- Target antiquarks and beam

$$\frac{d^2\sigma}{dx_b dx_t} = \frac{4\pi\alpha^2}{x_b x_t s} \sum_{q \in \{u, d, s, \dots\}} e_q^2 [\bar{q}_t(x_t) q_b(x_b) + \bar{q}_b(x_b) q_t(x_t)]$$

- u-quark dominance
(2/3)² vs. (1/3)²

$$+ \bar{q}_b(x_b) q_t(x_t)$$



Acceptance limited

- In leading order w/above assumptions

$$\frac{\sigma^{pd}}{2\sigma^{pp}} = \frac{1}{2} \left[1 + \frac{\bar{d}(x)}{\bar{u}(x)} \right]$$

Next-to-leading order terms have minimal effect on the ratio

- Suggested by Martin, Stirling and Roberts
Phys.Lett. B308 (1993) 377

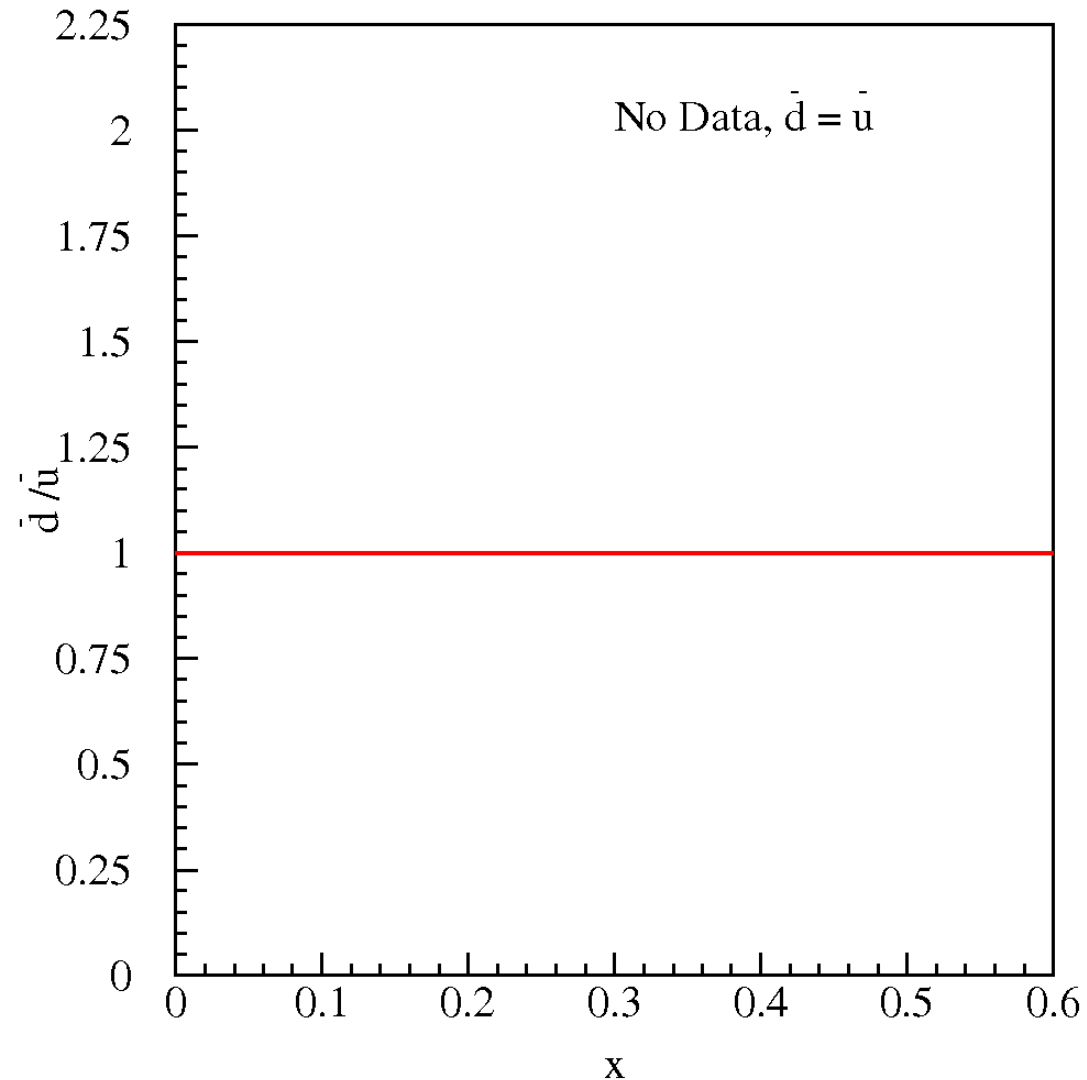
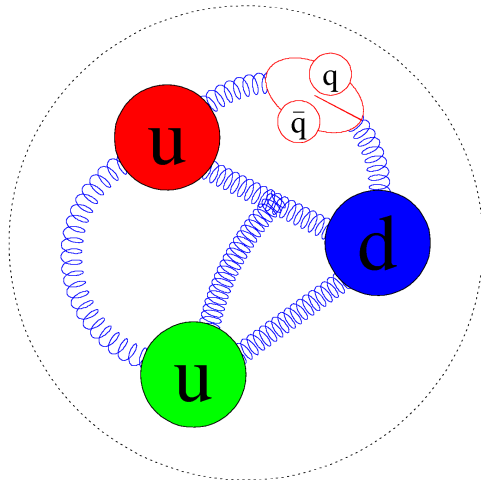
Light Antiquark Flavor Asymmetry: Brief History

Naïve Assumption:

$$\bar{d}(x) \equiv \bar{u}(x)$$

Why shouldn't it be

- the sea is generated by gluon splitting.
- Gluons couple to color, not flavor!



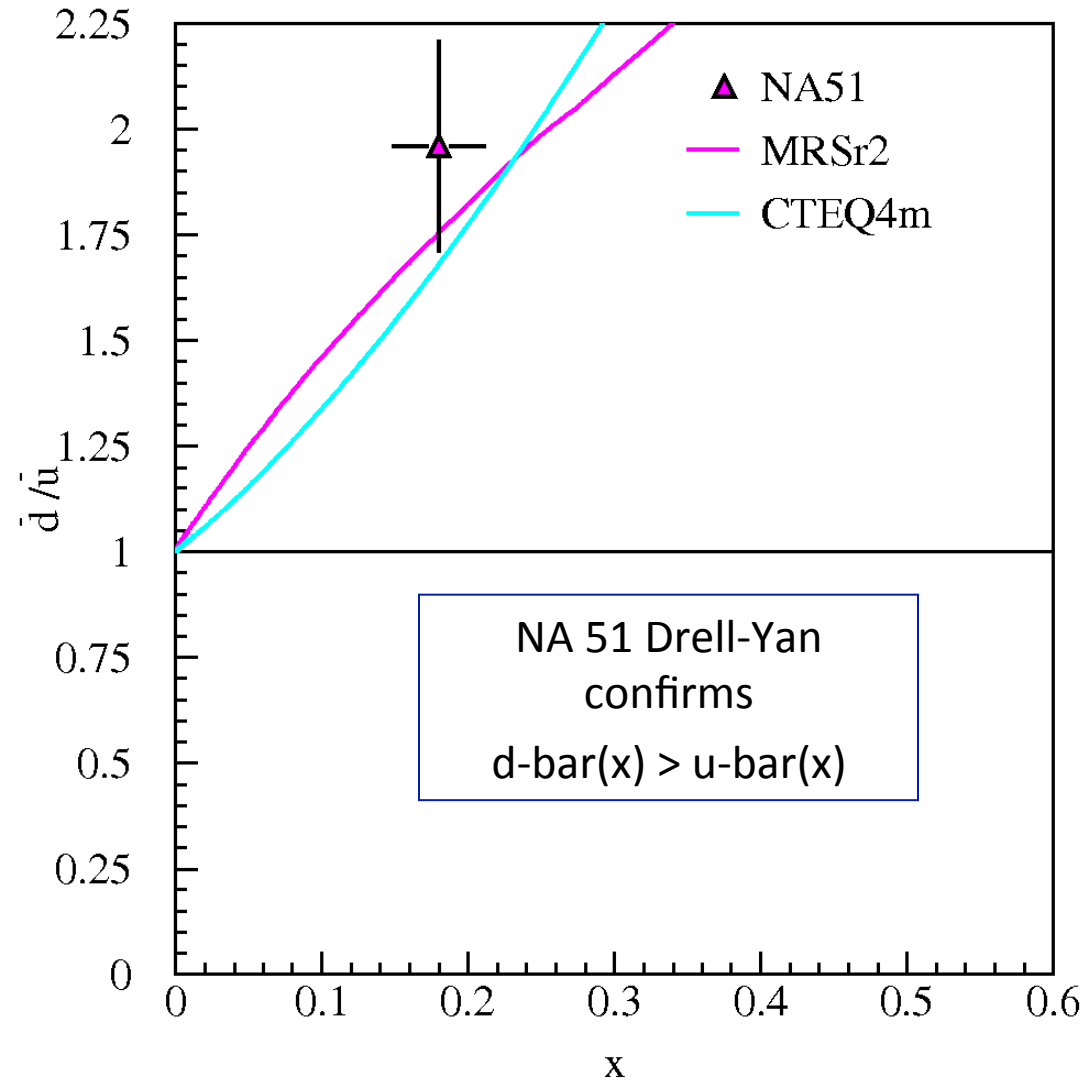
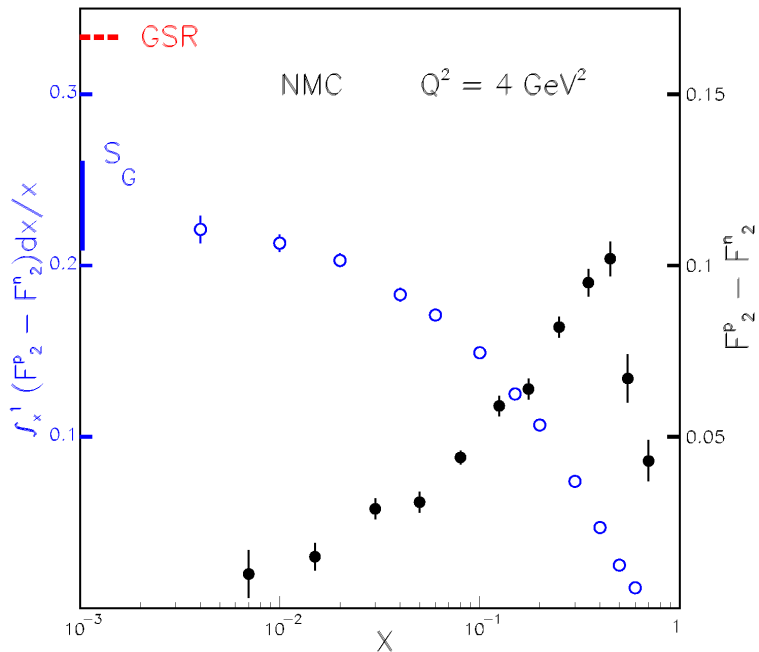
Light Antiquark Flavor Asymmetry: Brief History

- Naïve Assumption:

$$\bar{d}(x) = \bar{u}(x)$$

- CERN NMC (Gottfried Sum Rule)

$$\int_0^1 [\bar{d}(x) - \bar{u}(x)] dx \neq 0$$



Light Antiquark Flavor Asymmetry: Brief History

- Naïve Assumption:

$$\bar{d}(x) = \bar{u}(x)$$

- NMC (Gottfried Sum Rule)

$$\int_0^1 [\bar{d}(x) - \bar{u}(x)] dx \neq 0$$

- NA51 (Drell-Yan)

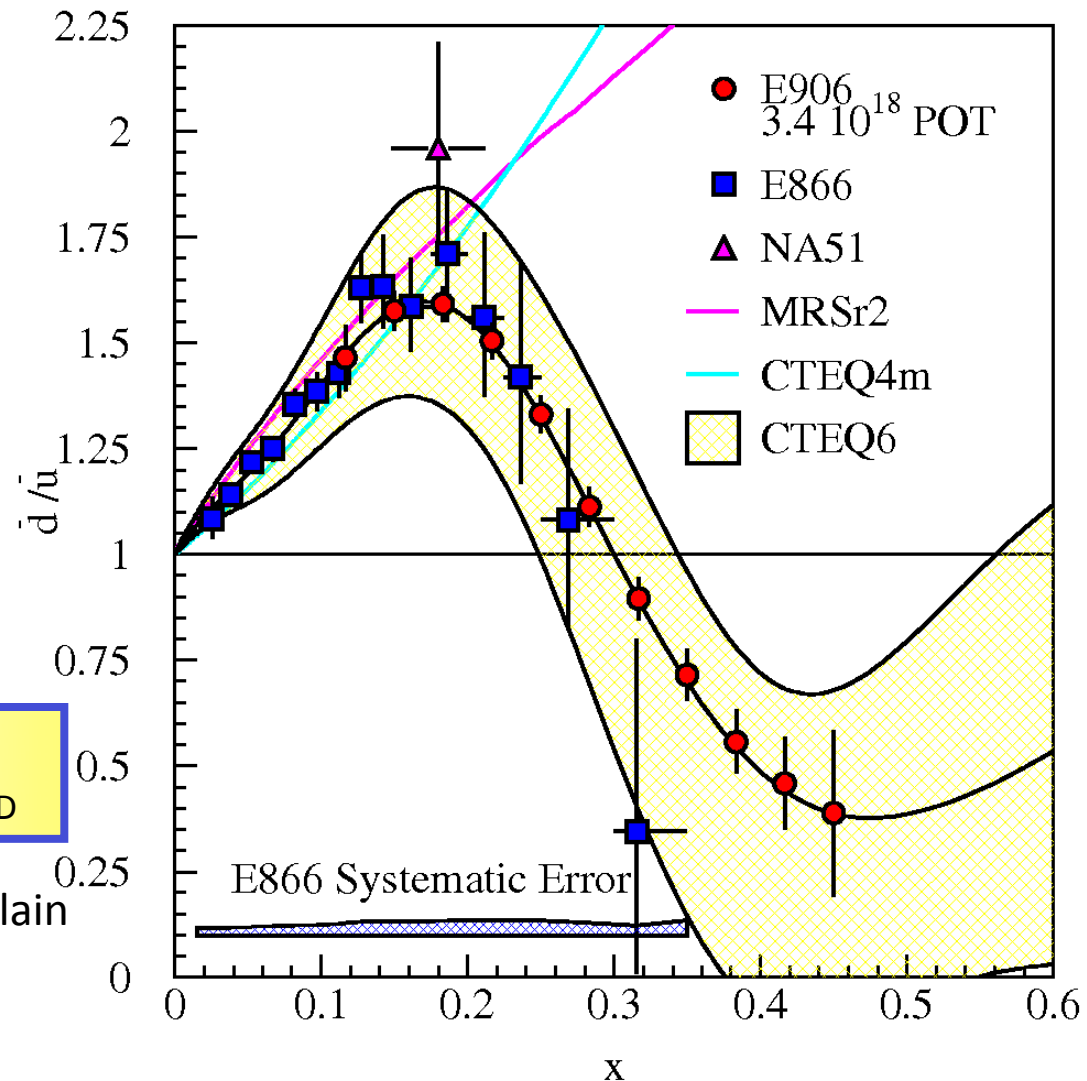
$$\bar{d} > \bar{u} \text{ at } x = 0.18$$

- E866/NuSea (Drell-Yan)

$$\bar{d}(x)/\bar{u}(x) \text{ for } 0.015 \leq x \leq 0.35$$

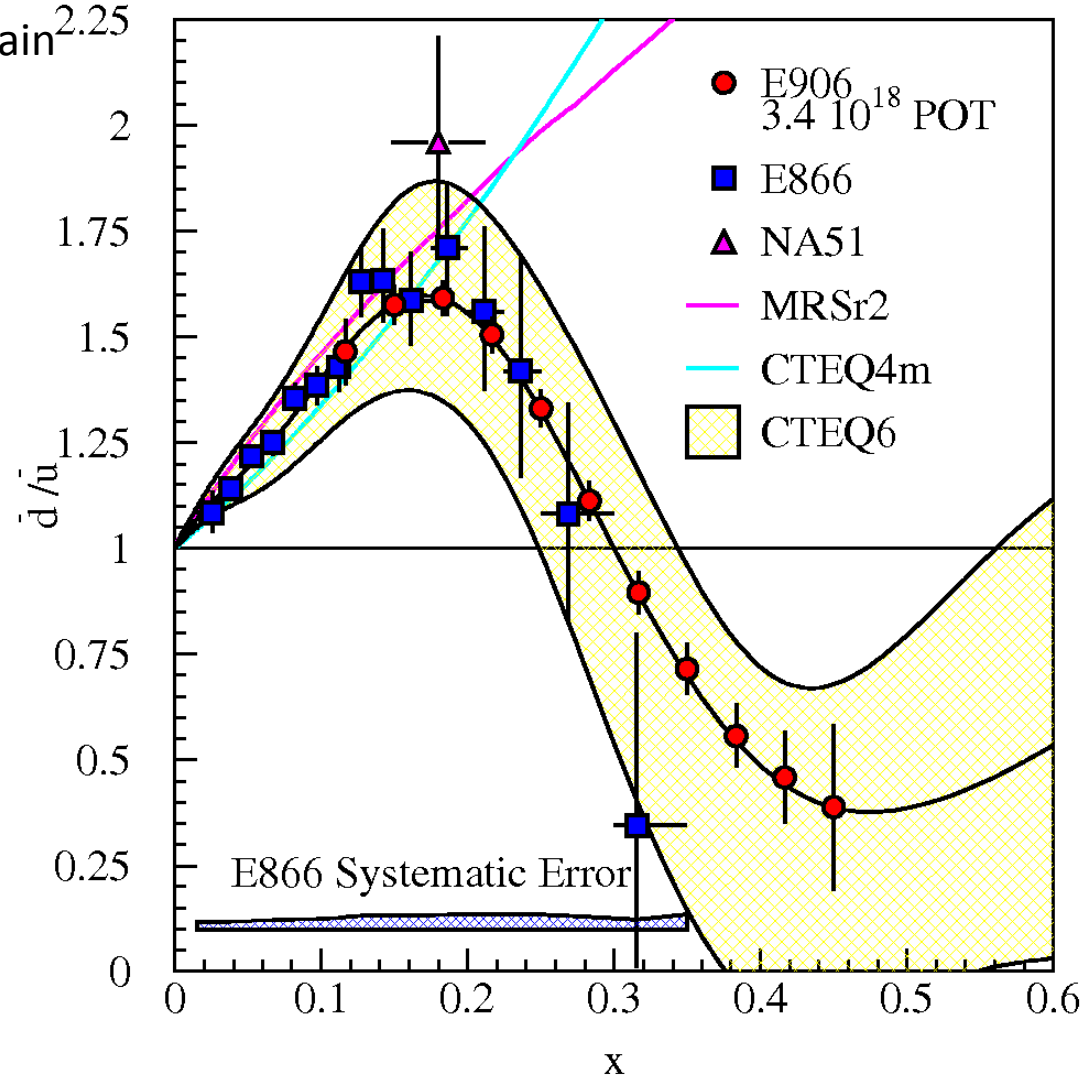
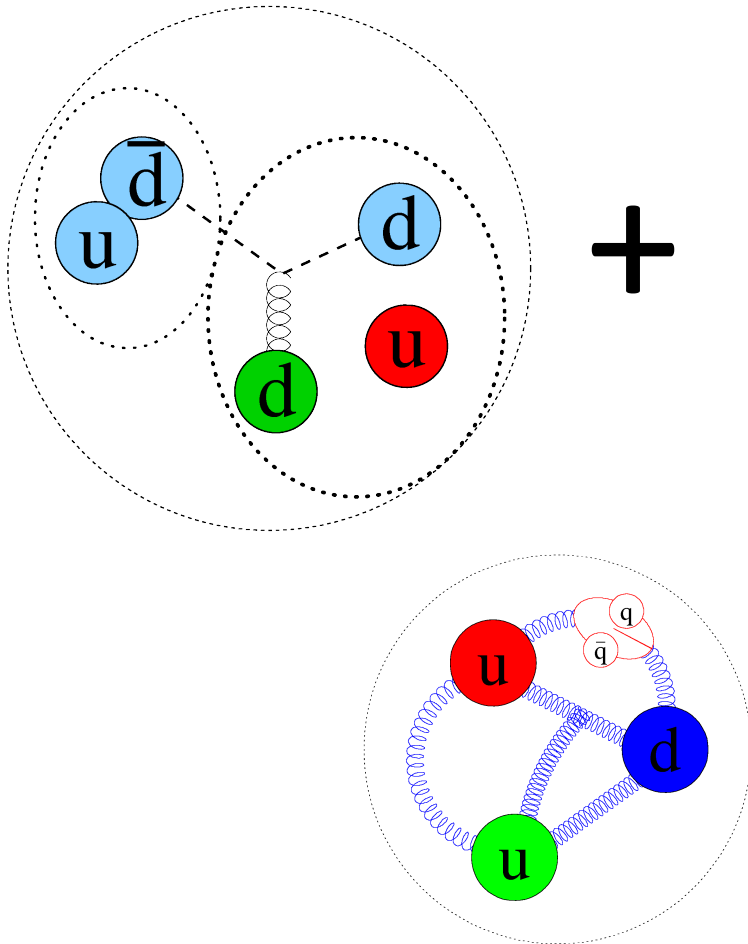
- Knowledge of sea dist. are data driven
- Sea quark distributions are difficult for Lattice QCD

- Non perturbative QCD models can explain excess d-bar quarks, but not return to symmetry or deficit of d-bar quarks



Light Antiquark Flavor Asymmetry: Brief History

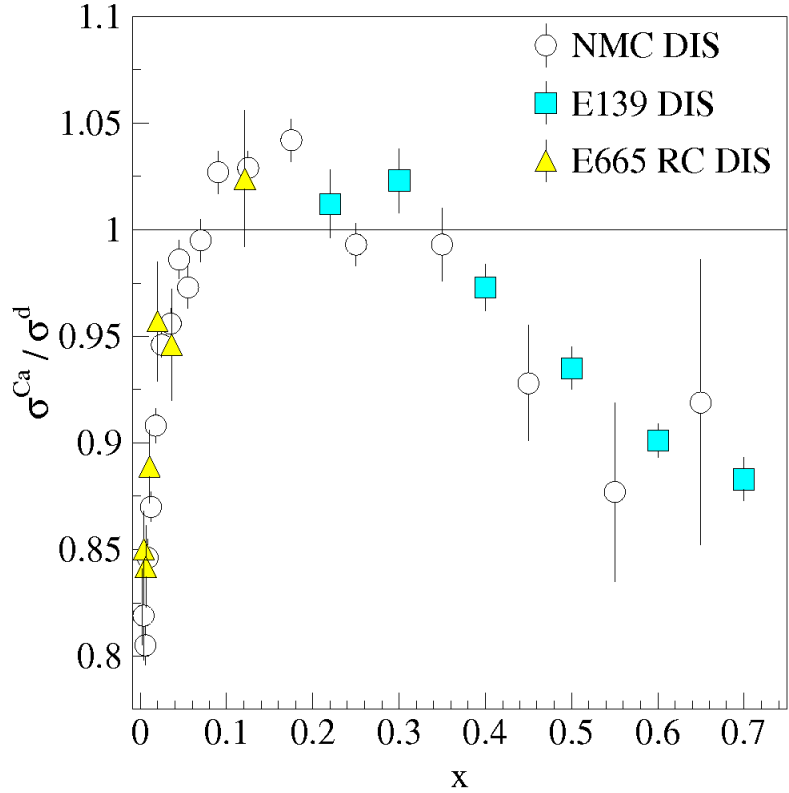
- Non perturbative QCD models can explain^{2.25} excess d-bar quarks, but not return to symmetry or deficit of d-bar quarks



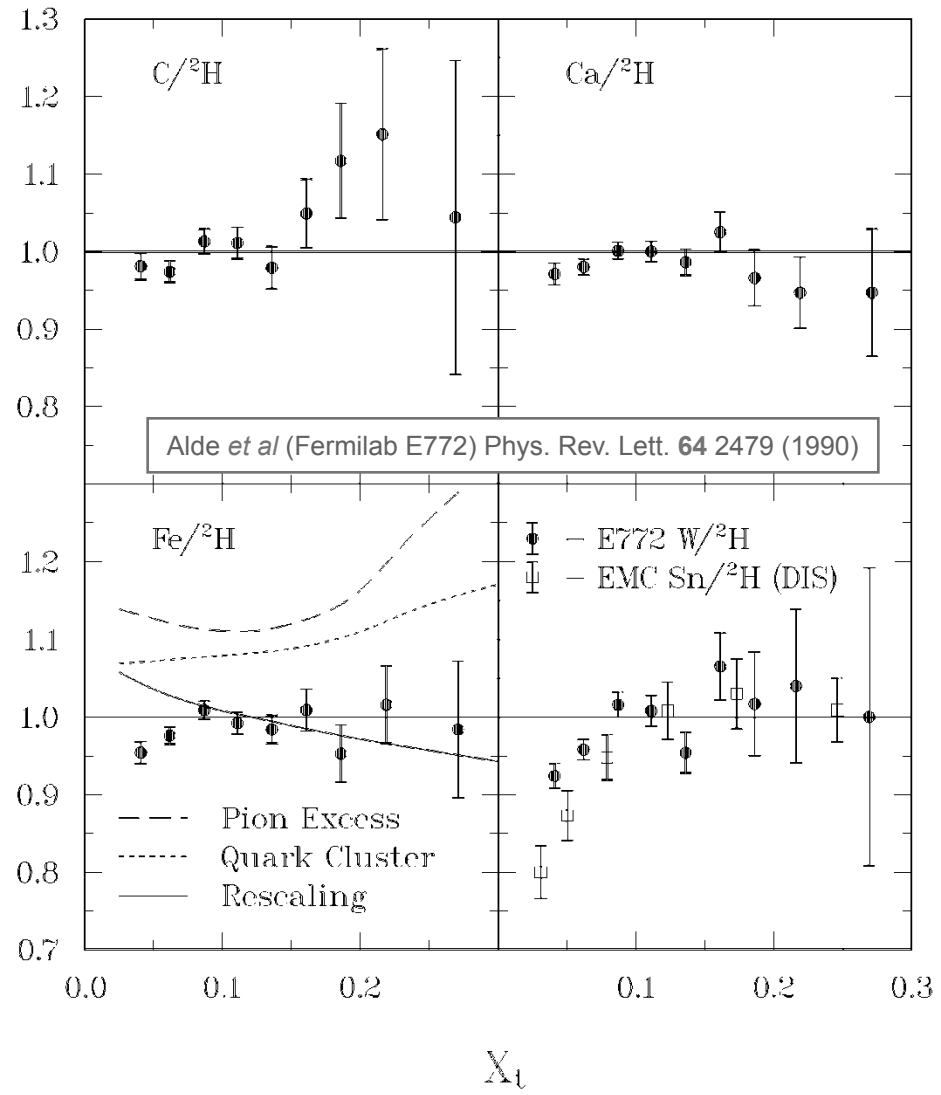
Structure of nucleonic matter: How do sea quark distributions differ in a nucleus?

Comparison with
Deep Inelastic Scattering (DIS)

- EMC: Parton distributions of bound and free nucleons are different.
- Antishadowing not seen in Drell-Yan — Valence only effect

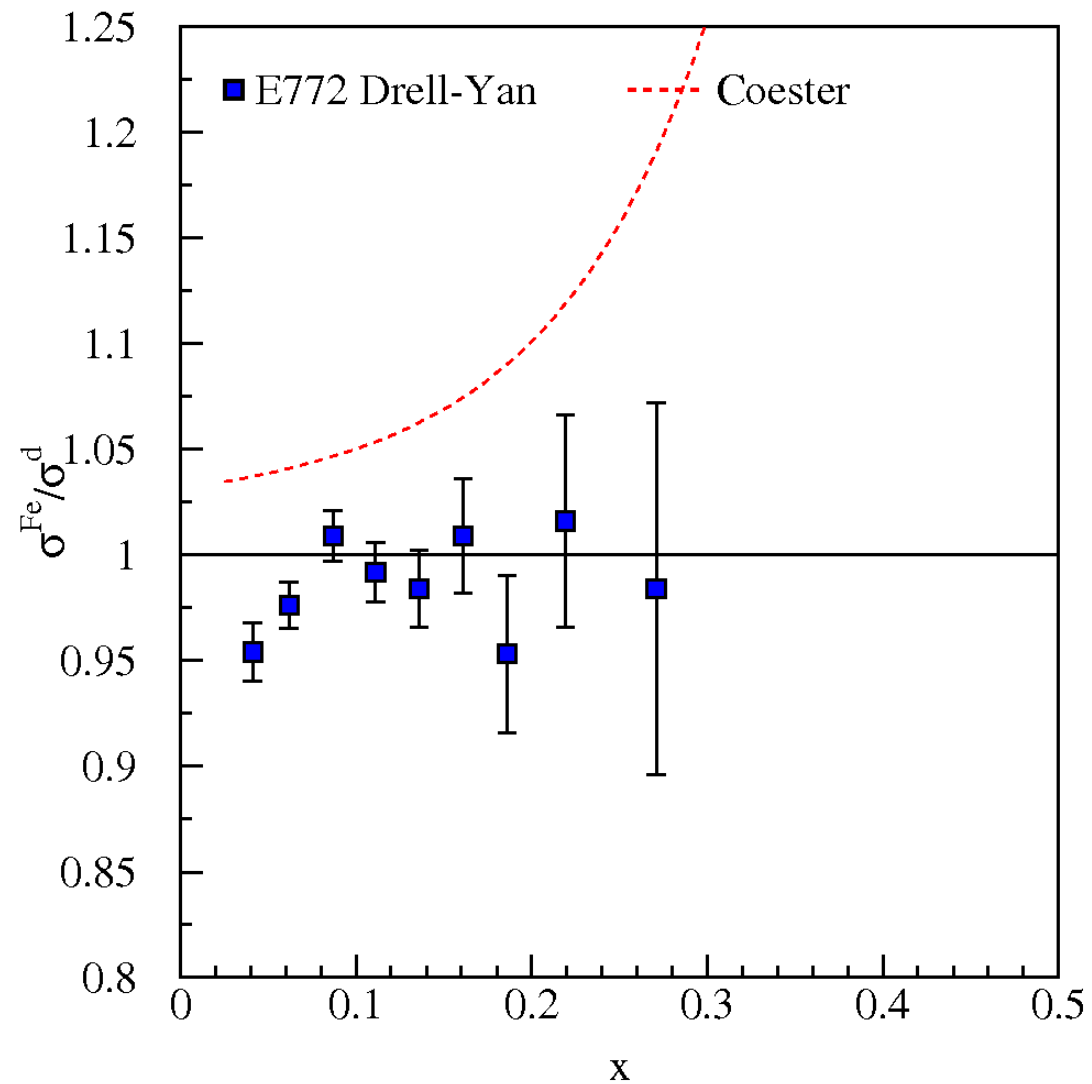


Drell-Yan Ratio



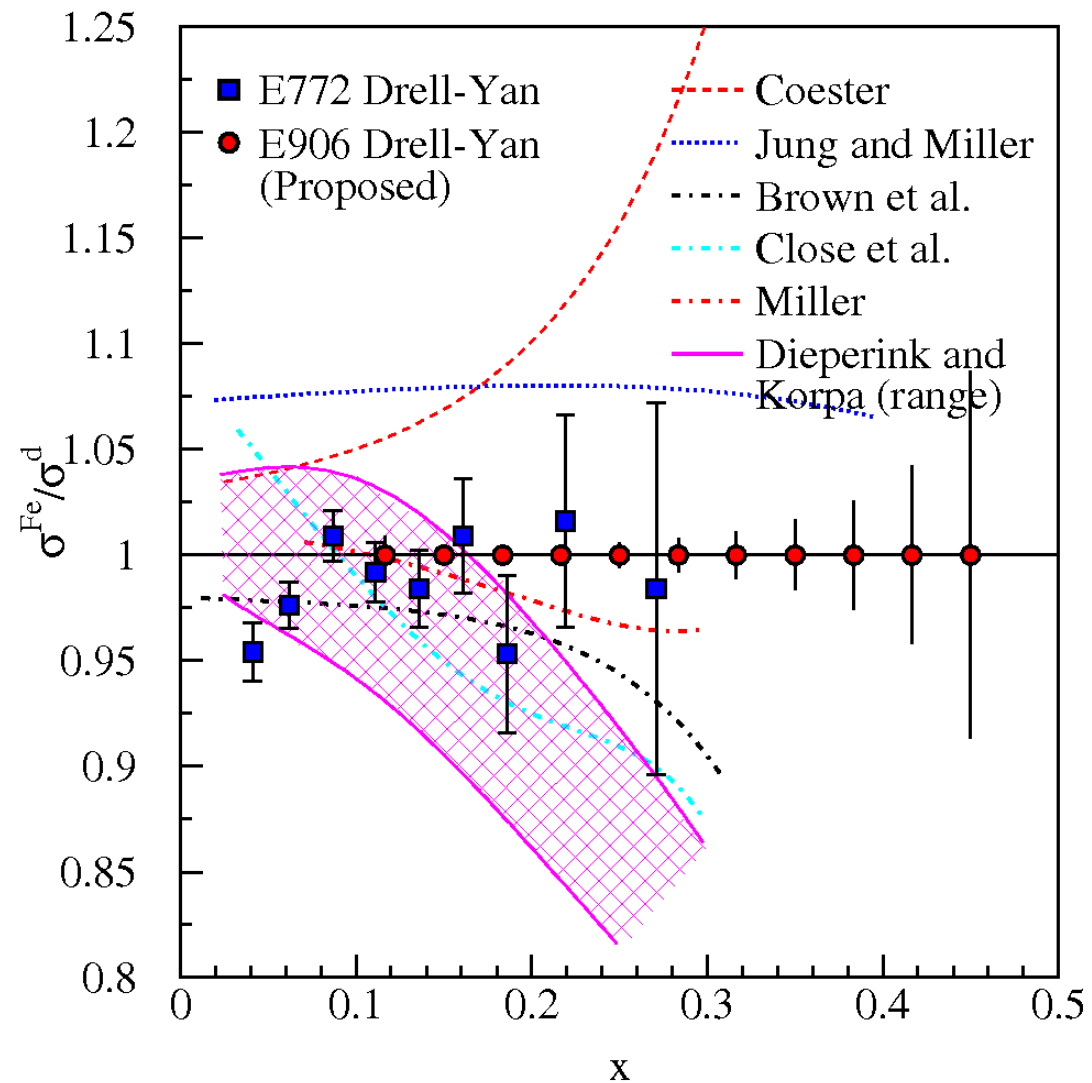
Structure of nucleonic matter: Where are the nuclear pions?

- The binding of nucleons in a nucleus is expected to be governed by the exchange of virtual “Nuclear” mesons (pions).
- No antiquark enhancement seen in Drell-Yan (Fermilab E772) data.



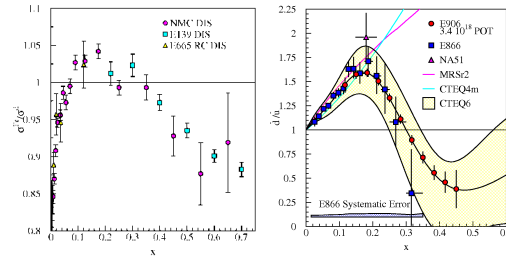
Structure of nucleonic matter: Where are the nuclear pions?

- The binding of nucleons in a nucleus is expected to be governed by the exchange of virtual “Nuclear” mesons.
- No antiquark enhancement seen in Drell-Yan (Fermilab E772) data.
- Contemporary models predict large effects to antiquark distributions as x increases.
- Models must explain both DIS-EMC effect and Drell-Yan



Pionic Parton Distributions: Why are should you be interested in the pion?

1. Explains quark asymmetry in nucleon sea.
2. Key role in nuclear binding and structure
3. Simple quark-antiquark system (valence).
 - *Should* have an easy theoretical interpretation.
 - What about mass—constituent quark mass ≈ 300 MeV each?
3. QCD's Goldstone Boson



How do we treat 3 and 4 in a unified way?

Models predict different behavior as $x \rightarrow 1$

At some base Q_0

$$\text{pQCD: } xq(x)/(1-x)^\beta \quad \beta = 2$$

$$\text{NJL: } xq(x)/(1-x)^\beta \quad \beta = 1$$

$$\text{DSE: } xq(x)/(1-x)^\beta \quad \beta \sim 1.9$$

Evolution to experimental Q increases β .

Pionic Parton Distributions: Why are should you be interested in the pion?

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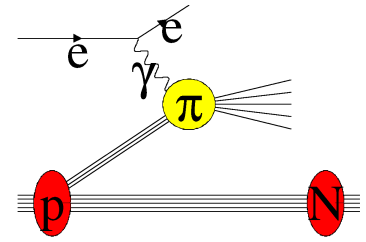
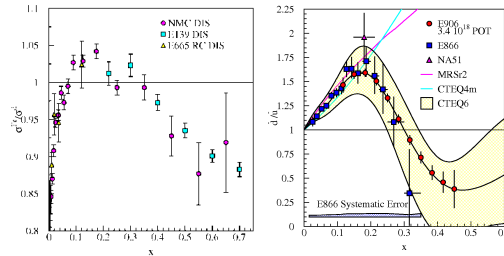
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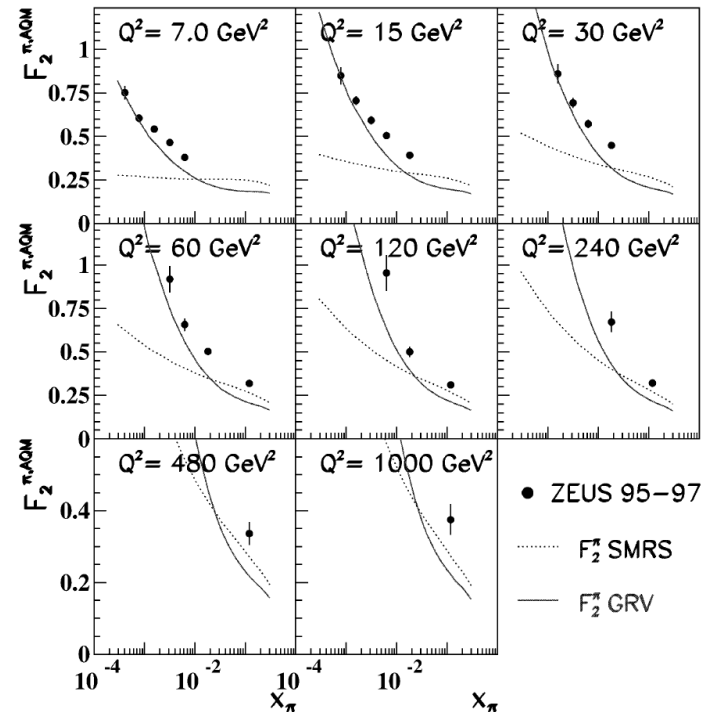


“Pion targets are not abundant” Hecht

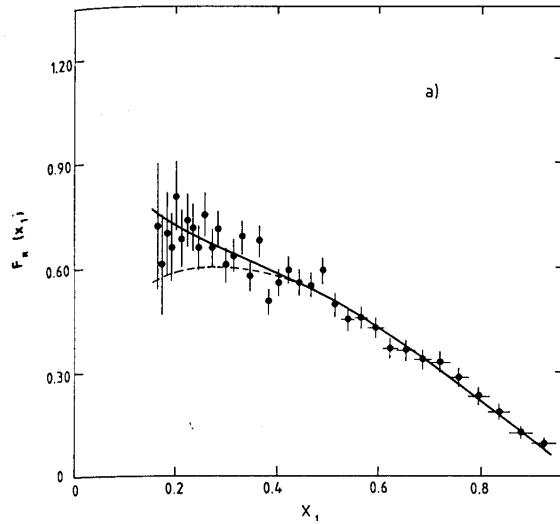
DIS on virtual pions:

$ep \rightarrow eN x$ HERA data [ZEUS, NPB637 3 (2002)]

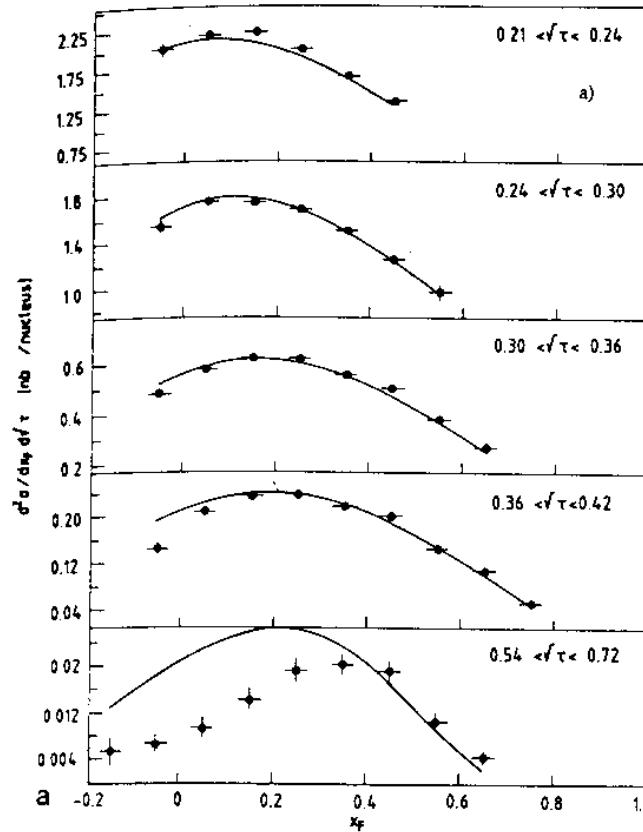
Possible JLab and EIC.



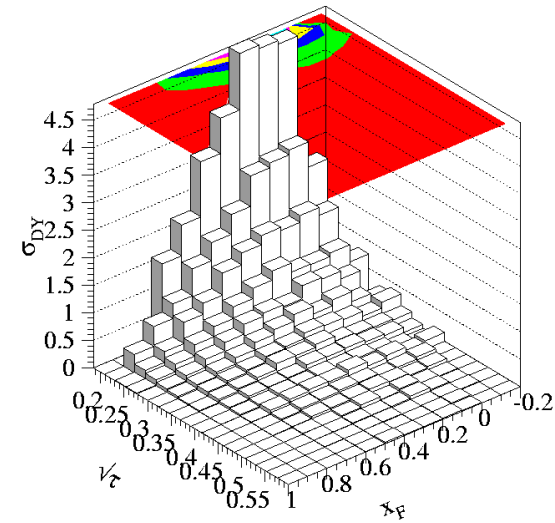
Pion Drell-Yan Data: CERN NA3 (π^+) NA10 (π^-)



CERN NA3 200 GeV π^- data
(also have 150 and 180 GeV π^-
and 200 GeV π^+ data).

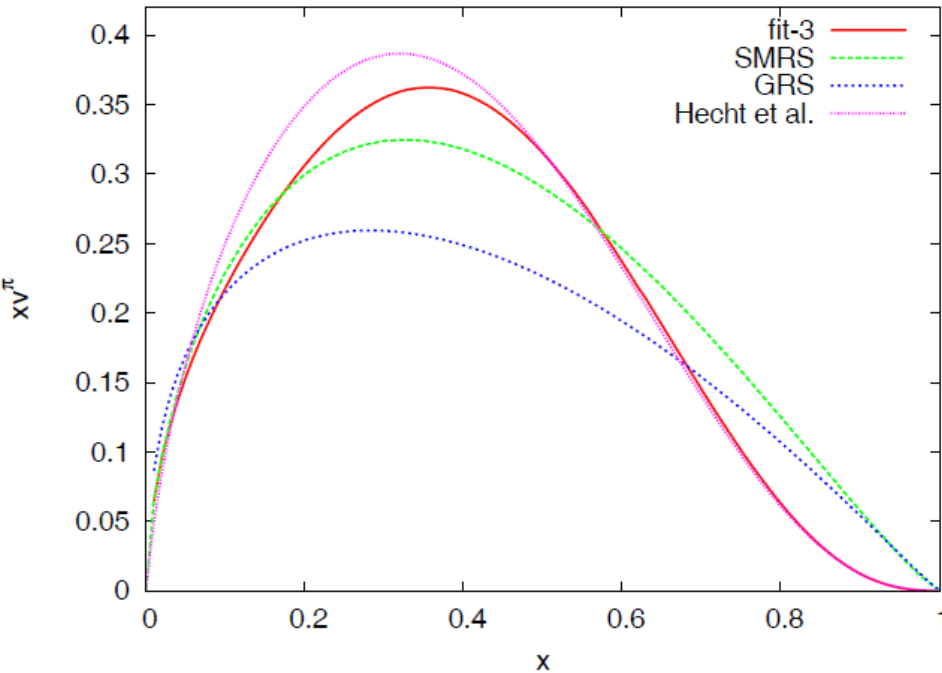


CERN NA10 194 GeV π^- data

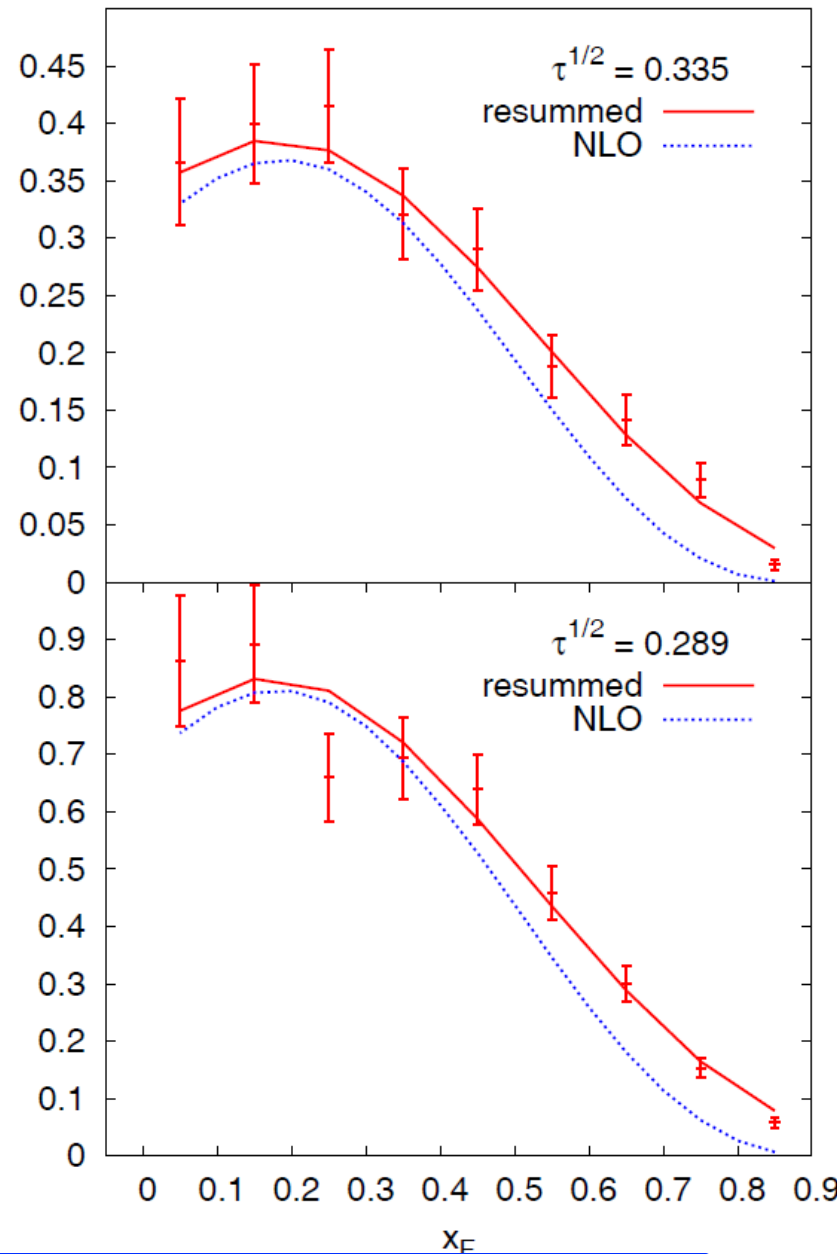


Fermilab E-615 π^- data

Soft Gluon Resummation



$d\sigma/d\tau^{1/2} dx_F$ (nb/nucleon)



$$xq_V^\pi(x) = A_V^\pi x^\alpha (1-x)^\beta (1 + \gamma x^\delta)$$

$$\beta = 2.03 \pm 0.06$$

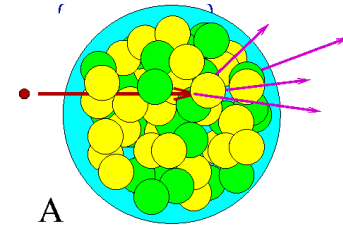
QCD and Dyson-Schwinger survive!

pQCD: $xq(x)/(1-x)^\beta$ $\beta = 2$

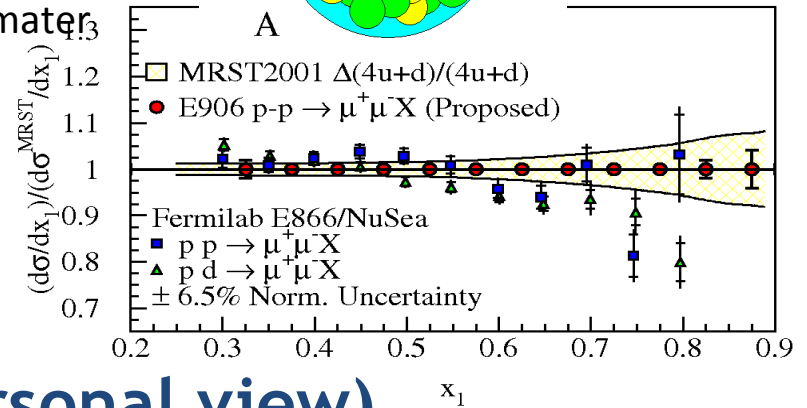
DSE: $xq(x)/(1-x)^\beta$ $\beta \sim 1.9$

Caveat: Better data would really help establish this.

Longitudinal Structure: Other topics



- Partonic energy loss
 - quarks traveling through cold, strongly interacting matter
 - Are antiquarks the same?
- Large-x proton parton distributions



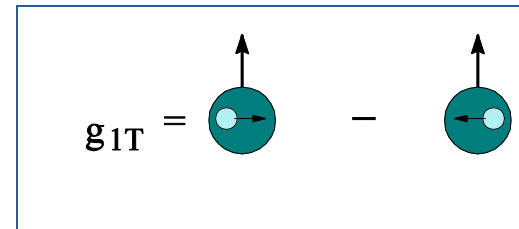
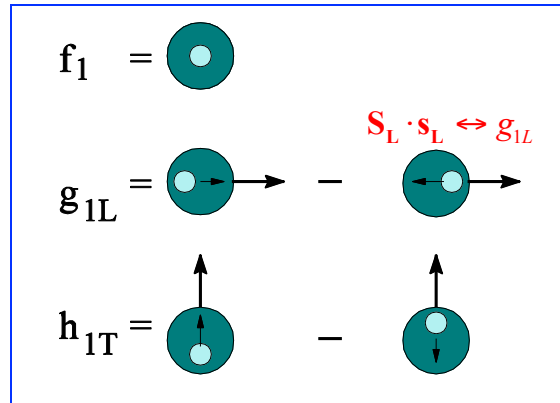
Longitudinal Structure: What data would I like to see? (My personal view)

- Proton deuterium Drell-Yan ratios $-\bar{d}/\bar{u}$
 - Fermilab E-906/SeaQuest
- Deuterium/A ratios—EMC and nuclear binding
 - Fermilab E-906/SeaQuest
- Unpolarized π -p Drell-Yan—nature of Goldstone Boson
 - CERN Compass
 - K-p Drell-Yan (another test of low energy QCD)
- Double Polarized Drell-Yan— $\Delta q(x)$ nice complement for SIDIS and W^\pm
 - CERN Compass, J-PARC, RHIC



Transverse Momentum Distributions: Introduction

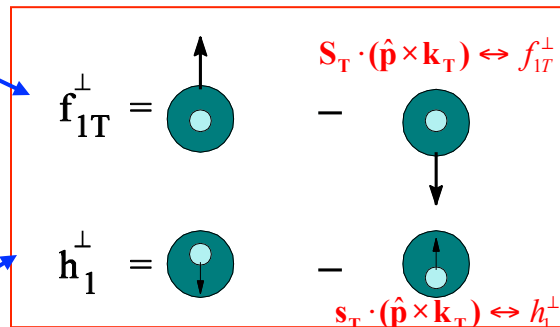
Survive k_T integration



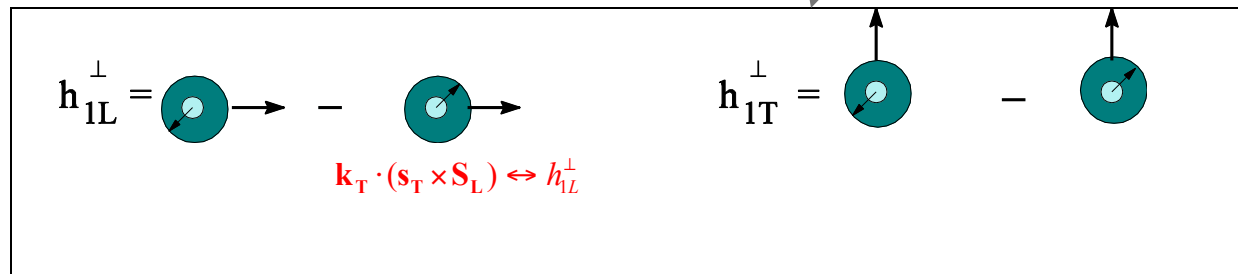
k_T - dependent, T-even

Sivers' Function

k_T - dependent,
Naive T-odd

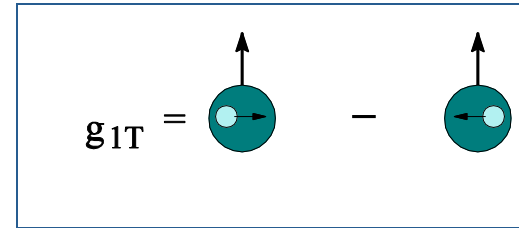
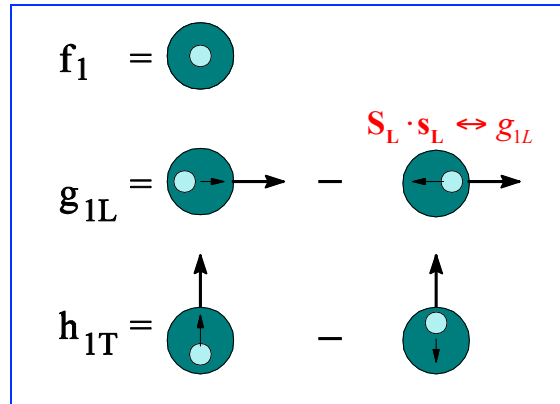


Boer-Mulders Function



Transverse Momentum Distributions: Introduction

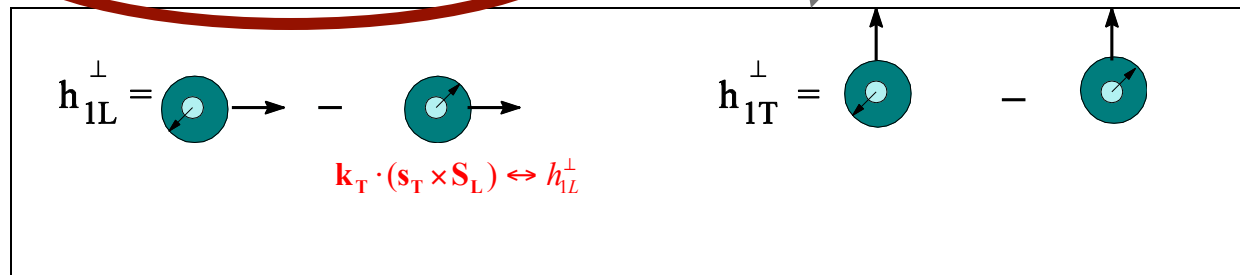
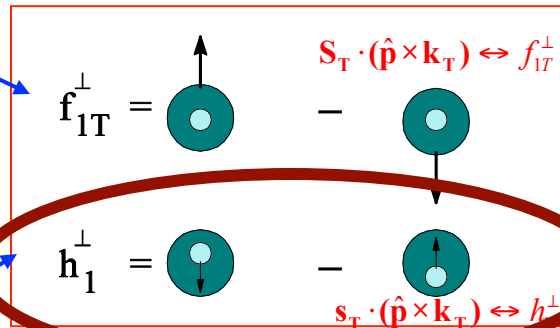
Survive k_T integration



k_T - dependent, T-even

Sivers' Function

k_T - dependent,
Naive T-odd



Boer-Mulders Function

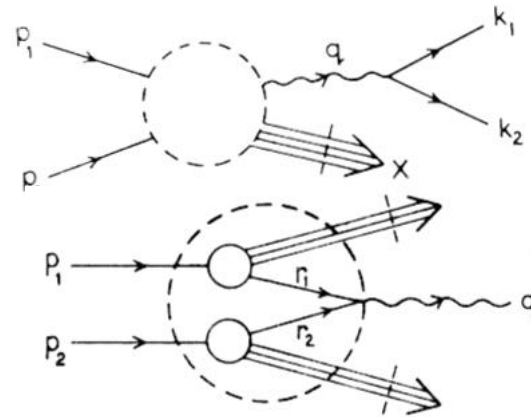
Generalized Angular Distributions & Lam Tung relation

Chi-Sing Lam and Wu-Ki Tung—basic formula for lepton pair production angular distributions PRD 18 2447 (1978)

$$\frac{d\sigma}{d^4q d\Omega_k^*} = \frac{1}{2} \frac{1}{(2\pi)^4} \left(\frac{\alpha}{M_S} \right)^2 \left[W_T (1 + \cos^2 \theta) + W_L (1 - \cos^2 \theta) \right. \\ \left. + W_{\Delta} \sin 2\theta \cos \phi + W_{\Delta\Delta} \sin^2 \theta \cos 2\phi \right]$$

Structure function formalism

- Derived in analogy to DIS
- Independent of Drell-Yan and parton “models”
- Showed same relations follow as a general consequence of the quark-parton model



Lam-Tung relation

- Derived in analogy to Colin-Gross relation of DIS

$$W_T = 2W_{\Delta\Delta}$$

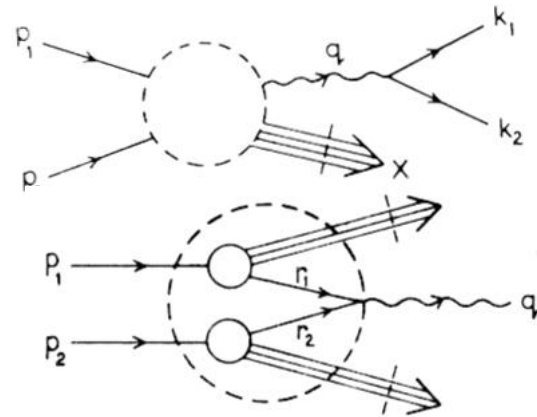
Generalized Angular Distributions & Lam Tung relation

Chi-Sing Lam and Wu-Ki Tung—basic formula for lepton pair production angular distributions PRD 18 2447 (1978)

$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

Structure function formalism

- Derived in analogy to DIS
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- Showed same relations follow as a general consequence of the quark-parton model

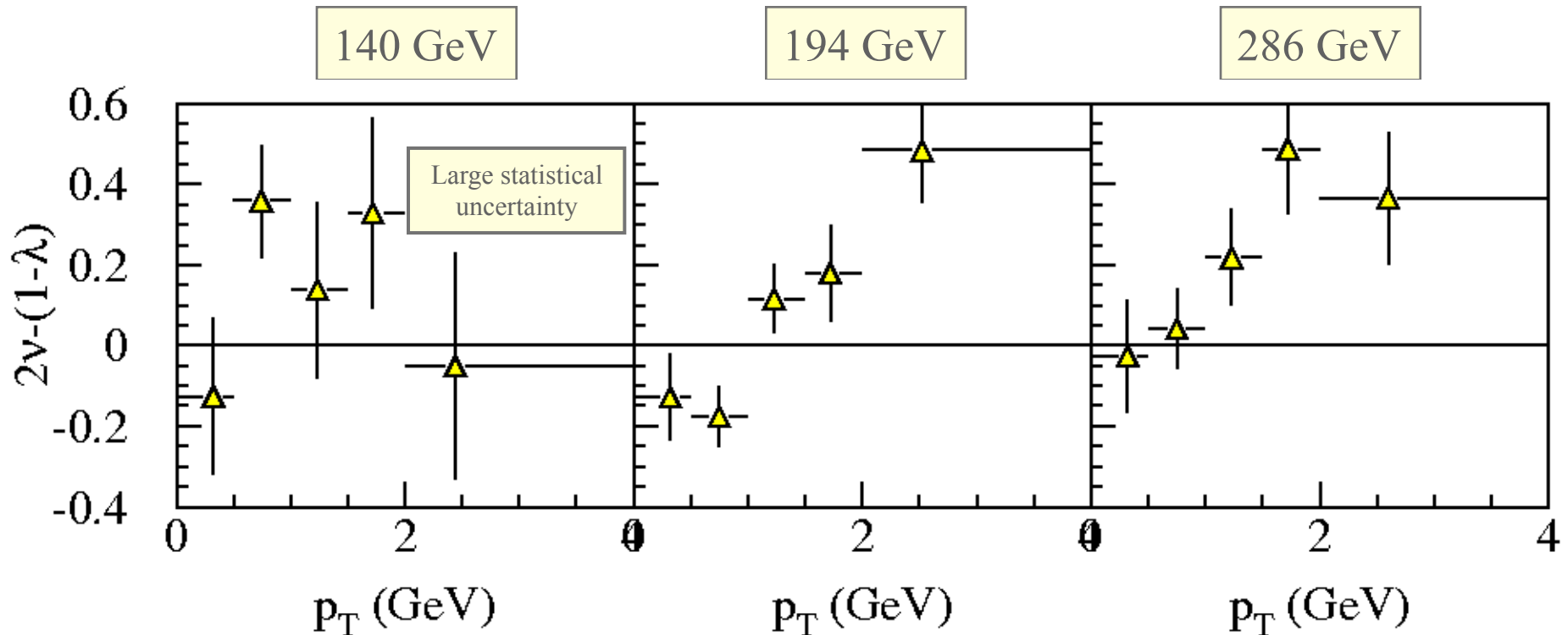


Lam-Tung relation

- Derived in analogy to Colin-Gross relation of DIS

$$1 - \lambda = 2\nu$$

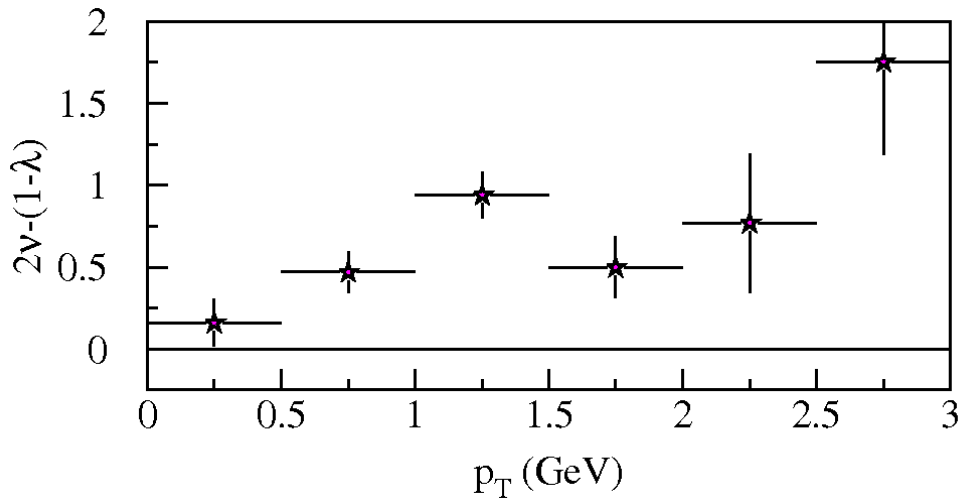
NA10 Lam-Tung Relation vs. p_T



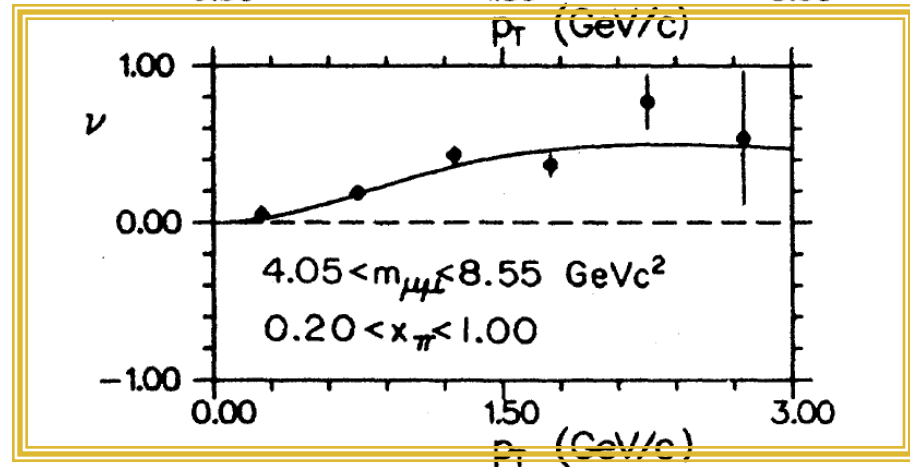
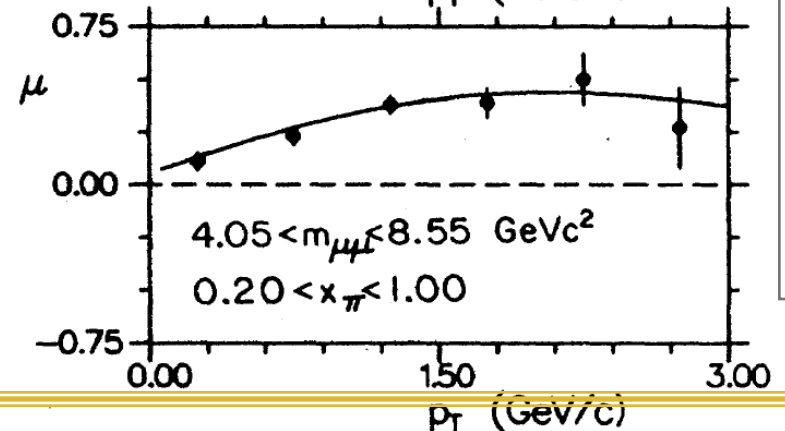
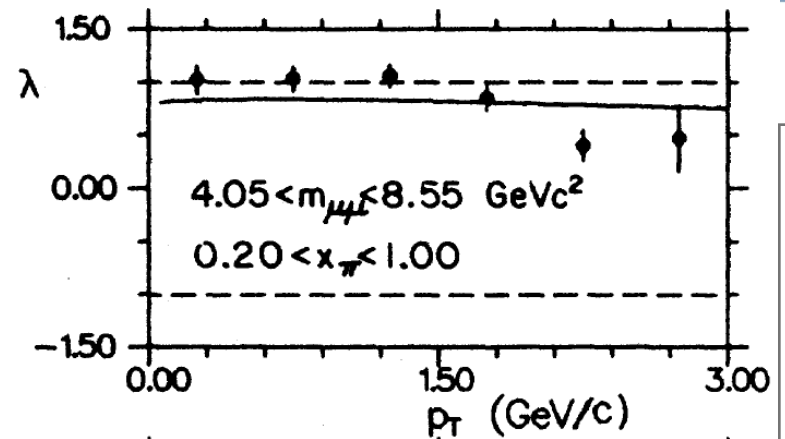
- Violation of Lam-Tung relation as p_T increases in higher momentum data. Statistics poor in 140 GeV data.
- Note: Correlation between λ and v uncertainties not considered.
- Since most data is at low p_T , *on average* the Lam-Tung relationship holds
- For all three energies $\lambda \sim 1$ and $v > 0$

Pionic Data Fermilab E615

- Significant non-zero ν coefficient
- Shows other kinematic dependencies



- Clear violation of Lam-Tung Relation vs. p_T .
- Violation larger than NA10



Conway et al. PRD 39 92 (1989)

Why is Lam-Tung relation important? What does it have to do with Transverse Momentum?

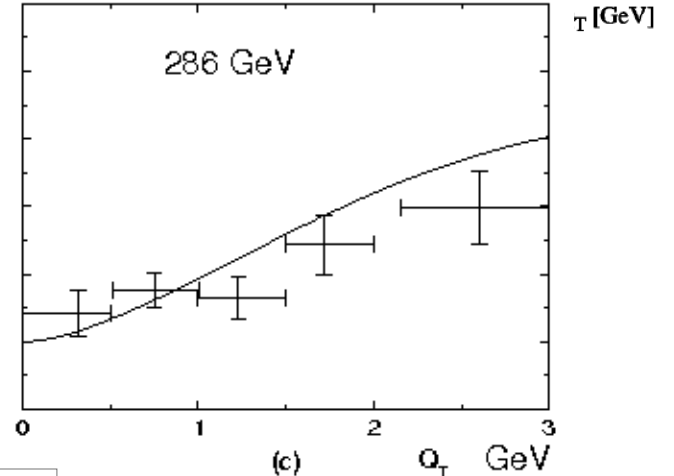
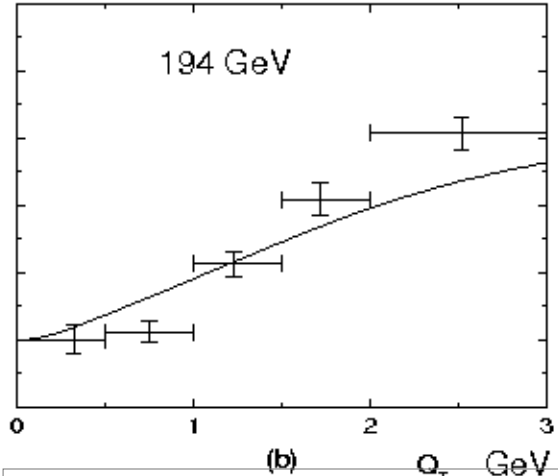
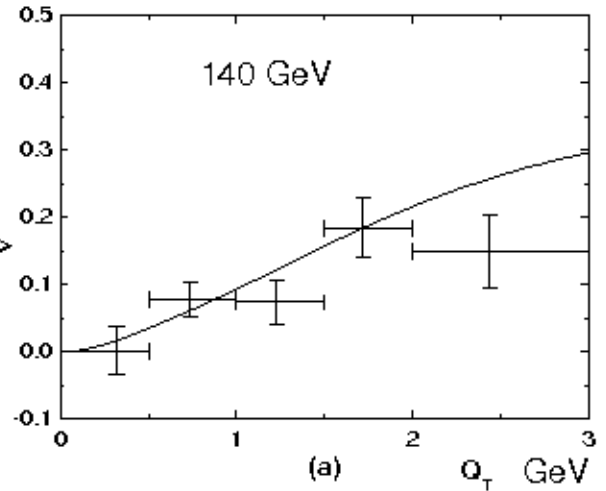
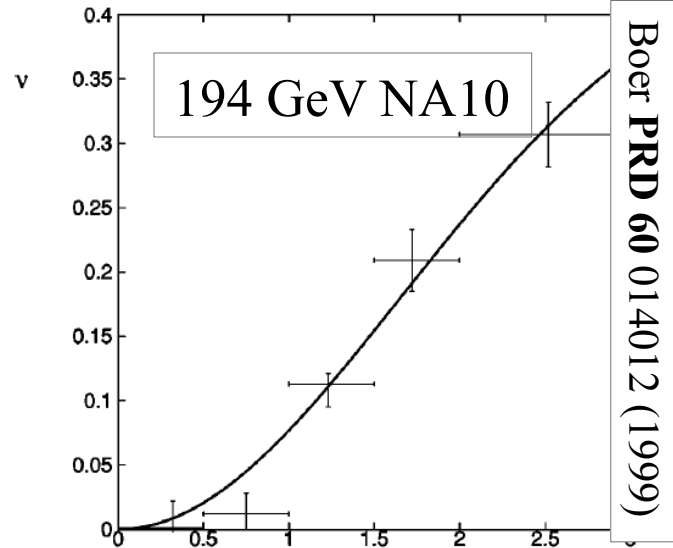
- Lam-Tug Relation is theoretically robust?
 - Unaffected by $O(\alpha_s)$ (NLO) corrections
 - NNLO [$O(\alpha_s^2)$] corrections also small Mirkes and Ohnemus, PRD 51 4891 (1995)
 - Soft Gluon Resummation–Berger, Qiu and Rodrigues-Pedraza showed that **the Lam-Tung relation is preserved** under resummation. arXiv:0707.3150, and PRD 76 074006 (2007)
- k_T dependent transverse momentum distribution (Boer Mulders h_1^\perp)
- Factorization breaking QCD Vacuum
- Alternatives:
 - Nuclear effects
 - Higher-Twist effects from quark-antiquark binding in pion
 - Neither seen in data

Boer-Mulders Structure Function

- Relates transverse spin and transverse momentum (k_T) in an unpolarized nucleon.
- Presence in both quark and antiquark in annihilation could form correlation contributing to $\cos(2\phi)$ distribution

$$\nu \propto h_{1q}^\perp(x_1) h_{1\bar{q}}^\perp(x_2)$$

- Reasonable fits to pionic data



Lu, Ma, PLB 615, 200 (2005)

Lam-Tung Relation—Alternative View: QCD effects

- Factorization breaking Brandenburg, Nachtmann and Mirkes, ZPC **60**, 679 (1993).
 - QCD Vacuum *may* correlate the spins and momenta of incoming partons
 - Effect could be instanton-induced Boer, Brandenburg, Nachtmann, Utermann, EPJC 40 55 (2005), Brandenburg, Ringwald, Utermann NPB 754, 107 (2006).

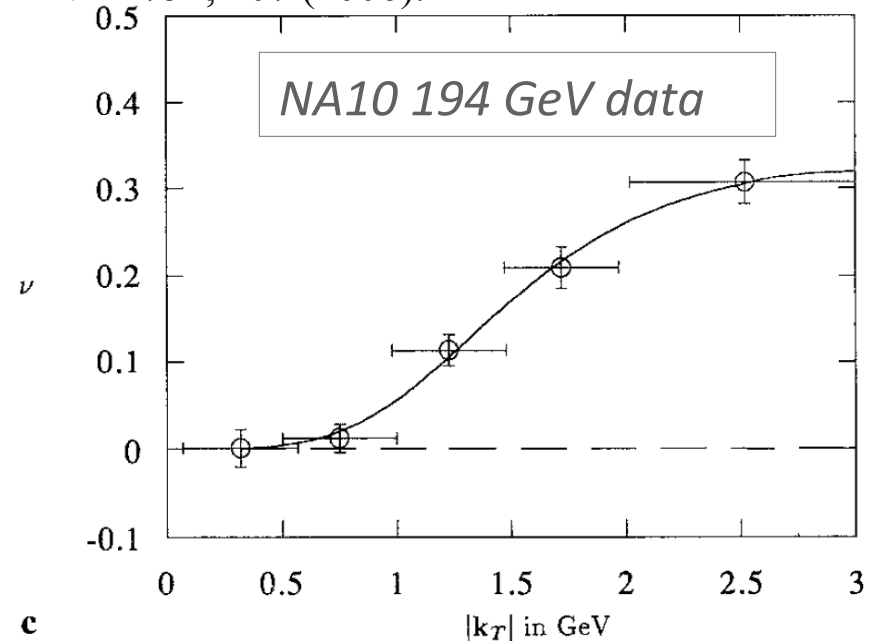
$$\nu \approx 2\mathcal{K} = 2\mathcal{K}_0 \frac{p_T^4}{p_T^4 + m_T^4}$$

$$\lambda \approx 1 \quad \mu \approx 0$$

- Fit NA10:

$$\kappa = 0.17$$

$$m_T = 1.5$$



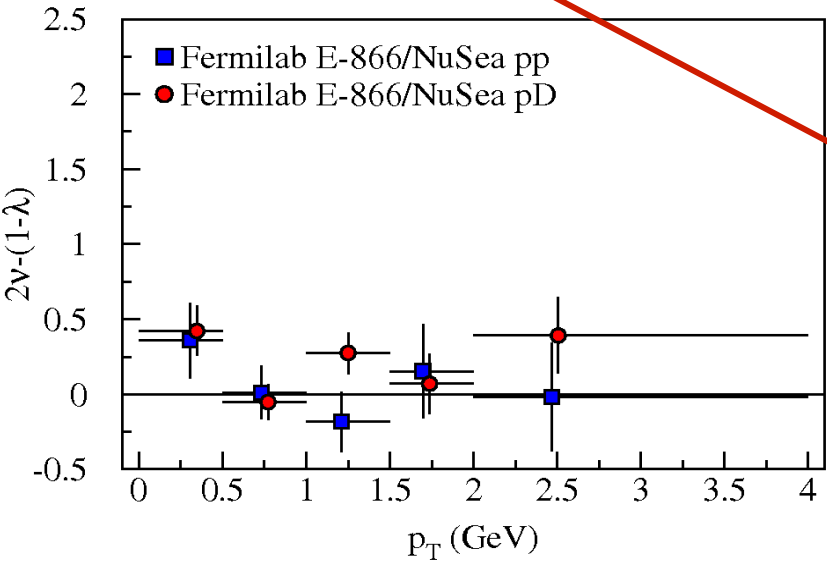
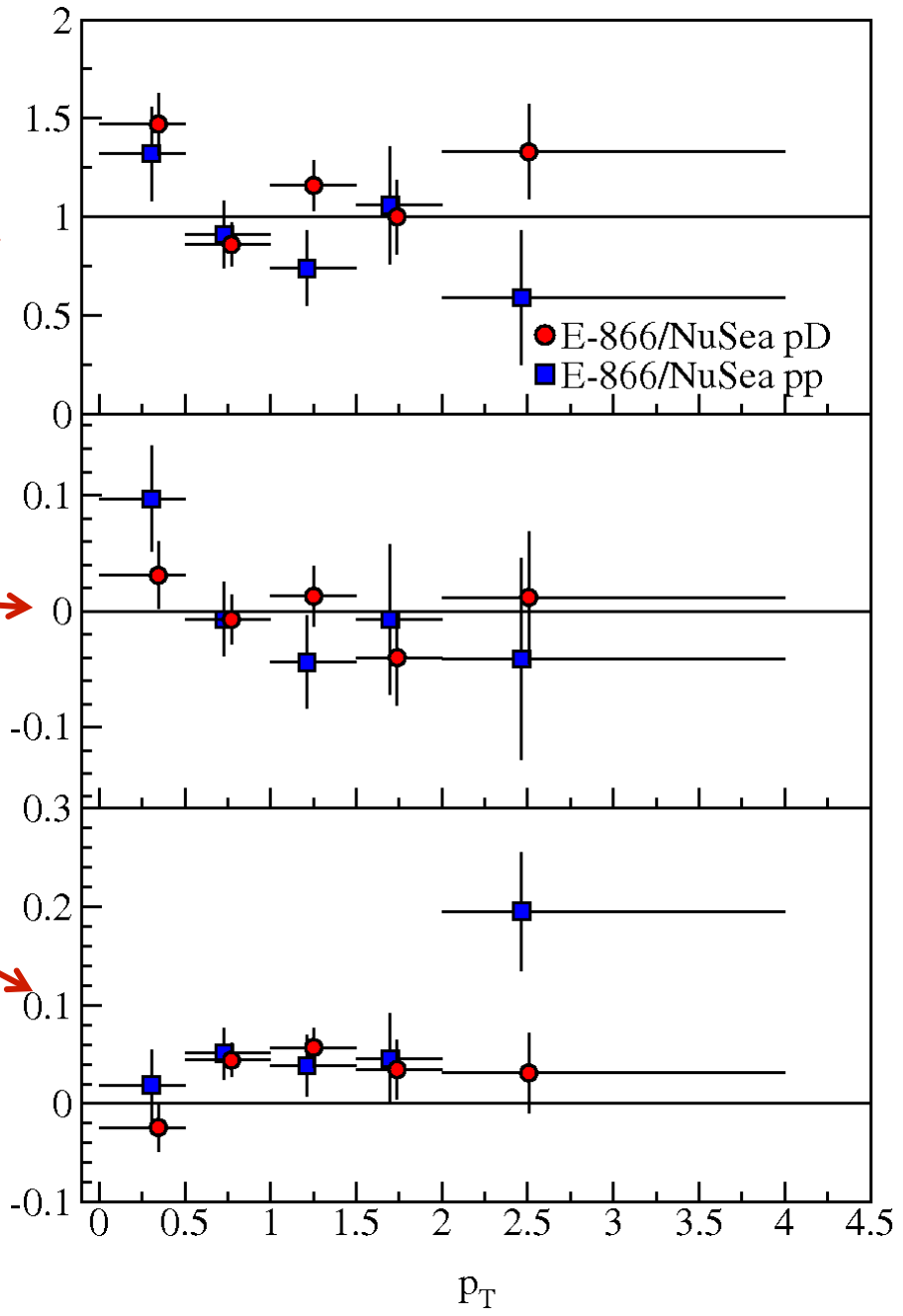
Should be flavor blind and seen in both sea and valence distributions

Protonic Drell-Yan

E-866/NuSea: Angular distrib.

$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

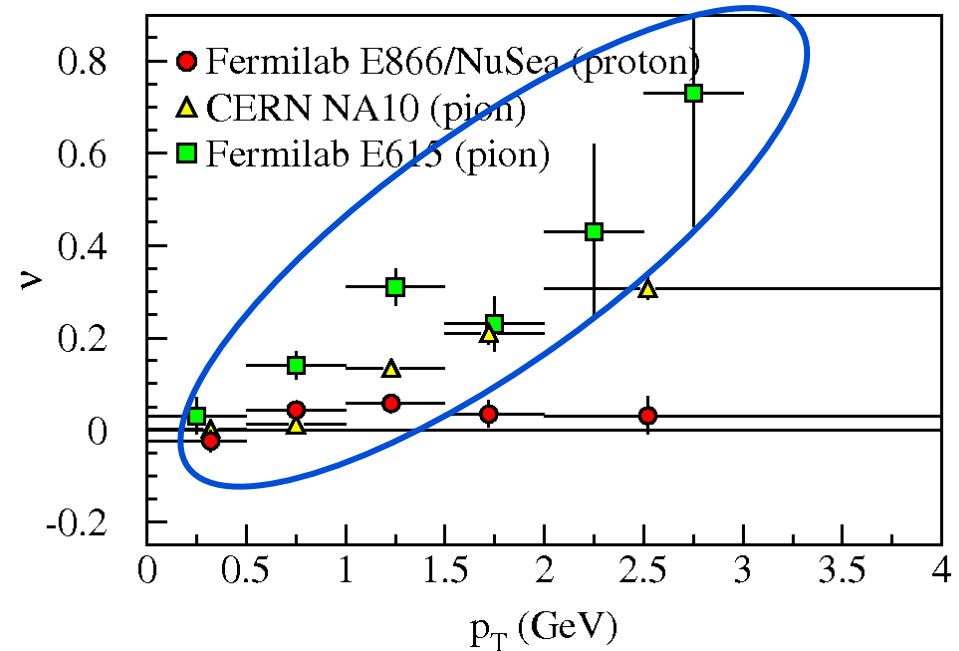
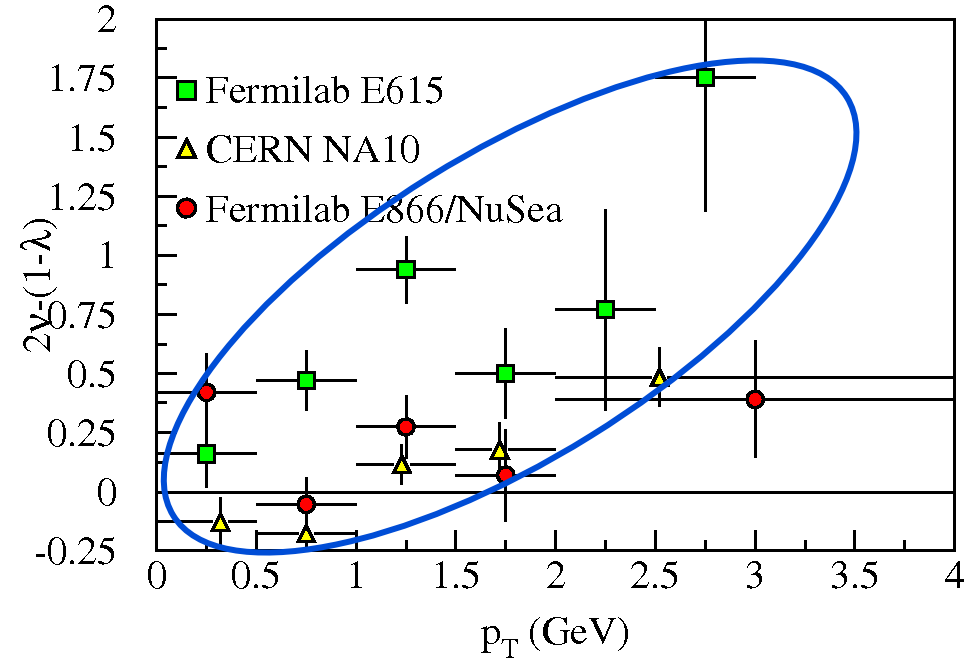
- λ consistent with 1
- μ consistent with 0
- ν consistent with 0 (or slightly positive)



Lam-Tung Relation

■ π^- Drell-Yan

- Violates L-T relation
- Large v ($\cos 2\phi$) dependence
- Strong with p_T



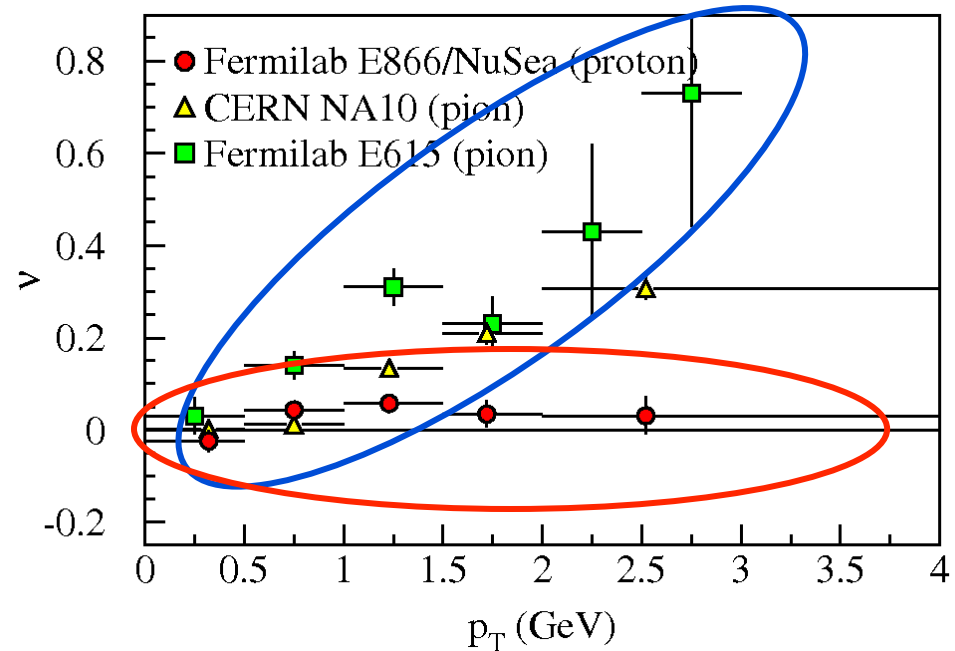
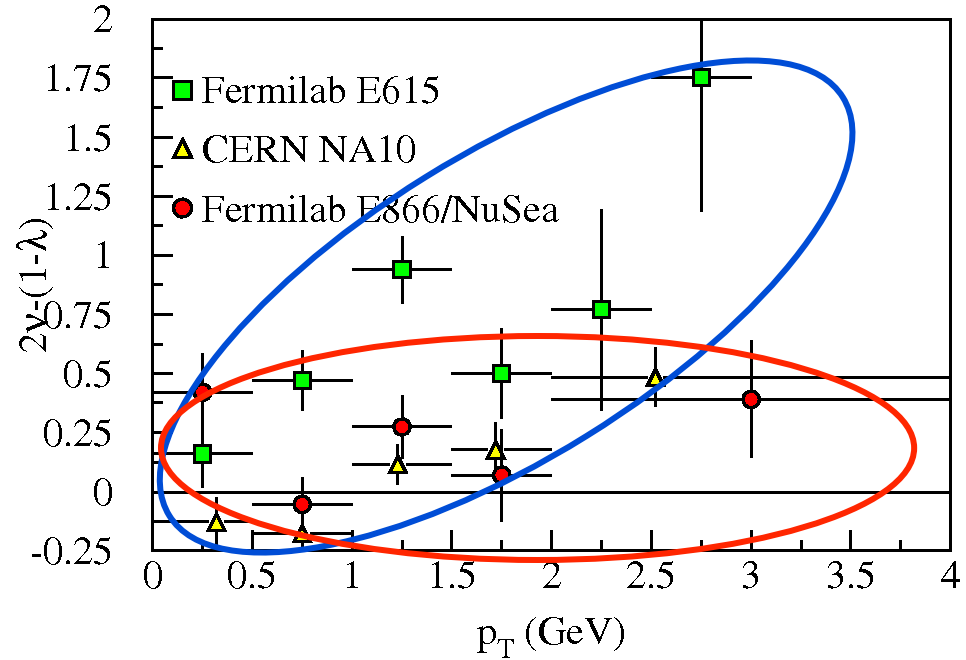
Lam-Tung Relation

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■ Proton Drell-Yan

- Consistent with L-T relation
- No ν ($\cos 2\phi$) dependence
- No p_T dependence



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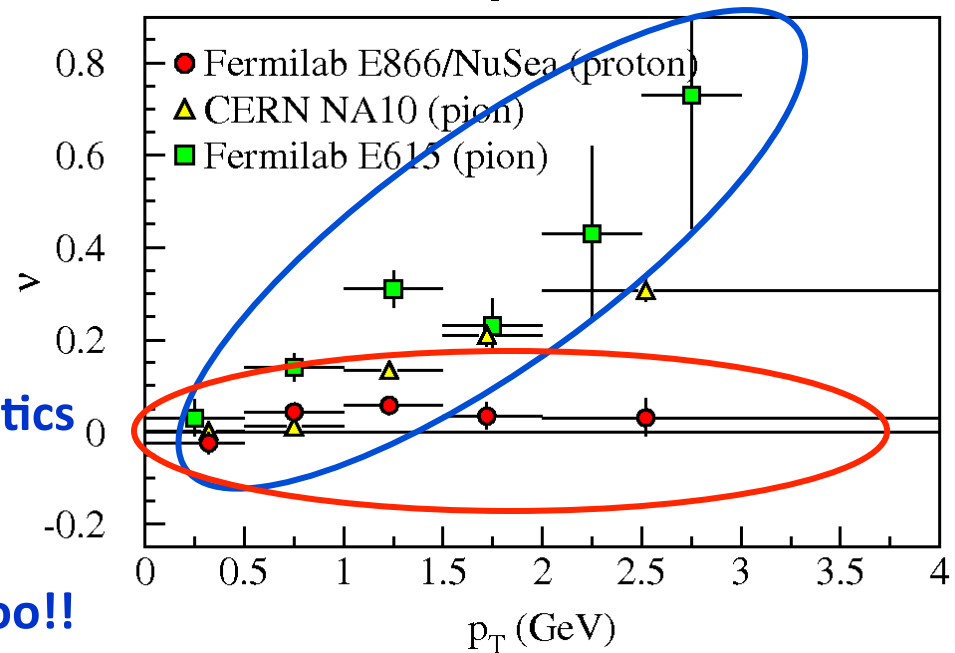
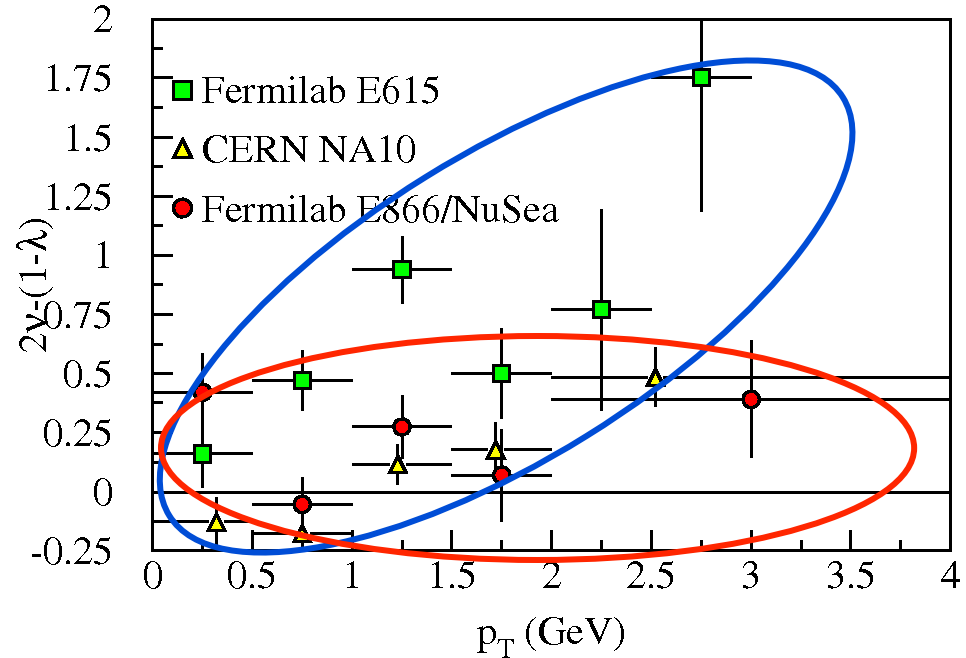
■ With Boer-Mulders function h_1^{\perp} :

- $v \pi$ -W $\rightarrow \mu^+ \mu^- X$
valence $h_1^{\perp}(\pi)$ * *valence* $h_1^{\perp}(p)$
- $v pd \rightarrow \mu^+ \mu^- X$
valence $h_1^{\perp}(p)$ * *sea* $h_1^{\perp}(p)$

■ Consistent story but need better statistics

- π -p $\rightarrow \mu^+ \mu^- X$ —CERN Compass
- pd $\rightarrow \mu^+ \mu^- X$ —Fermilab E906/SeaQuest

– Remember QCD effects are important too!!



Spin Asymmetries in Drell-Yan reactions



Single Spin Leading Order Drell-Yan Cross Section

$$\begin{aligned} \frac{d\sigma^{\text{LO}}}{d^4q d\Omega} = & \frac{\alpha^2}{Fq^2} \hat{\sigma}_{\text{U}}^{\text{LO}} \left[1 + D_{\sin^2 \theta}^{\text{LO}} A_{\text{U}}^{\cos 2\phi} \cos 2\phi \right. \\ & + S_{\text{L}} D_{\sin^2 \theta}^{\text{LO}} A_{\text{L}}^{\sin 2\phi} \sin 2\phi \\ & + \left| \vec{S}_{\text{T}} \right| A_{\text{T}}^{\sin \phi_{\text{S}}} \sin \phi_{\text{S}} \\ & + \left| \vec{S}_{\text{T}} \right| D_{\sin^2 \theta}^{\text{LO}} A_{\text{T}}^{\sin(2\phi + \phi_{\text{S}})} \sin(2\phi + \phi_{\text{S}}) \\ & \left. + \left| \vec{S}_{\text{T}} \right| D_{\sin^2 \theta}^{\text{LO}} A_{\text{T}}^{\sin(2\phi - \phi_{\text{S}})} \sin(2\phi - \phi_{\text{S}}) \right] \end{aligned}$$

$A_{\text{U}}^{\cos 2\phi}$

Boer-Mulders of beam hadron

$A_{\text{T}}^{\sin \phi_{\text{S}}}$

Sivers' for target nucleon

$A_{\text{T}}^{\sin(2\phi + \phi_{\text{S}})}$

Boer-Mulders of beam hadron and $h_{1\perp}$ and pretzelosity of target nucleon

$A_{\text{T}}^{\sin(2\phi - \phi_{\text{S}})}$

Boer-Mulders of beam hadron and h_1 and transversity of target nucleon

Formula: from Aram Kotzinian
Slide: content Oleg Denisov's
Transversity 2011 talk

Transverse Momentum Distributions: The Sivers' function

1. What is the Sivers' function?

$$f_{1T}^\perp = \begin{array}{c} \uparrow \\ \text{○} \\ \text{---} \\ \text{○} \\ \downarrow \end{array} \quad \text{---} \quad \text{S}_T \cdot (\hat{\mathbf{p}} \times \mathbf{k}_T) \leftrightarrow f_{1T}^\perp$$

k_T - dependent,
Naïve T-odd

2. When will we have Drell-Yan data?



Transverse Momentum Distributions: The Sivers' function

1. What is the Sivers' function?

$$f_{1T}^\perp = \begin{array}{c} \uparrow \\ \odot \end{array} - \begin{array}{c} \odot \\ \downarrow \end{array} \quad S_T \cdot (\hat{\mathbf{p}} \times \mathbf{k}_T) \leftrightarrow f_{1T}^\perp$$

k_T - dependent,
Naïve T-odd



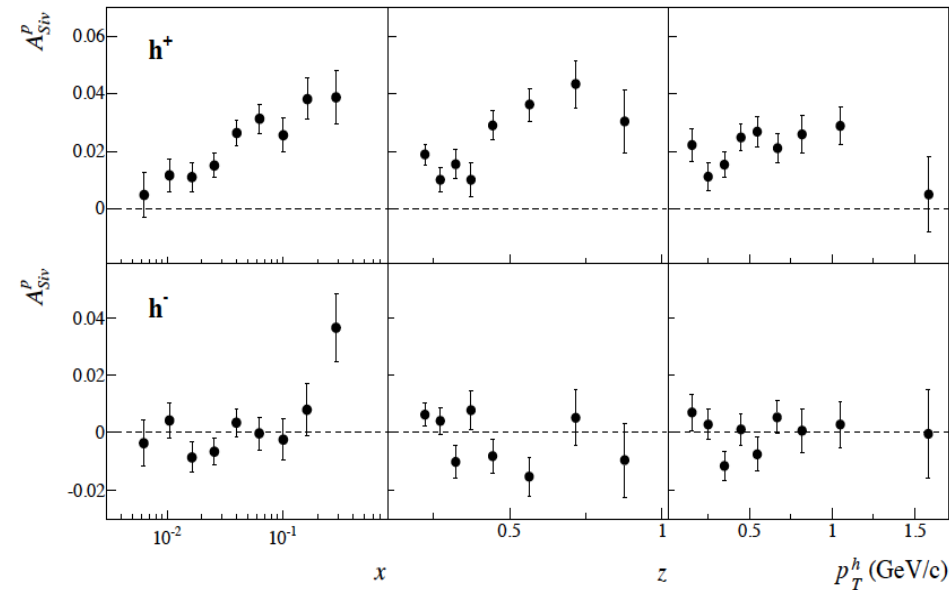
2. When will we have Drell-Yan

$$f_1^\perp(x, k_T) \Big|_{\text{DIS}} \stackrel{?}{=} - f_1^\perp(x, k_T) \Big|_{\text{D-Y}}$$

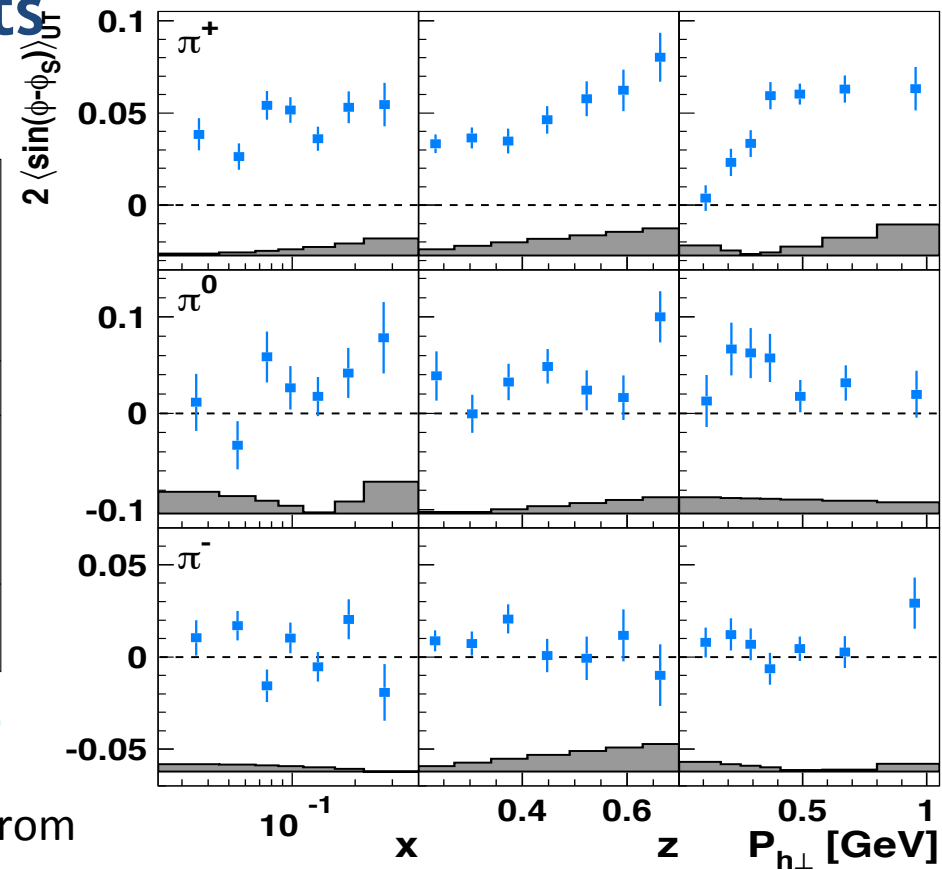


SIDIS Sivers' measurements

COMPASS



HERMES



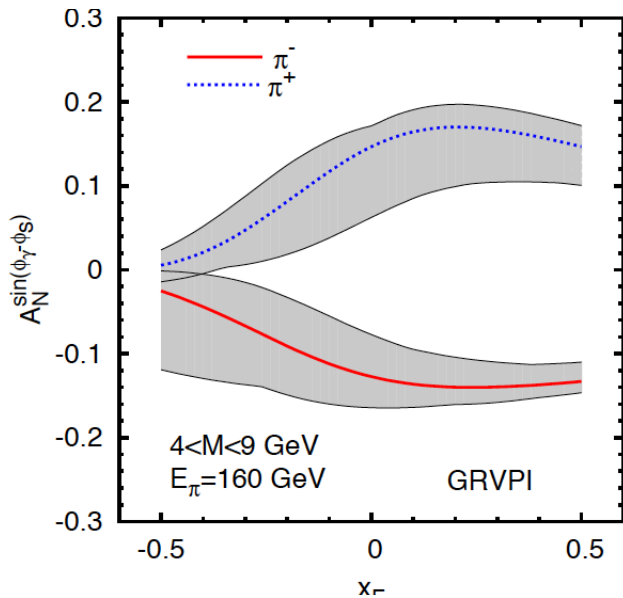
- Global fit to $\sin(\phi_h - \phi_S)$ asymmetry in SIDIS from HERMES, and COMPASS
- Comparable measurements needed for single spin asymmetries in Drell-Yan process
- Caution:
 - Must cover same kinematics with Drell-Yan and SIDIS (or have robust QCD evolution)
 - Right now, nothing excludes a node

Future Polarized Drell-Yan Measurements

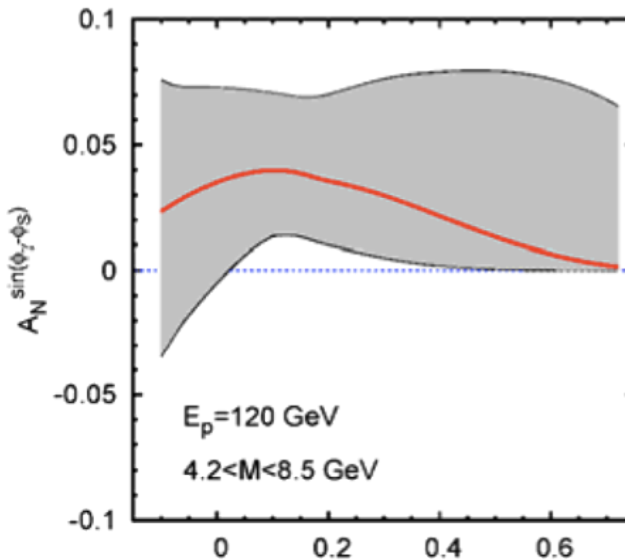
experiment	particles	energy	x_b or x_t	Luminosity	timeline
COMPASS (CERN)	$\pi^\pm + p^\uparrow$	160 GeV $\sqrt{s} = 17.4$ GeV	$x_t = 0.2 - 0.3$	$2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$	2014
PAX (GSI)	$p^\uparrow + p_{\text{bar}}$	collider $\sqrt{s} = 14$ GeV	$x_b = 0.1 - 0.9$	$2 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$	>2017
PANDA (GSI)	$p_{\text{bar}} + p^\uparrow$	15 GeV $\sqrt{s} = 5.5$ GeV	$x_t = 0.2 - 0.4$	$2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$	>2016
NICA (JINR)	$p^\uparrow + p$	collider $\sqrt{s} = 20$ GeV	$x_b = 0.1 - 0.8$	$1 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$	>2014
PHENIX (RHIC)	$p^\uparrow + p$	collider $\sqrt{s} = 500$ GeV	$x_b = 0.05 - 0.1$	$2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$	>2018
SeaQuest (unpol.) (FNAL)	$p + p$	120 GeV $\sqrt{s} = 15$ GeV	$x_b = 0.35 - 0.85$ $x_t = 0.1 - 0.45$	$3.4 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$	2012
pol. SeaQuest [§] (FNAL)	$p^\uparrow + p$	120 GeV $\sqrt{s} = 15$ GeV	$x_b = 0.35 - 0.85$	$1 \times 10^{36} \text{ cm}^{-2} \text{ s}^{-1}$	>2015

Actual luminosity at Fermilab limited to $2 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ by proton needs of other experiments

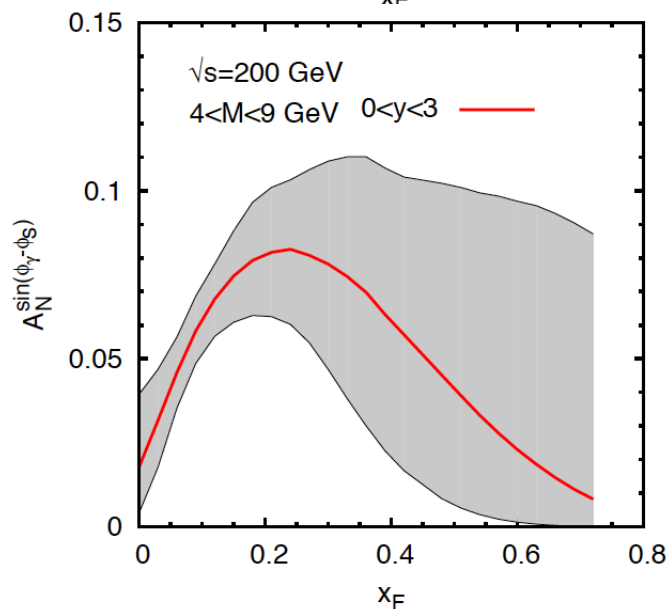
$A_N^{\sin(\phi_\gamma - \phi_S)}$ prediction Anselmino *et al*



COMPASS



Fermilab

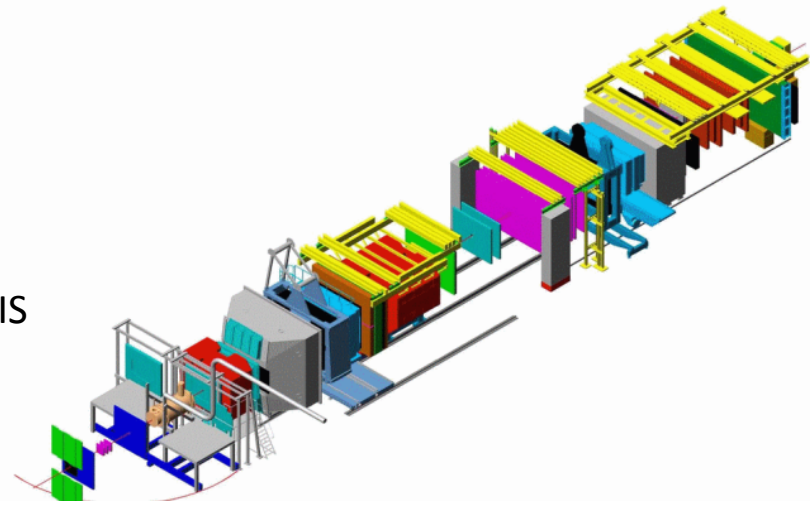


Anselmino, *et al* Phys. Rev. D79 (2009) 054010
and private communication (for Fermilab)

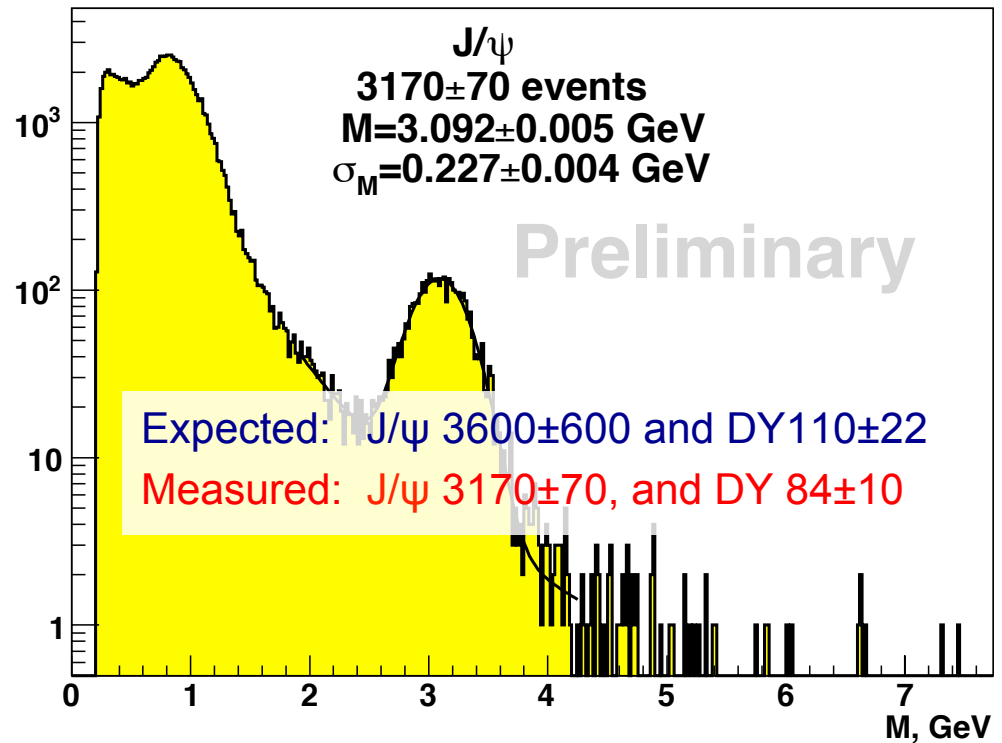
- Uncertainties in predictions dominated by uncertainties in SIDIS data
- Also predictions for GSI PAX, GSI PANDA, J-PARC, JINR NCIA, IHEP SPASCHARM

COMPASS@CERN

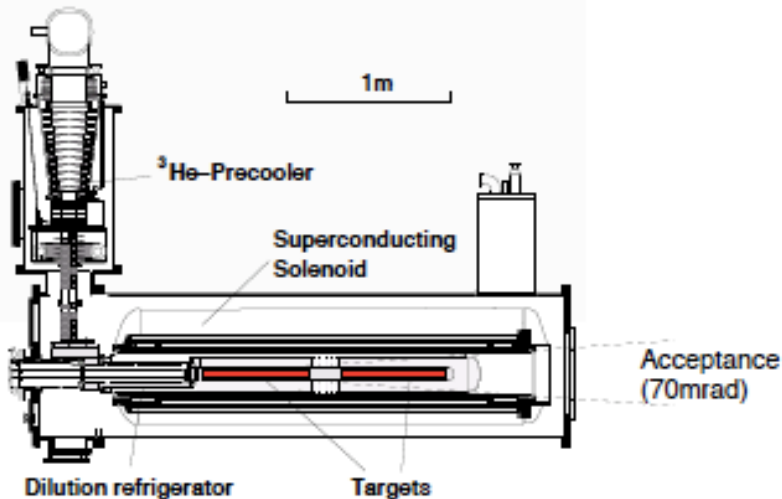
- See talk by S. Takekawa (next)
- Very well understood apparatus
 - A_N angular measurements already completed for SIDIS
 - MC and test data match
- Polarized target
- Data in 2014



COMPASS DY beam test 2009



Polarized target: ^6LiD or NH_3





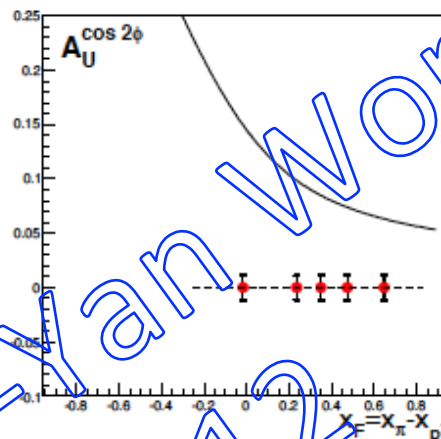
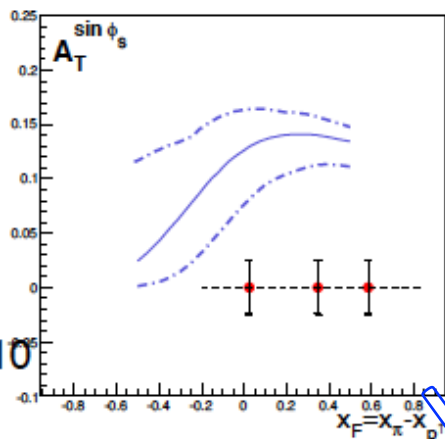
Comparing with theory predictions

...but: Q^2 evolution not properly accounted for in the predictions...

$$DY: 4 < M_{\mu\mu} < 9 \text{ GeV}/c^2$$

$f_1^\pi \otimes \text{Sivers}^P$

Anselmino et al,
PRD79(2009)054010

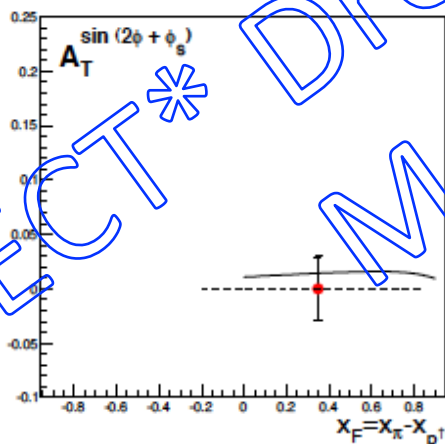


$BM^\pi \otimes BM^P$

B. Zhang et al,
PRD77(2008)054011

$BM^\pi \otimes \text{pretzel}^P$

Zhun Lu et al,
arXiv:1101.2702v2



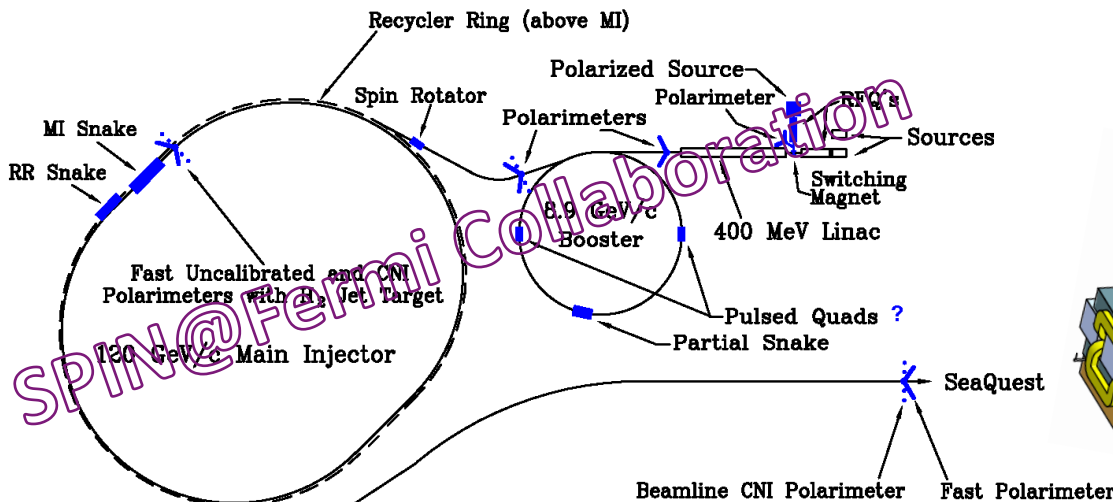
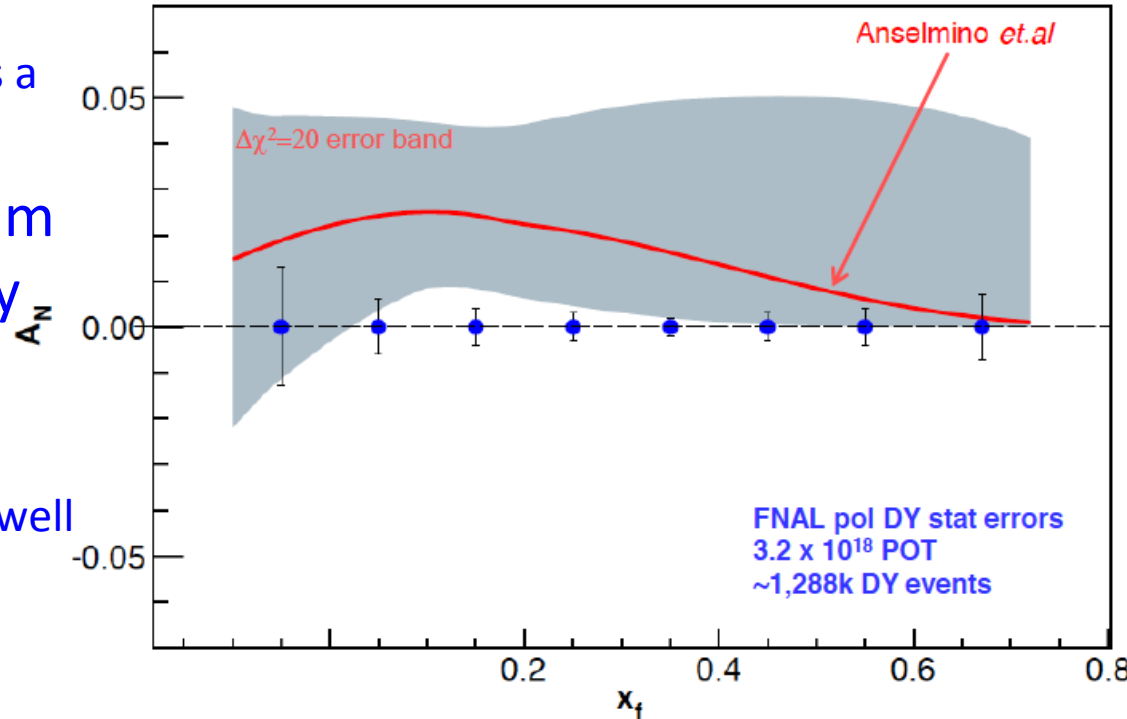
$BM^\pi \otimes \text{transv}^P$

A.N. Sissakian et al,
Phys.Part.Nucl.41:
64-100,2010

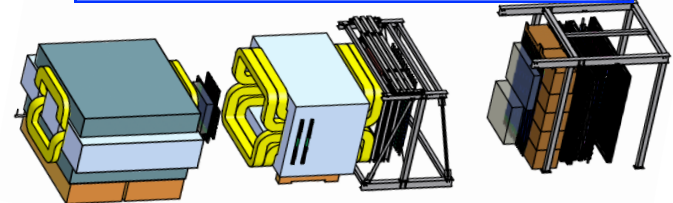
From ECT* Drell-Yan Workshop May 2012

Polarized Drell-Yan@Fermilab Main Injector

- Polarized beam
 - Major advantage—the beam is a blow torch—Luminosity
 - Major disadvantage—the beam polarization is presently *virtual*—only a proposal
- Spectrometer:
 - By 2014, spectrometer will be well understood, including angular acceptance



Exploring polarized target
 ■ sea quark Sivers'



Drell-Yan Data:

What data would I like to see? (My personal view)

- Siver's function relation to DIS tested

$$f_{1T}^\perp|_{\text{DIS}} = - f_{1T}^\perp|_{\text{DY}}$$

- CERN COMPASS, GSI PAX, J-PARC, RHIC

- Unpolarized Angular Distributions—Boer-Mulders Structure Function, QCD

- Fermilab E-906/SeaQuest, J-PARC, RHIC (pp, pA)
- CERN COMPASS (πp)

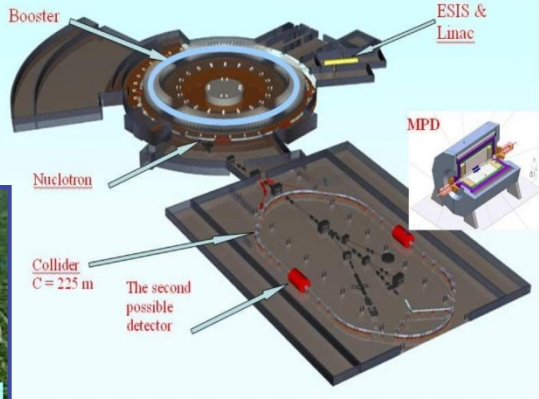
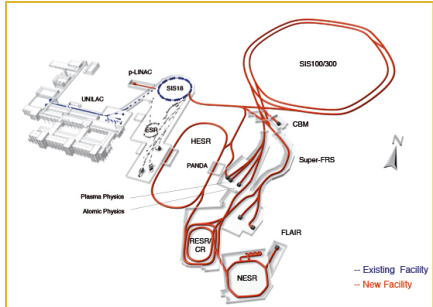
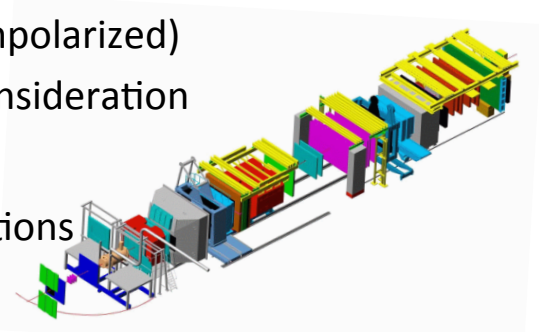
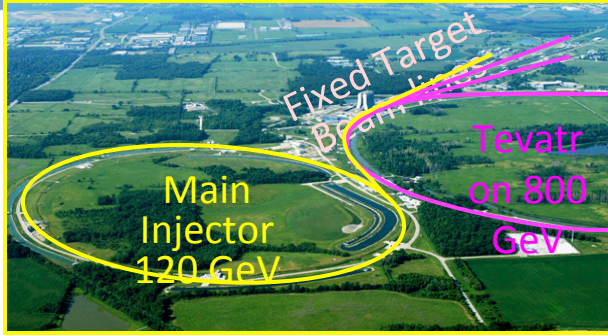
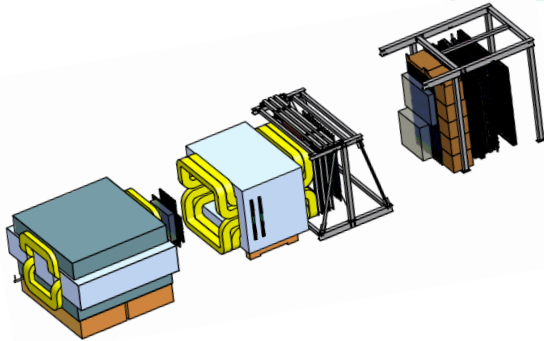
- Double polarized Drell-Yan— $h_1(x)$

$$\begin{aligned} A_{TT}^{\text{DY}} &= \frac{d\sigma^{\uparrow\uparrow} - d\sigma^{\uparrow\downarrow}}{d\sigma^{\uparrow\uparrow} + d\sigma^{\uparrow\downarrow}} \\ &= \alpha_{TT} \frac{\sum_i e_i^2 [h_{1i}(x_1)\bar{h}_{1i}(x_2) + \bar{h}_{1i}(x_1)h_{1i}(x_2)]}{\sum_i e_i^2 [f_i(x_1)\bar{f}_i(x_2) + \bar{f}_i(x_1)f_i(x_2)]} \\ &\approx \alpha_{TT} \frac{h_{1u}(x_1)h_{1u}(x_2)}{u(x_1) + u(x_2)} \end{aligned}$$

- GSI-FAIR PAX antiproton-proton
- Fermilab polarized beam and target

Future Drell-Yan Experiments

- Fermilab E-906/Drell-Yan
 - Better statistical precision (unpolarized)
 - Polarized extension under consideration
- COMPASS
 - Pion beam—valence distributions
- GSI FAIR—PAX experiment
 - Antiproton beam will sample valence distributions of targets
- JINR Dubna-NICA
- J-PARC
- RHIC



Epilogue:



This is a “simulated” (that is, equivalent to Monte Carlo?) picture of what the effect might look like done with GIMP software.

NA10 angular distributions vs. p_T

140 GeV/c

194 GeV/c

286 GeV/c

