



Karlsruher Institut für Technologie



# Higgs from Loops

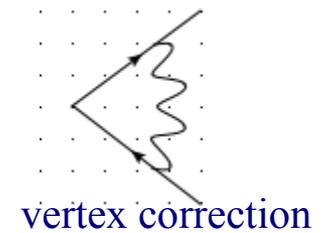
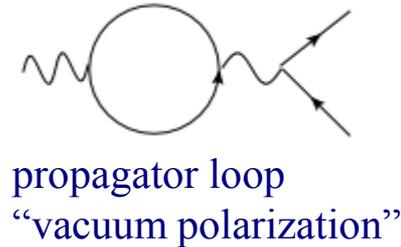
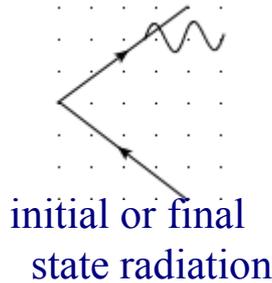
*G. Quast, A. Raspereza*

**Course “Higgs Physics”**

*Lecture 7, 14/06/2012*

## To take home from last lecture

### Higher orders modify tree-level diagrams:



- running coupling constants:  $\alpha \rightarrow \alpha(s)$  (holds for strong, em and weak interactions)
- modified decay widths:  $\Gamma \rightarrow \Gamma (1+c \alpha(s) + \dots)$  (for em and strong interactions)
- Z couplings  $\rightarrow$  “effective couplings”, or  
 $\sin^2(\Theta_w) \rightarrow \sin^2(\Theta_w)^{\text{eff}}$  and  $\rho=1 \rightarrow \rho^{\text{eff}}=(1+\Delta\rho)$
- relation between  $\sin^2(\Theta_w)$  from W and Z masses and  $\sin^2(\Theta_w)^{\text{eff}}$ :  $\Delta\rho$
- large vertex corrections in diagrams involving b and t -quarks:  $\Delta\rho_b$

### Corrections depend

- **quadratically** on top quark mass
- **logarithmically** on Higgs boson mass

Need precision calculations and precision measurements  
to become sensitive to Higgs in loops



Precision tests of ew theory in the 90ties  
(LEP, SLC, Tevatron)

# Measurements in $e^+ e^-$ at LEP – the four experiments

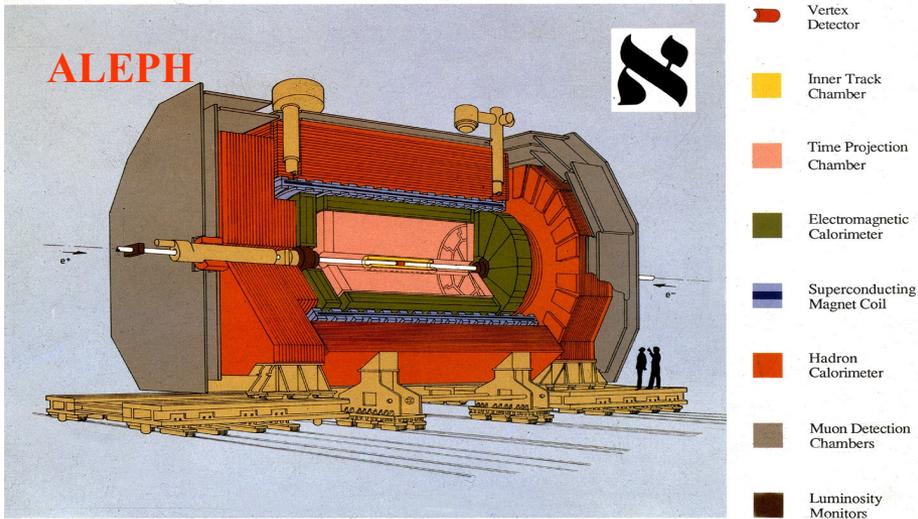
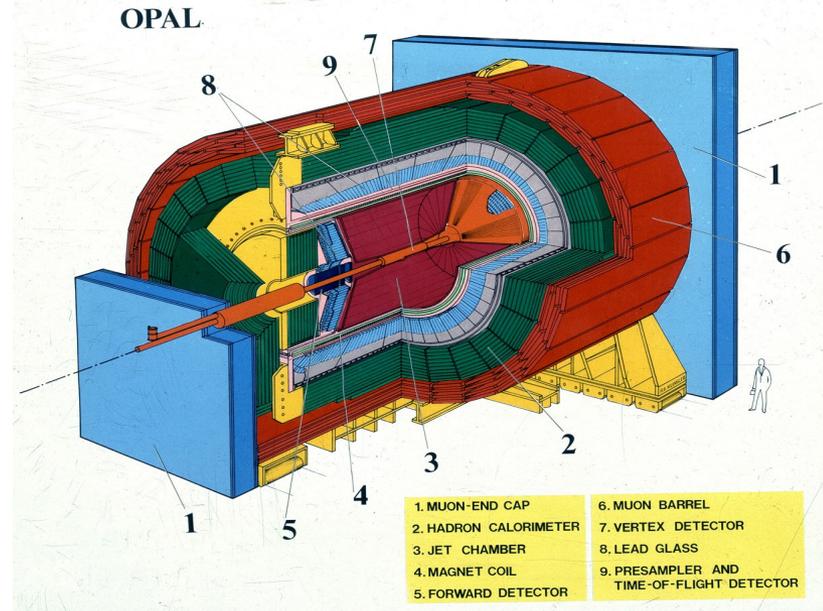
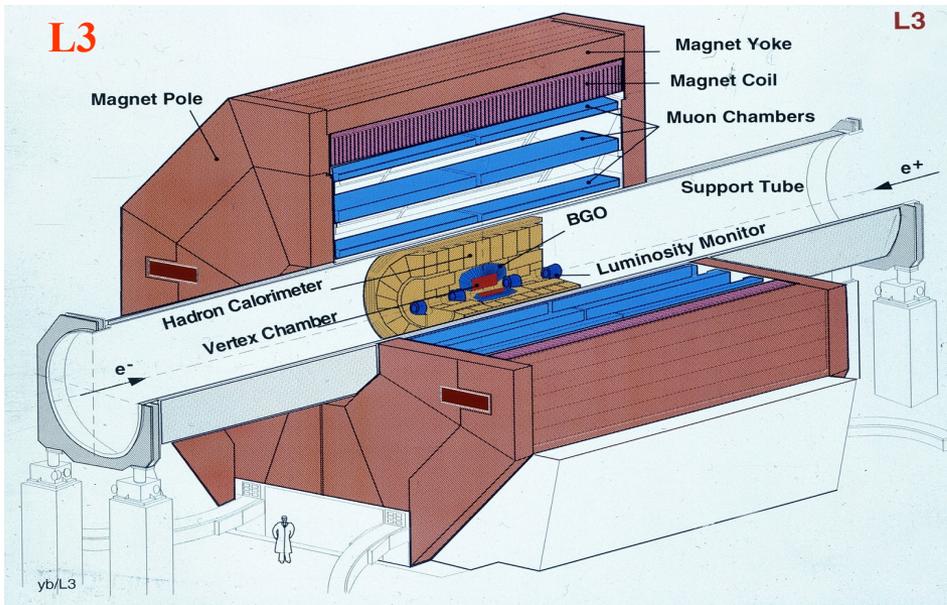
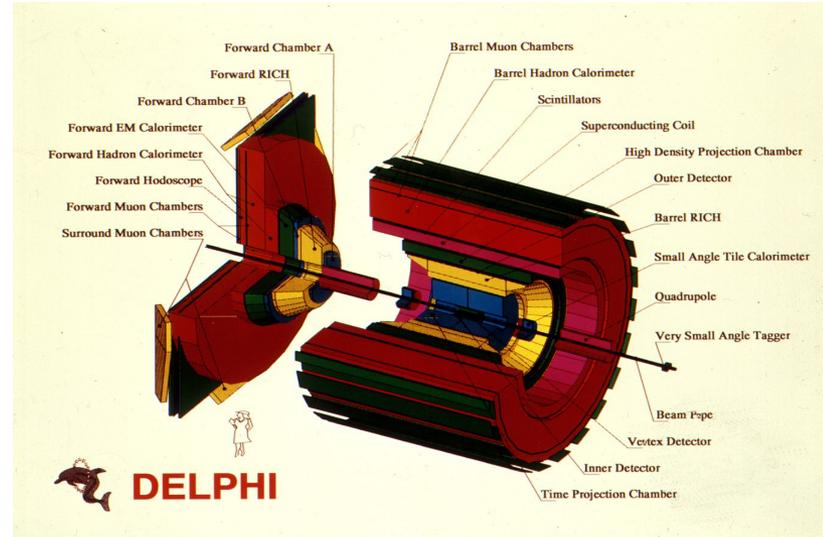
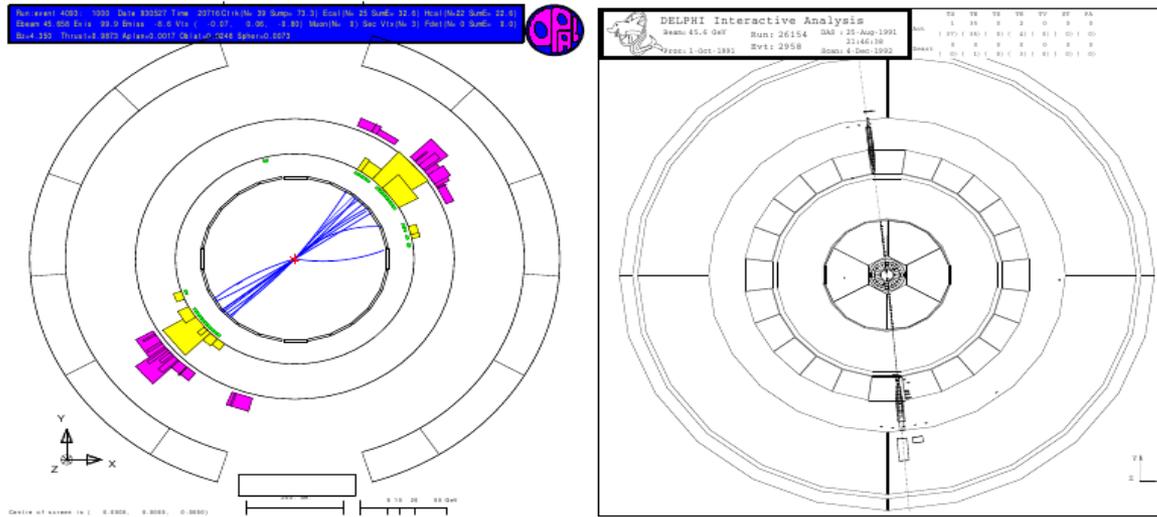


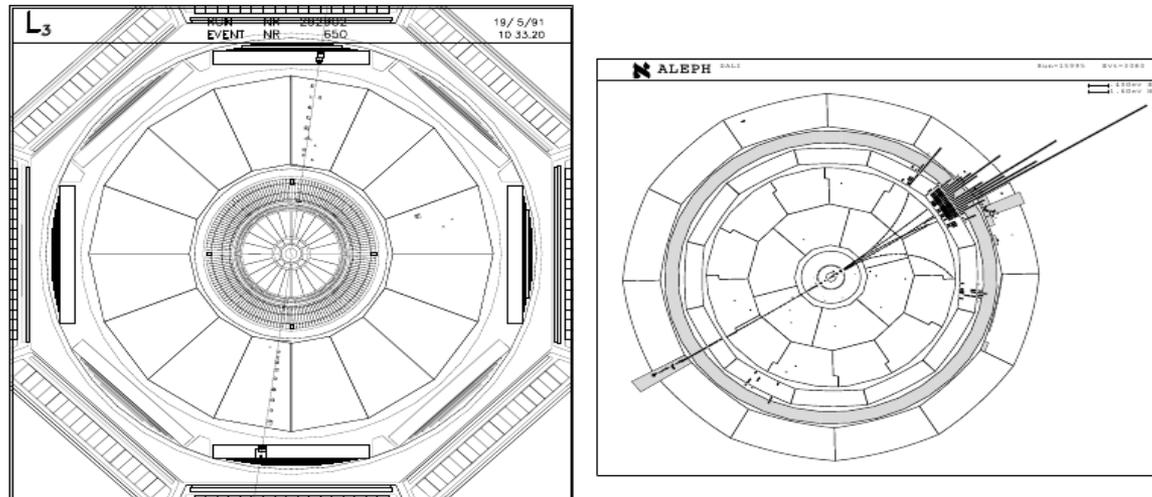
Fig. 1 - The ALEPH Detector



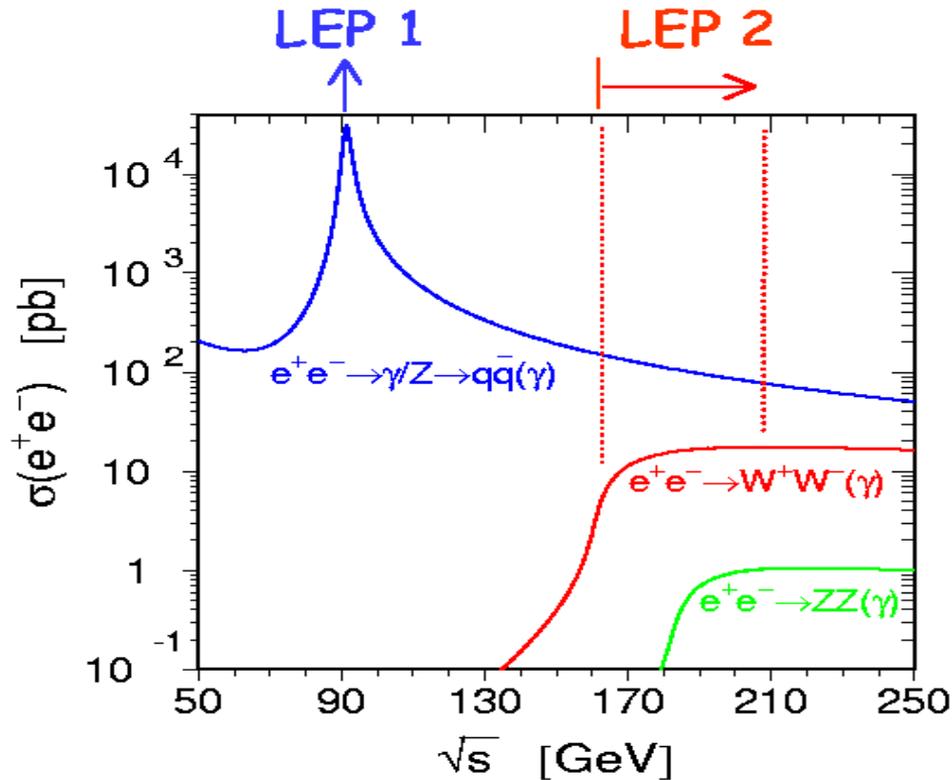
# Measurements in $e^+ e^-$ at LEP



## Event Displays from LEP experiments



# Measurements in $e^+ e^-$ at LEP



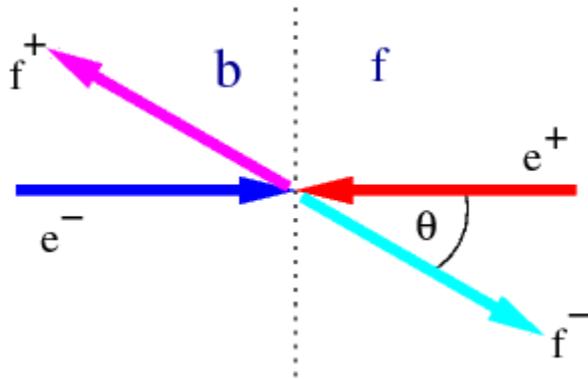
**LEP 1:** measurement of Z boson parameters (16 million Z's)

**LEP 2:** measurement of W and Z boson pair production,  
W boson Parameters (40'000 W pairs)

## total cross section

$$\sigma_{tot} \equiv \int_{-1}^1 \frac{d\sigma}{d\cos\theta} d\cos\theta$$

depends on  $(1+\cos^2 \Theta)$  terms only



## Forward-backward asymmetry

$$A_{FB} \equiv \frac{\int_0^1 \frac{d\sigma}{d\cos\theta} d\cos\theta - \int_{-1}^0 \frac{d\sigma}{d\cos\theta} d\cos\theta}{\sigma_{tot}}$$

depends on  $\cos^2 \Theta$  terms only

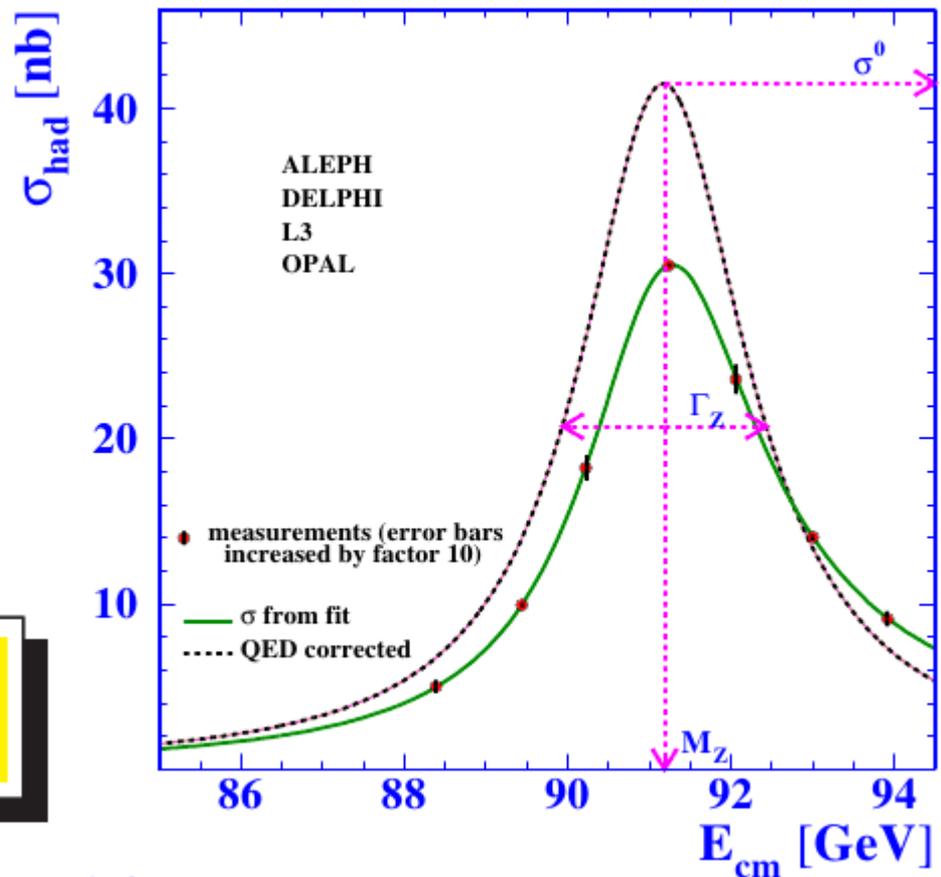
## Principle:

Extract combinations of Z-couplings to fermions from measurements of  $\sigma_{tot}$  and  $A_{FB}$  in different channels and at different energies

## Measurement of cross sections:

- Determine candidate events for signal process
- Subtract background passing selection criteria
- Determine selection efficiency: from Monte-Carlo simulation or using “data-driven” methods
- Normalize corrected number of signal events to luminosity

$$\sigma(E_i) = \frac{N_{\text{ff}}^{\text{cand}}(E_i) - N_{\text{ff}}^{\text{bkg}}(E_i)}{\varepsilon_{\text{ac}}(E_i)} \frac{1}{\int L(E_i)},$$



## Luminosity determination:

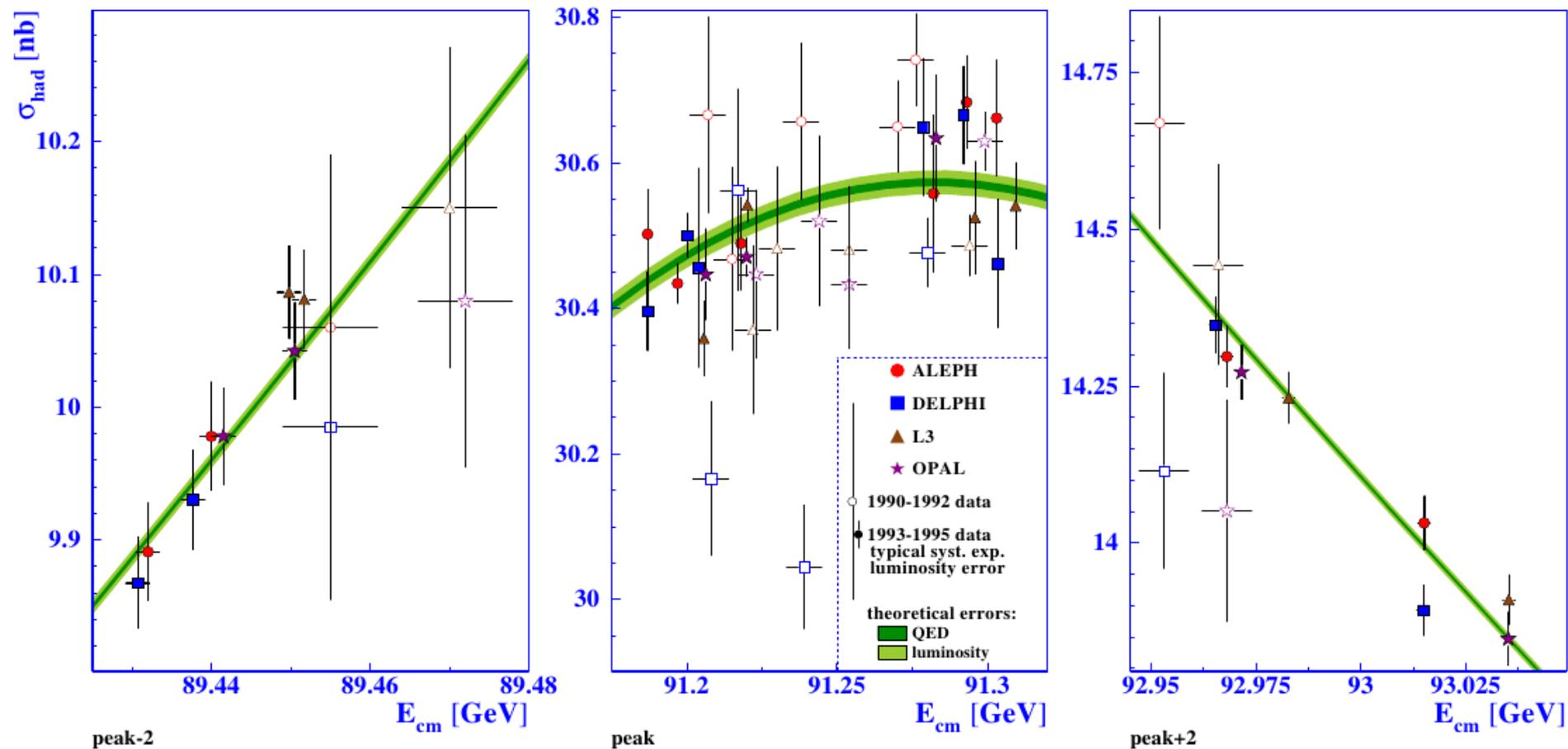
count number of events in reference reaction

with well-known cross section:  $\int L = n_{\text{ref}} / \sigma_{\text{ref}}$

reference reaction @ LEP:  $e^+e^- \rightarrow e^+e^-$  at small angles (“Bhabha scattering”)

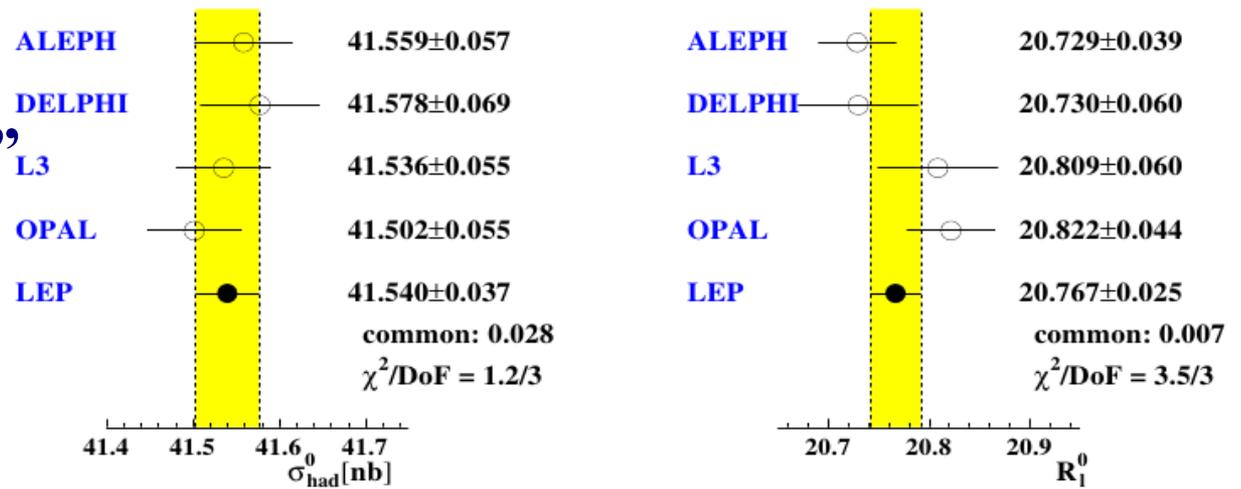
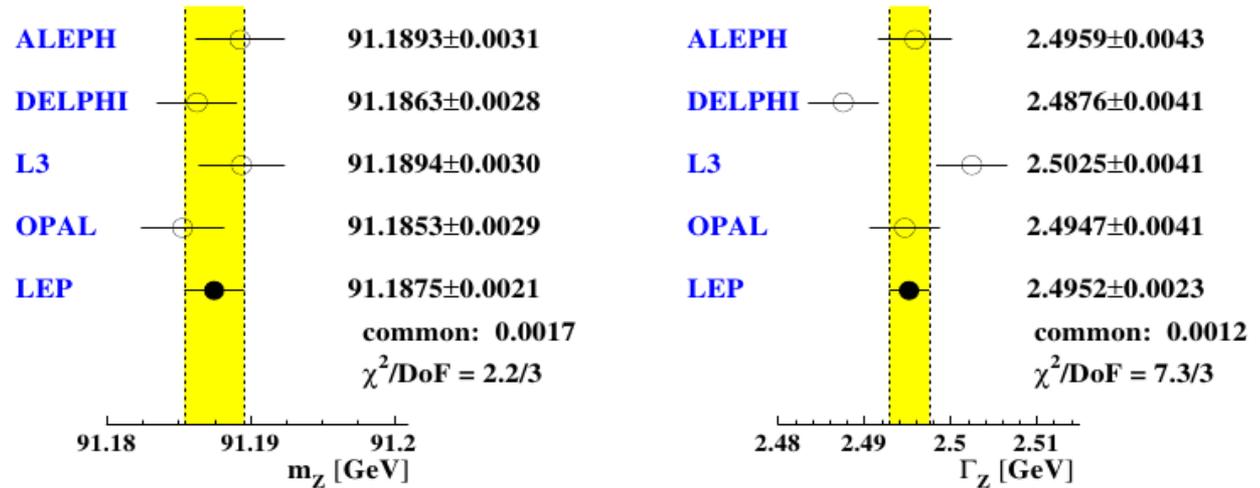
*photonic corrections are large,  
but well known theoretically*

# Precision measurements: Z line shape



**The challenge:** combining more than 800 individual measurements (different channels, CM-energies and data taking periods)

# Precision measurements: Z line shape for leptons and hadrons



Parametrization  
of cross section  
using  
“pseudo-observables”

- $m_Z$
- $\Gamma_Z$
- $\sigma_{\text{had}}^o = \frac{12\pi}{m_Z^2} \frac{\Gamma_{ee}\Gamma_{\text{had}}}{\Gamma_Z^2}$
- $R_e = \Gamma_{\text{had}} / \Gamma_{ee}$
- $R_\mu = \Gamma_{\text{had}} / \Gamma_{\mu\mu}$
- $R_\tau = \Gamma_{\text{had}} / \Gamma_{\tau\tau}$

partial decay width

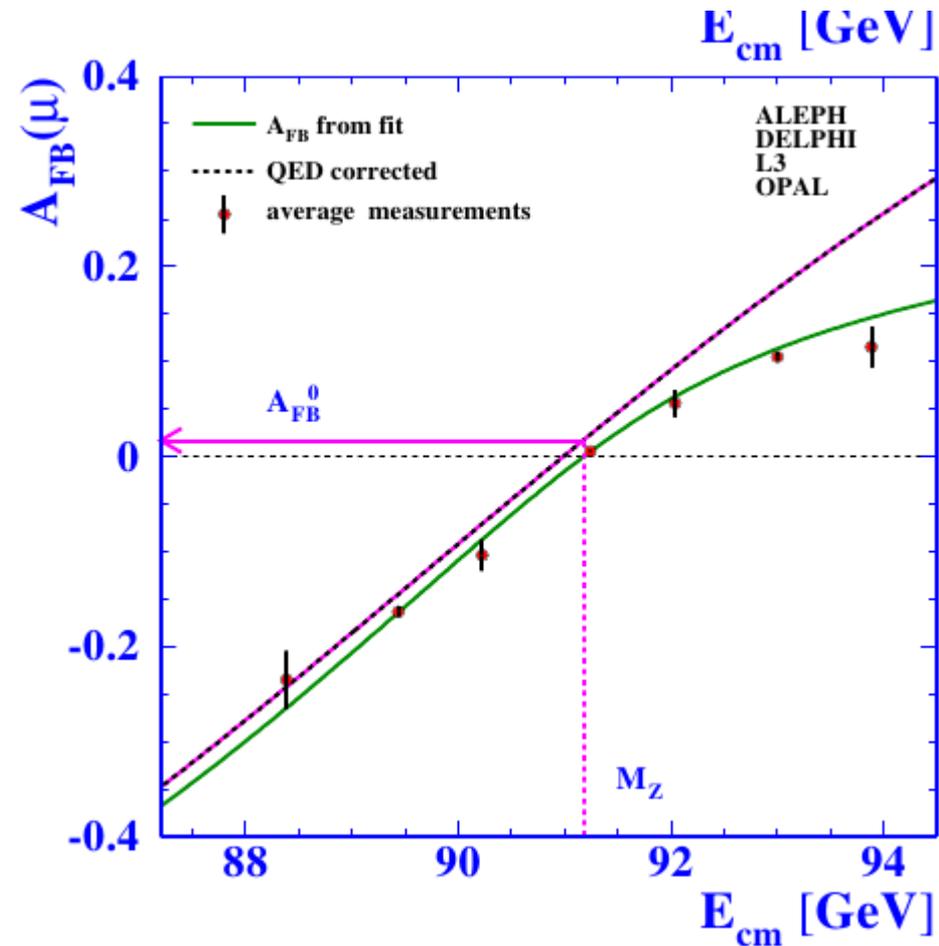
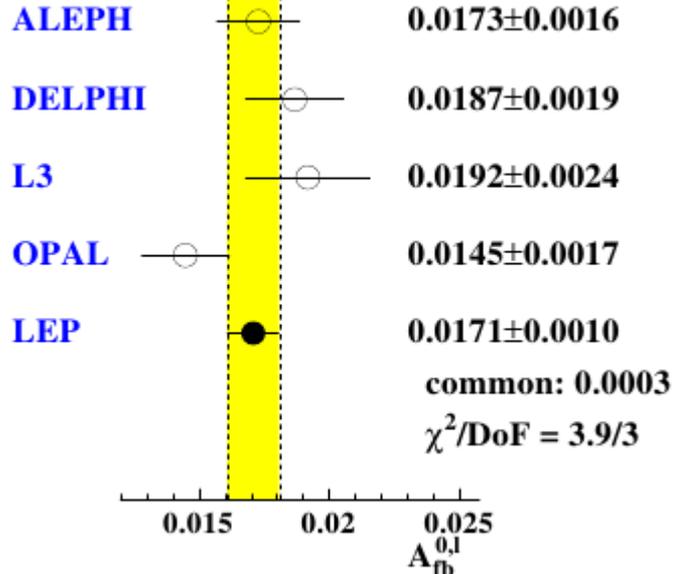
$$\Gamma_{ff} \propto (g_{Vf}^2 + g_{Af}^2) \text{ for } f=e, \mu, \tau$$

# Precision measurements: forward-backward asymmetries of lepton pairs

$$A_{FB} = \frac{N_{\text{forw}} - N_{\text{back}}}{N_{\text{tot}}} = \frac{\int_0^1 \frac{d\sigma}{d\cos\theta} d\cos\theta - \int_{-1}^0 \frac{d\sigma}{d\cos\theta} d\cos\theta}{\sigma_{\text{tot}}}$$

## pseudo-observables:

$$A_{FB}^{0,f} \equiv \frac{3}{4} \mathcal{A}_e \mathcal{A}_f \quad \text{with} \quad \mathcal{A}_f = \frac{2g_{Vf}/g_{Af}}{1 + (g_{Vf}/g_{Af})^2}$$



# More Precision measurements – tau lepton polarisation

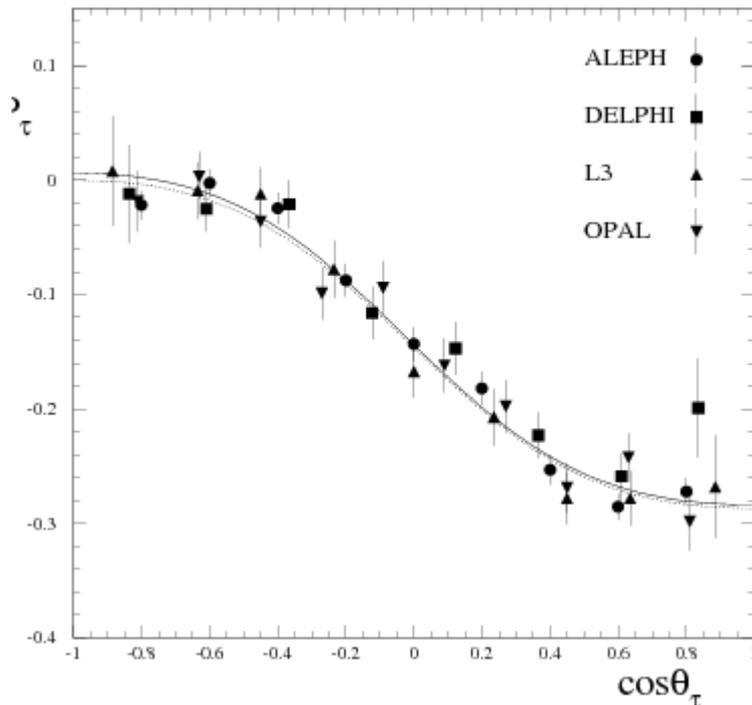
in  $e^+e^- \rightarrow \tau^+\tau^-$ :

Spin in final state can be measured assuming V-A structure in  $\tau$  decay

average  $\tau$  polarization depends on e and  $\tau$  couplings

$$\mathcal{P}_\tau(\cos\theta) = -\frac{\mathcal{A}_\tau(1 + \cos^2\theta) + 2\mathcal{A}_e \cos\theta}{1 + \cos^2\theta + 2\mathcal{A}_e \mathcal{A}_\tau \cos\theta}$$

Measured  $\mathcal{P}_\tau$  vs  $\cos\theta_\tau$



allows precise measurement  
of vector and axial vector  
couplings of  $\tau$  to  $Z$ !

# More Precision measurements – with polarized e-beam

Measurements at SLAC linear collider:

**polarized**  $e^-$  colliding with unpolarized  $e^+$  at  $\sqrt{s}=M_Z$

measurements analogous to LEP, but

can determine  $\sigma$  and  $A_{FB}$  for left- and right-handed  $e^-$

observables at SLAC:

$$A_{LR} = \frac{1}{\mathcal{P}_e} \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} \propto \mathcal{A}_e$$

$$A_{fb,LR} = \frac{1}{\mathcal{P}_e} (A_{fb,L} - A_{fb,R}) \propto \mathcal{A}_f$$

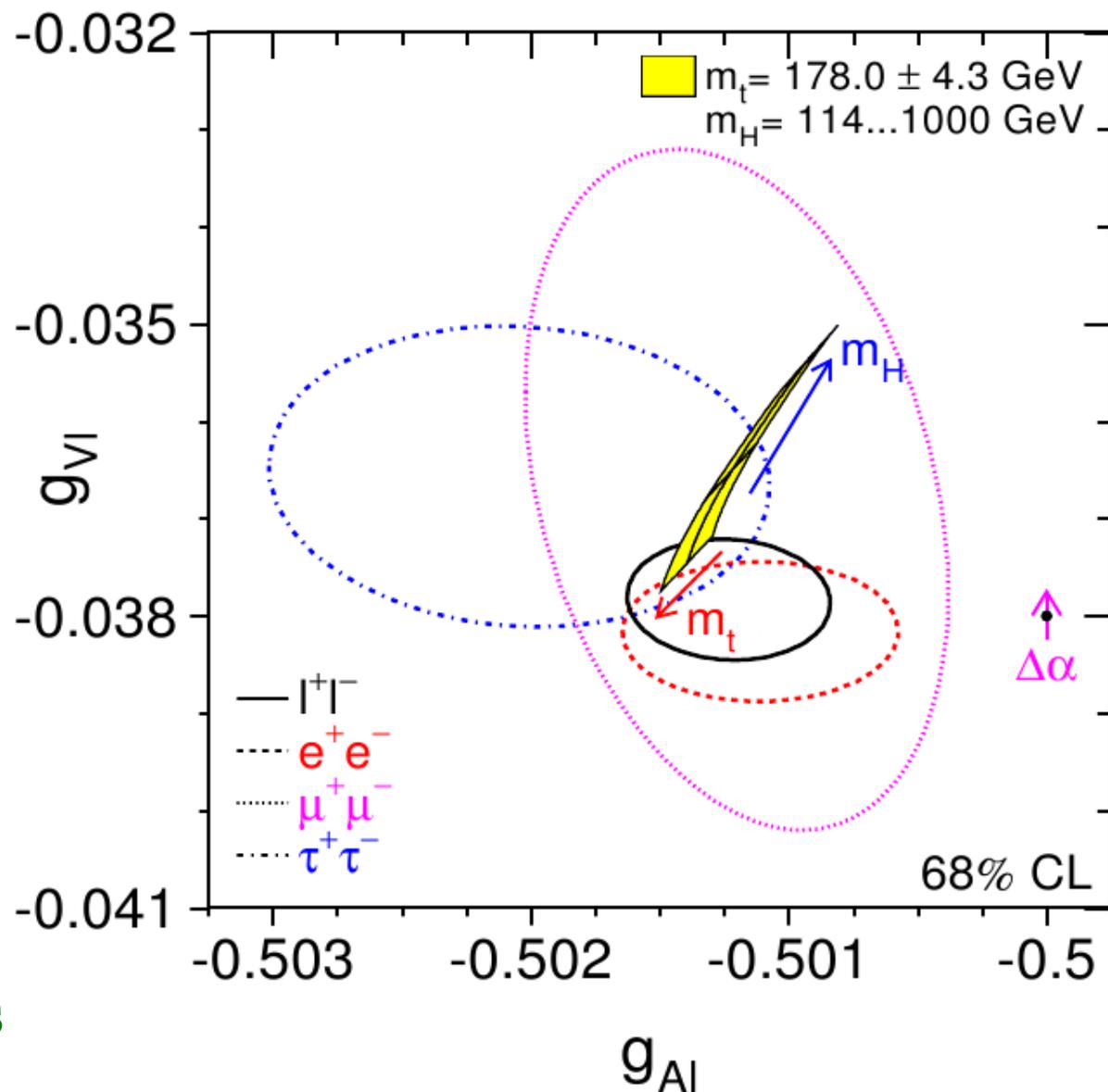
$A_{LR}$  is most sensitive single measurement to  $\sin^2\theta_W^{\text{eff}}$

# Determination of lepton couplings to Z boson

Combination of previously shown measurements gives most precise values of lepton couplings.

Consistent with “lepton universality”  
→ combine all into leptonic  $g_V$  and  $g_A$

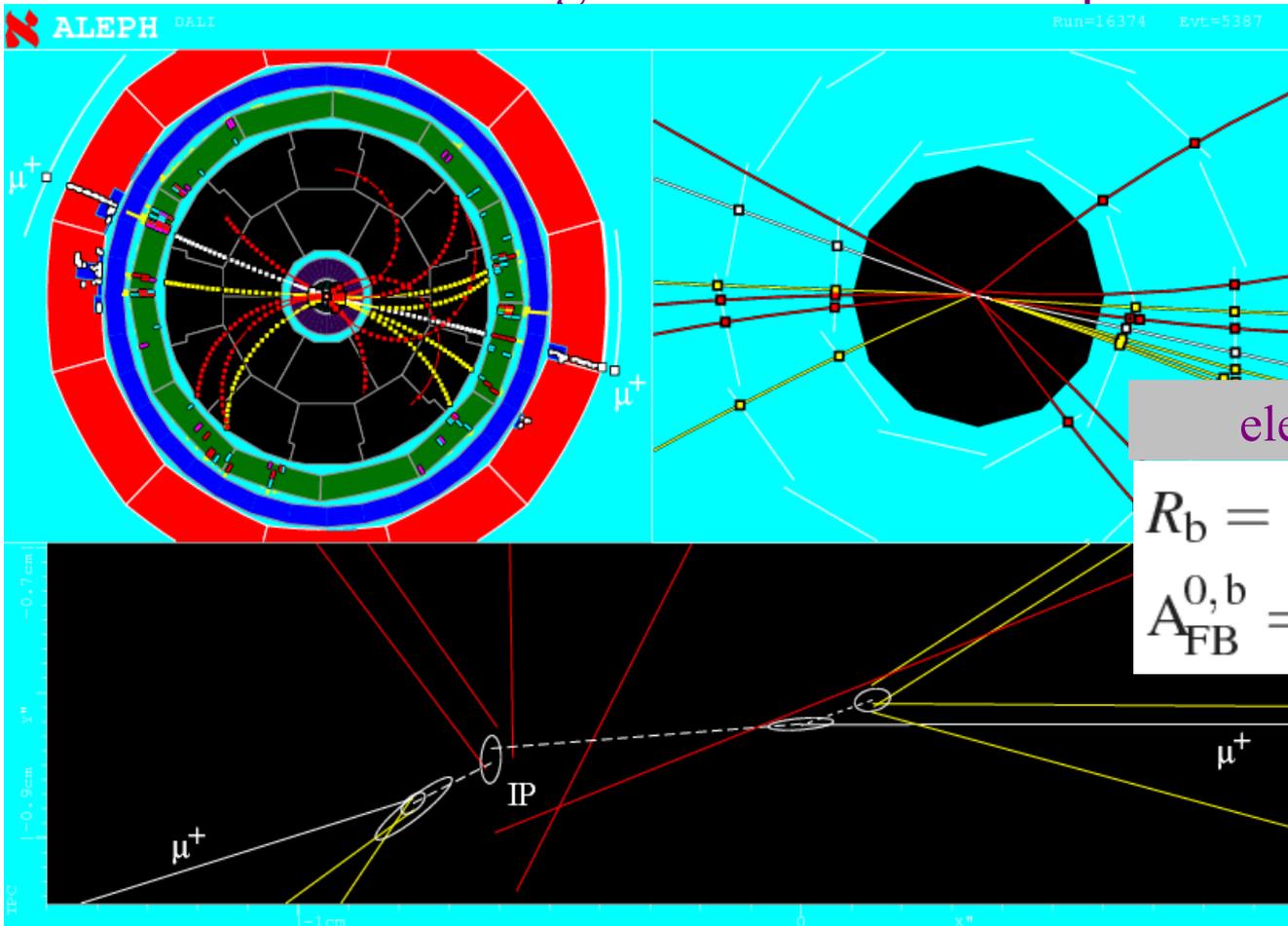
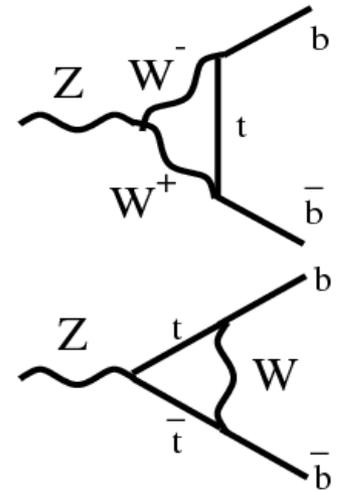
**Measurements are sensitive to top and Higgs boson masses**



# More Precision measurements – heavy quarks

- b-quarks** – are (relatively) long-lived  
 – are heavy  
 – decay in cascades → can be distinguished from other quarks

special relation to top quark:



electroweak observables

$$R_b = \frac{\Gamma_{b\bar{b}}}{\Gamma_{had}} \Rightarrow \rho_b$$

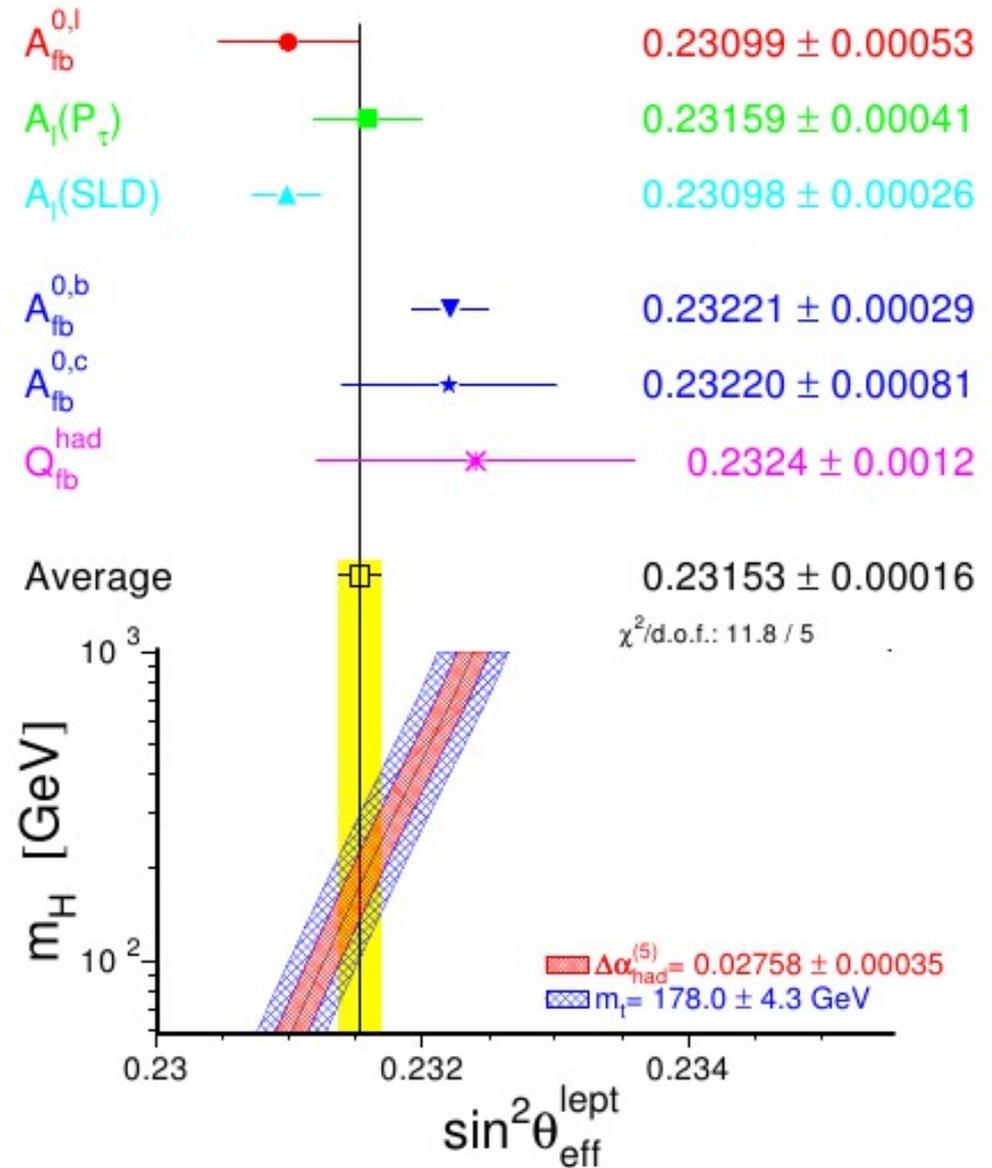
$$A_{FB}^{0,b} = \frac{4}{3} \mathcal{A}_e \mathcal{A}_b \Rightarrow \sin^2 \theta_{eff}^{lept}$$

# Summary of asymmetry-type measurements

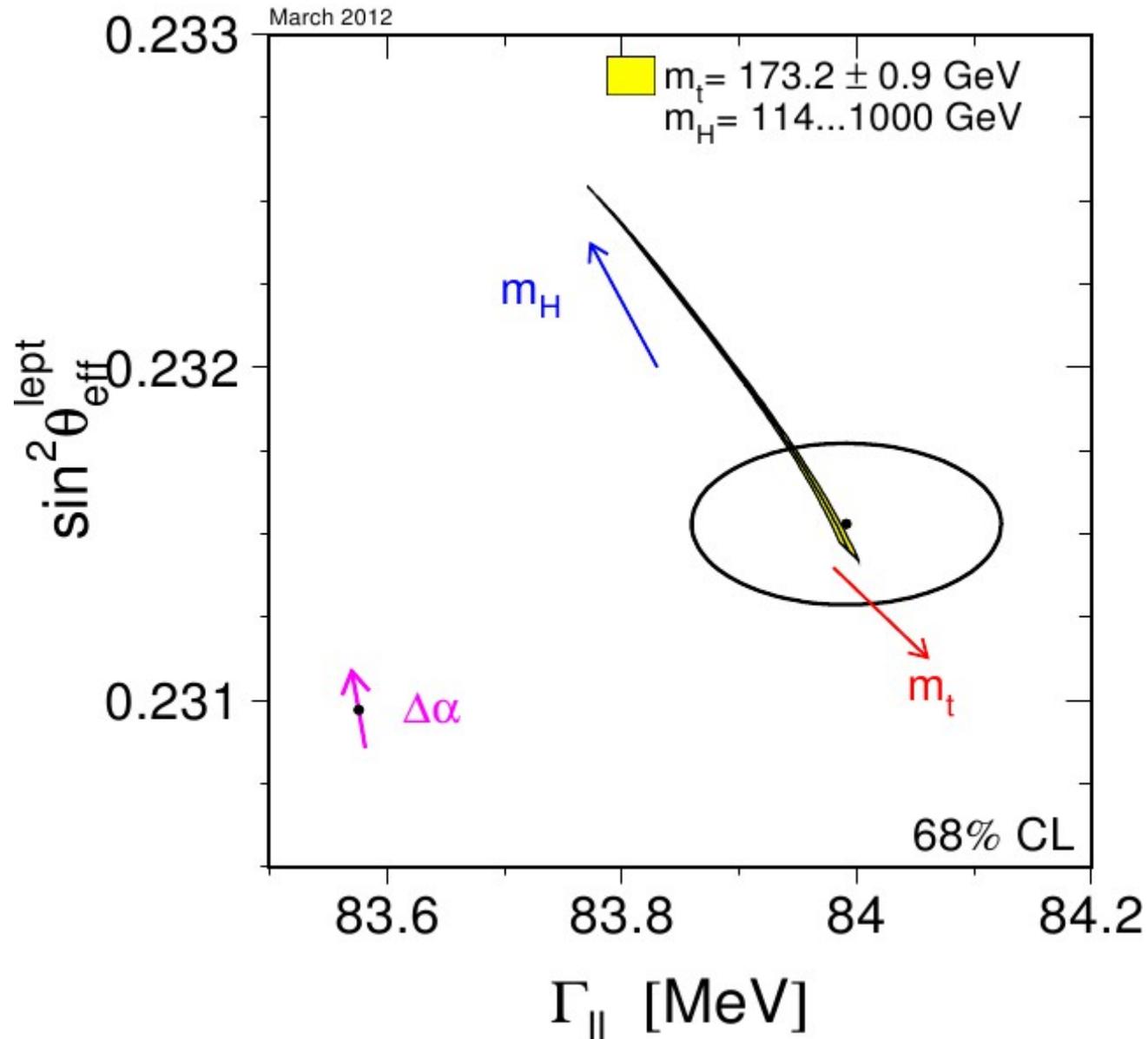
Asymmetry-type measurements depend on vector and axial vector couplings,  $g_V$  and  $g_A$ , and hence on  $\sin^2\theta_W^{\text{eff}}$

Average of  $\sin^2\theta_W^{\text{eff}}$  very sensitive to Higgs mass

There is, however, some “tension” between the most precise measurements !

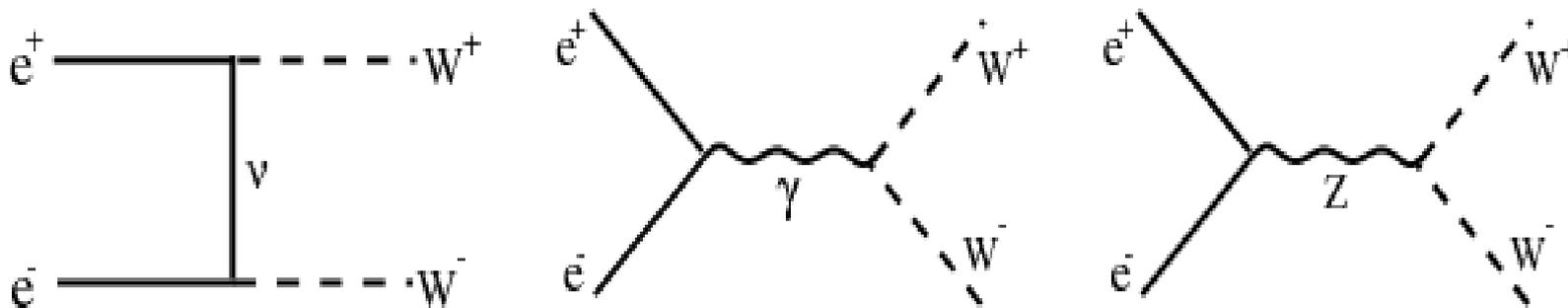


# Combination of precision measurements: $\Gamma_{\text{lept}}$ and $\sin^2\theta_{\text{W}}^{\text{eff}}$



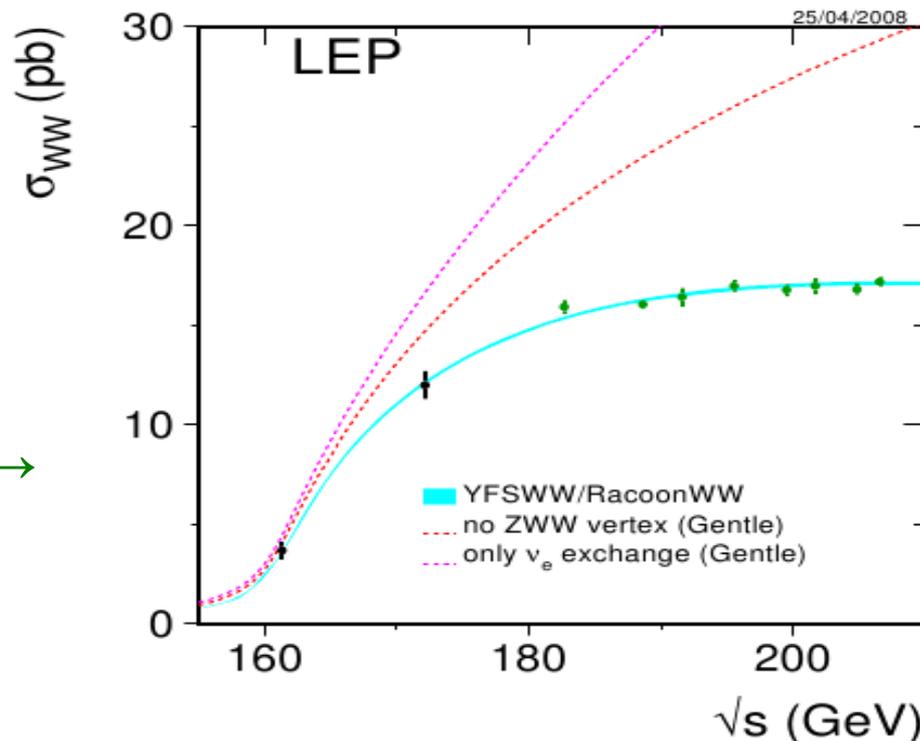
$\Gamma_{\text{lept}}$  and  $\sin^2\theta_{\text{W}}^{\text{eff}}$   
 together provide  
 strong constraint  
 on Higgs mass !

# W pair production at LEP II



$e^+e^- \rightarrow W^+W^-$  cross section in EW theory remains finite only if photon, Z and neutrino exchange are all taken into account -

**experimentally well confirmed** →



**Remark:** there is also a Higgs diagram, but negligible in  $e^+e^-$  !

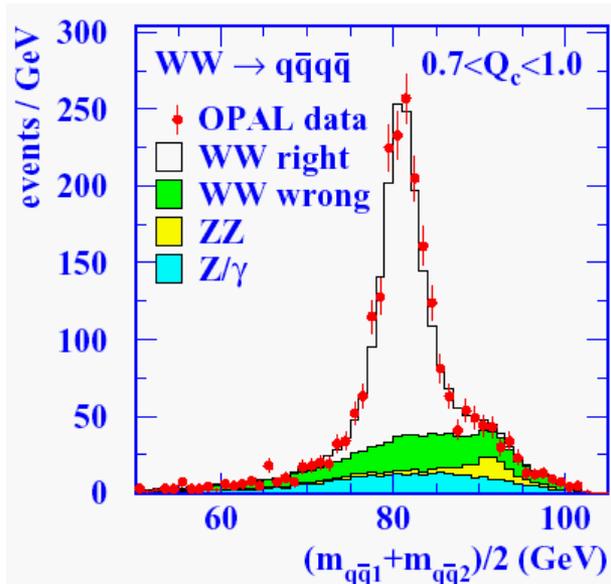
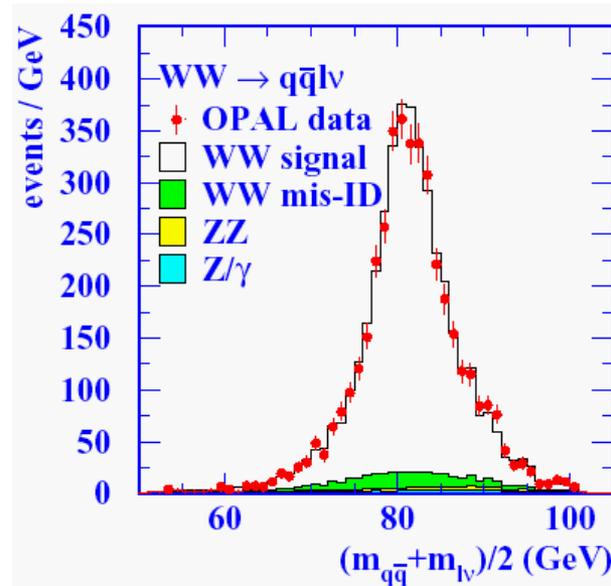
# W Boson mass

## W-Boson mass

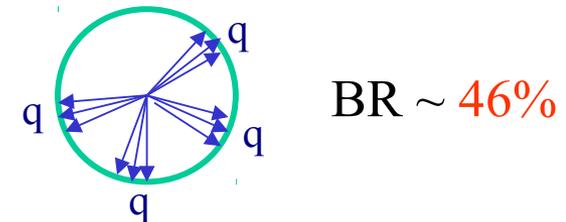
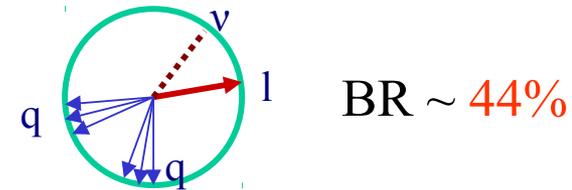
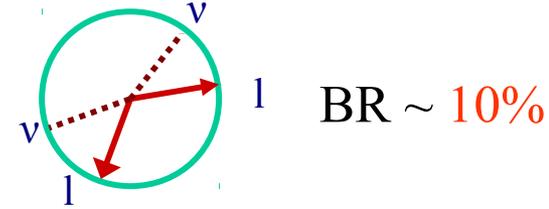
from reconstructed objects in detectors by LEP experiments ALEPH, DELPHI, L3 and OPAL and by Tevatron experiments CDF & D0

(LHC not yet ...)

very difficult,  
dominated by  
systematic errors

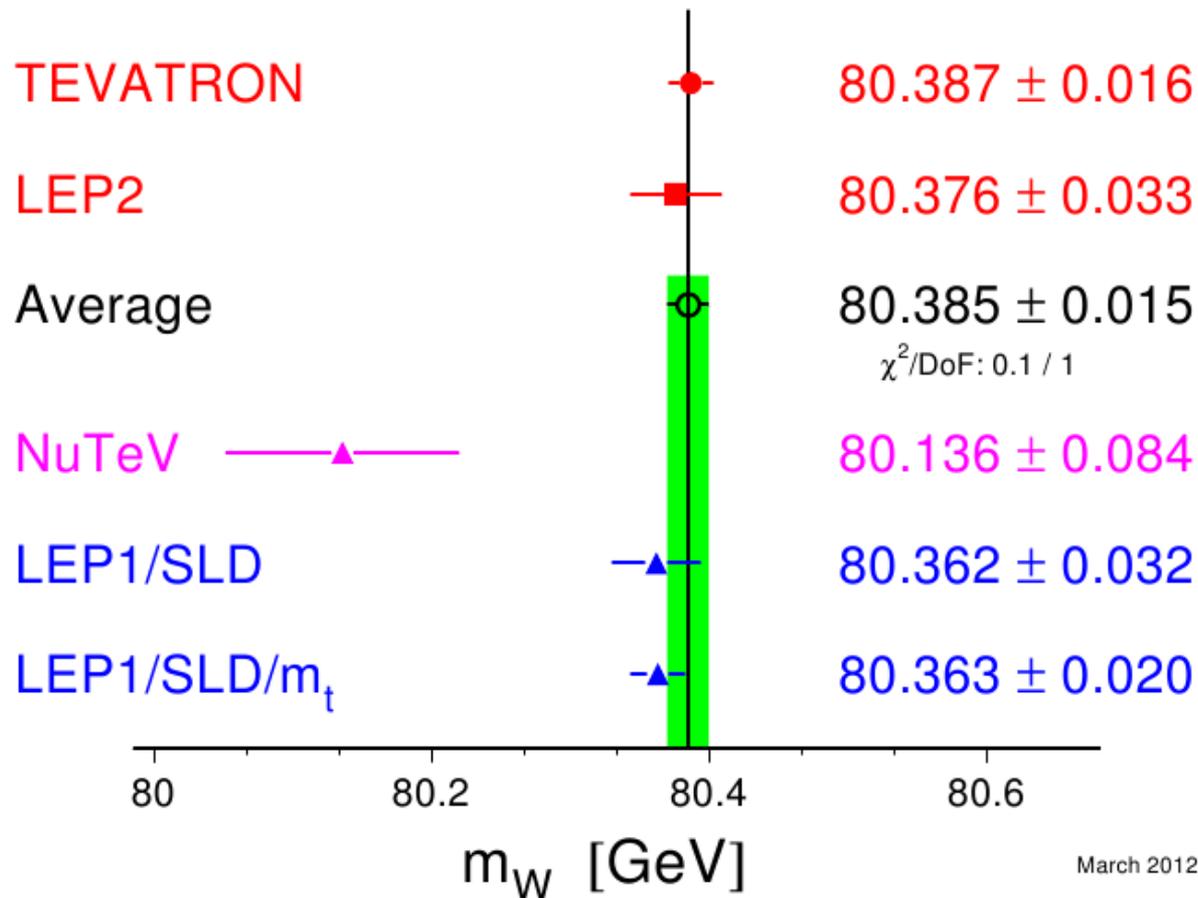


@LEP: W pair events:



# W Boson mass – world average 2012

W-Boson Mass [GeV]



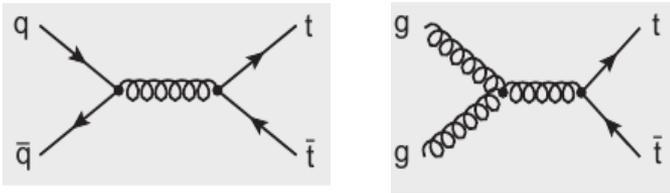
**Remember:** Relevance of W mass:  
determines on-shell weak angle

# One more ingredient: need the top quark mass

## top quark mass

measured from reconstructed objects in detectors by Tevatron experiments CDF & D0 (LHC on the way ...)

- top quarks (mostly) produced in pairs via  $q\bar{q}$  or  $gg$



- dominant decay:

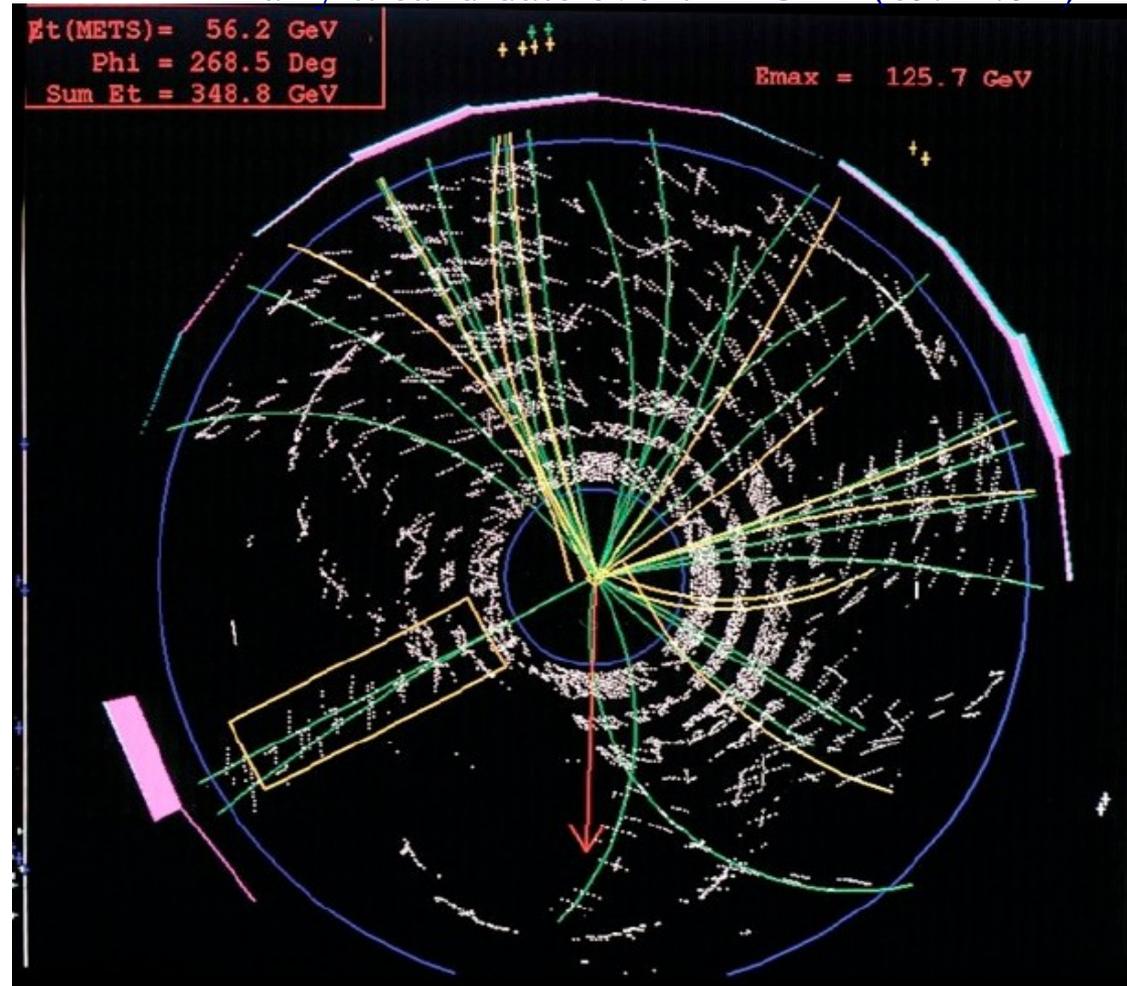
$$t \rightarrow b + W^+$$

$$W \rightarrow \bar{q}q' \text{ or } l\nu$$

signatures studied:

- fully hadronic
- lepton + jets
- di-lepton

Early  $t\bar{t}$  candidate event in CDF (09/24/92)



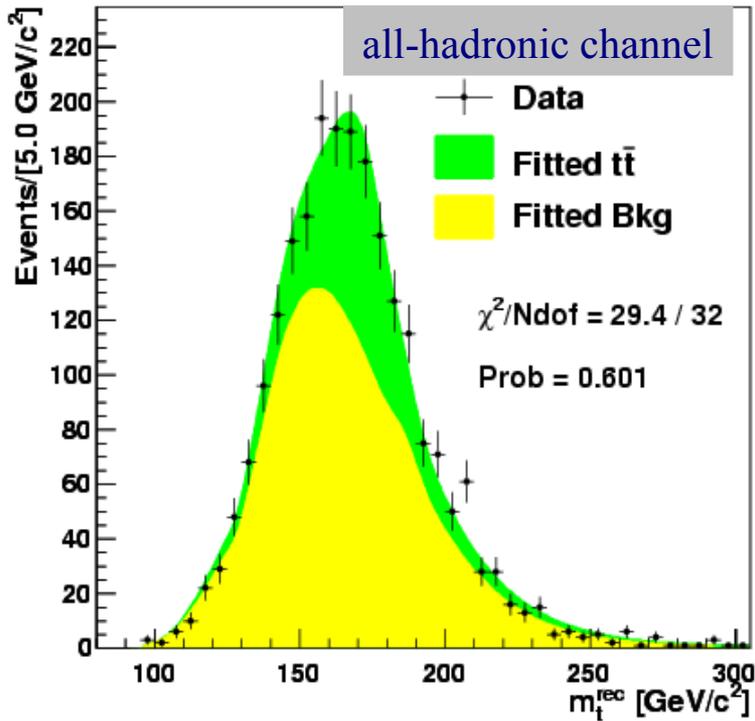
**Signature:** “lepton + jets”

- 1 lepton
- four jets
- missing transverse energy

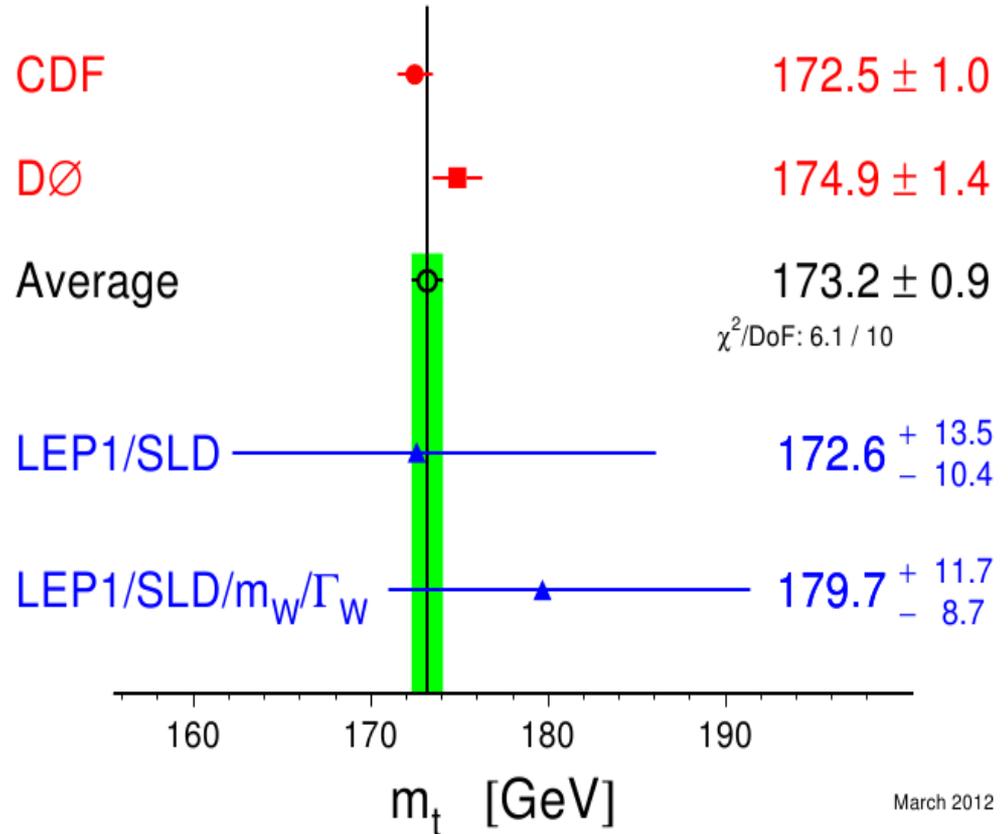
# One more ingredient: top quark mass

Example of observed mass distribution:

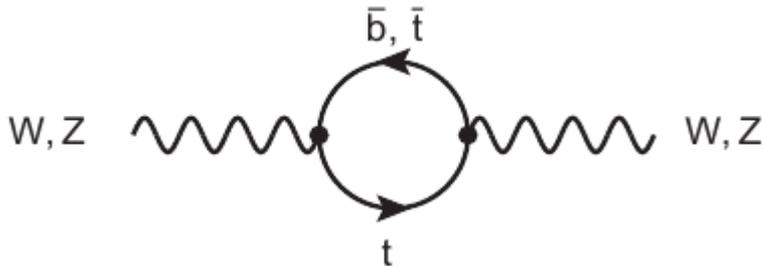
CDF Run II Preliminary (5.8 fb<sup>-1</sup>)



Top-Quark Mass [GeV]



March 2012



**Remember:** dominant radiative corrections arise from top ( $\sim G_F m_t^2$ ) !

# Top quark mass from loop corrections ...

... at Moriond conference  
in March 1994:

direct search for top at Tevatron:

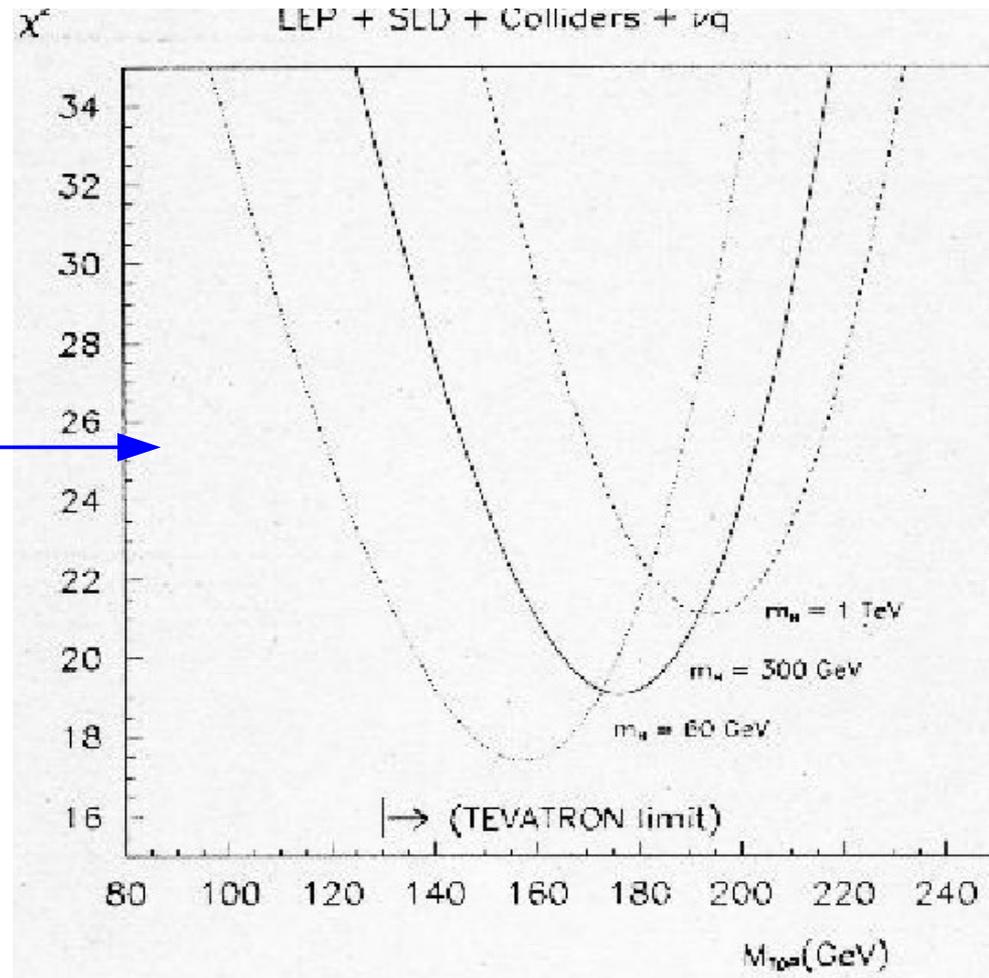
$$m_t > 130 \text{ GeV}/c^2$$

from radiative corrections:

$$m_t = 177 \pm 11_{-19}^{+18} \text{Higgs} \text{ GeV}/c^2$$

a little later in summer, direct  
observation of top quark (CDF):

$$m_t = 174 \pm 10_{-12}^{+13} \text{syst} \text{ GeV}/c^2$$



Excellent agreement between  
top from loops and from direct measurement!

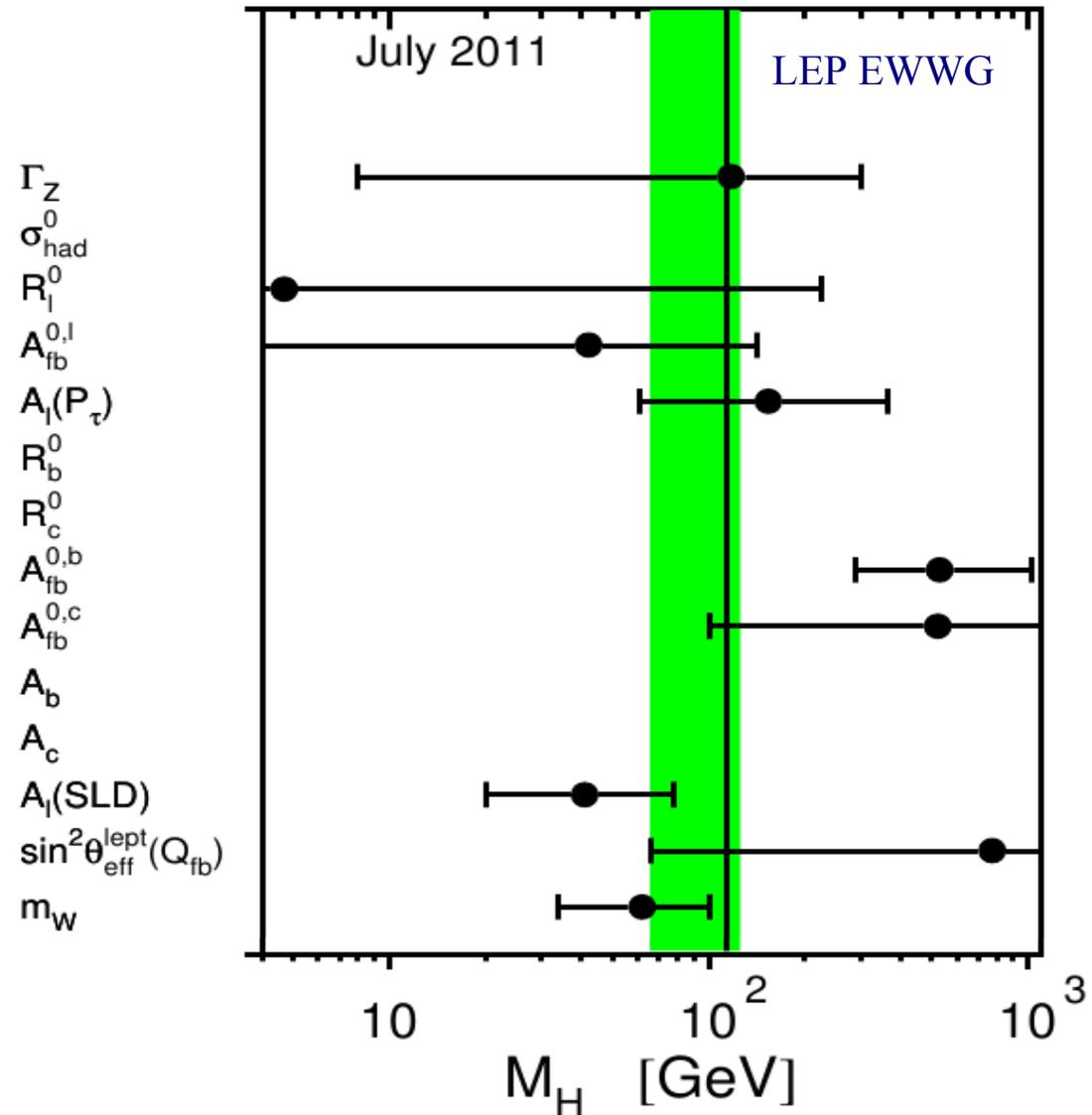
# $M_H$ from individual precision measurements

With known top mass,  
can now disentangle  
radiative effects from  
top and Higgs

( provided nothing else  
is in the loops ...)

→

$M_H$  from precision  
measurements  
with large errors



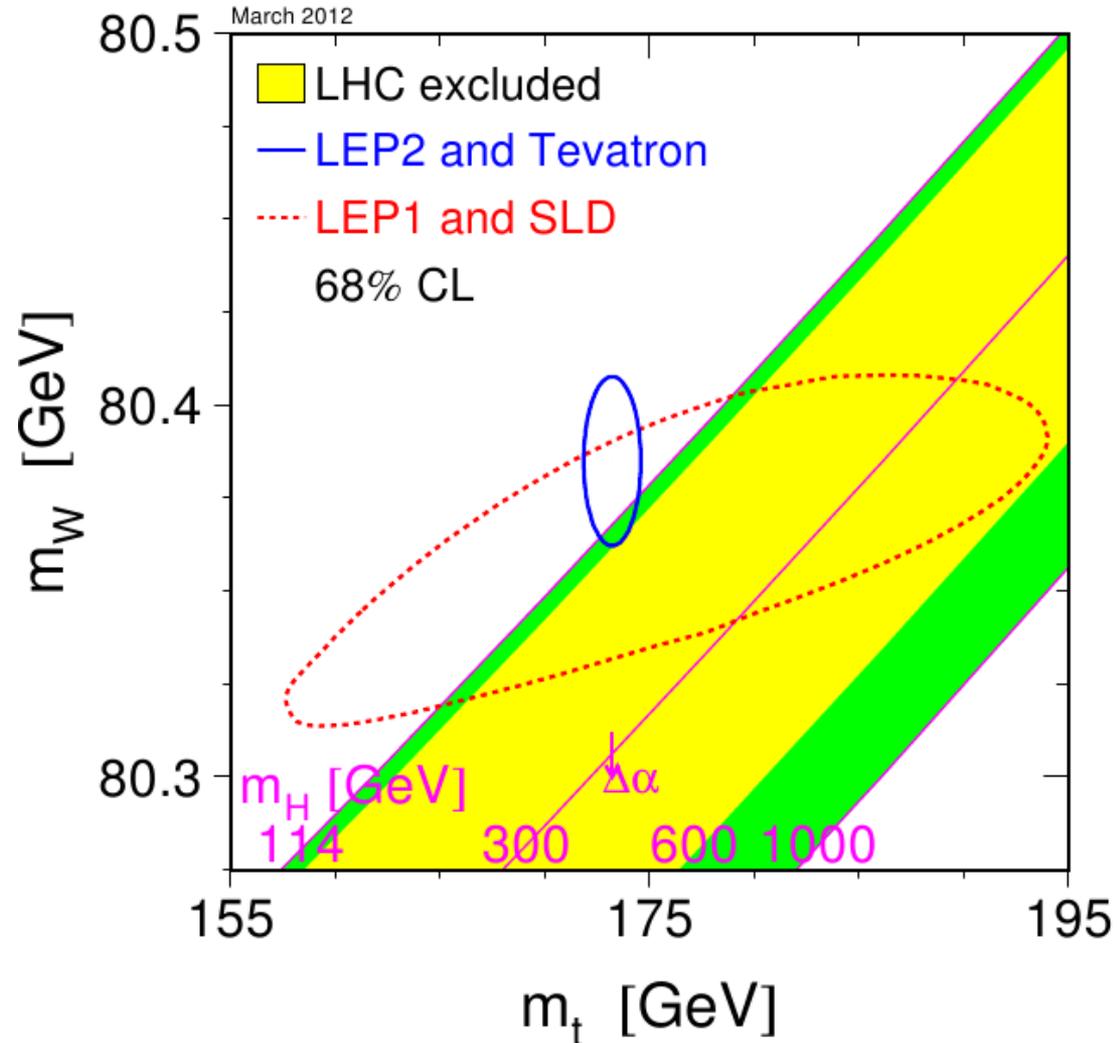
# Putting it all together ....

... in a rather complicated plot:

top quark vs. W boson mass

1. direct measurements
2. indirect determination from precision variables:  
 $O := O(\alpha, G_F, M_Z, m_t, M_H)$   
solved numerically  
(“Newton's method”) for  $m_t$  and  $M_H$  (using SM)
3. W mass predicted for different values of  $M_H$

**nice consistency check,**  
prefers low value for  
Higgs mass !



# Putting it all together ....

... in an overall fit

of Higgs boson mass  
within the minimal  
Standard Model

“pull plot” and overall  
value of  $\chi^2$  to identify  
potential problems

**good overall consistency,**  
but there is one “outlier”;  
the measurement of the  
b-quark forward backward  
asymmetry.



July 2011

<http://lepewwg.web.cern.ch/LEPEWWG>

# Putting it all together ....

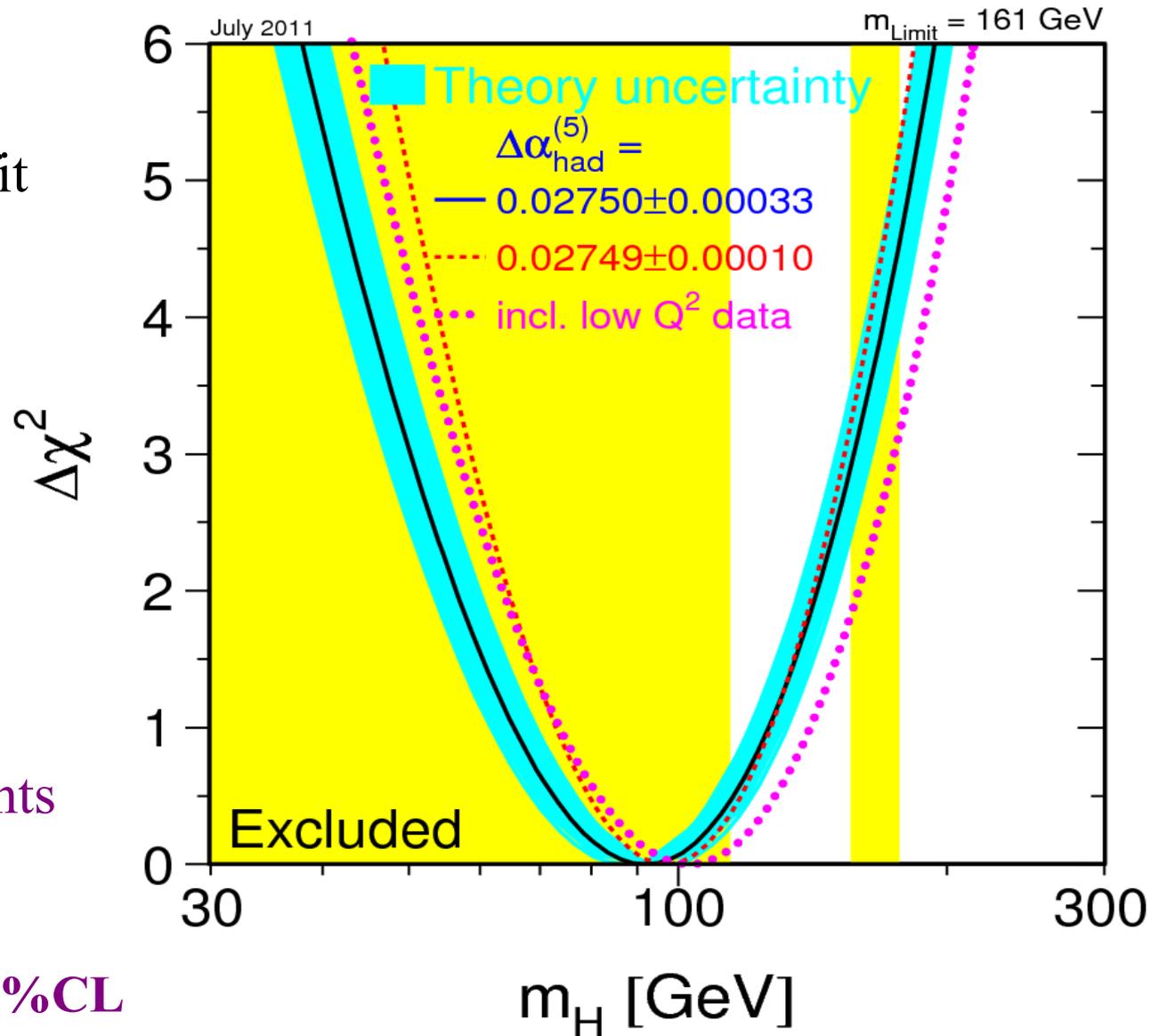
... profiled  $\chi^2$

for  $M_H$  from overall fit  
within the minimal  
Standard Model

**“Blue Band Plot”**

Precision measurements  
predict mSM Higgs  
boson to be light,

$M_H < 161 \text{ GeV}/c^2 @ 95\% \text{ CL}$



Loops affect theory itself

# Conclusions on results from precision measurements

## Standard Model established as a renormalizable field theory:

- in agreement with (almost) all experimental results, prec.  $\sim 0.1\%$
- quantum corrections seen and well established
- indirect determination of top quark mass
- triple gauge boson couplings confirmed with prec. 1%
- precision measurement constrain Higgs to be light

Theoretical work on proof of renormalizability awarded with Nobel prize in 1999



Gerardus 't Hooft



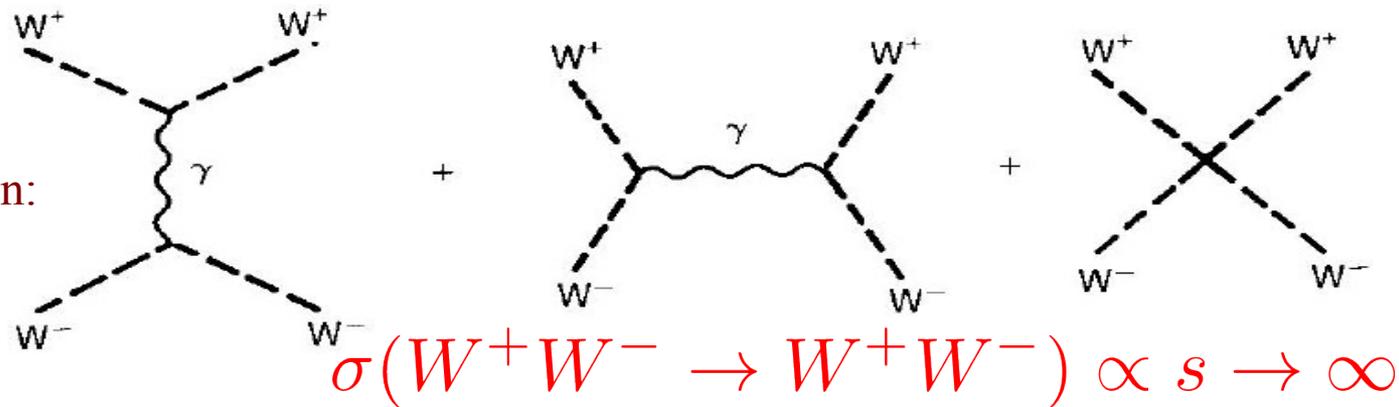
Martinus J.G. Veltman

The Nobel Prize in Physics 1999 was awarded jointly to Gerardus 't Hooft and Martinus J.G. Veltman "for elucidating the quantum structure of electroweak interactions in physics"

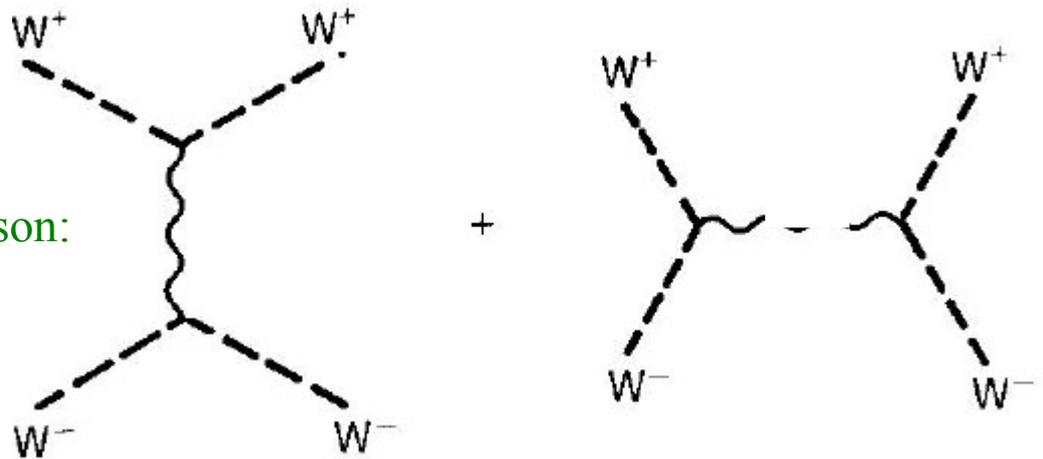
# Theoretical constraints on Higgs boson mass

Diagrams with Higgs-Boson prevent divergencies:

Without Higgs boson:



additional diagrams involving Higgs boson:

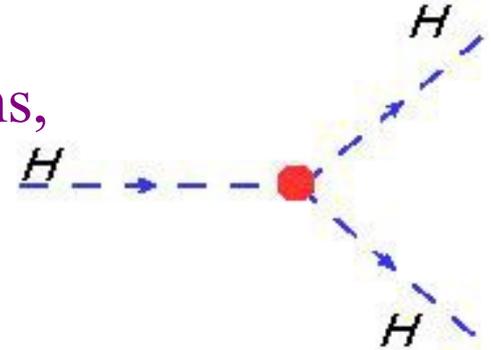


Cross section for WW scattering remains finite

Higgs must not be too heavy::  $< \sim 1 \text{ TeV}/c^2$

# Theoretical constraints $M_H$ loop corrections

Higgs propagator affected by higher order corrections,  
in particular from Higgs self couplings



$$m^2(p^2) = m_0^2 + \frac{\text{Diagram 1}}{p} + \text{Diagram 2} + \text{Diagram 3}$$

The equation shows the Higgs mass squared  $m^2(p^2)$  as a function of momentum  $p$ . It is equal to the tree-level mass  $m_0^2$  plus three loop corrections:

- Diagram 1:** A loop of spin-1 particles (J=1), represented by a wavy line. It is labeled with  $J=1$  and  $H$ .
- Diagram 2:** A loop of spin-1/2 particles (J=1/2), represented by a fermion loop. It is labeled with  $J=1/2$ .
- Diagram 3:** A loop of spin-0 particles (J=0), represented by a scalar loop. It is labeled with  $J=0$  and  $H$ .

Higgs in loops couples to itself with strength proportional to Higgs mass  $\rightarrow$

large corrections to Higgs mass and Higgs potential

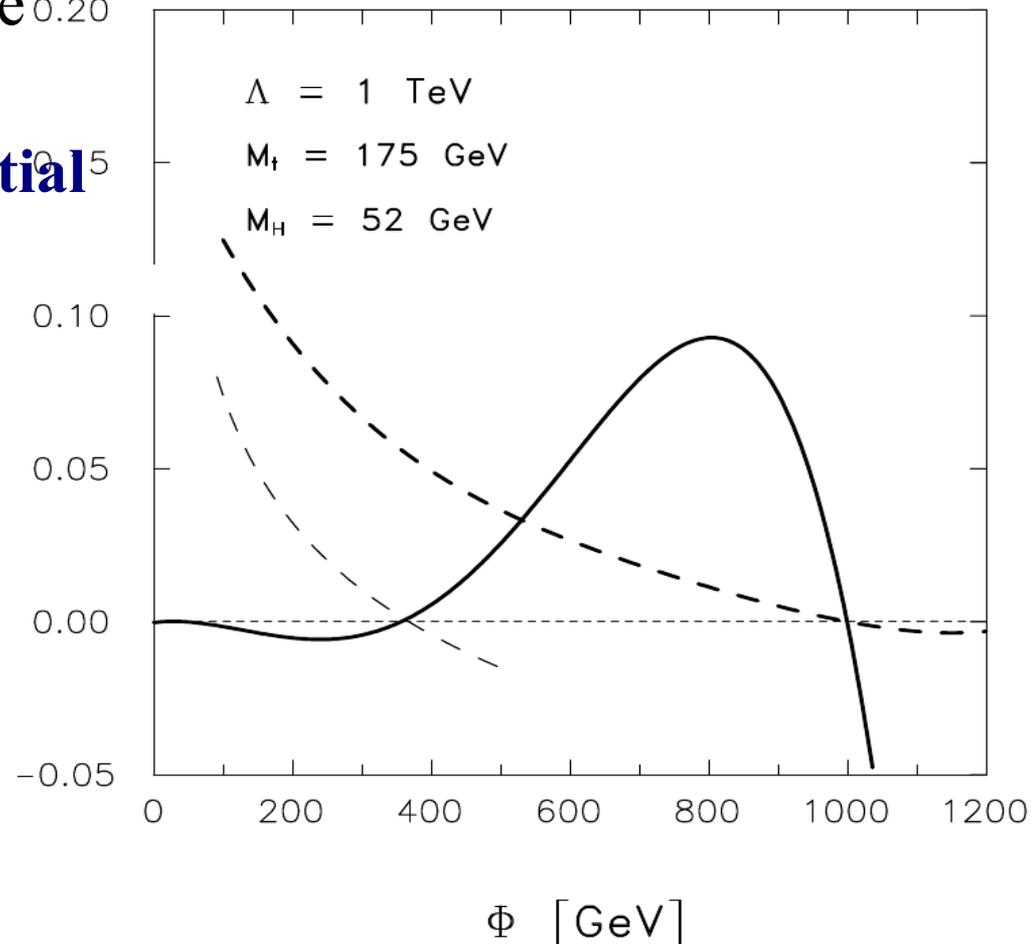
- high Higgs mass: Higgs self-coupling becomes strongly interacting, breakdown of perturbation theory
- small Higgs mass: Higgs self-coupling becomes negative  $\rightarrow$  unstable vacuum  
( in contradiction to age of universe)

# Theoretical constraints on $M_H$ – vacuum stability

## Example of renormalized Higgs Potential

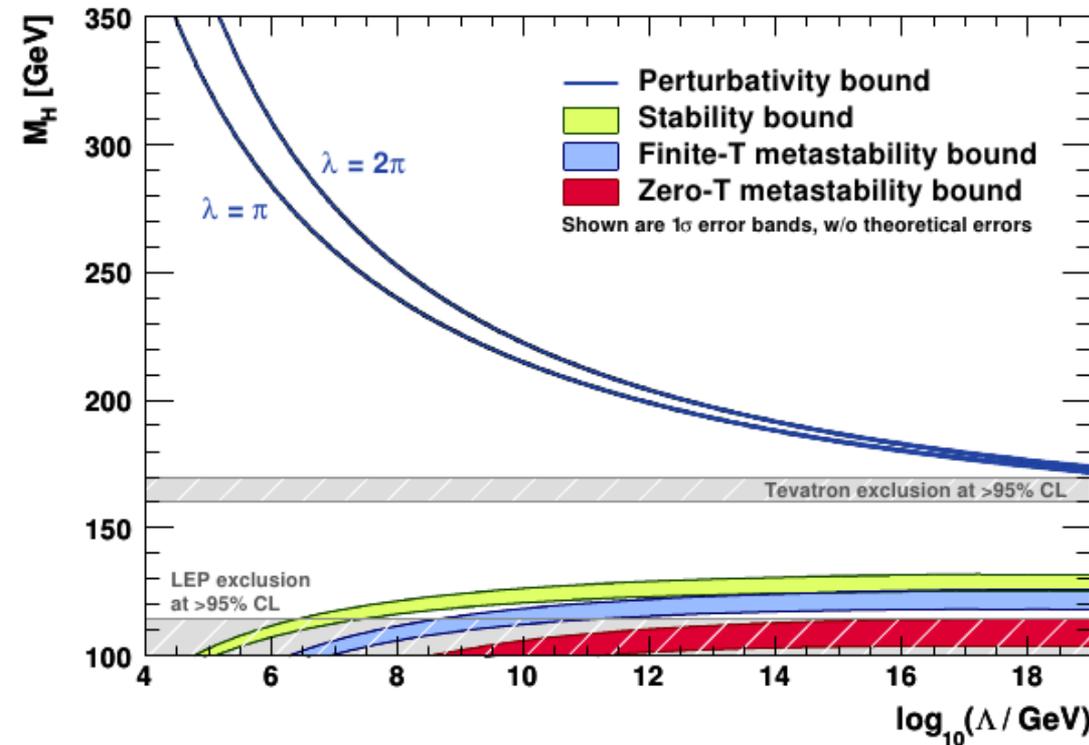
Casas, Espinosa, Quiros (1996)

**Loops** (in particular top) drive quartic coupling negative above some scale  $\Lambda$  for small  $M_H$   
→ renormalized **Higgs potential** is either **unbounded** (inconsistent with the existence of the universe) or in local minimum “false vacuum”, (assume tunneling time  $> 14$  billion years to derive limit on  $M_H$ )

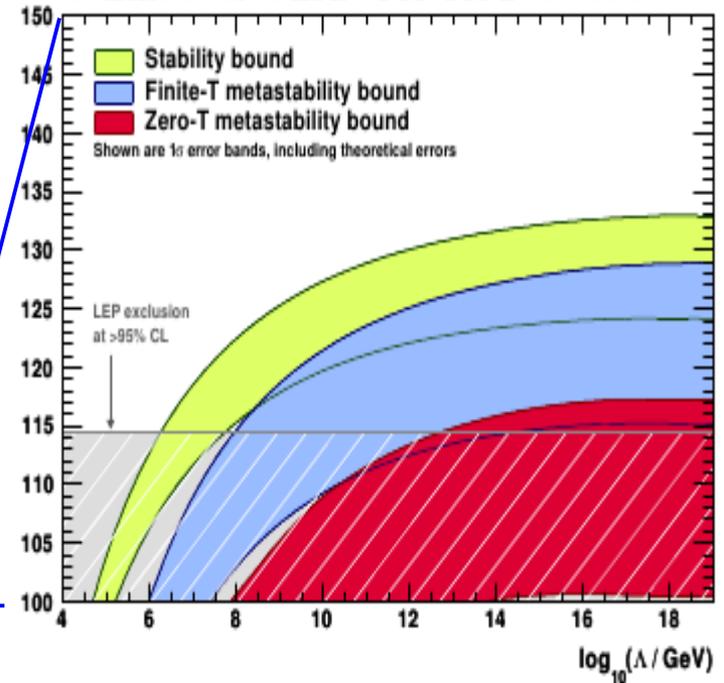


# Theoretical constraints on $M_H$ – vacuum stability

→ beyond energy scale  $\Lambda$  new Physics must be present



J. Ellis et al. CERN-PH-TH/2009-058



→ corridor of allowed values of Higgs mass within Standard Model

# Final remarks

We know that the Standard Model is incomplete:

- no candidate for dark matter  
(would have to be heavy, stable and only weakly interacting)
- no explanation of “dark energy”
- inclusion of gravity causes problems  
(Why are Higgs mass and Boson masses so “small” compared to the relevant scale, the Planck mass?)

There must be new physics, even if the Higgs boson will be found one day at a mass making the Standard Model self-consistent).

# Summary:

Standard Model looks fine and gauge couplings and Higgs mechanism appear to work, although the model is clearly incomplete.

## To be done:

- discovery of Higgs boson
- determination of Yukawa couplings to fermions
- test Higgs potential via measurement of Higgs self couplings

Will the model show inconsistencies and thus point to way to extend it, or will it remain to be a self-consistent theory for the part of reality it describes ?