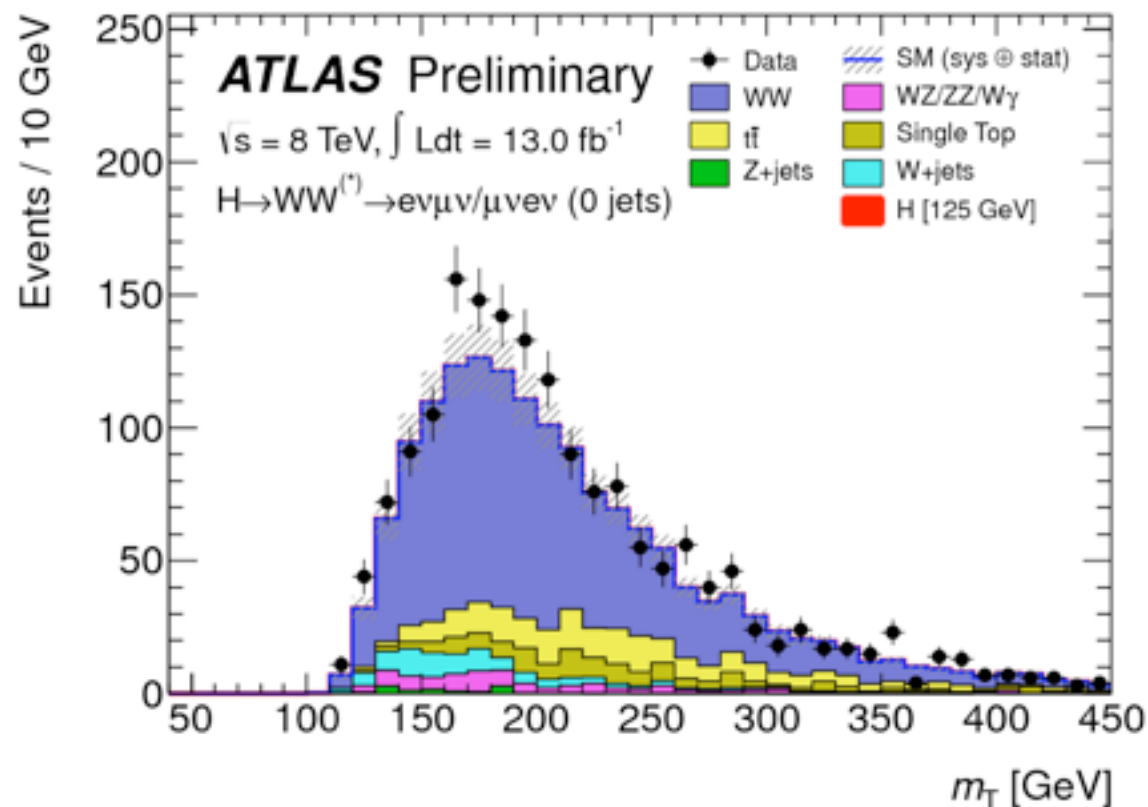
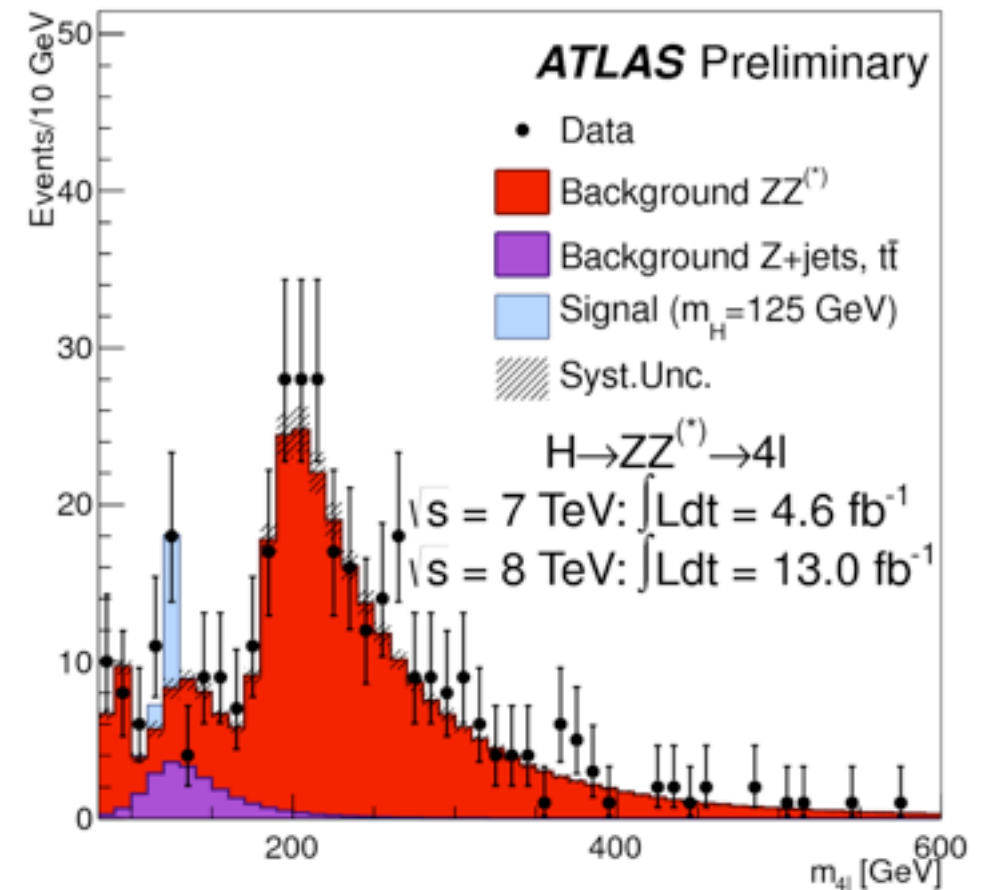
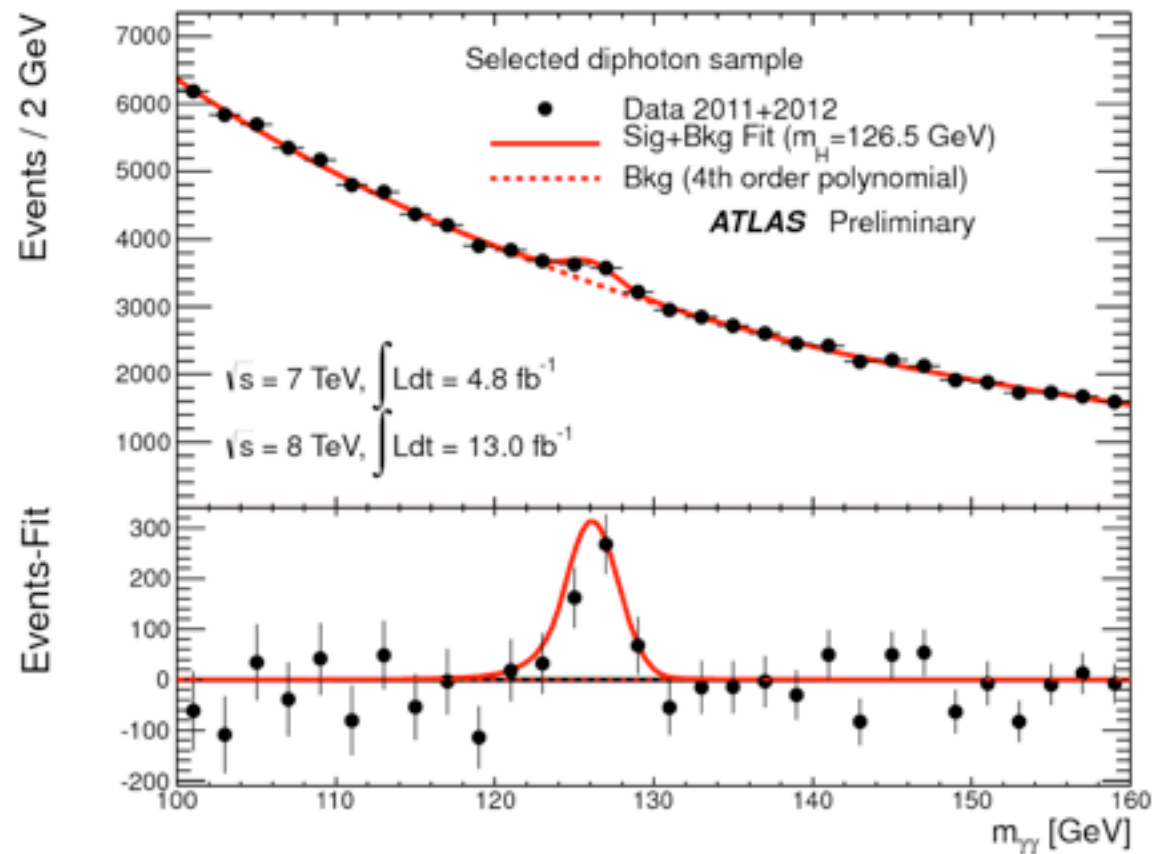


Higgs and Jets at the LHC

Michael Spannowsky

University of Durham

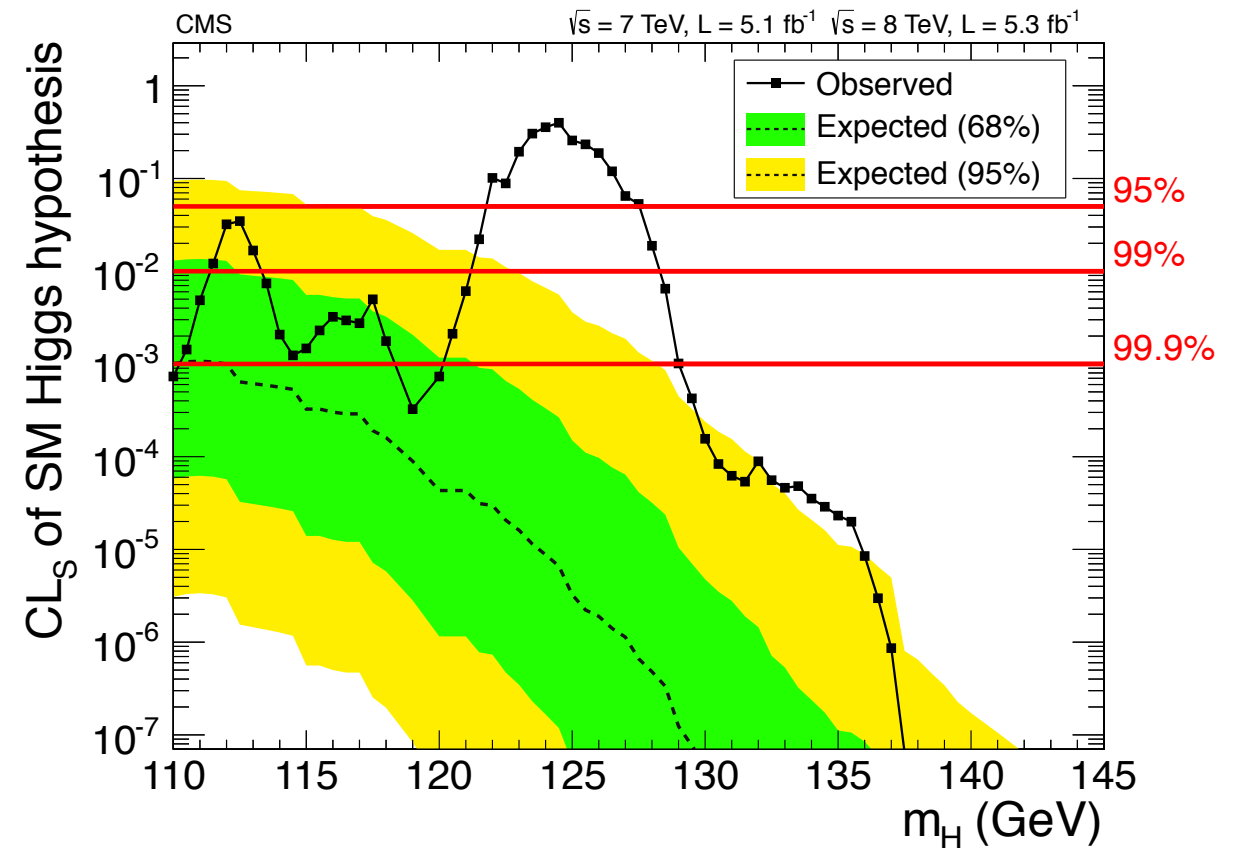
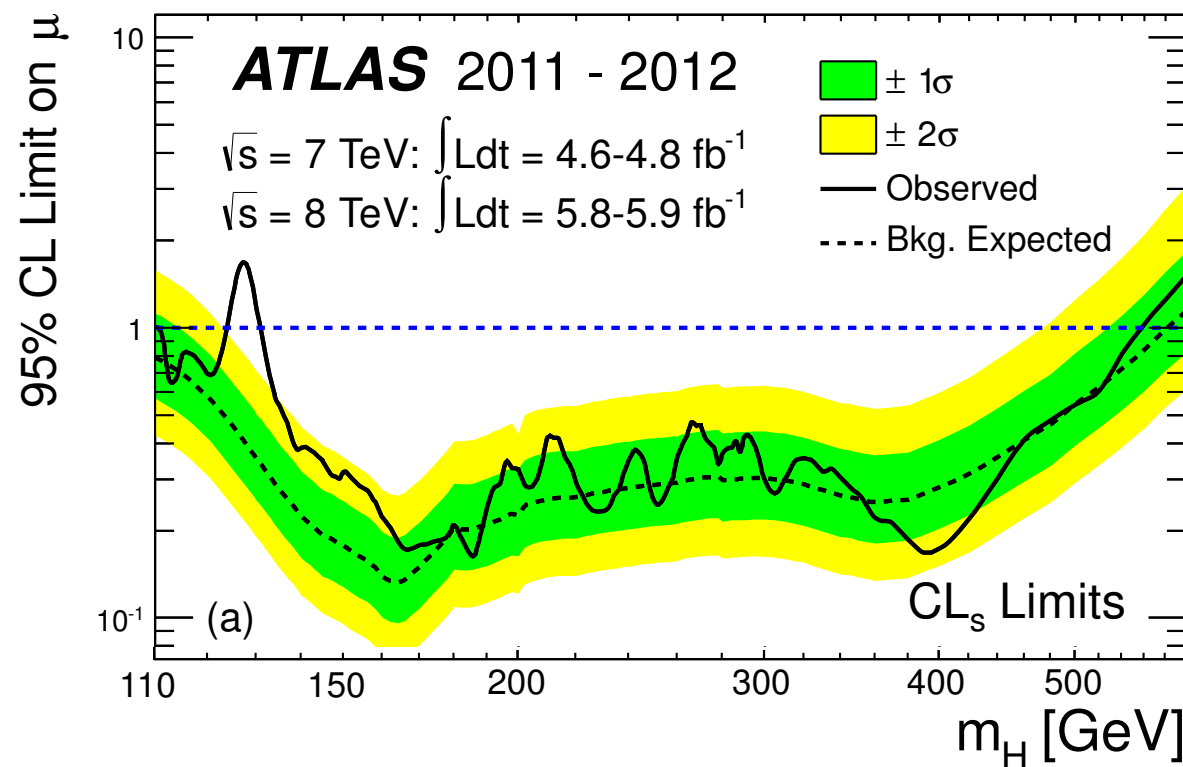
7 + 8 TeV data:



Recently observed clear excesses in 3 independent channels:

- 2 Photons
- 4 Leptons (electrons/muons)
- 2 Leptons (electrons/muons) + MET

Combined results for each experiment



Discoveries by ATLAS and CMS

- Particle consistent with Higgs boson, predicted 50 years ago
- Huge **international** and **intergenerational success!**



Whats next?

- Is what we observe "THE HIGGS BOSON"?
- Is minimal SSB mechanism realized in nature?

Only one $SU(2)$ doublet

Couplings to Fermions

Couplings to Gauge Bosons

Higgs potential and selfcoupling

CP property of Resonance

Spin of Resonance

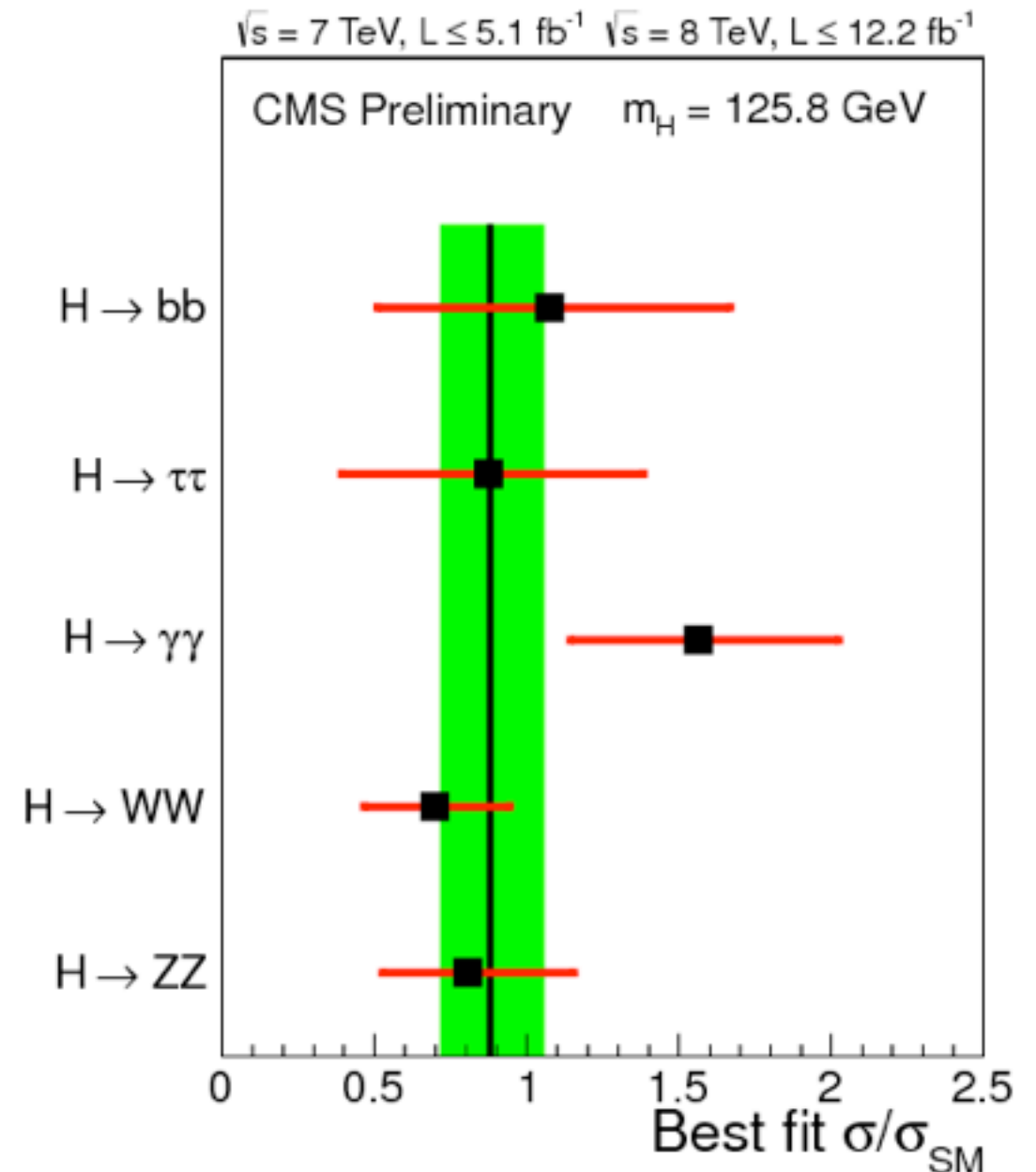
After discovery the Higgs couplings have to be measured:

Present status:

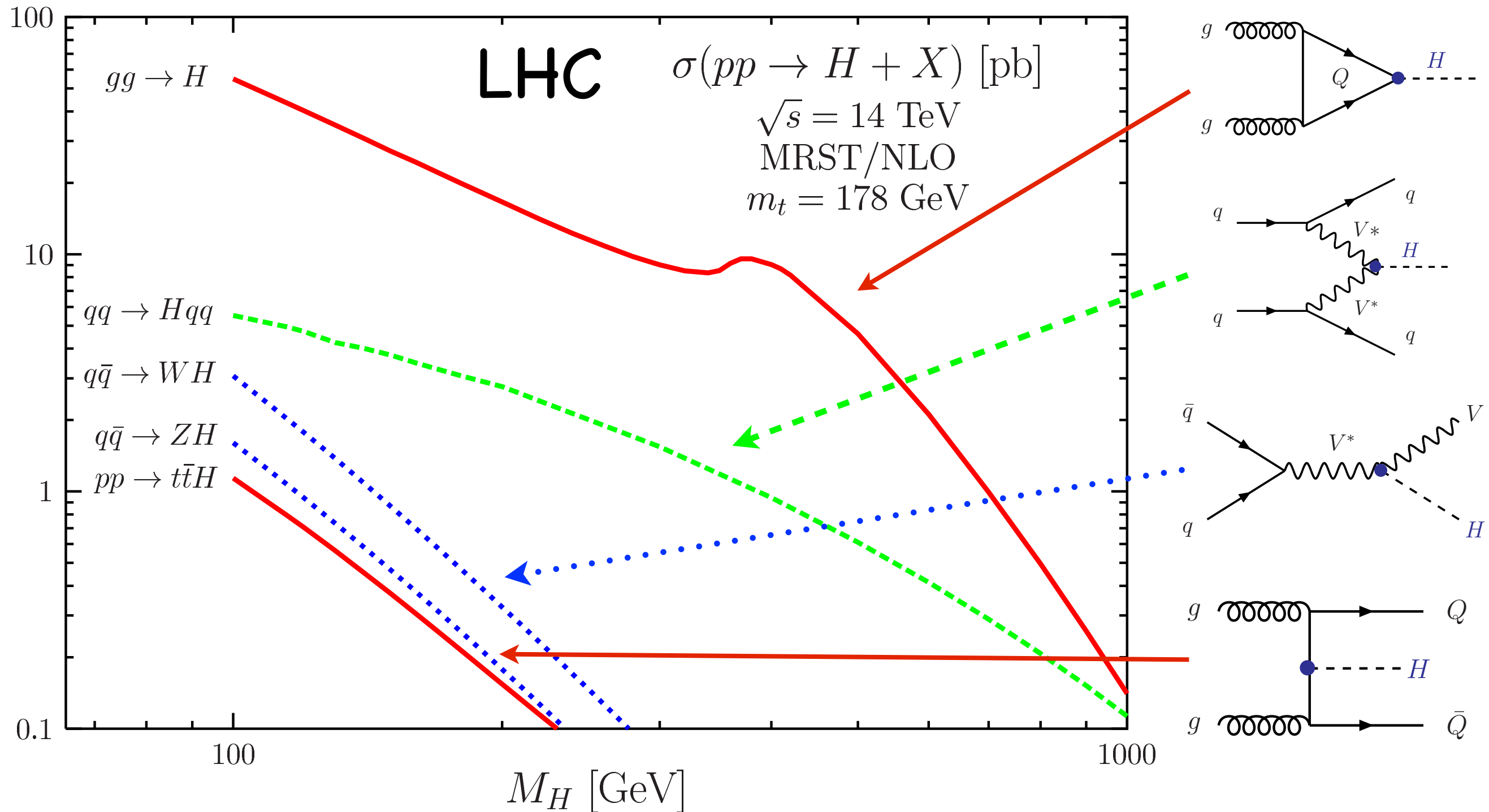
- CMS and ATLAS did fit for couplings already
- For the overall CS CMS has

$$\sigma/\sigma_{\text{SM}} = 0.87 \pm 0.23$$

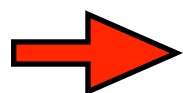
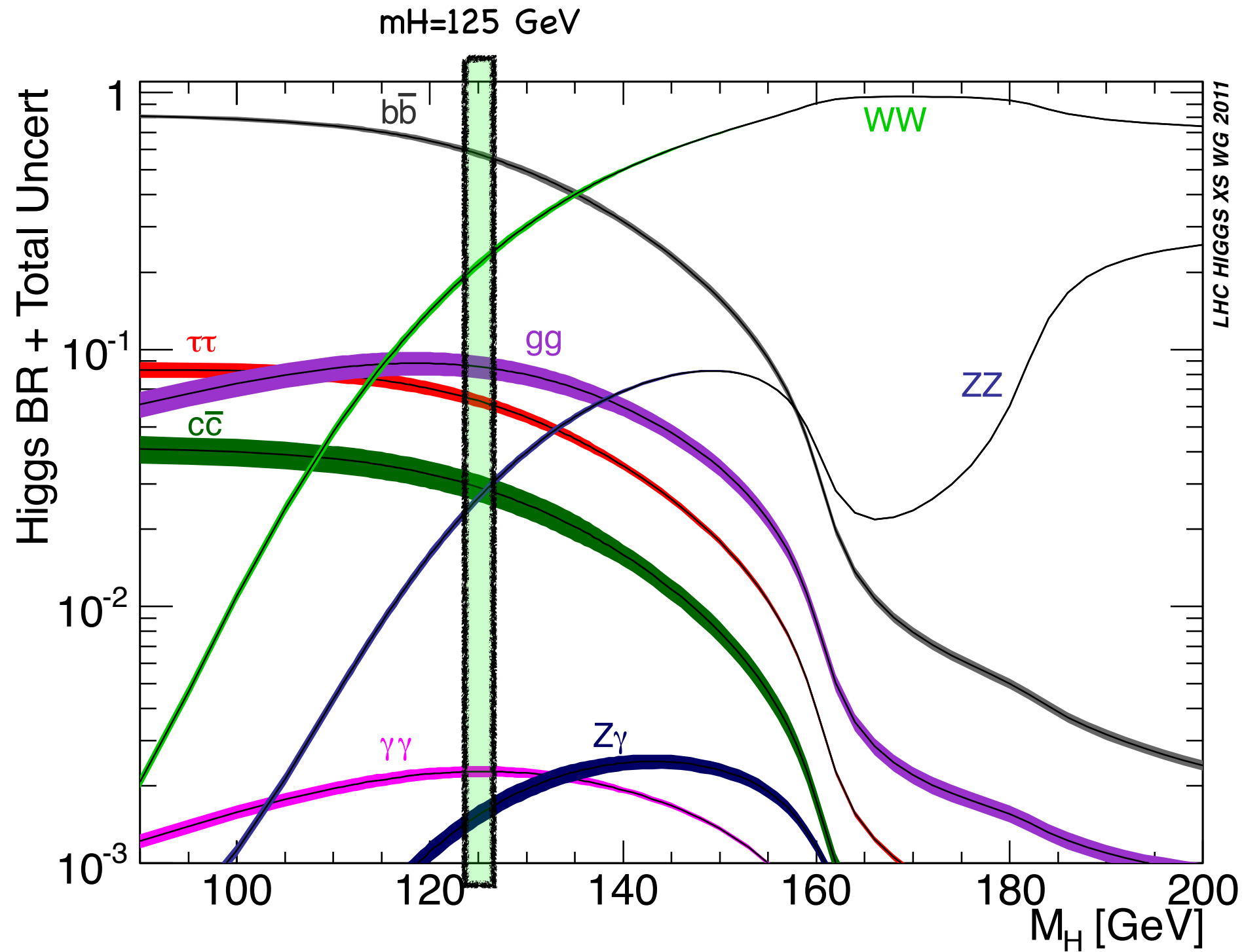
- Green band indicates $\pm 1\sigma$ uncertainty including stat. and sys. uncertainties
- Decay to photons a bit high to taus a bit low, but so far all in all good agreement with SM
- Channels are mutually related!



Production cross section at hadron colliders



Predicted decay channels



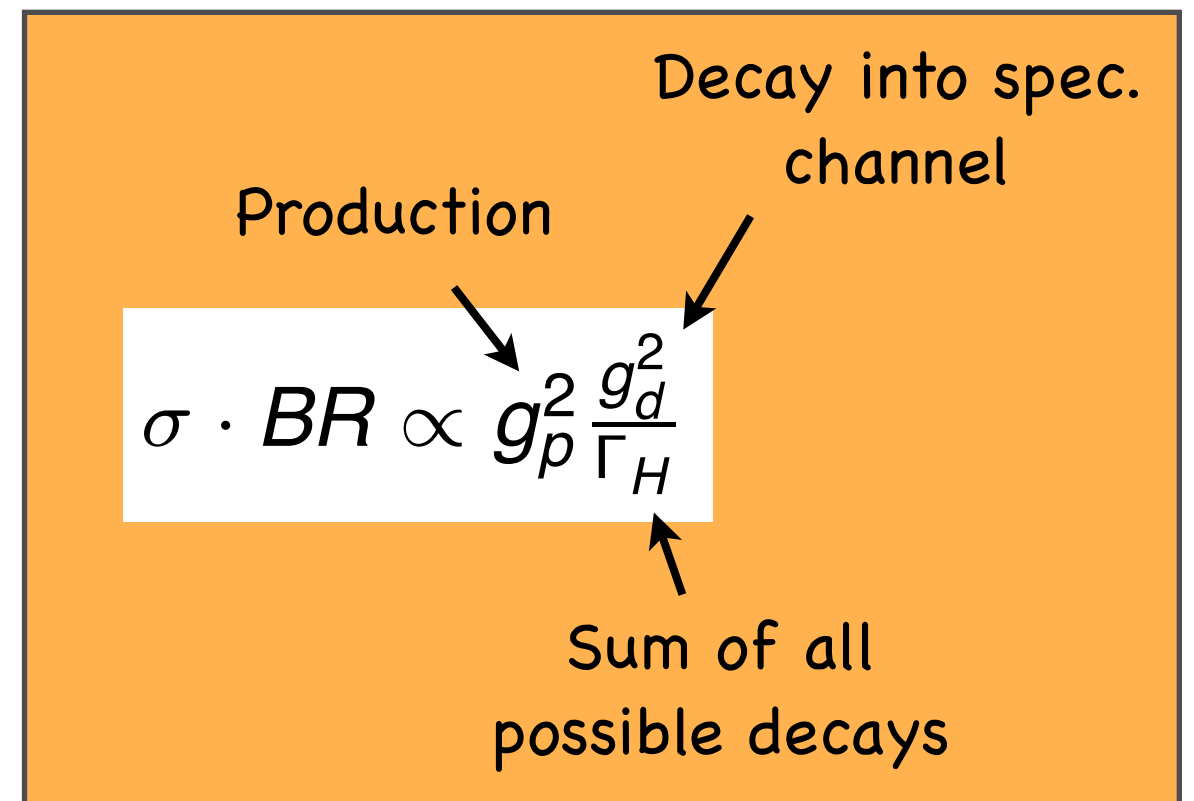
$m_H = 125 \text{ GeV}$ is smooth spot. Important couplings accessible

- If the Higgs is SM-like it has to show up in several channels

Need cross correlation between many channels:

	production	decay
	$gg \rightarrow H$	ZZ
	qqH	ZZ
	$gg \rightarrow H$	WW
	qqH	WW
	$t\bar{t}H$	$WW(3\ell)$
	$t\bar{t}H$	$WW(2\ell)$
	inclusive qqH	$\gamma\gamma$
	$t\bar{t}H$	$\gamma\gamma$
	WH	$\gamma\gamma$
	ZH	$\gamma\gamma$
	qqH	$\tau\tau(2\ell)$
	qqH	$\tau\tau(1\ell)$
	$t\bar{t}H$	$b\bar{b}$
	WH/ZH	$b\bar{b}$ (subject)

[Lafaye, Plehn, Rauch, Zerwas, Duehrssen (2009)]



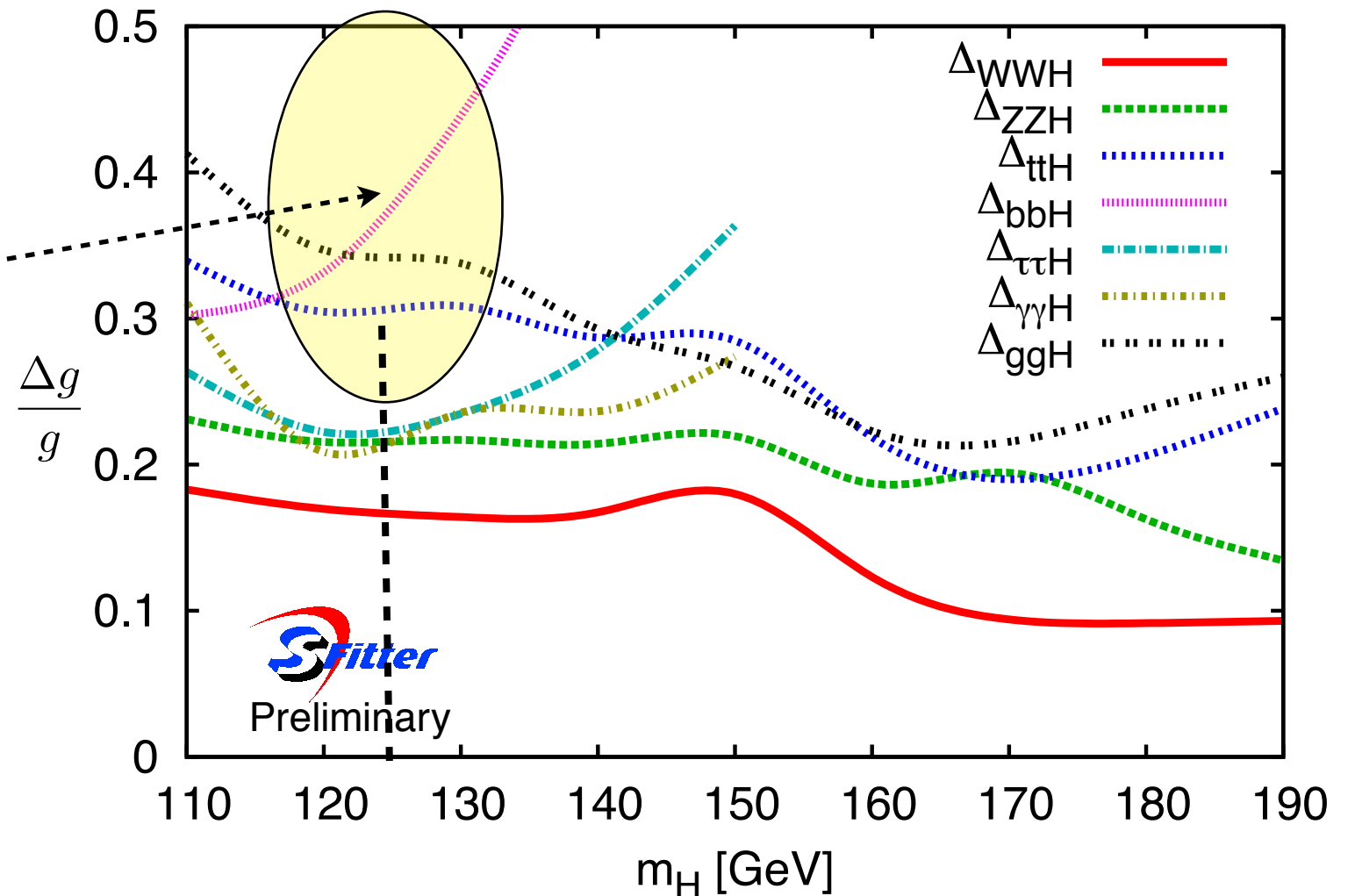
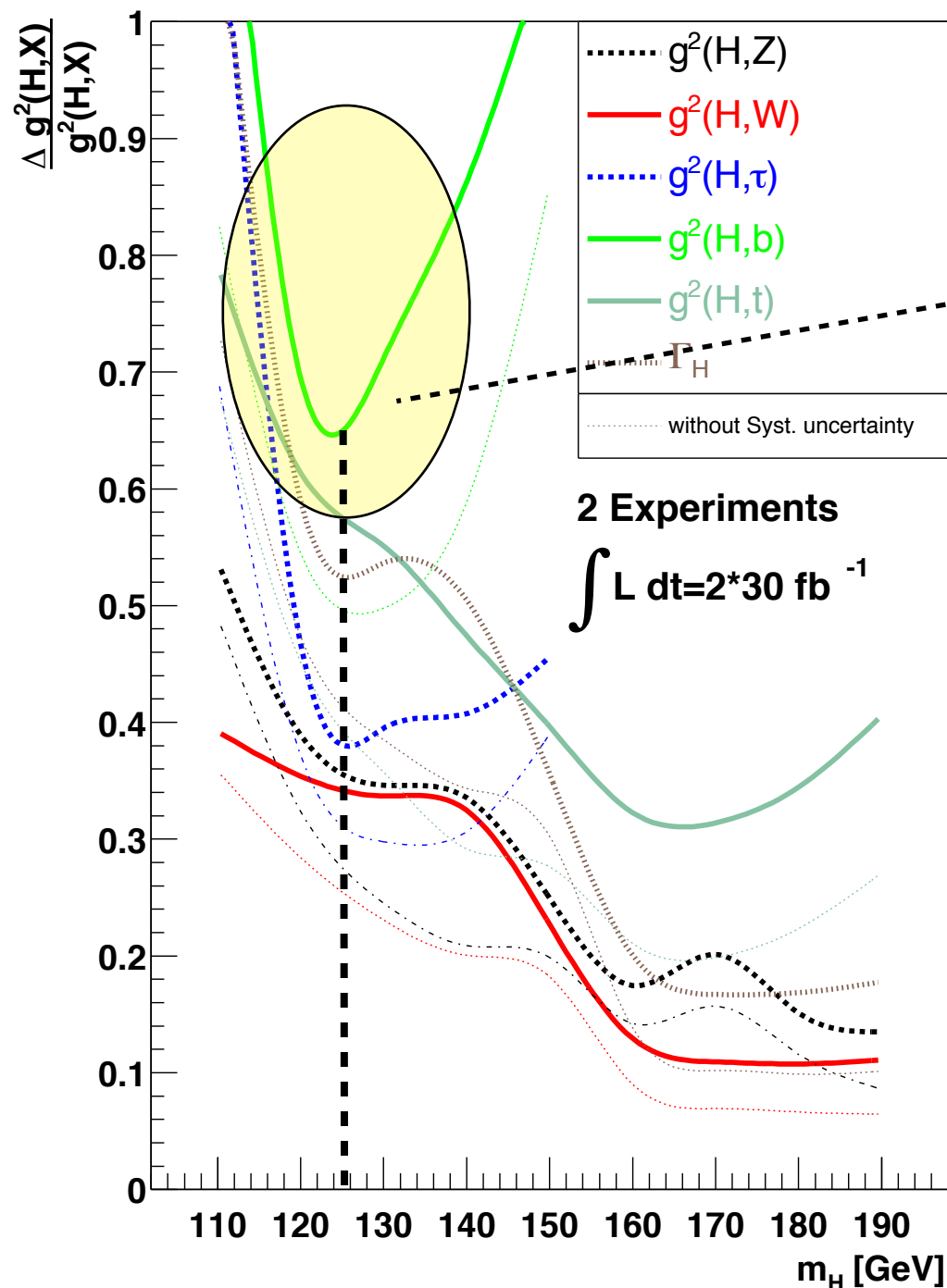
- Every measurement affected by production and decay

➔ Uncertainty of all coupling measurements driven by total width, i.e. channel with largest BR: $H \rightarrow b\bar{b}$

$Hb\bar{b}$ difficult but can use new techniques, i.e. **Jet substructure!**

[Zeppenfeld, Kinnunen, Nikitenko, Richter-Was PRD 62 (2000);
Duehrssen (2005)]

[Lafaye, Plehn, Rauch, Zerwas, Duehrssen (2009)]

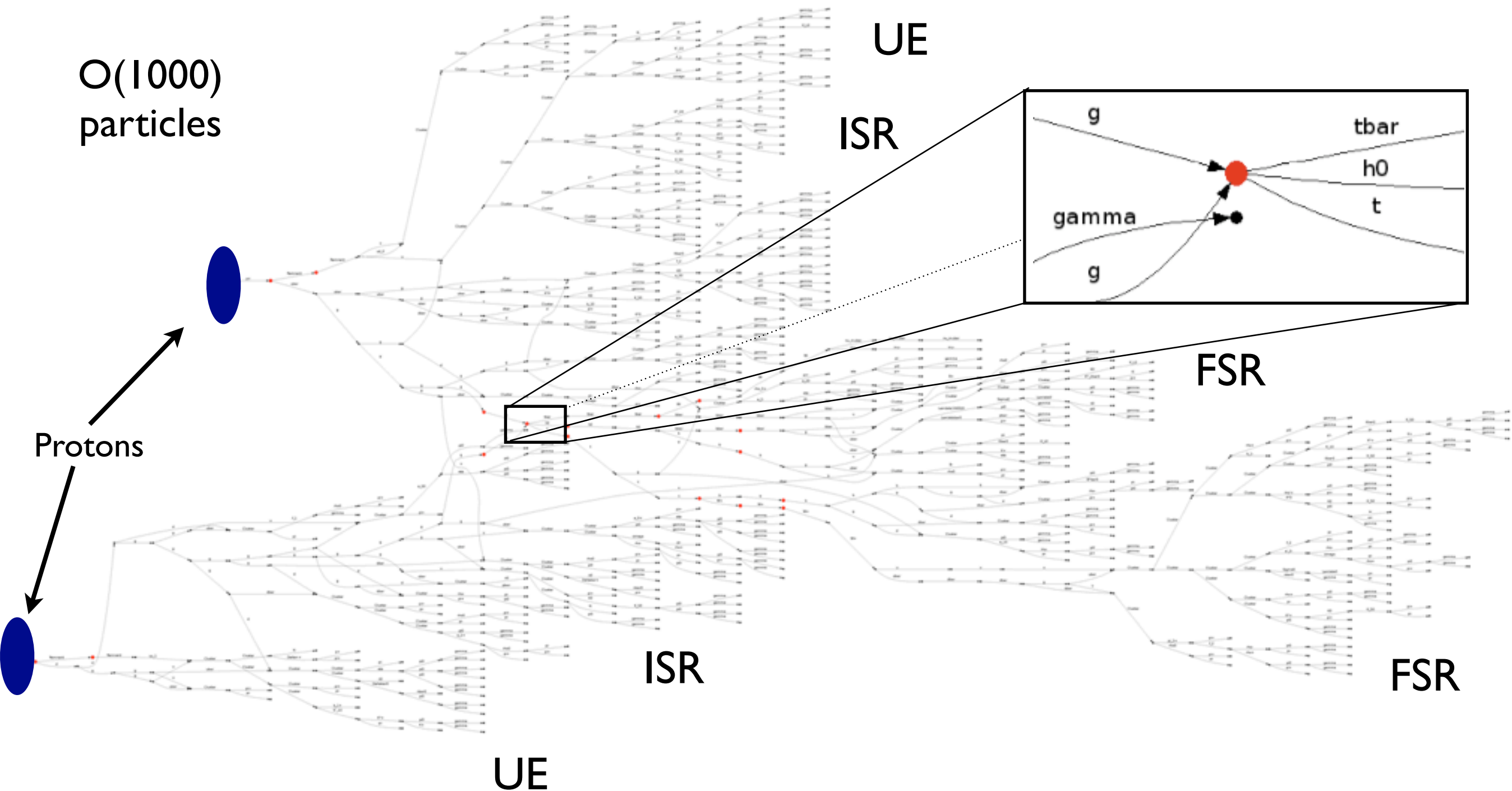


- Huge improvement from boosted Higgs analysis
- also for non-b decay modes due to better knowledge of total width



To reduce uncertainty for
all coupling, need to
measure b and t coupling

New techniques face difficult environment:



Tedious for theorists and experimentalists

Using jet substructure at LHC requires understanding of all radiation in event

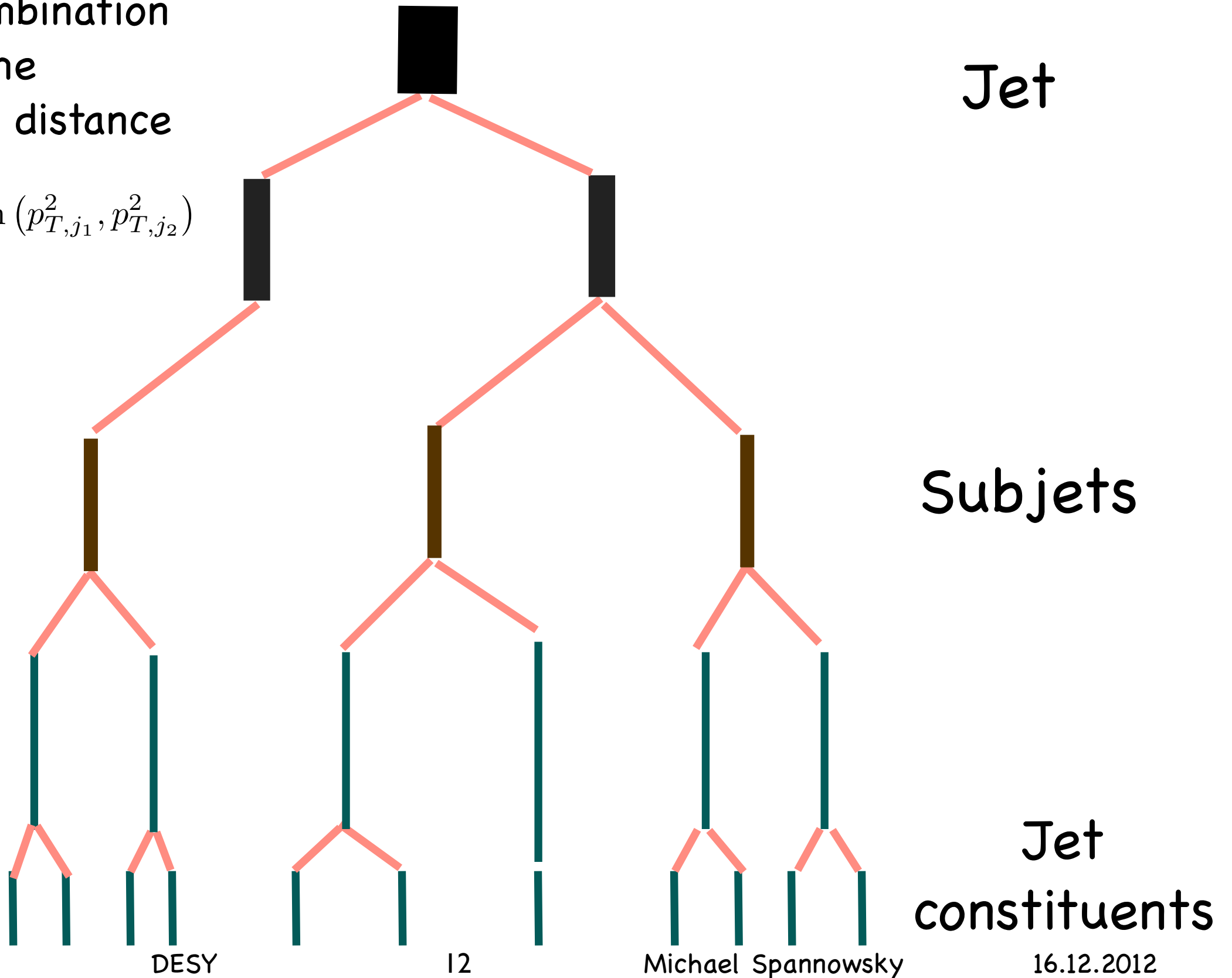
At the LHC many sources of radiation:

- Pileup $\rightarrow$ Can add up to 100 GeV of soft radiation per unit rapidity
[Cacciari, Salam, Sapeta JHEP 1004]
 - Underlying Event $\rightarrow \langle \delta m_j^2 \rangle \simeq \Lambda_{\text{UE}} p_{T,j} \left(\frac{R^4}{4} + \frac{R^8}{4608} + \mathcal{O}(R^{12}) \right)$ with $\Lambda_{\text{UE}} \sim \mathcal{O}(10)$ GeV
[Dasgupta, Magnea, Salam JHEP 0802]
 - Initial state radiation (ISR)
 - Hard radiation from many resonances in event
- $\rightarrow$ Need methods to separate final state radiation (FSR) from rest of event

Mass not a good discriminator -> look into jet substructure

order of recombination
defined by some
metric, e.g. kT distance

$$d_{j_1 j_2} = \frac{\Delta R_{j_1 j_2}^2}{D^2} \min(p_{T,j_1}^2, p_{T,j_2}^2)$$



Jet

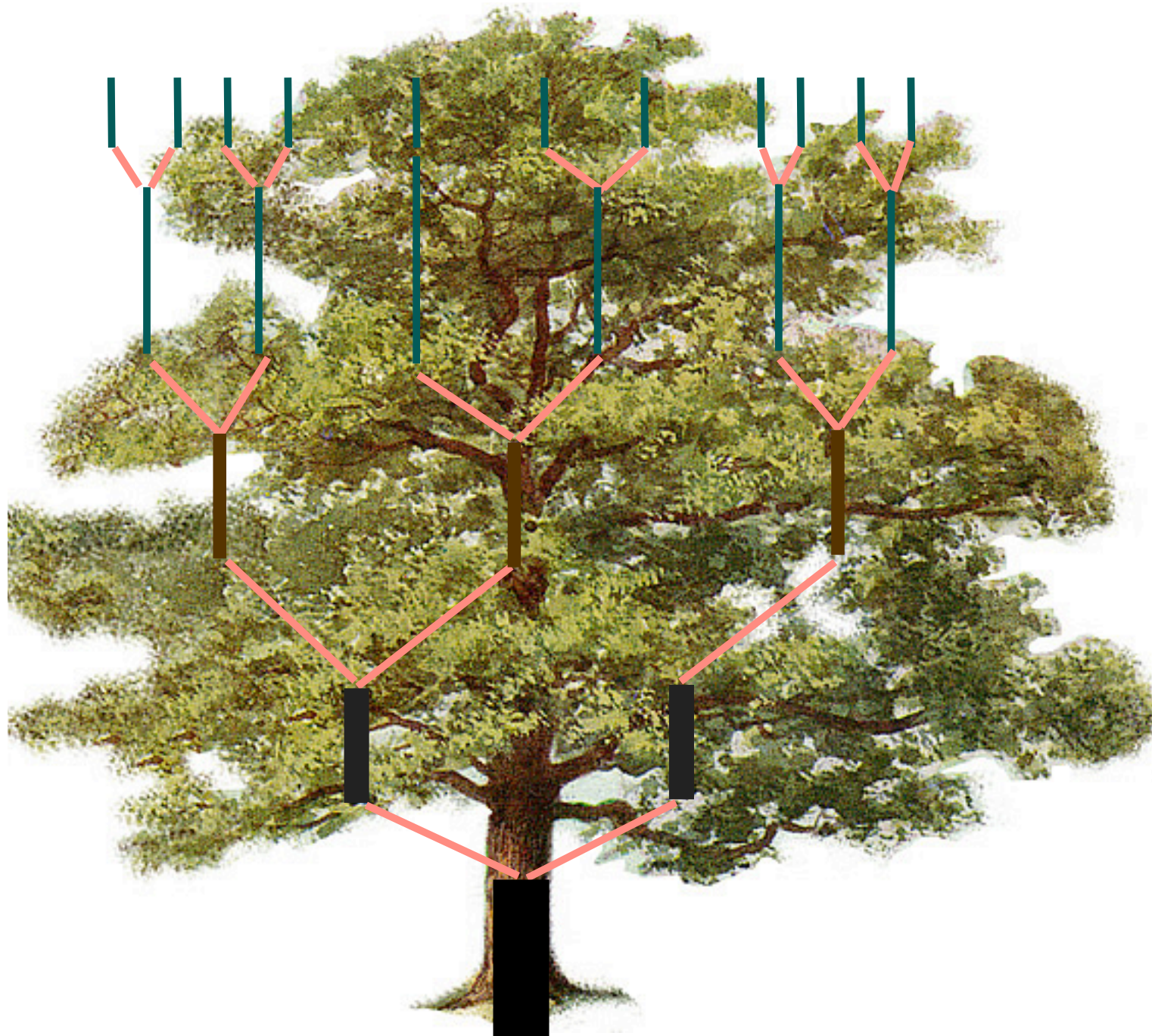
Subjects

Jet
constituents

Intrinsic scales and energy sharing of jet substructure different

→ Can be used to discriminate QCD jet from resonance jet

Higgs-boson jet



QCD jet



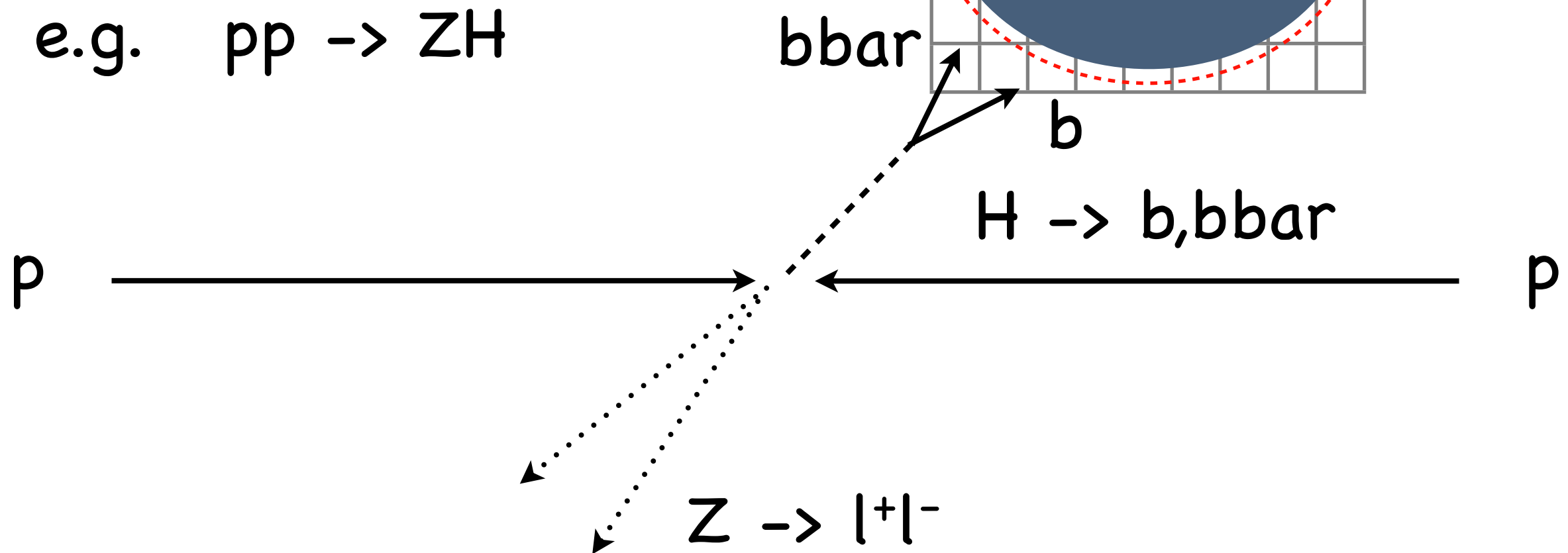
I. Measuring the Higgs-bottom coupling using jet substructure

[Butterworth, Davison, Rubin, Salam PRL 100 (2008)]

Collect FSR

Reject ISR and UE

e.g. $pp \rightarrow ZH$



Idea: [M. H. Seymour, Z. Phys. C 62, 127 (1994)]

HV - Higgs discovery channel

[Butterworth, Davison, Rubin, Salam PRL 100 (2008)]

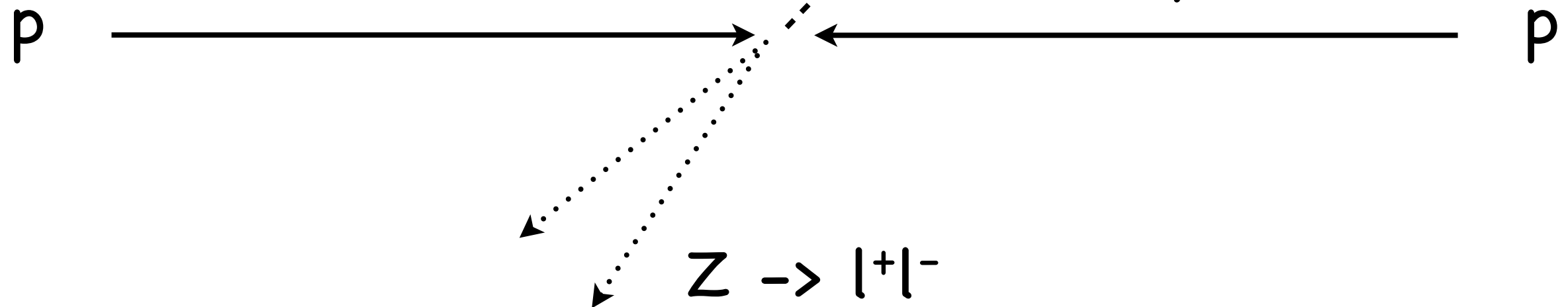
mass drop:

1) check for mass drop

$$m_{j1} < 0.66 m_j$$

2) check "asymmetry"

$$y = \frac{\min(p_{tj1}^2, p_{tj2}^2)}{m_j^2} \Delta R_{j1,j2}^2 > y_{\text{cut}}$$

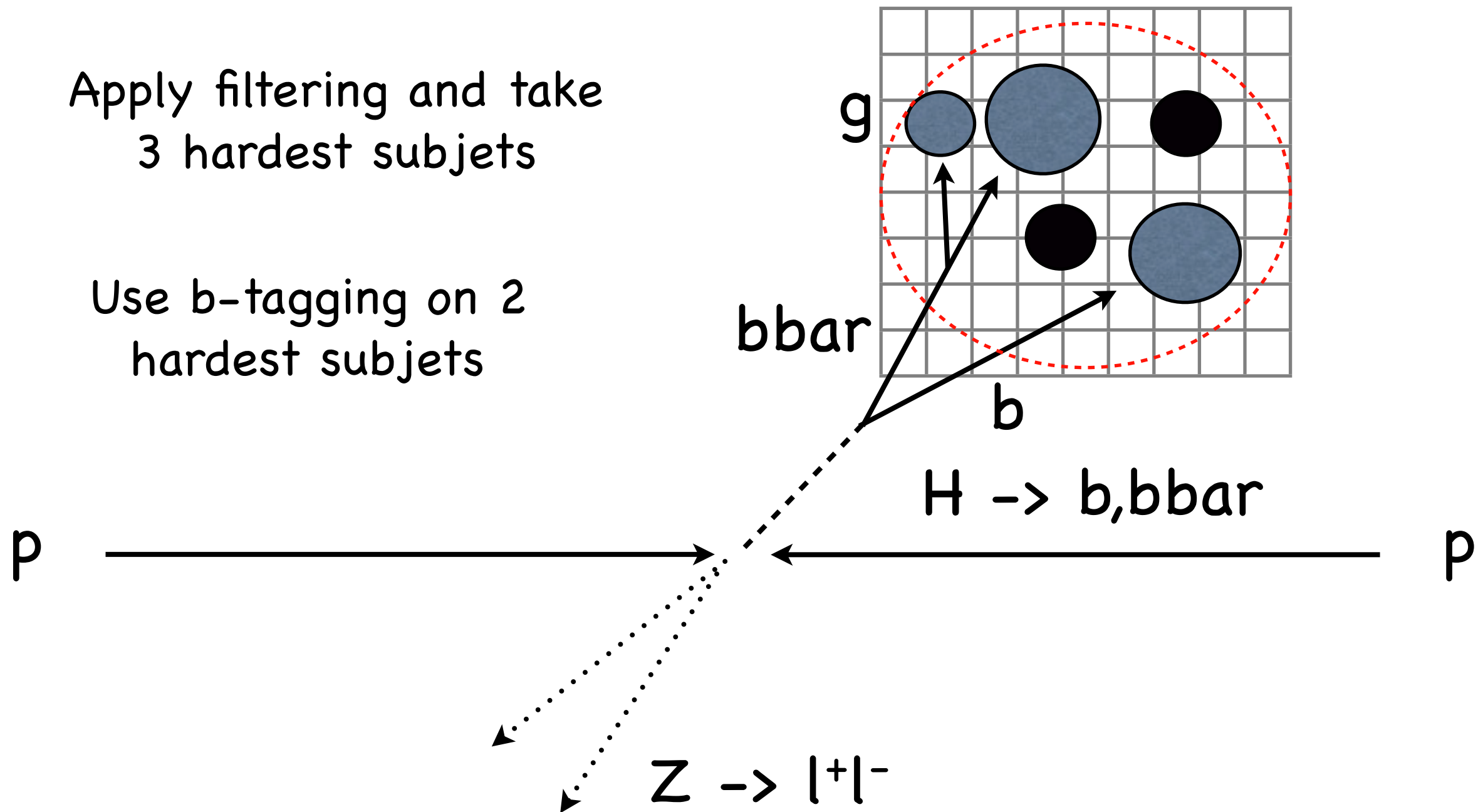


HV - Higgs discovery channel

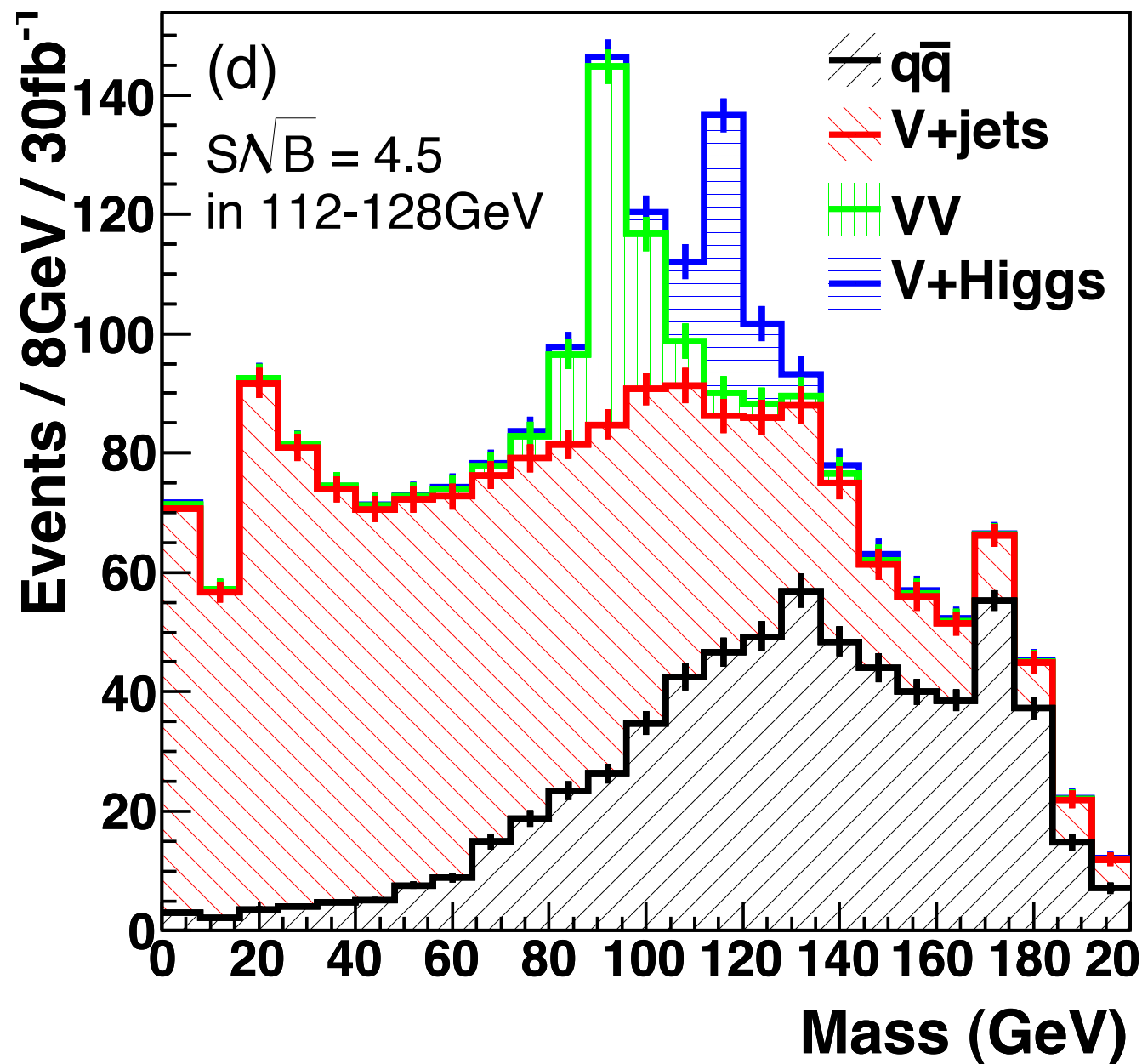
[Butterworth, Davison, Rubin, Salam PRL 100 (2008)]

Apply filtering and take
3 hardest subjets

Use b-tagging on 2
hardest subjets



BDRS Result

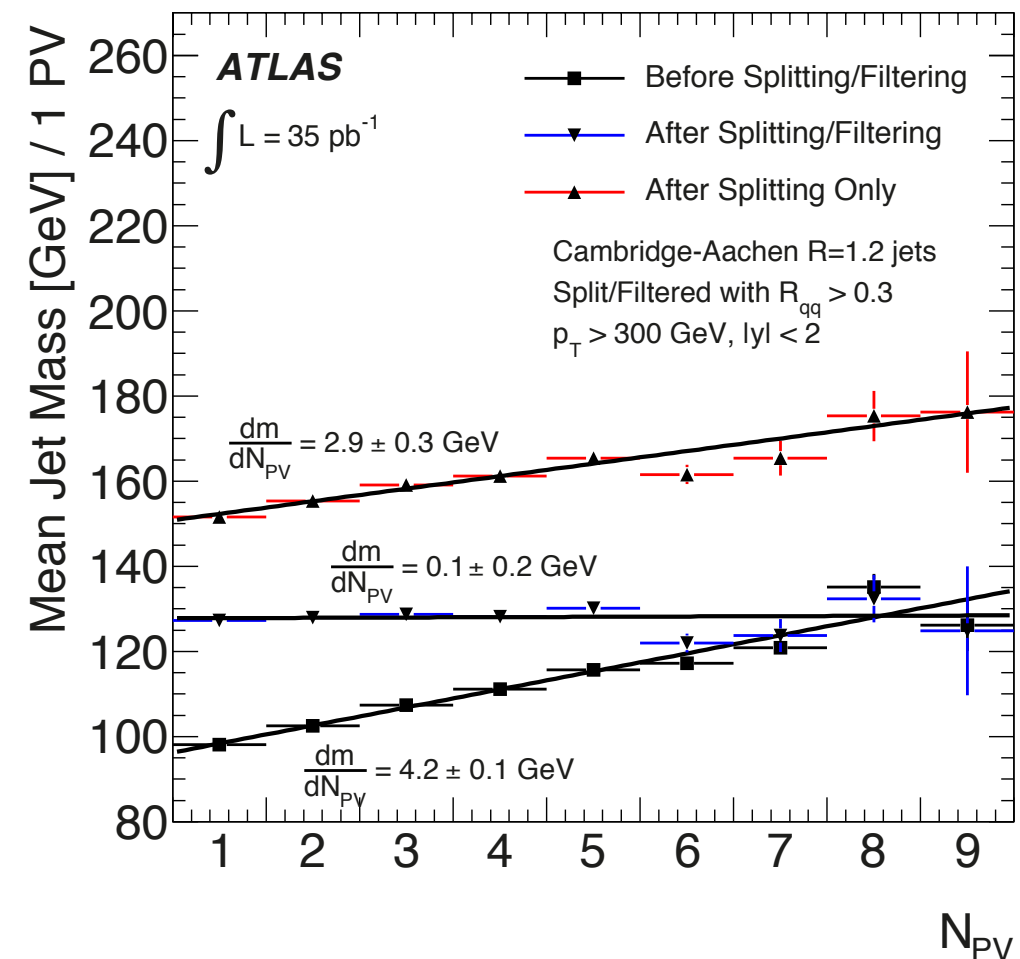
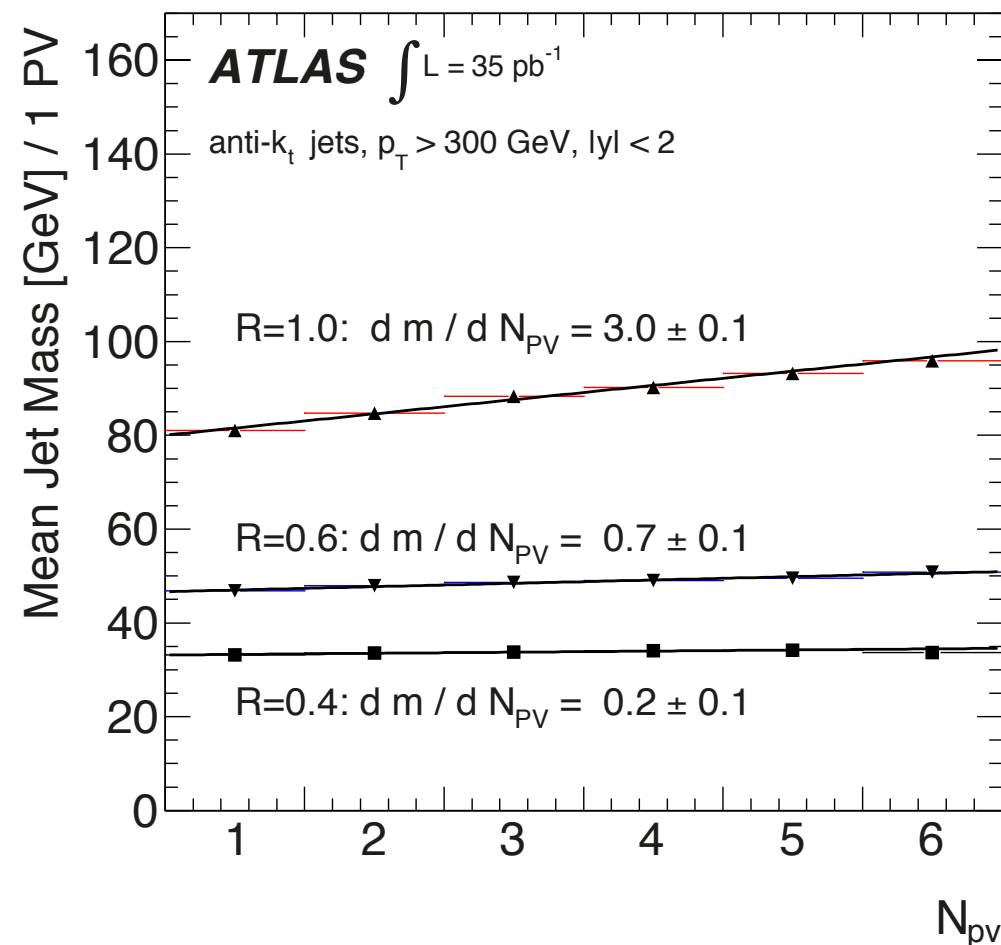


- LHC 14 TeV; 30 fb⁻¹
- HERWIG/JIMMY/Fastjet cross-checked with PYTHIA with "ATLAS tune"
- 60% b-tag; 2% mistag
- Combination of HZ and HW channels

Confirmed in ATLAS full detector simulation

First studies of method using data: (see ATLAS 1203.4606)

- Jet mass in good agreement with MC
- γ -splitter observable in good agreement with MC
- Massdrop + Filtering as predicted by MC
- Pileup under control so far:



All measurements indicate large potential for jet substructure techniques and good agreement with MC

Jet substructure methods can still be improved a lot!

Our approach:

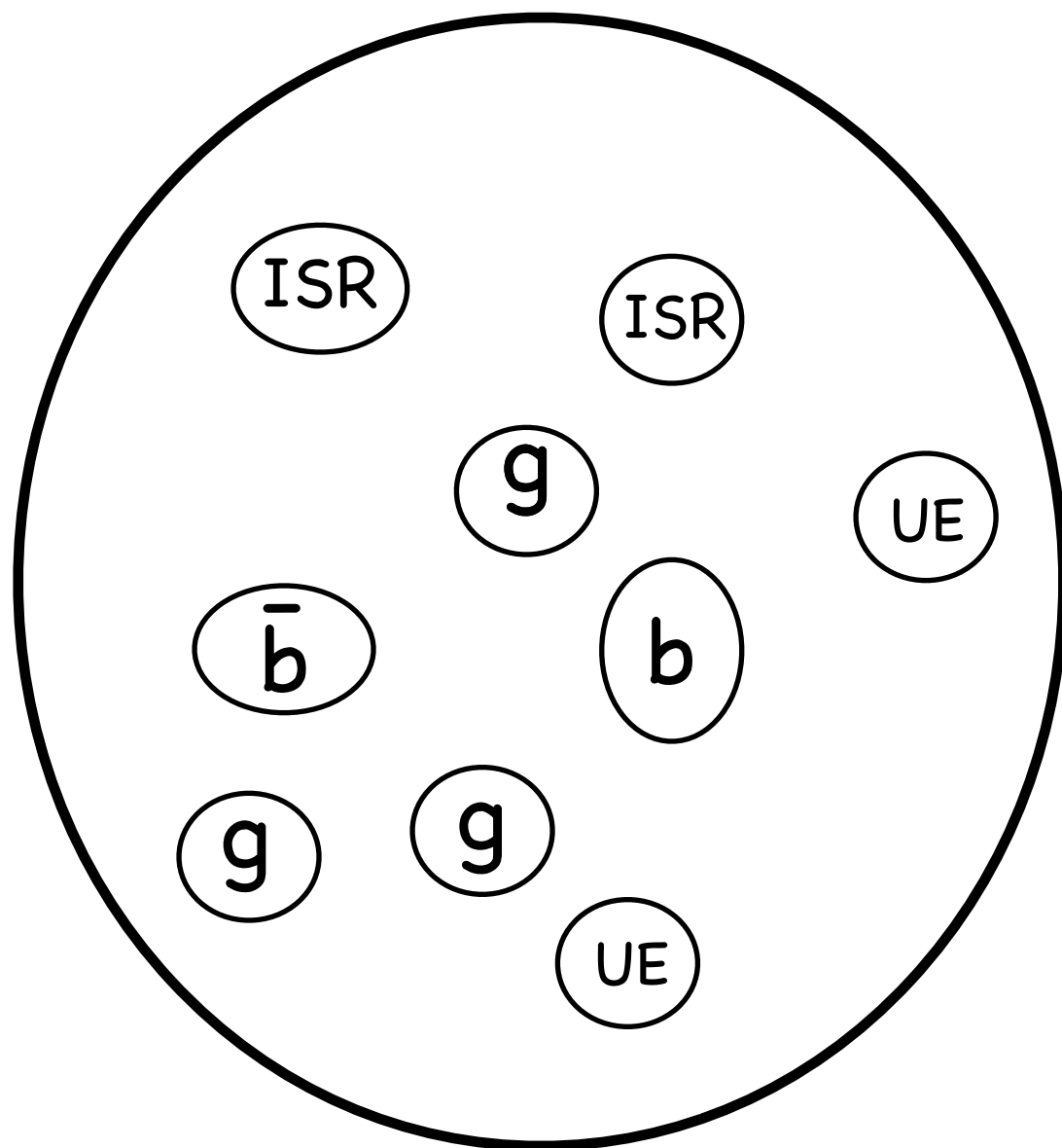
Shower deconstruction

- Maximal information approach to discriminate signal from backgrounds
→ UE, ISR, FSR, hard process
- We want one discriminating analytic function
- Have to respect experimental limitations

[Soper, MS PRD 84 (2011)]

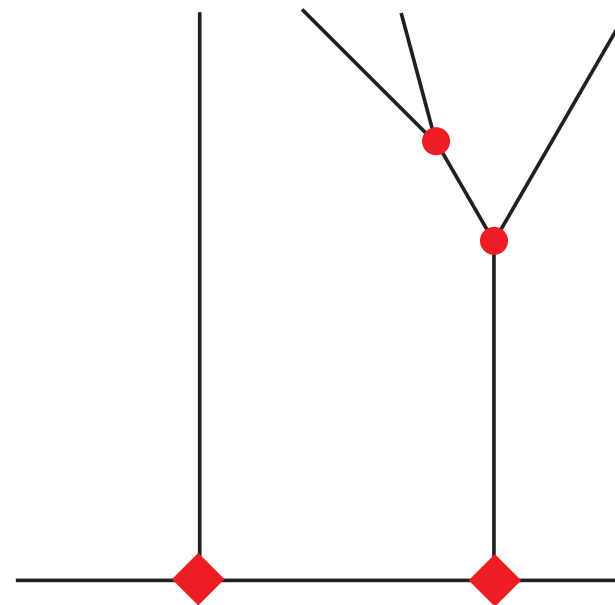
Playground: Boosted HZ final state

Fat jet: $R=1.2$, anti- k_T

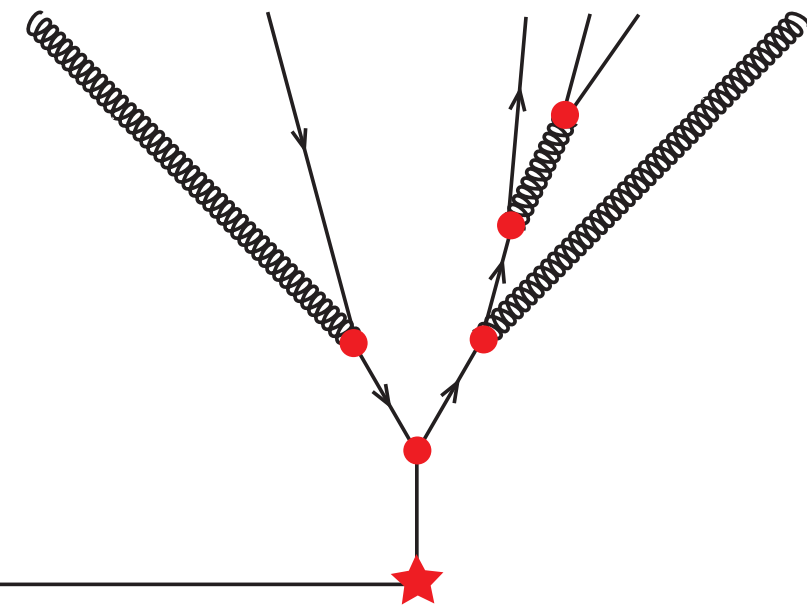


microjets
 $R=0.15$, k_T

ISR/UE



hard interaction

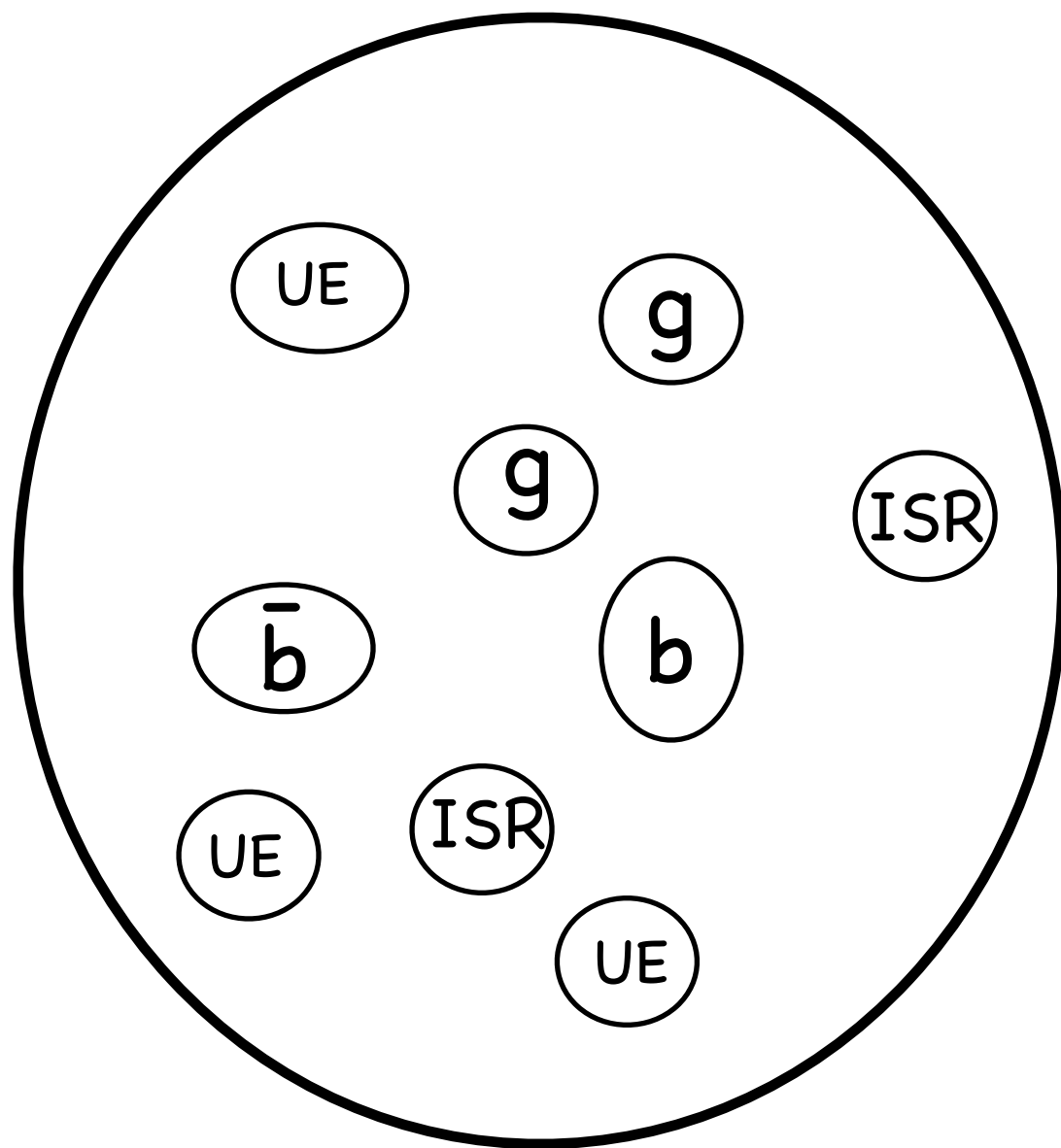


Build all possible shower histories

signal vs background hypothesis
based on:

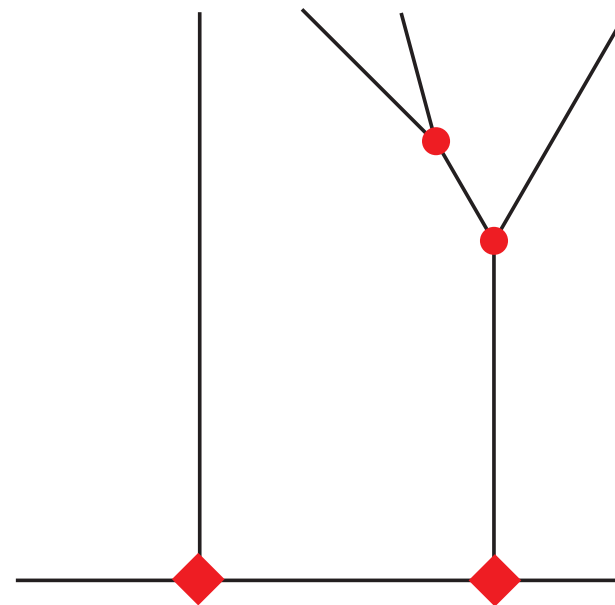
- Emission probabilities
- Color connection
- Kinematic requirements
- b -tag information

Fat jet: $R=1.2$, anti- k_T

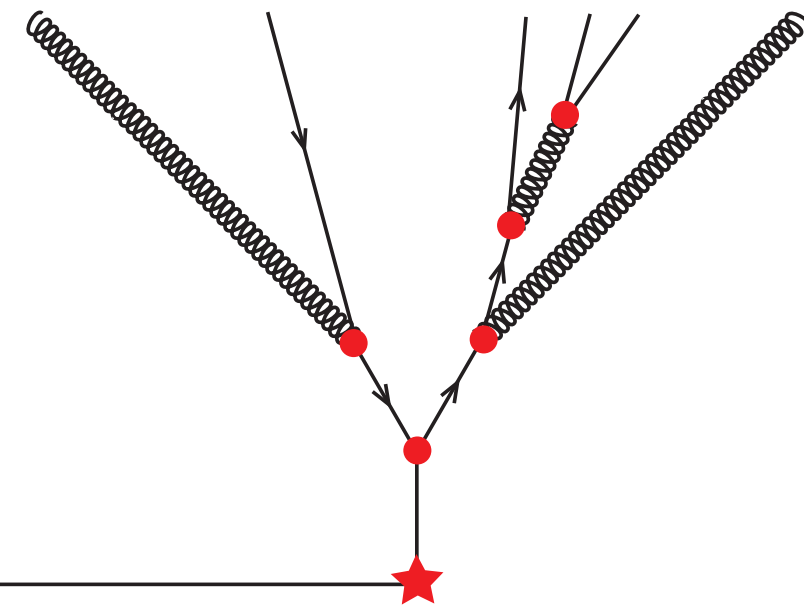


microjets
 $R=0.15$, k_T

ISR/UE



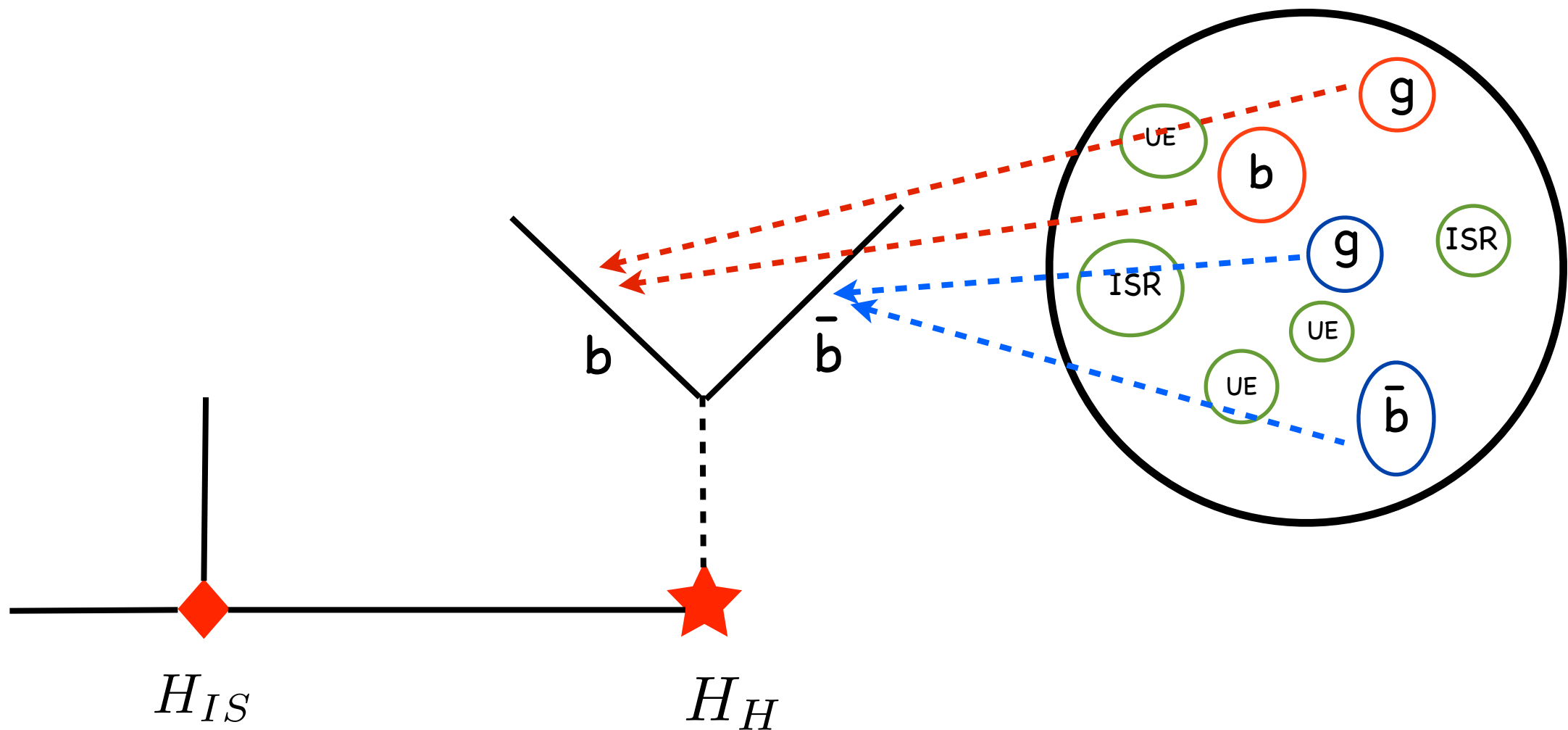
hard interaction



Build all possible shower histories

signal vs background hypothesis
based on:

- ▶ Emission probabilities
- ▶ Color connection
- ▶ Kinematic requirements
- ▶ b -tag information



Higgs has to decay:

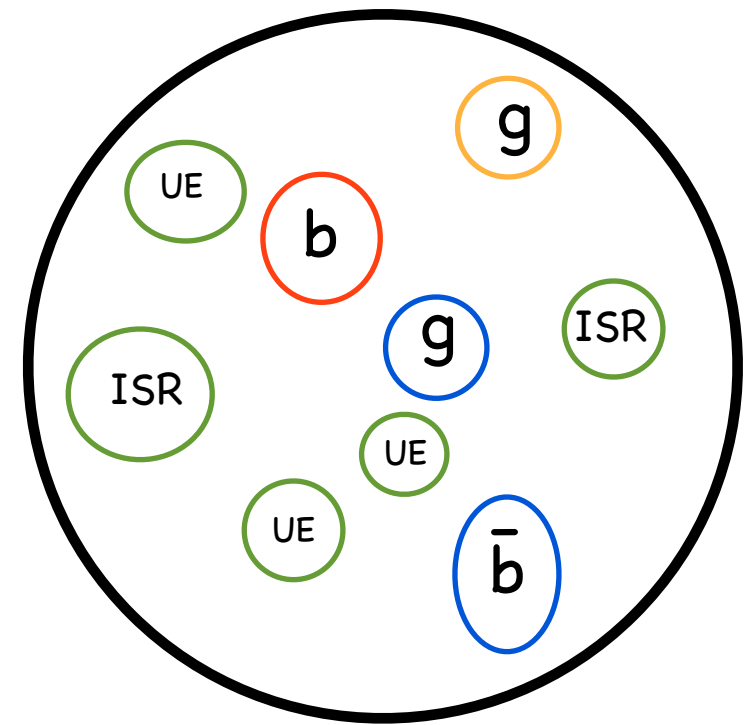
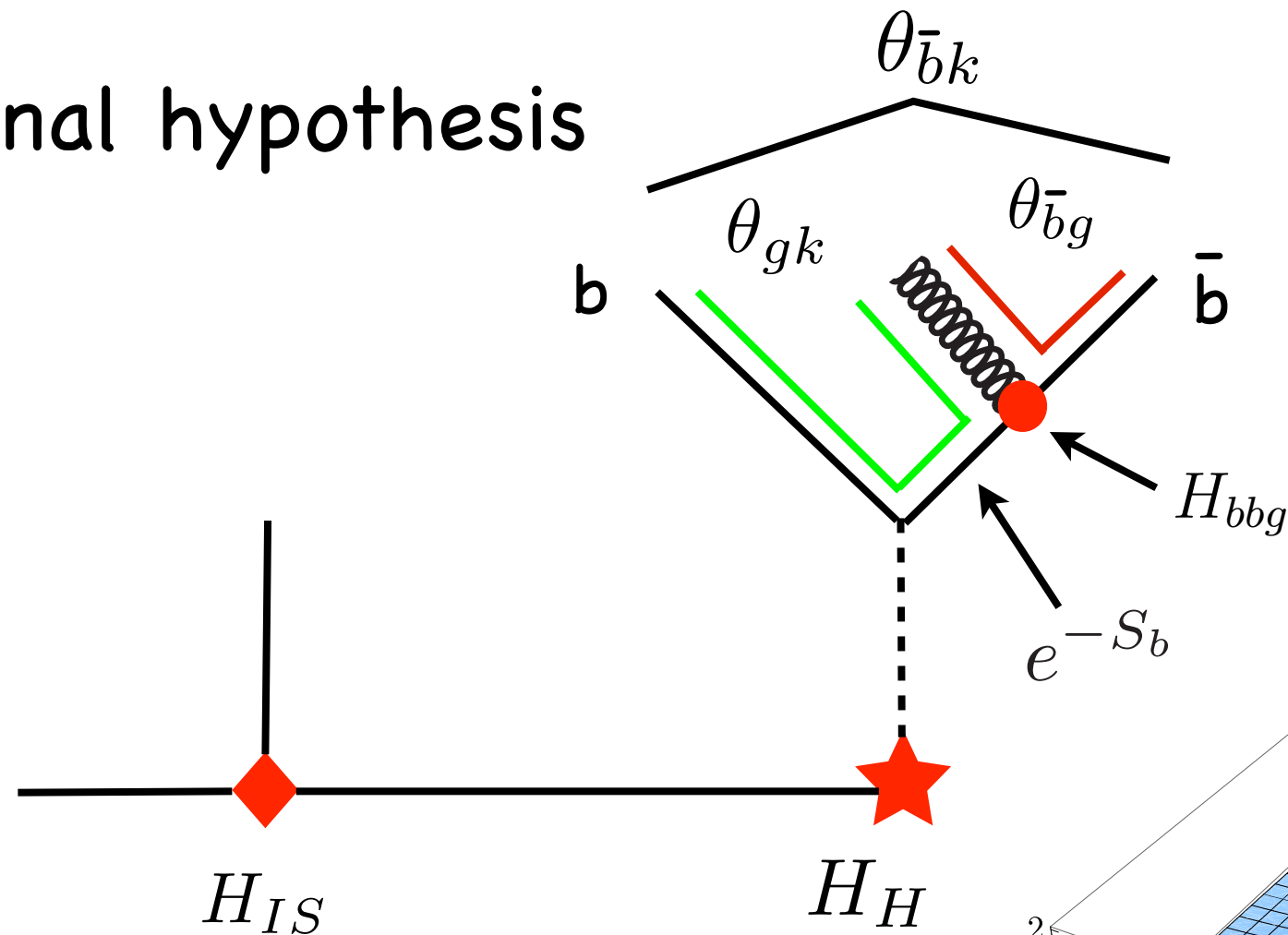
$$He^{-S} = 16\pi^2 \frac{\Theta(|m_{b\bar{b}} - m_H| < \Delta m_H)}{4m_H \Delta m_H}$$

$$\frac{1}{4(2\pi)^3} \int dm_{b\bar{b}}^2 \int dz \int d\varphi He^{-S} = 1$$

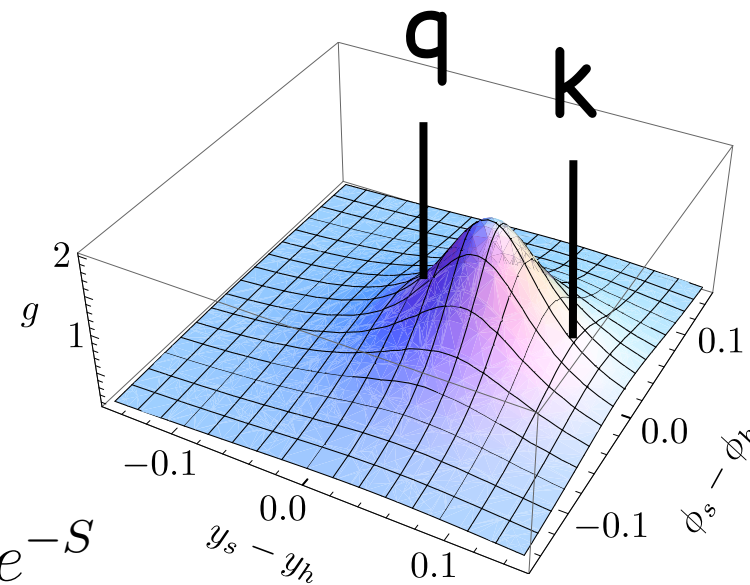
Mass window

$$\Delta m_H = 10 \text{ GeV}$$

Signal hypothesis



[Marchesini, Webber, (1984)]



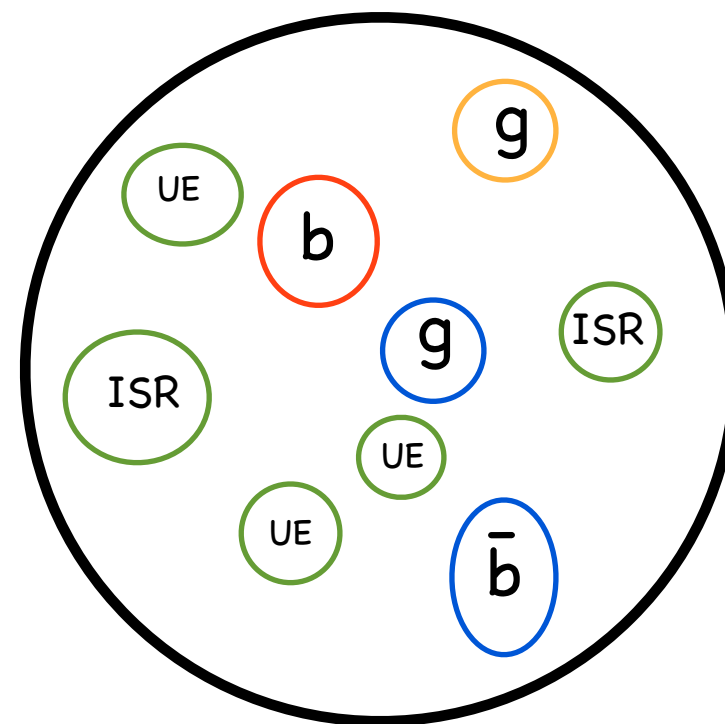
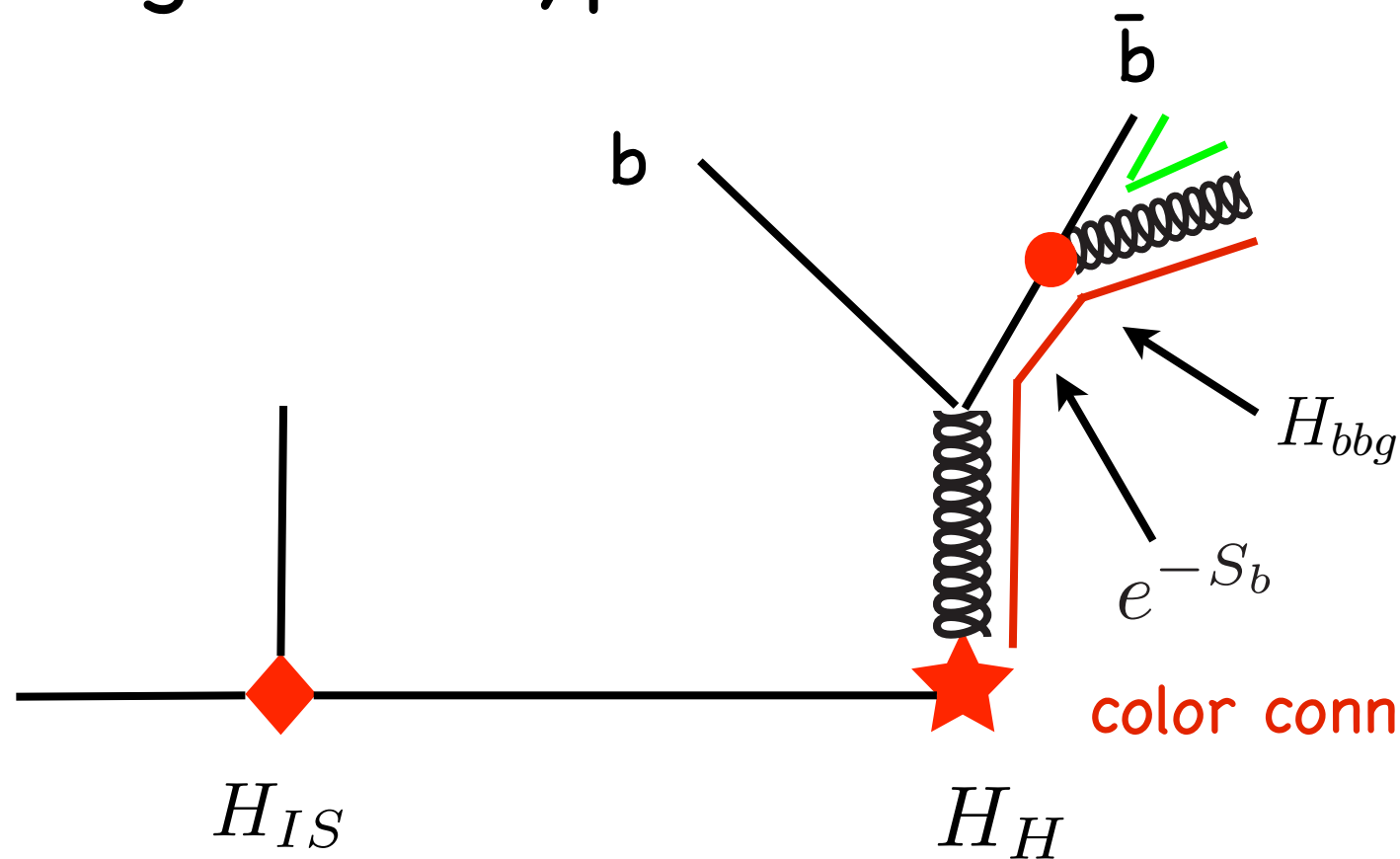
b-quarks radiate gluons

$$\int d\mathcal{P} = \int d\mu_J^2 \int d\Delta\phi \int d\Delta y \sum_s J H e^{-S}$$

$$S \approx \int d\bar{\mu}_J^2 \Theta(\mu_J^2 < \bar{\mu}_J^2) \int d\Delta\bar{y} \int d\Delta\bar{\phi} \sum_{\bar{s}} J(\bar{p}_A, \bar{p}_B) H(\bar{p}_A, \bar{p}_B) \Theta(\{\bar{p}_A, \bar{p}_B\} \in \text{fat jet})$$

$$H_{qqg} = H_{\bar{q}g\bar{q}} = 8\pi C_F \frac{\alpha_s(\mu_J^2)}{\mu_J^2} \frac{k_J}{k_g} \left[1 + \left(\frac{k_q}{k_J} \right)^2 \right] \frac{\theta_{qk}^2}{\theta_{gq}^2 + \theta_{gk}^2} \Theta\left(2 \frac{\mu_J^2}{k_J} < \frac{\mu_K^2}{k_K} \right)$$

Background hypothesis



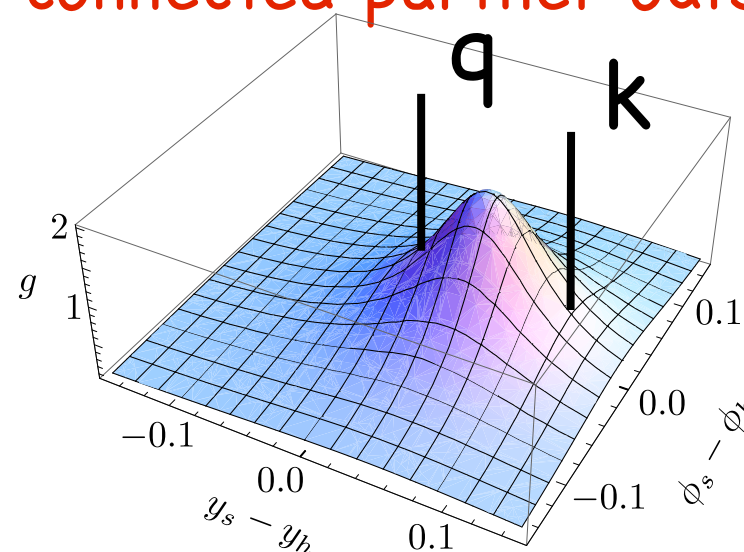
color connected partner outside fat jet

b-quarks radiate gluons

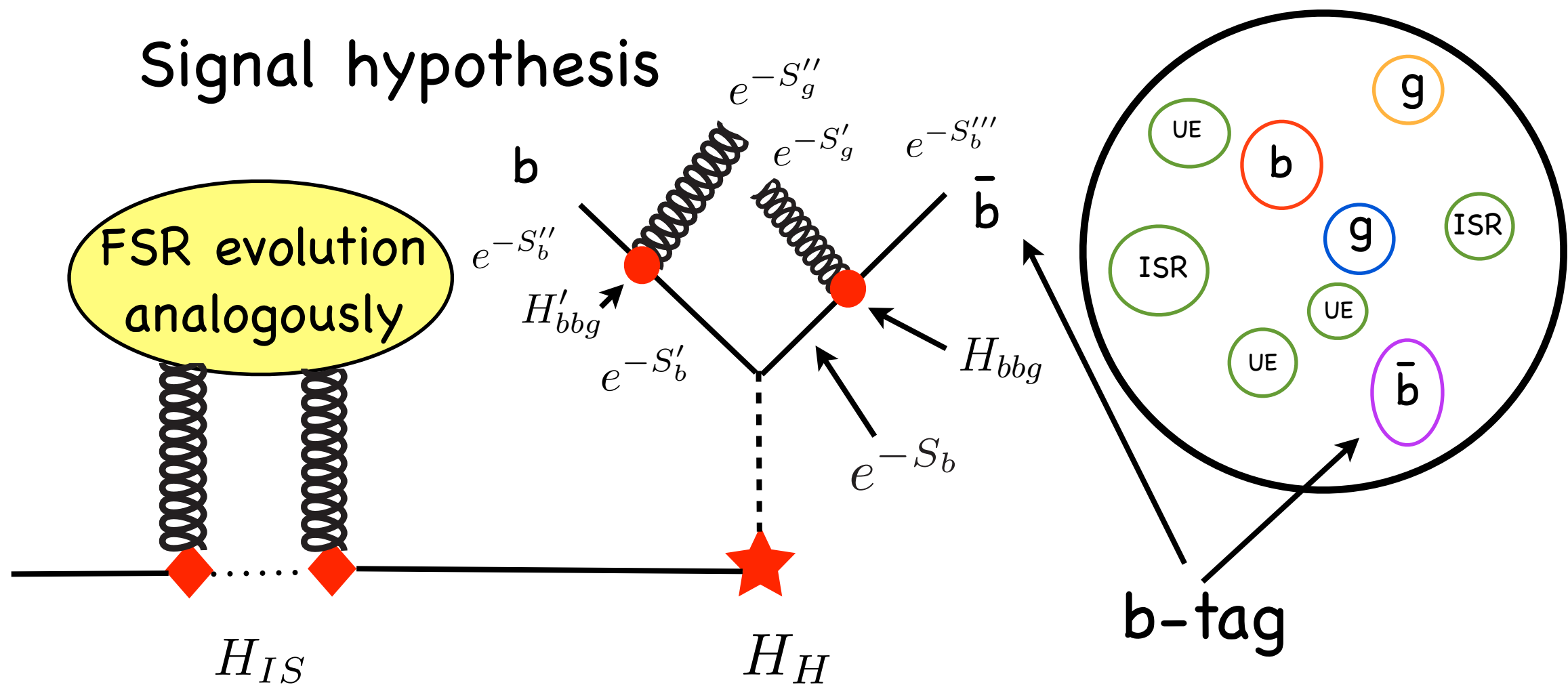
$$\int d\mathcal{P} = \int d\mu_J^2 \int d\Delta\phi \int d\Delta y \sum_s J H e^{-S}$$

$$S \approx \int d\bar{\mu}_J^2 \Theta(\mu_J^2 < \bar{\mu}_J^2) \int d\Delta \bar{y} \int d\Delta \bar{\phi} \sum_{\bar{s}} J(\bar{p}_A, \bar{p}_B) H(\bar{p}_A, \bar{p}_B) \Theta(\{\bar{p}_A, \bar{p}_B\} \in \text{fat jet})$$

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Signal hypothesis



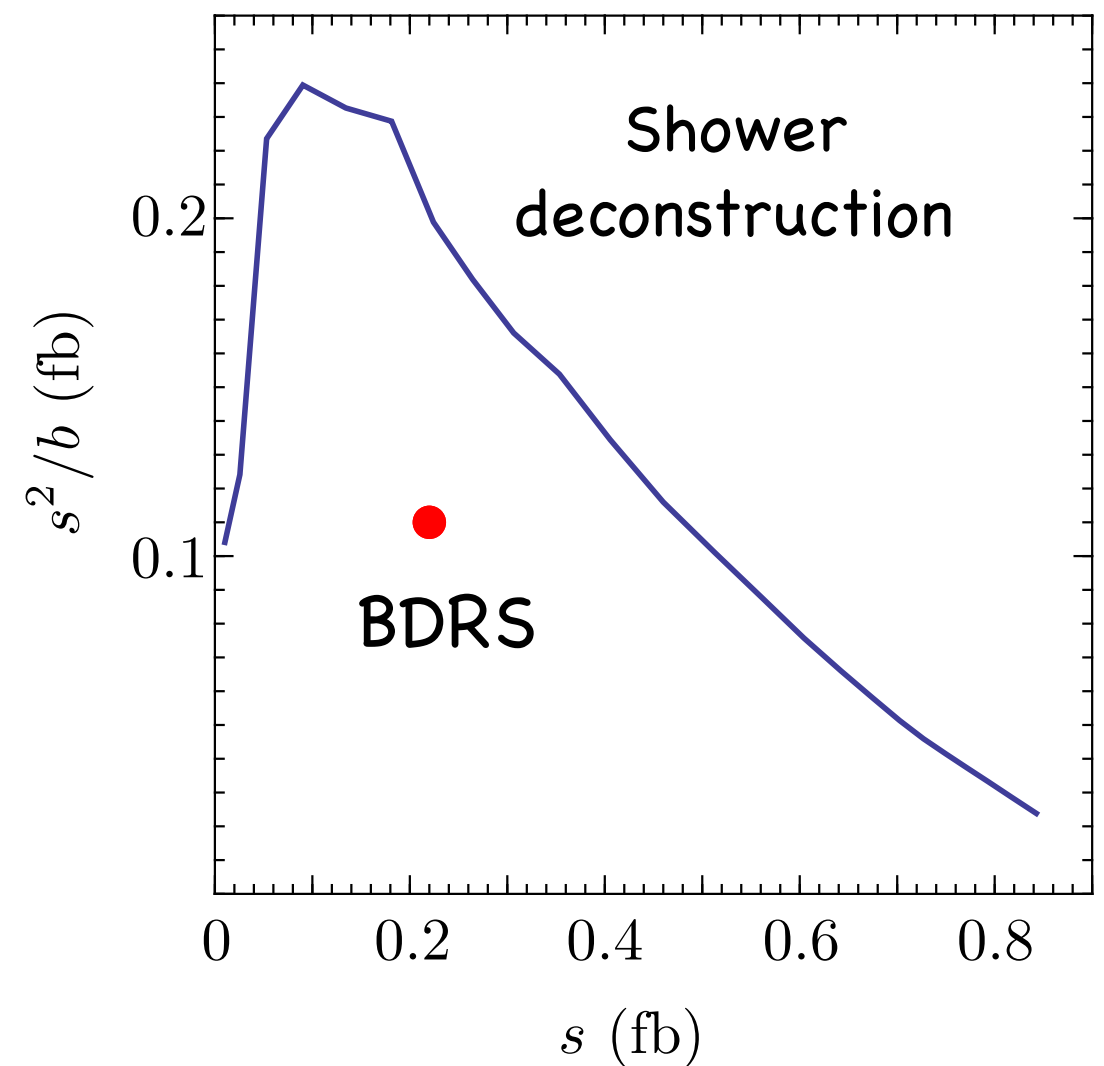
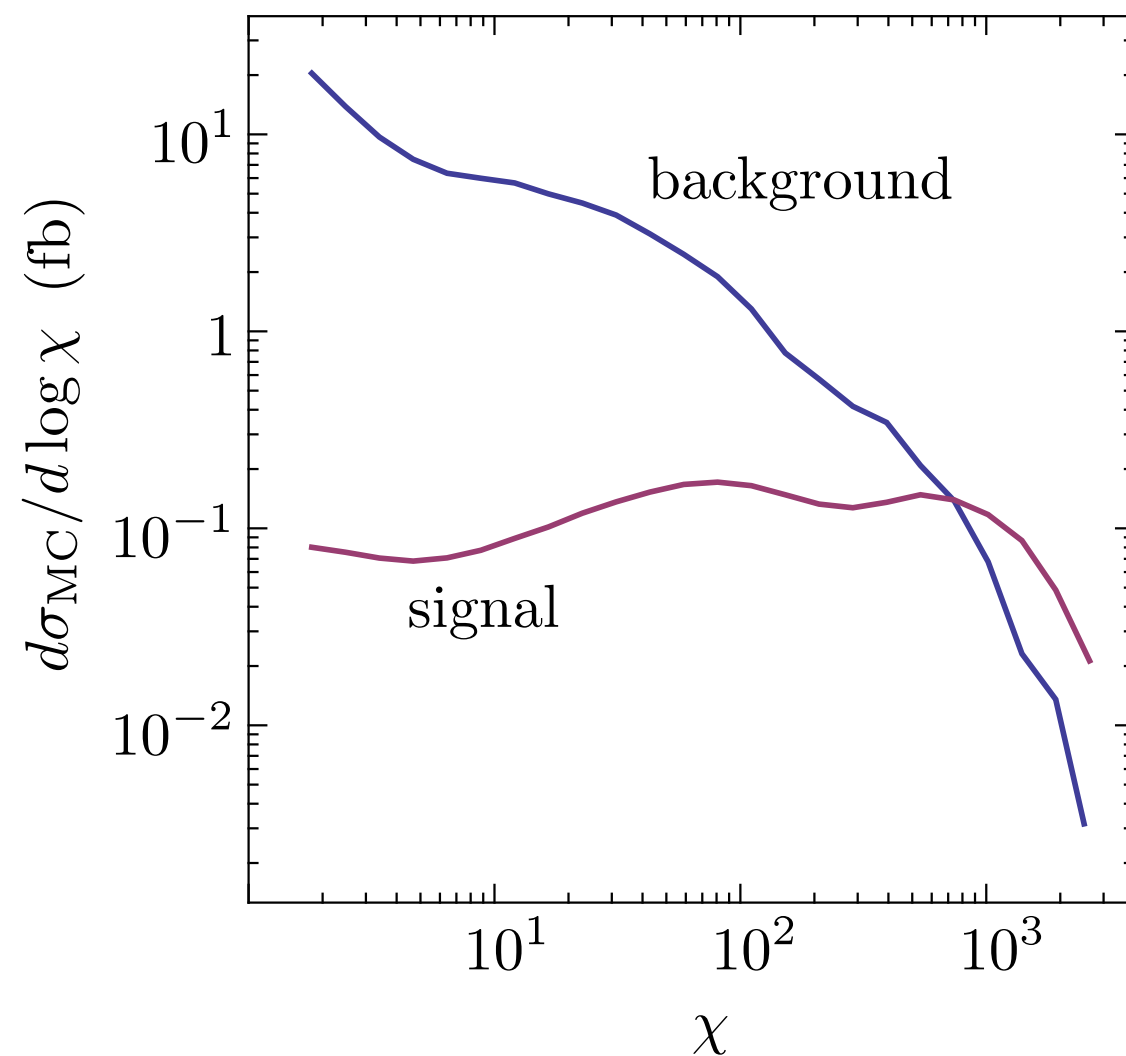
Wrapping up all factors gives weight for shower history

$$\chi = \frac{\sum_{ISR/Hard} \left(\sum_i ISR_i \times \sum_j Signal_j \right)}{\sum_{ISR/Hard} \left(\sum_i ISR_i \times \sum_j Backg_j \right)}$$

Here $Signal_1 = H_H H_{split} e^{-S_{split}} H_{bbg} e^{-S'_b} e^{-S''_b} e^{-S'_g} H'_{bbg} e^{-S''_b} e^{-S_b} e^{-S''_g}$

Results of shower deconstruction (SD)

$$\chi(\{p, t\}_N) = \frac{P(\{p, t\}_N | S)}{P(\{p, t\}_N | B)}$$



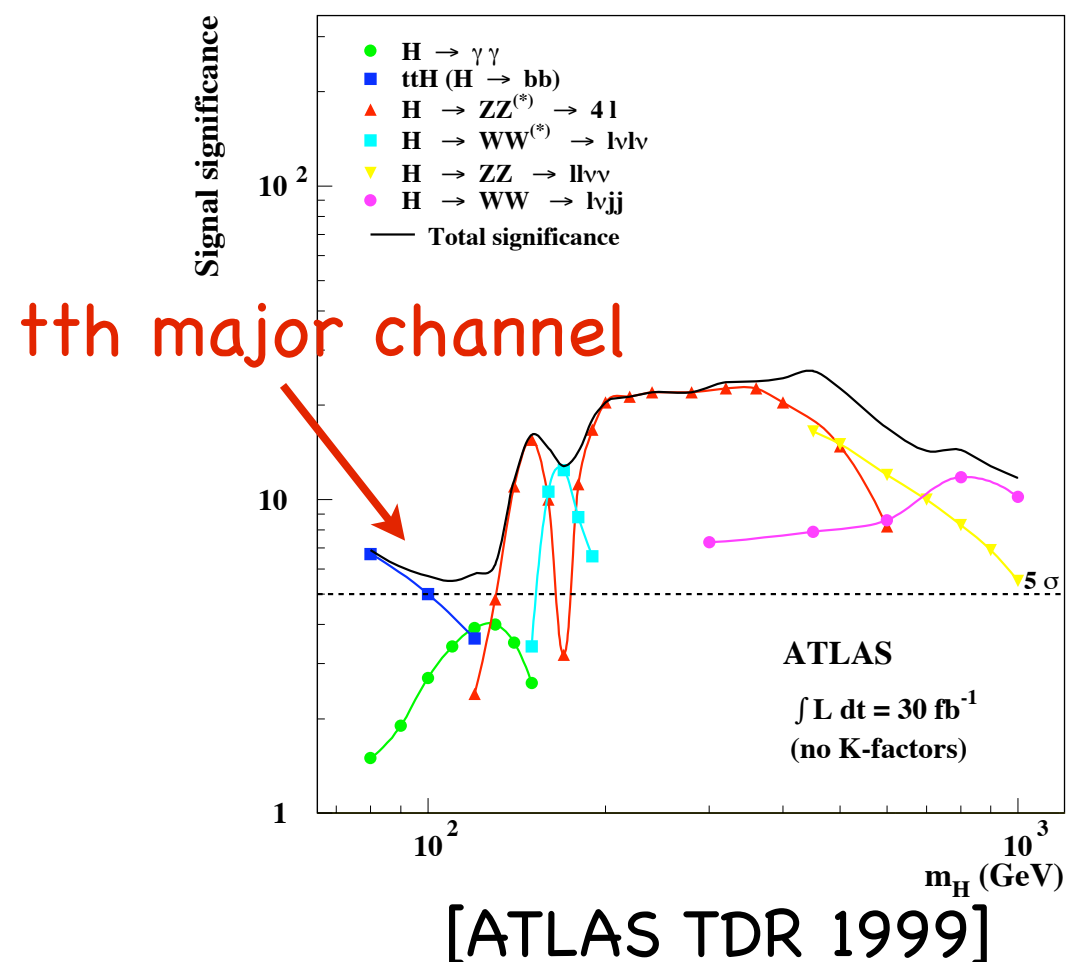
imperfect b-tagging (60%,2%) no b-tag required

II. Measuring the Higgs-top coupling using boosted techniques

[Plehn, Salam, MS PRL 104 (2010)]

- Motivation:
- sizable cross-section
 - Higgs discovery contribution in low mass range
 - access to t - and b -Yukawa couplings

High expectations:



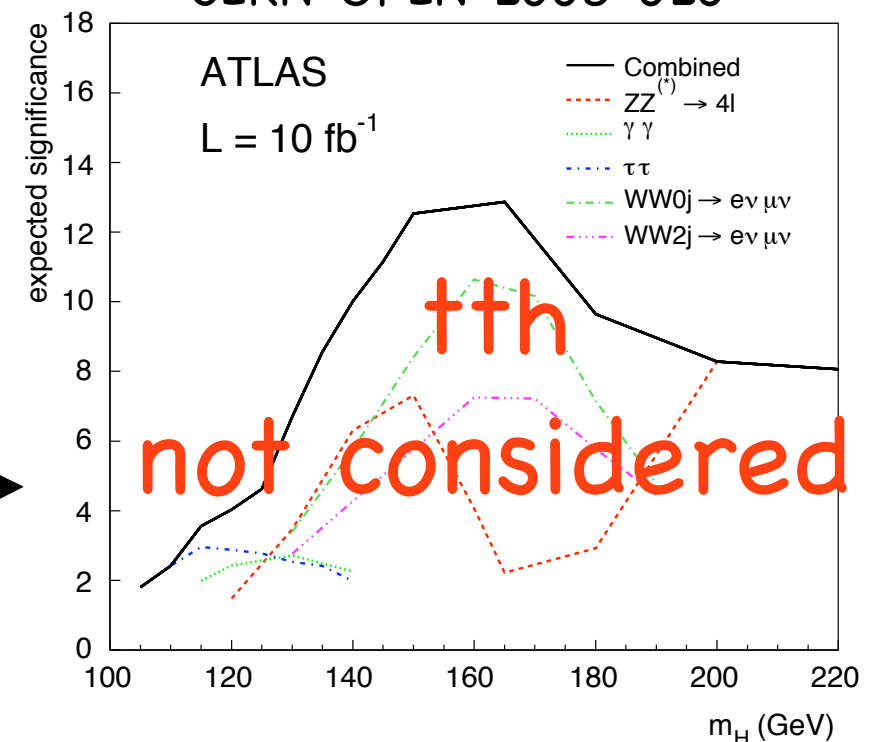
Cammin
and
Schumacher
(ATLAS)

$$S/B \simeq 1/9$$

$$S/\sqrt{B} \simeq 2.2$$



Expected Performance of the
ATLAS Experiment,
CERN-OPEN-2008-020



$t\bar{t}h$ (Signal)

Beenakker et al.,
PRL 87 2001;
Reina et al.,
PRD 65 2002

→ $K=1.57$

$t\bar{t}b\bar{b}$

Bredenstein et al.,
PRL 103 2009;
Belivacqua et al.,
JHEP 0909 2009

→ $K=2.3$

$t\bar{t}+\text{jets}$

Dittmaier et al.,
PRL 98 2007
Bevilacqua et al.,
PRL 104 2010

→ $K=1.0$

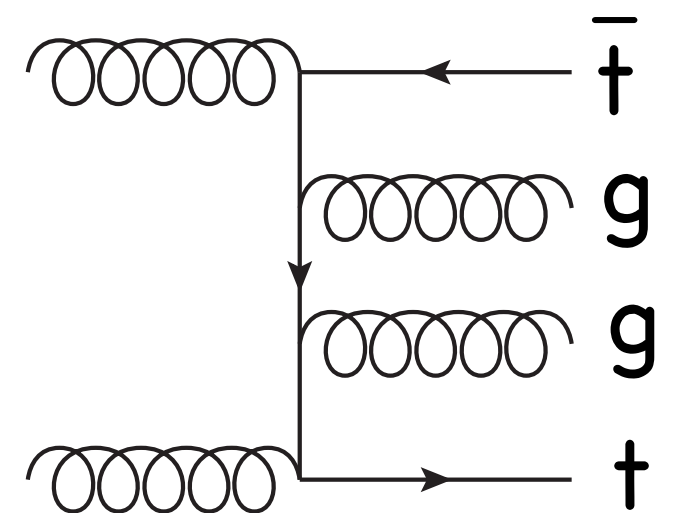
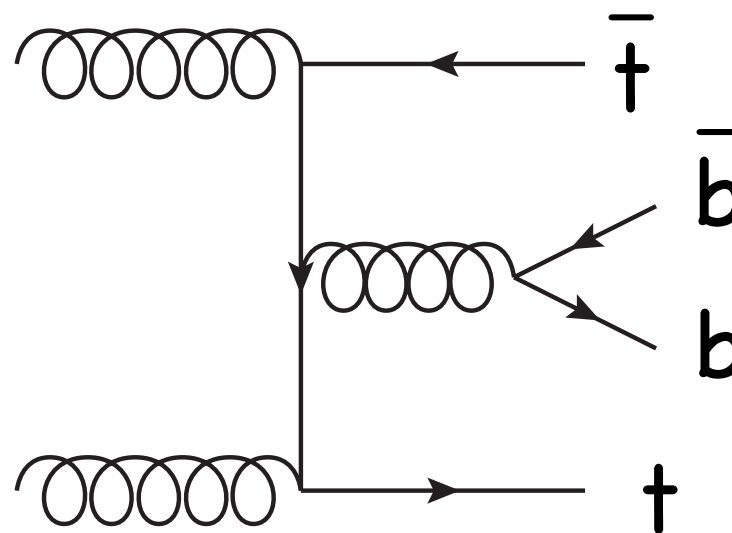
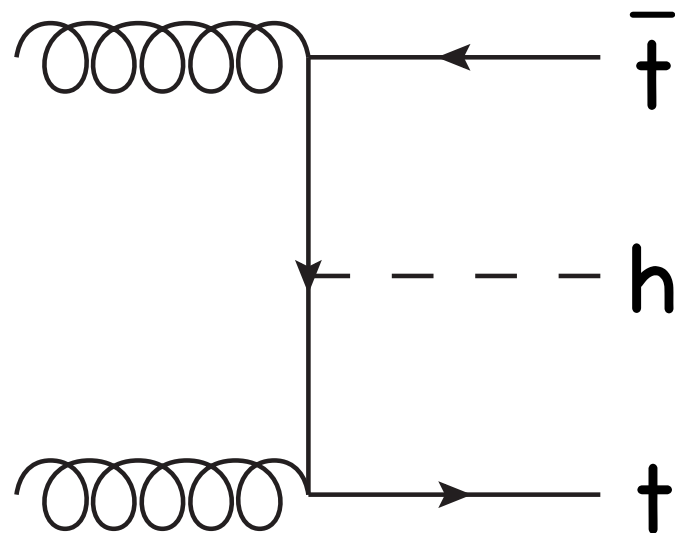
$t\bar{t}z$

Lazopoulos et al.,
PLB 666 2008

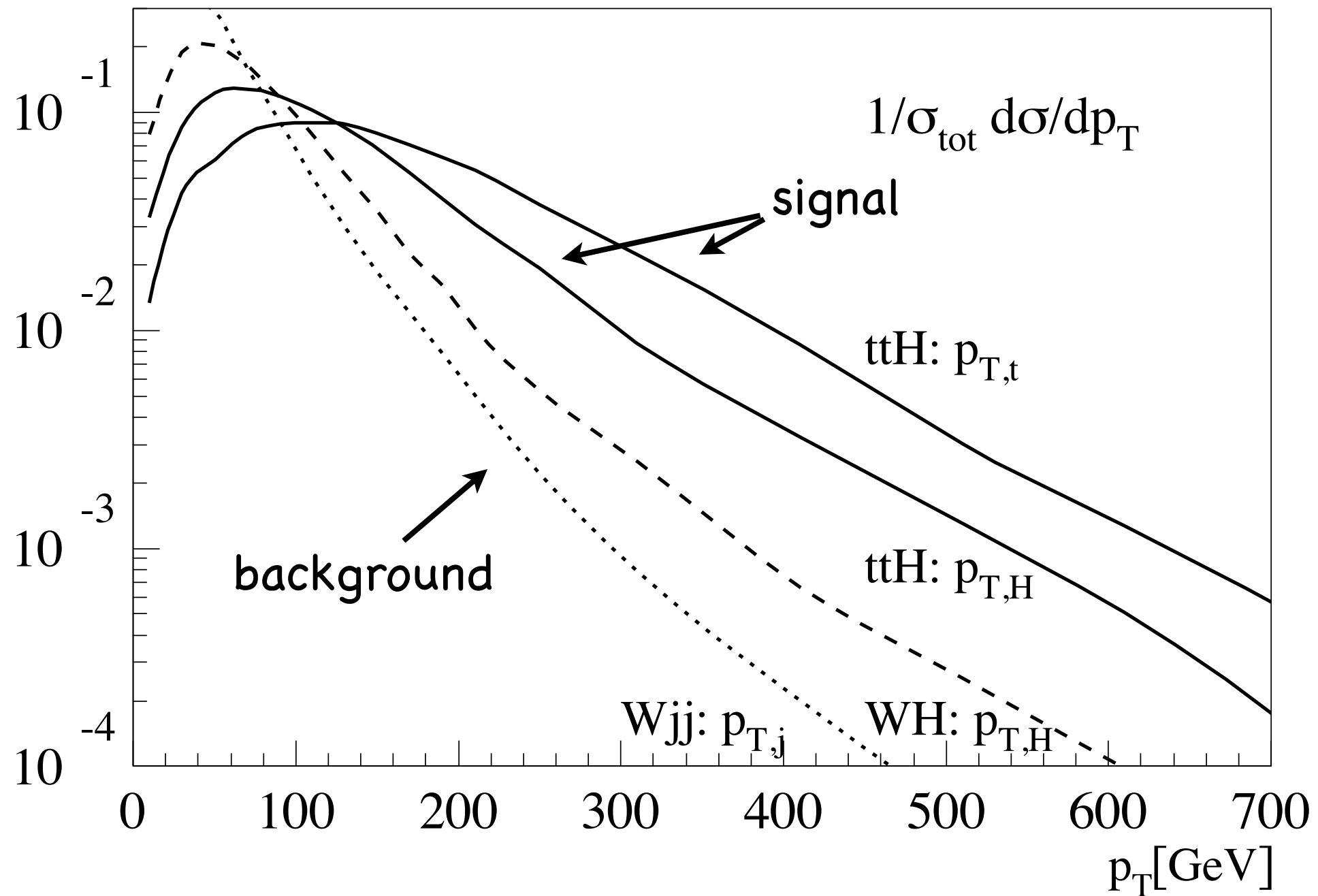
→ $k=1.53$

$w+\text{jets}$

negligible after
b-tags and
taggers



pT distributions relevant for tth

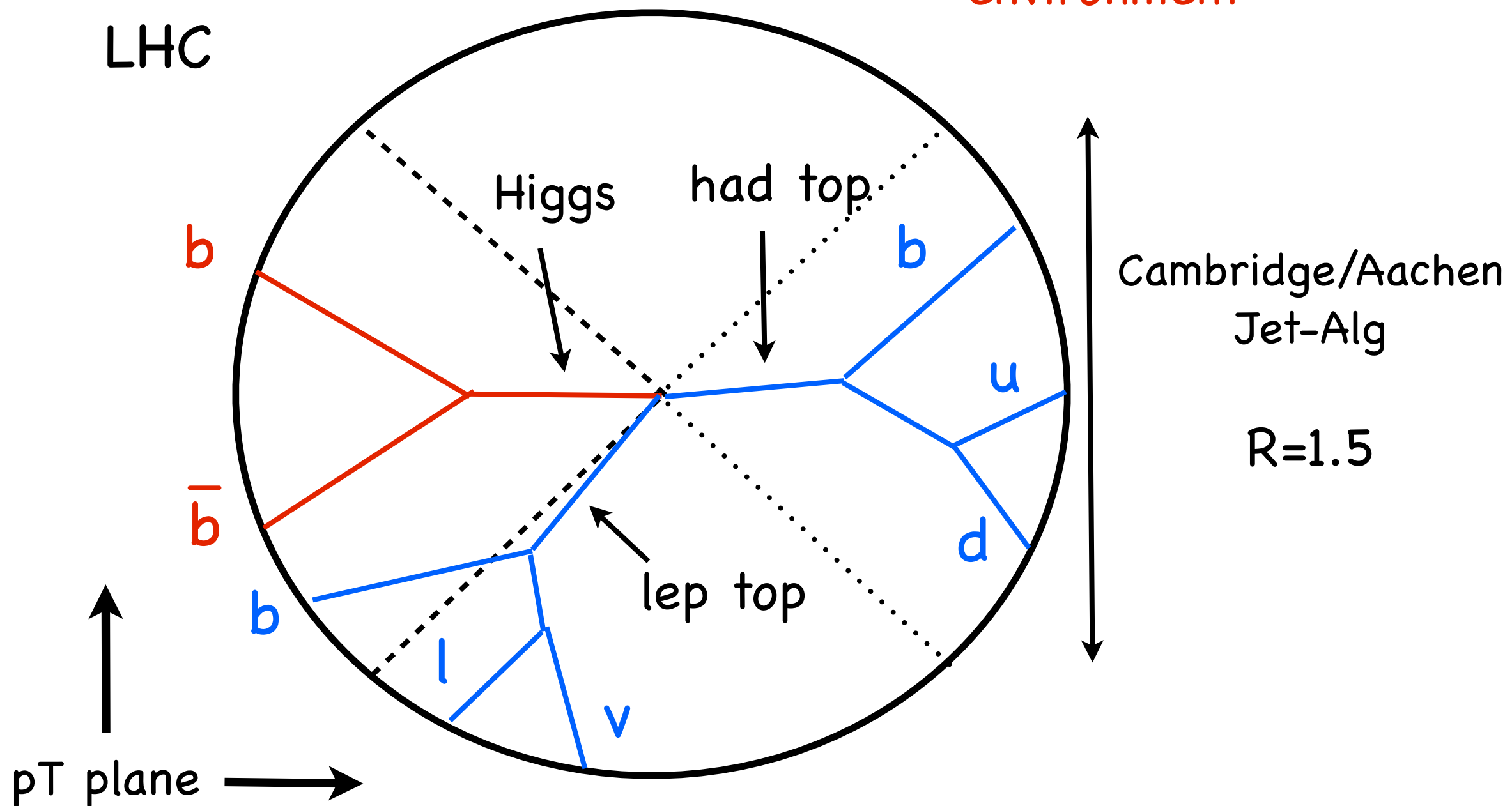


Problems in event reconstruction:

- (b-)jet multiplicity
- reconstruction efficiency

Boost should help
but

need tagger for this
environment

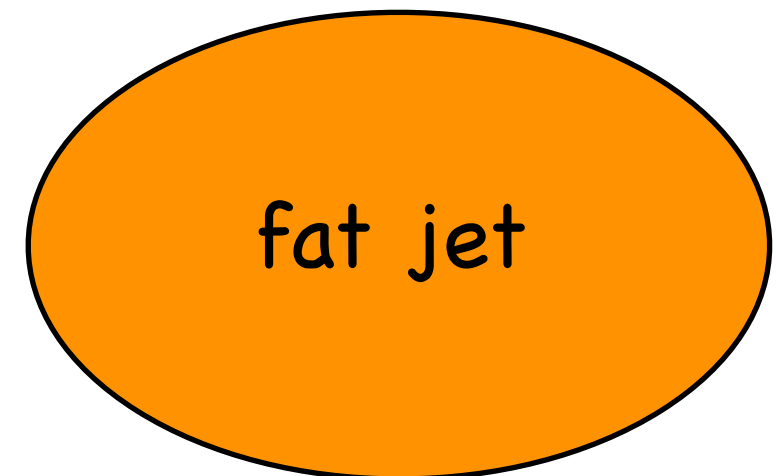


How does the HEPTopTagger work?

I. Find fat jets (C/A, $R=1.5$, $p_T > 200$ GeV)

II. Find hard substructure using mass drop criterion

Undo clustering, $m_{\text{daughter}_1} < 0.8 m_{\text{mother}}$ to keep both daughters

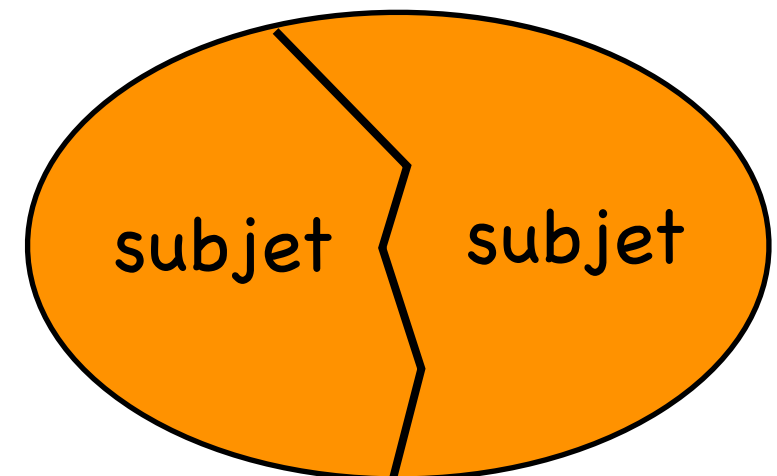


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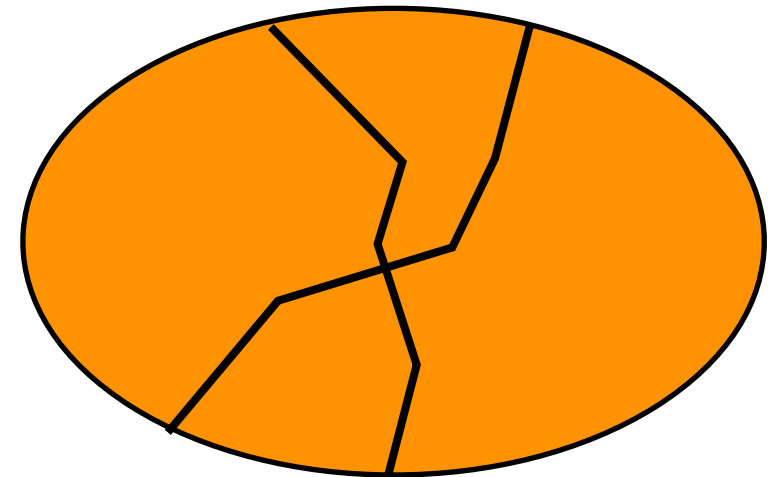


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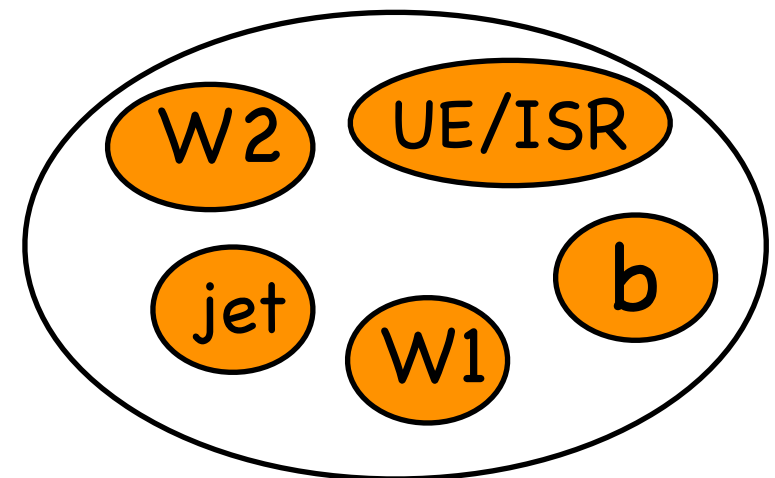
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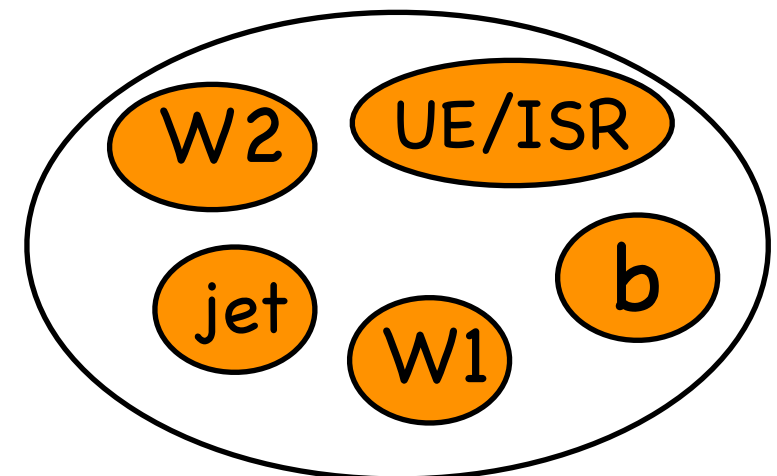
III. Apply jet grooming to get top decay candidates



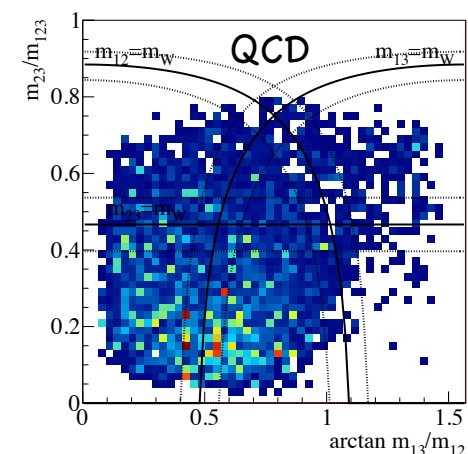
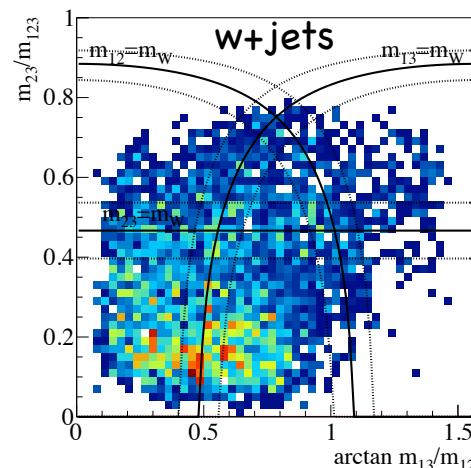
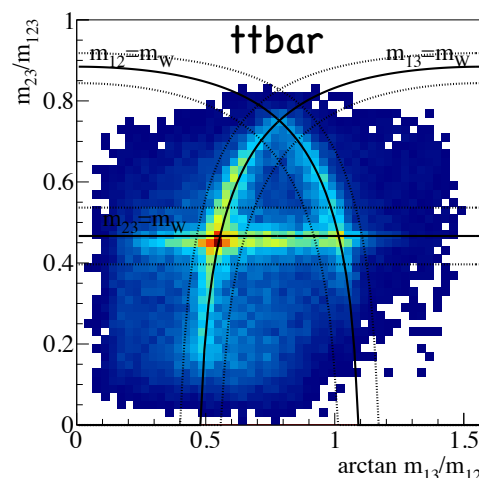
How does the HEPTopTagger work?

- I. Find fat jets (C/A, R=1.5, $p_T > 200$ GeV)
- II. Find hard substructure using mass drop criterion
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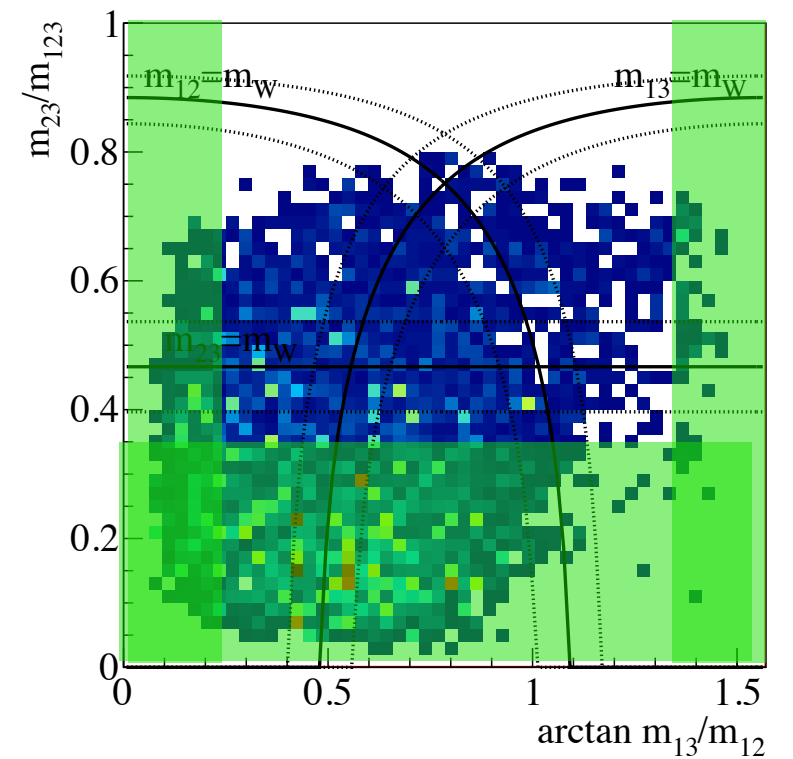
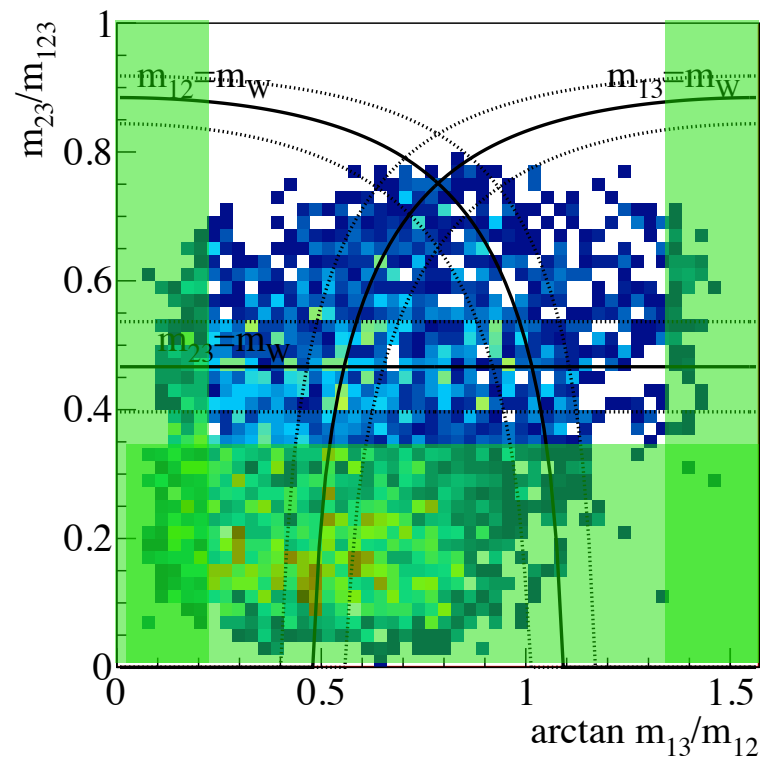
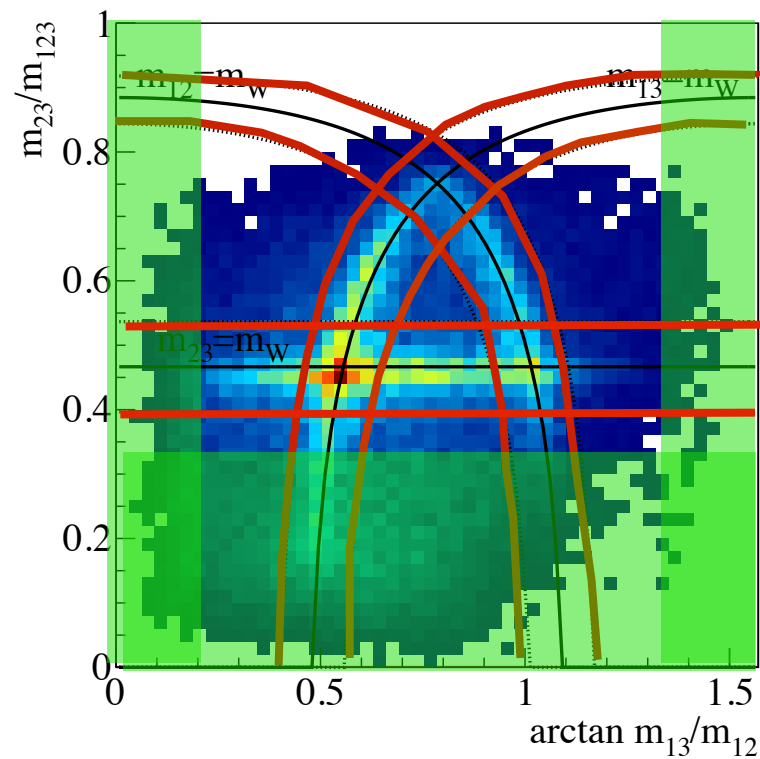


IV. Choose pairing based on kinematic correlation, e.g. top mass, W mass and invariant subjet masses



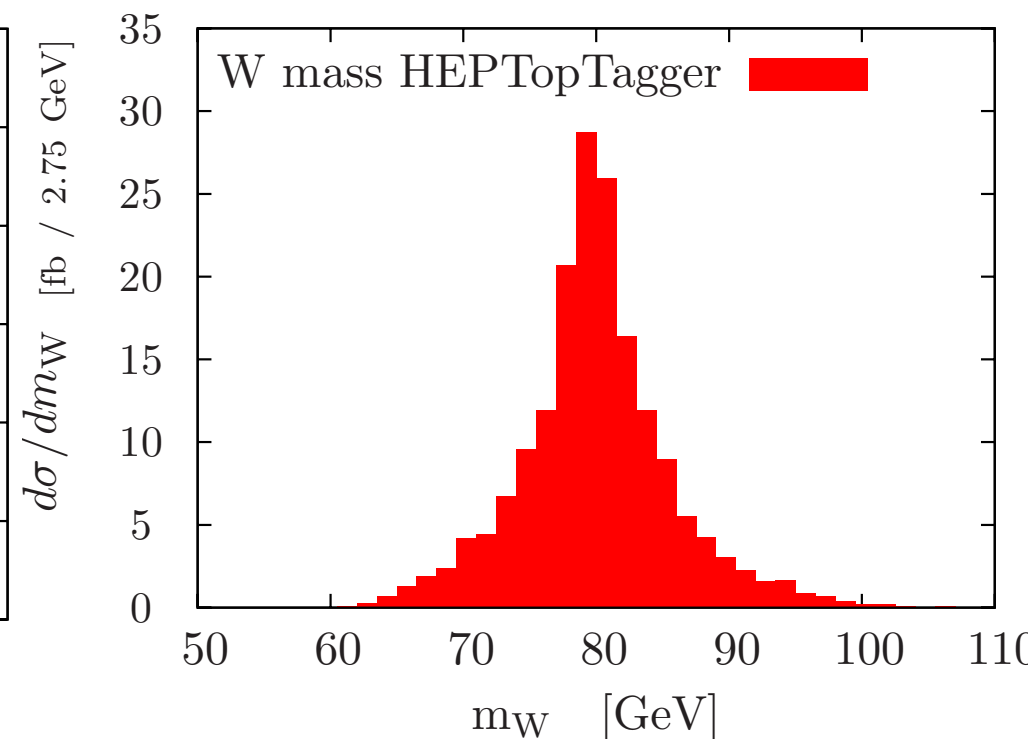
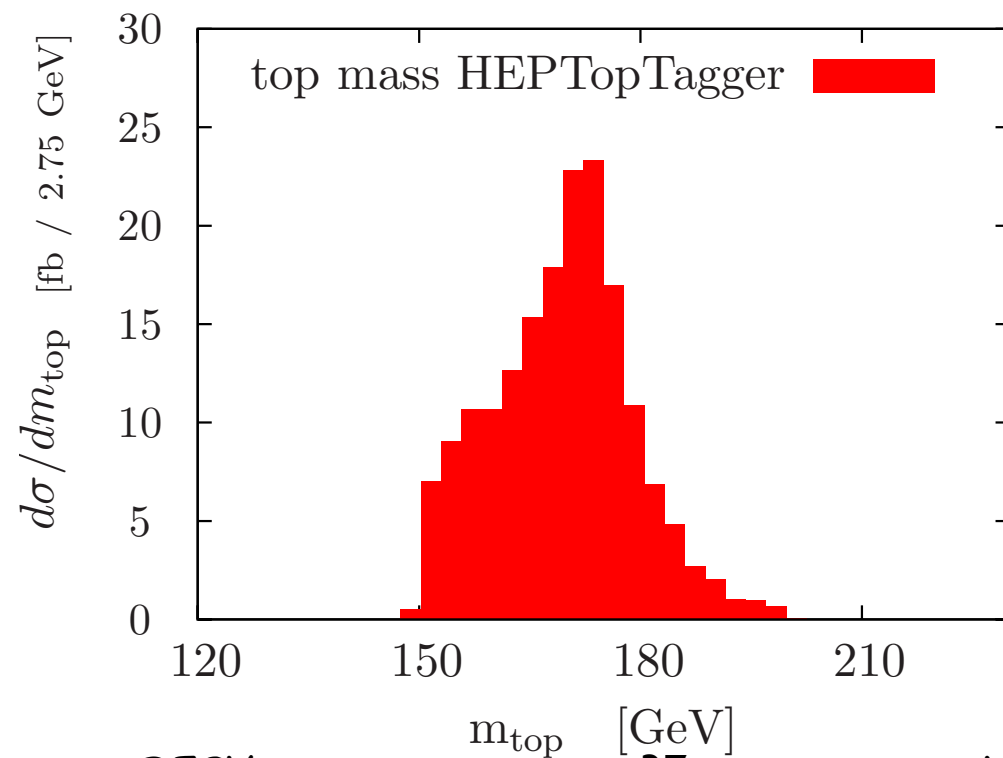
IV. check mass ratios

Cluster top candidate into 3 subjets j_1, j_2, j_3

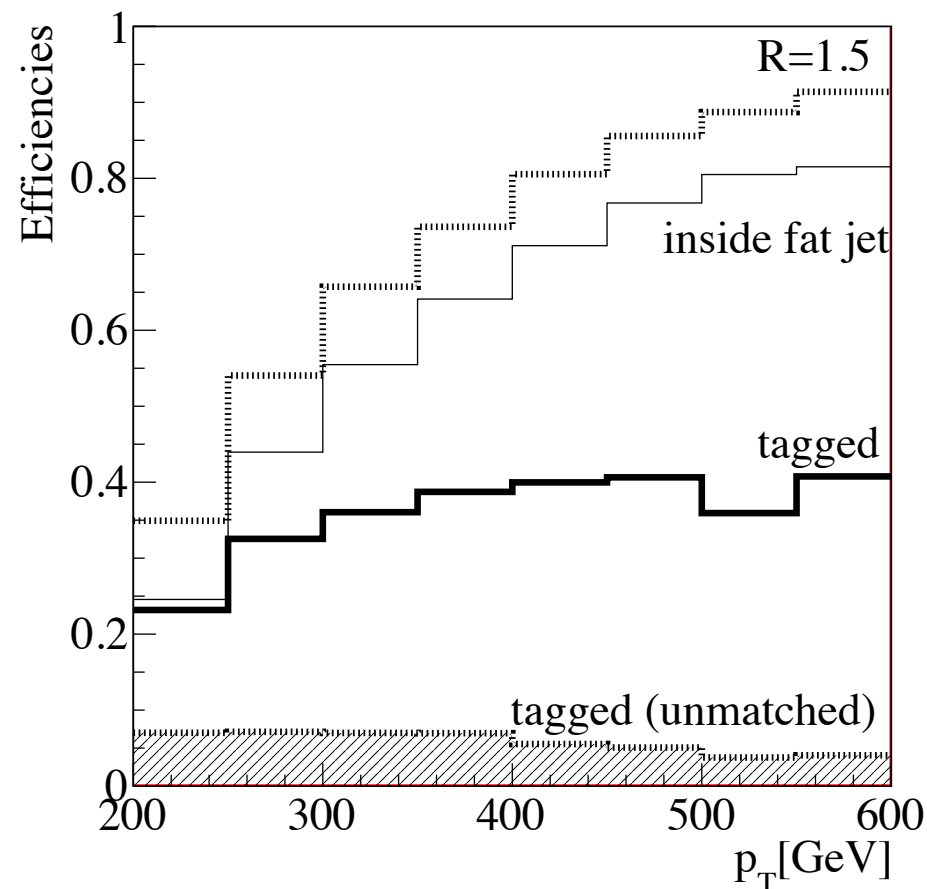
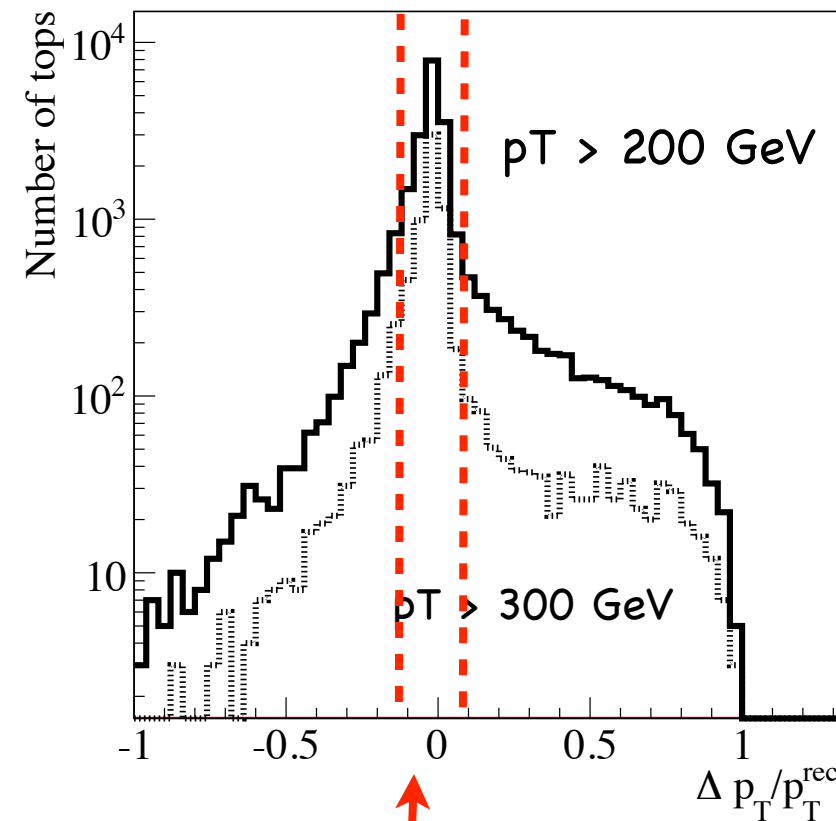
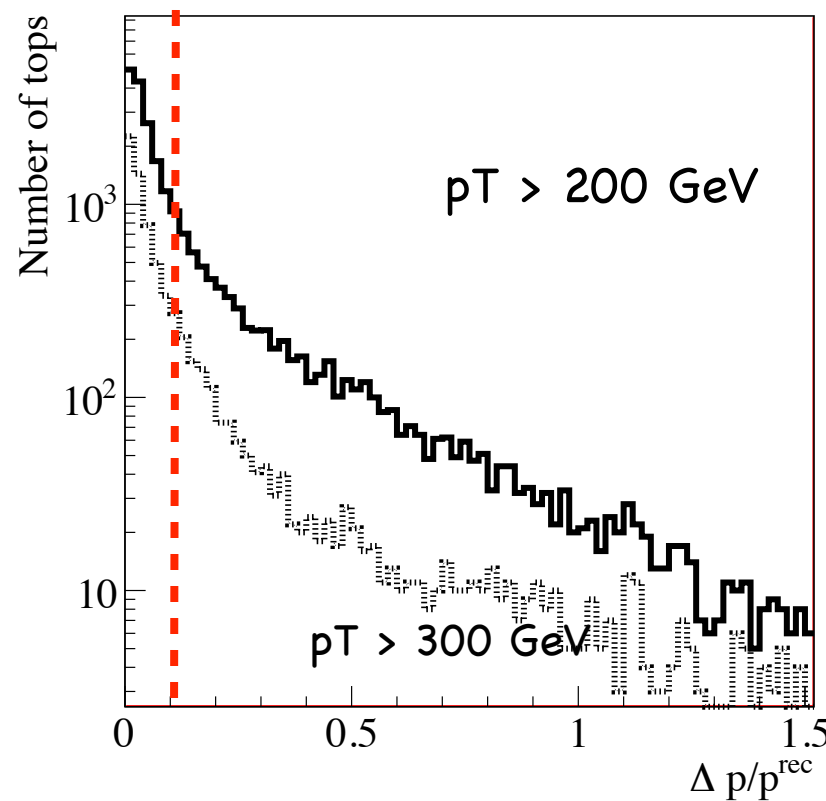


No fix pairing
for W mass
reconstruction

Only invariants for
reconstruction



Top quark momentum reconstruction

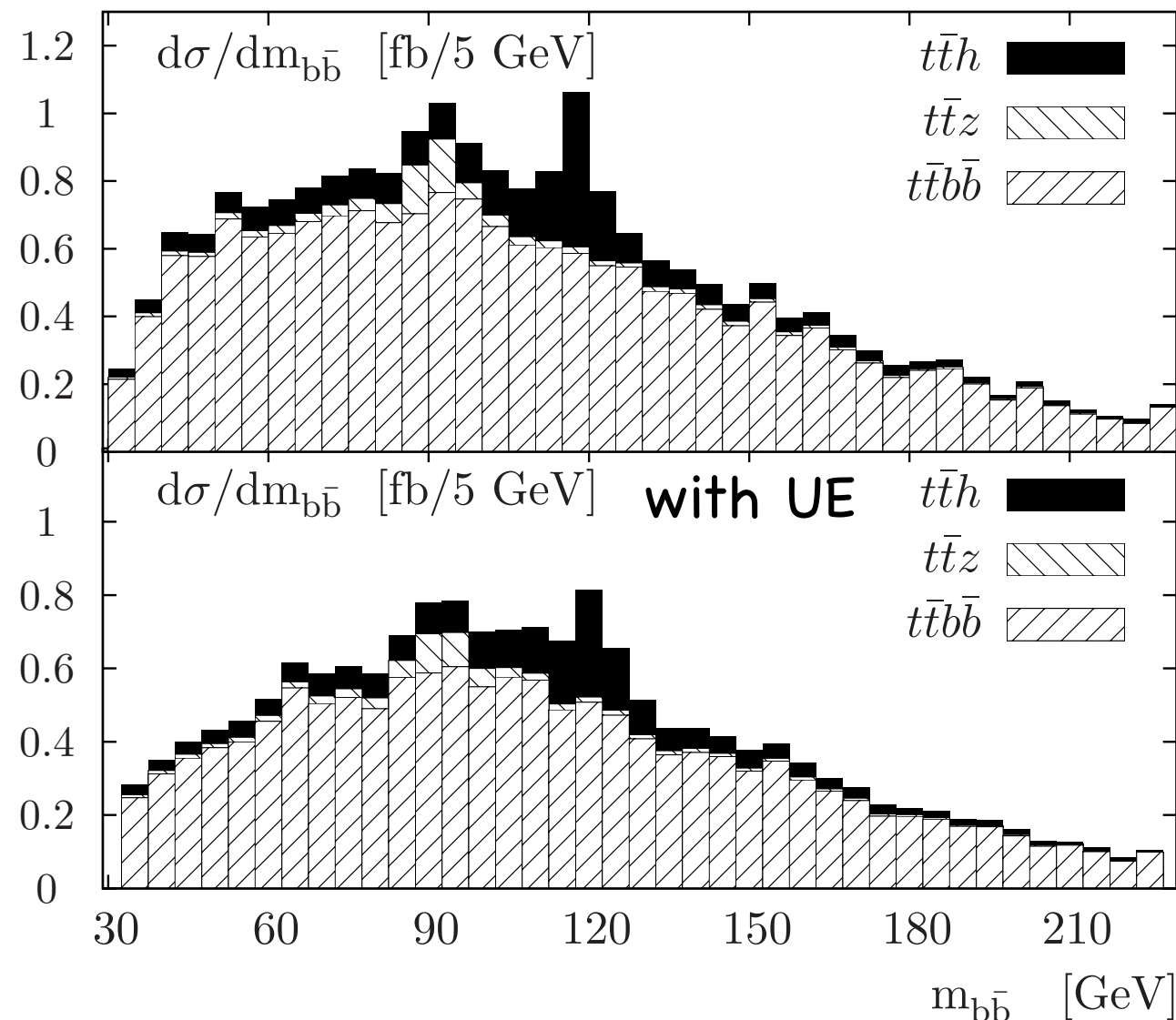


► Great reconstruction of top quark momentum

► 35% tagging efficiency
2% W+jets fake rate

► Tagger used in resonance searches in ATLAS: 1207.2409

Results for $t\bar{t}h$



- 5 sigma sign. with 100 1/fb

- Development of Higgs and top tagger for busy final state

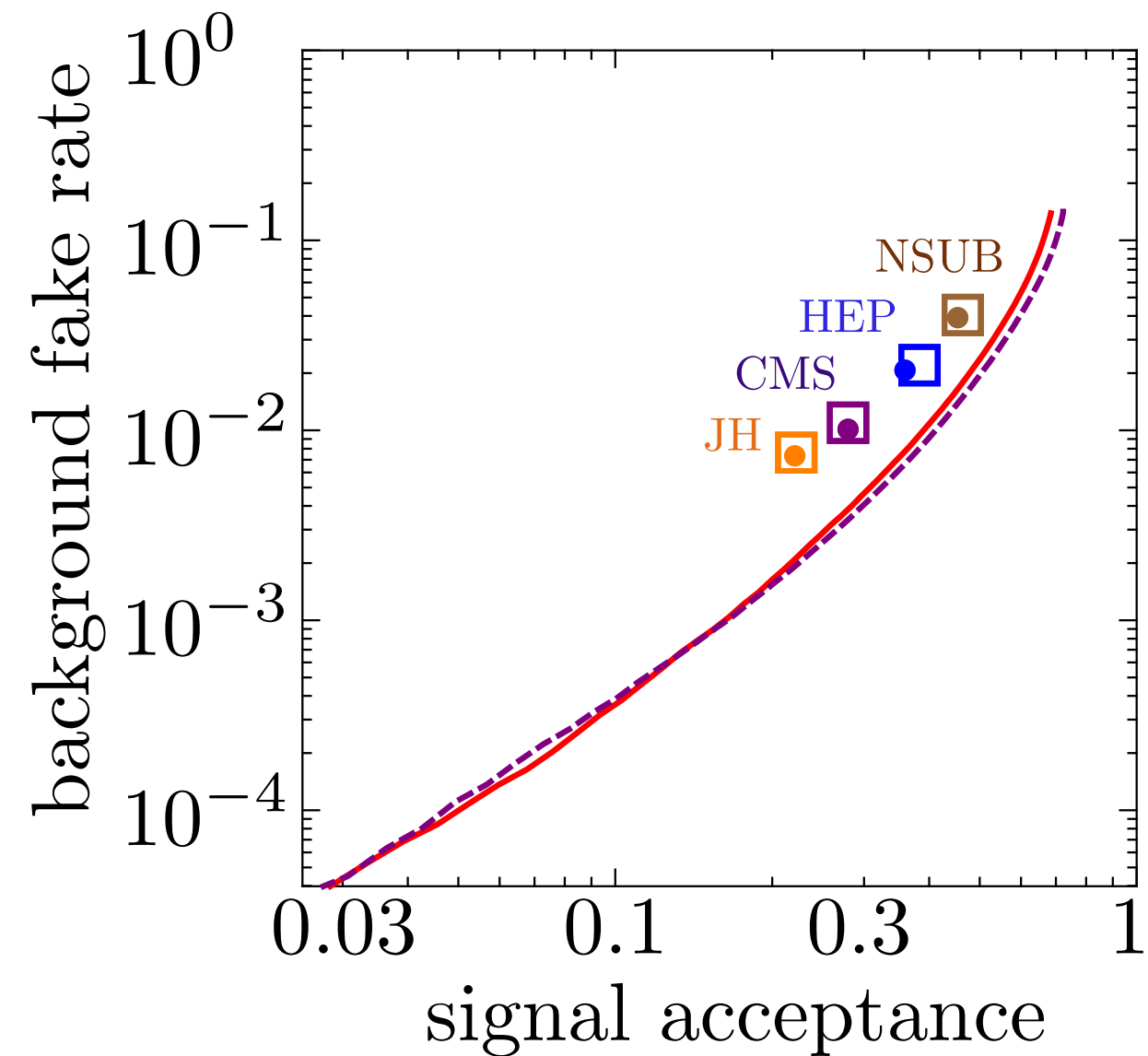
- Improvement of S/B from 1/9 to 1/2

→ $t\bar{t}h$ might be a window to Higgs-top coupling

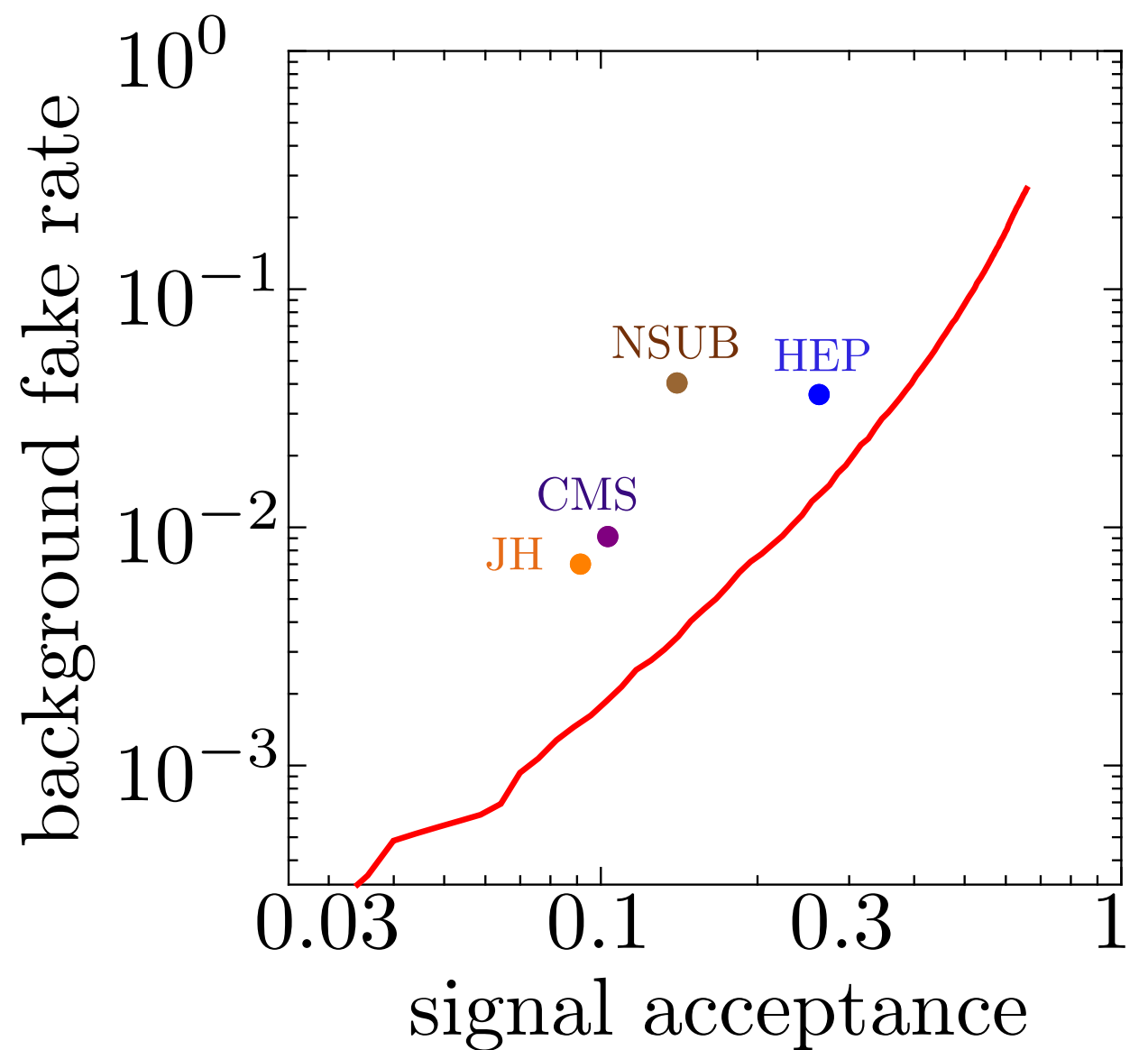
Shower deconstruction based top tagger can do even better:

Input for SD: small jets with $R=0.2$ and $p_T > 5.0$ GeV

fat jet $p_T > 500$ GeV and $R=1.0$:



fat jet $p_T > 200$ GeV and $R=1.5$:



III. Measuring the CP of the Higgs boson

- For light Higgs with 125 GeV CP can be measured using angular correlations of tagging jets in Gluon Fusion with 2 additional jets

[Plehn, Rainwater, Zeppenfeld PRL 88 (2002)]

- Event shape observables can be used to measure CP of Higgs

[Englert, MS, Takeuchi 1203.5788]

Interaction:

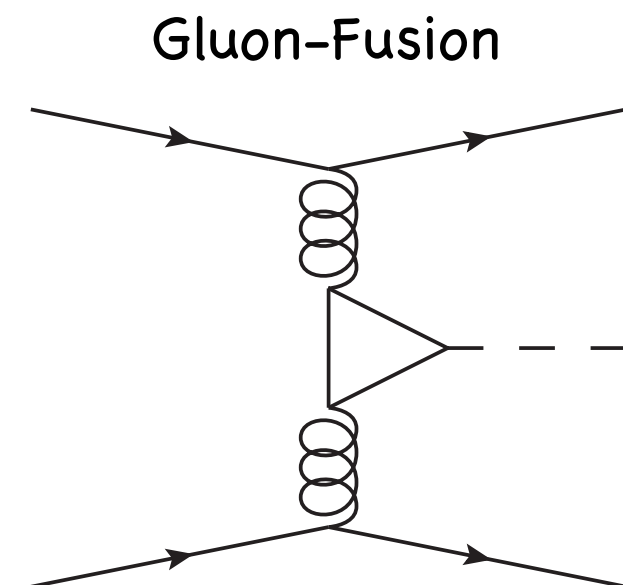
$$\mathcal{L} = \frac{\alpha_s}{12\pi v} H G_{\mu\nu}^a G^{a\mu\nu} + \frac{\alpha_s}{16\pi v} A G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$$

For tagging jets with $|p_z^J| \gg |p_{x,y}^J|$

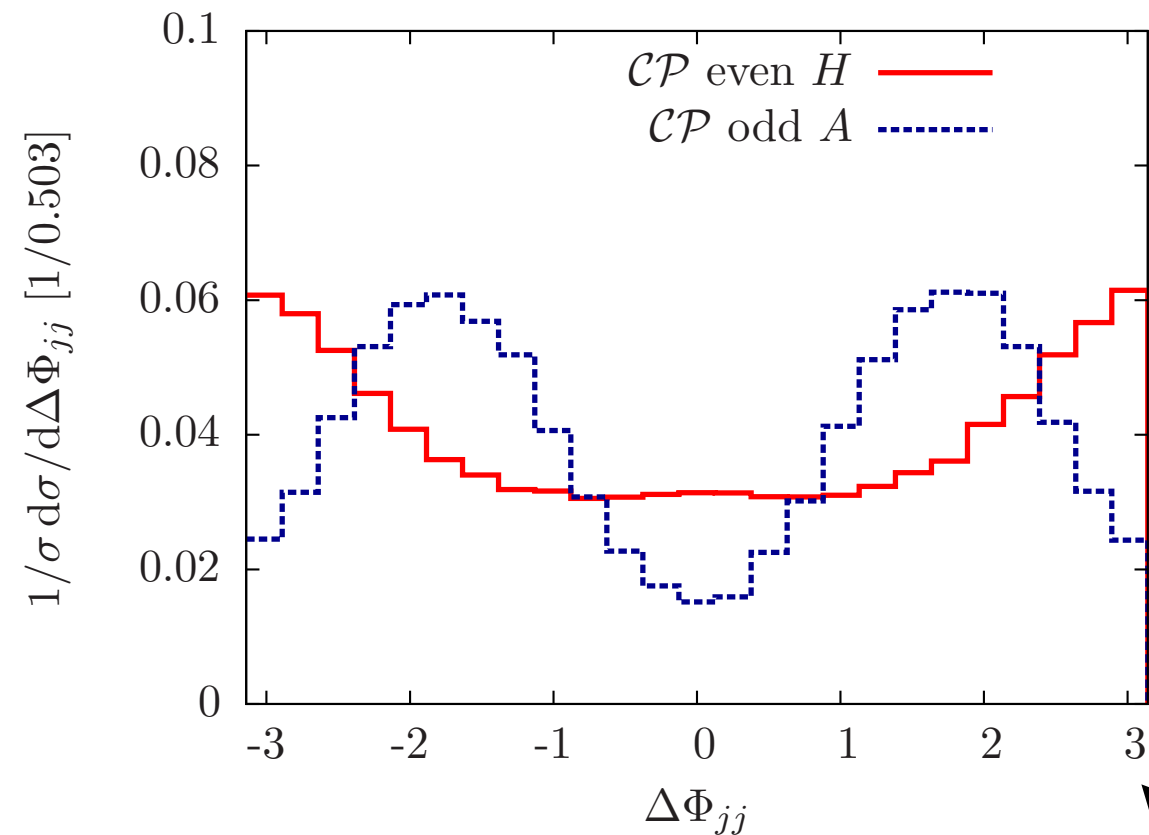
$$\mathcal{M}_{\text{even}} \sim J_1^\mu J_2^\nu [g_{\mu\nu}(q_1 \cdot q_2) - q_{1\nu} q_{2\mu}]$$

$$\sim [J_1^0 J_2^0 - J_1^3 J_2^3] \mathbf{p}_T^{J_1} \cdot \mathbf{p}_T^{J_2} \sim 0 \text{ for } \Delta\phi_{jj} = \pi/2$$

$\mathcal{M}_{\text{odd}}$ contains Levi-Civita tensor which is 0 if two of momenta linearly dependent, i.e. if $\Delta\phi_{jj} = 0$ or $\Delta\phi_{jj} = \pi$



Tagging jets approach:

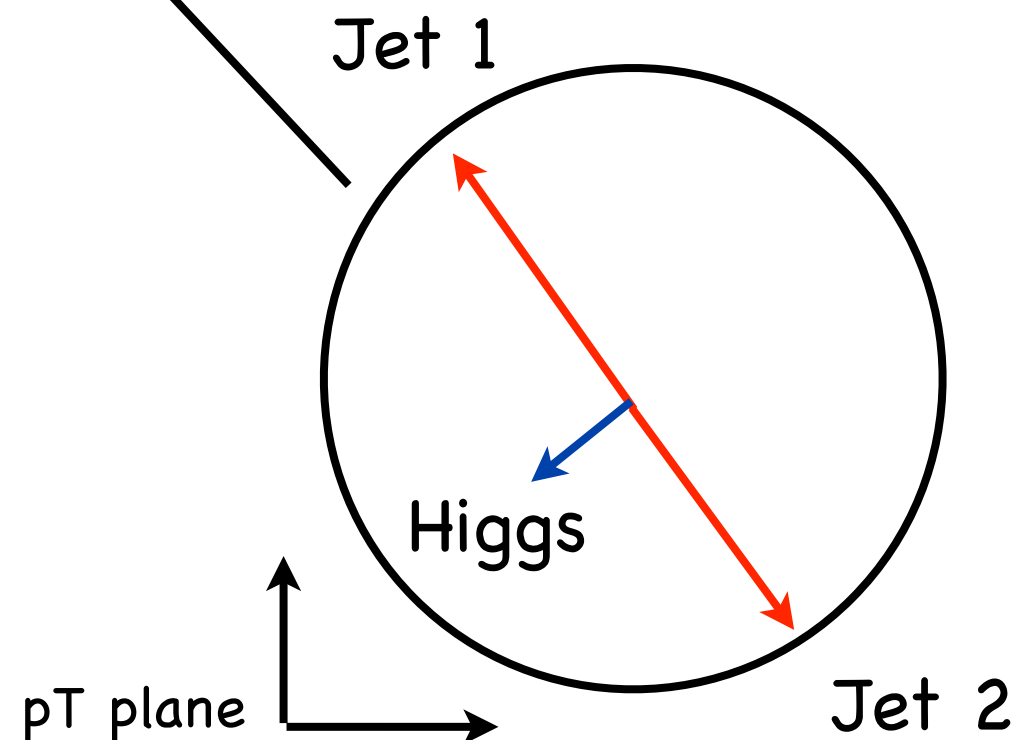
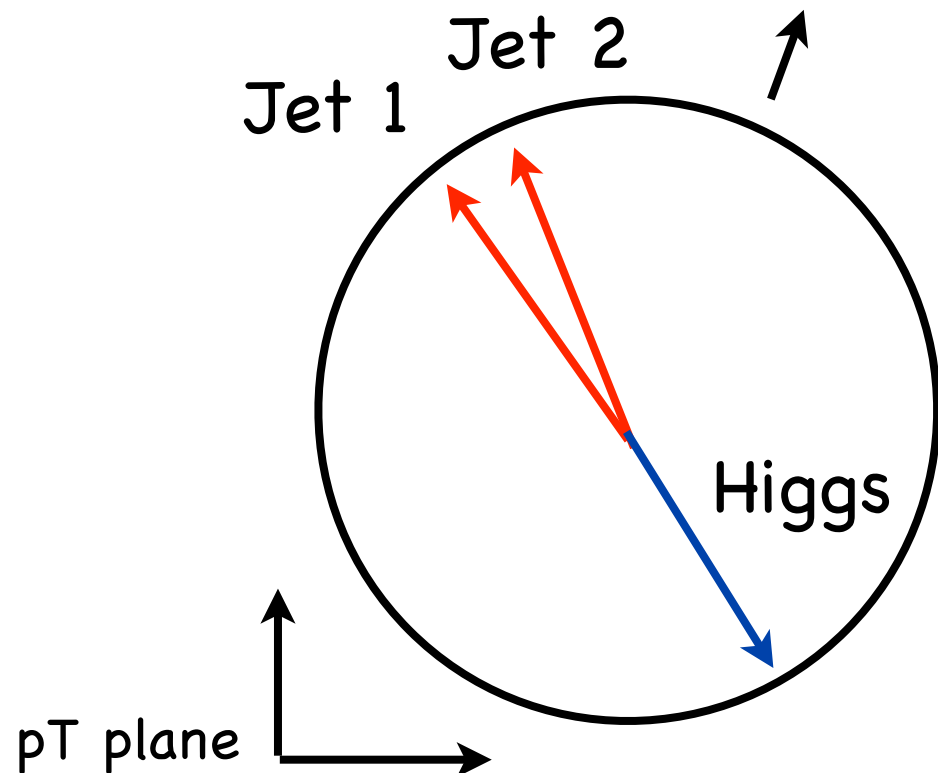


azimuthal angle between all jets
with larger or smaller rapidity
wrt Higgs

$$p_{<}^{\mu} = \sum_{j \in \{\text{jets: } y_j < y_h\}} p_j^{\mu}$$

$$p_{>}^{\mu} = \sum_{j \in \{\text{jets: } y_j > y_h\}} p_j^{\mu}$$

$$\Delta\Phi_{jj} = \phi(p_{>}) - \phi(p_{<})$$



Event shapes

- Event shapes well studied experimentally and theoretically

[Bethke, Nucl.Phys.Proc.Suppl. 121 (2003)]

[Kluth. et al, EPJC 21 (2011)]

[Banfi et al., JHEP 0408]

[Gehrmann-De Ridder et al., JHEP 0712]

- Avoids vetoing of jets which can induces large theo. uncertainties

[Stewart, Tackmann PRD 85 (2012)]

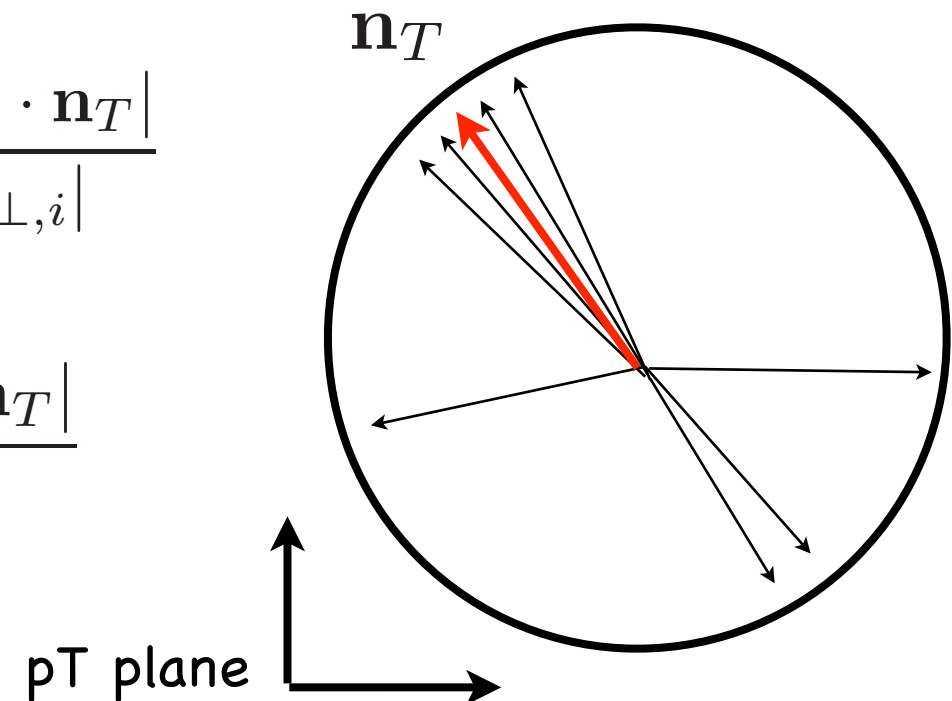
- Event shape measurements established in experimental collaborations already now
e.g.

transverse thrust

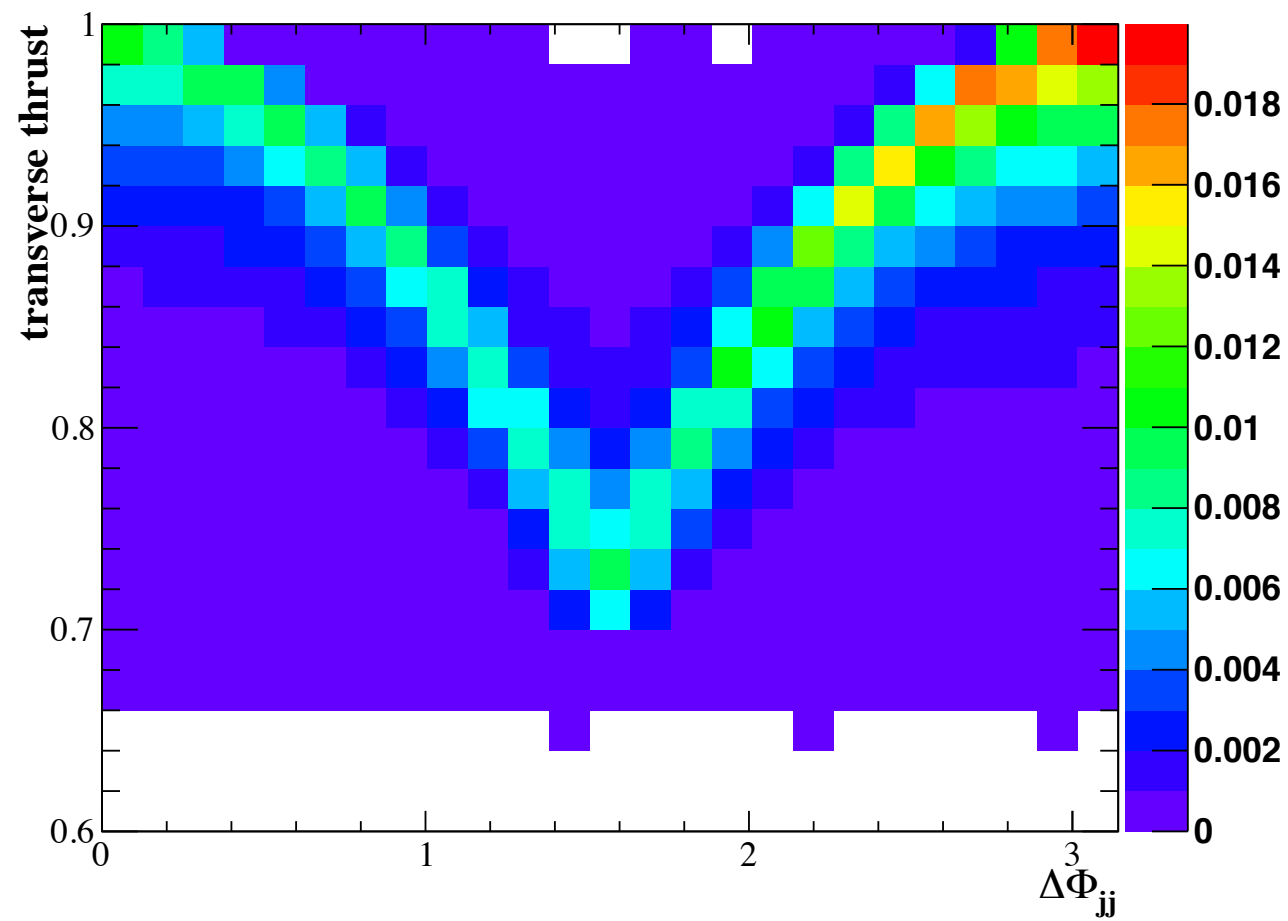
$$T_{\perp,g} = \max_{\mathbf{n}_T} \frac{\sum_i |\mathbf{p}_{\perp,i} \cdot \mathbf{n}_T|}{\sum_i |\mathbf{p}_{\perp,i}|}$$

transverse thrust
minor

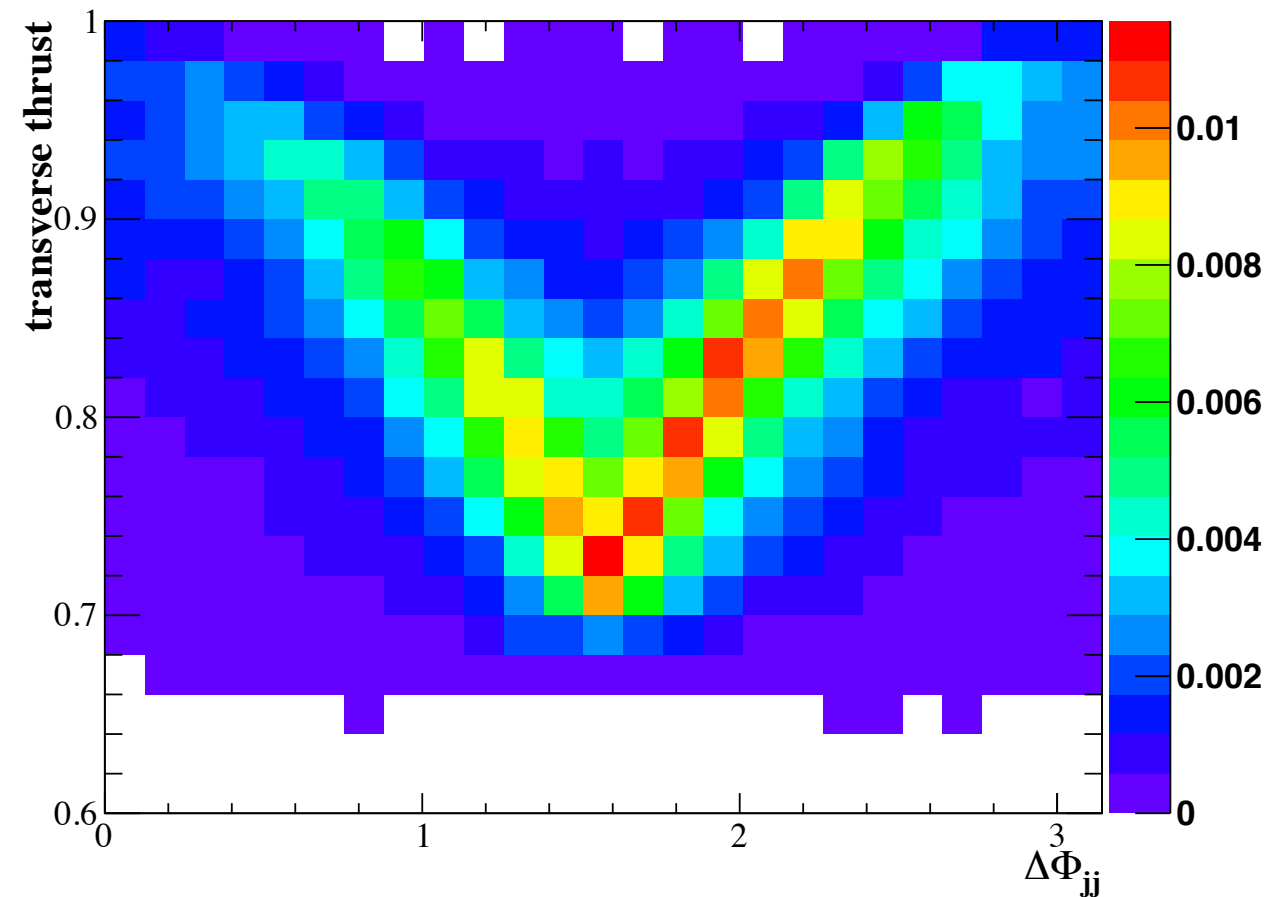
$$T_{m,g} = \frac{\sum_i |\mathbf{p}_{\perp,i} \times \mathbf{n}_T|}{\sum_i |\mathbf{p}_{\perp,i}|}$$



Obvious correlation between thrust and $\Delta\Phi_{jj}$



(a) $\mathcal{CP}$ even Higgs



(b) $\mathcal{CP}$ odd Higgs

Event selection cuts

two tagging jets: $p_{T,j} \geq 40 \text{ GeV}$, and $|y_j| \leq 4.5$

$$m_{jj} = \sqrt{(p_{j,1} + p_{j,2})^2} \geq 600 \text{ GeV}$$

two taus, hard and central: $p_{T,\tau} \geq 20 \text{ GeV}$, and $|y_\tau| \leq 2.5$

$$|m_{\tau\tau} - m_H| < 20 \text{ GeV}$$

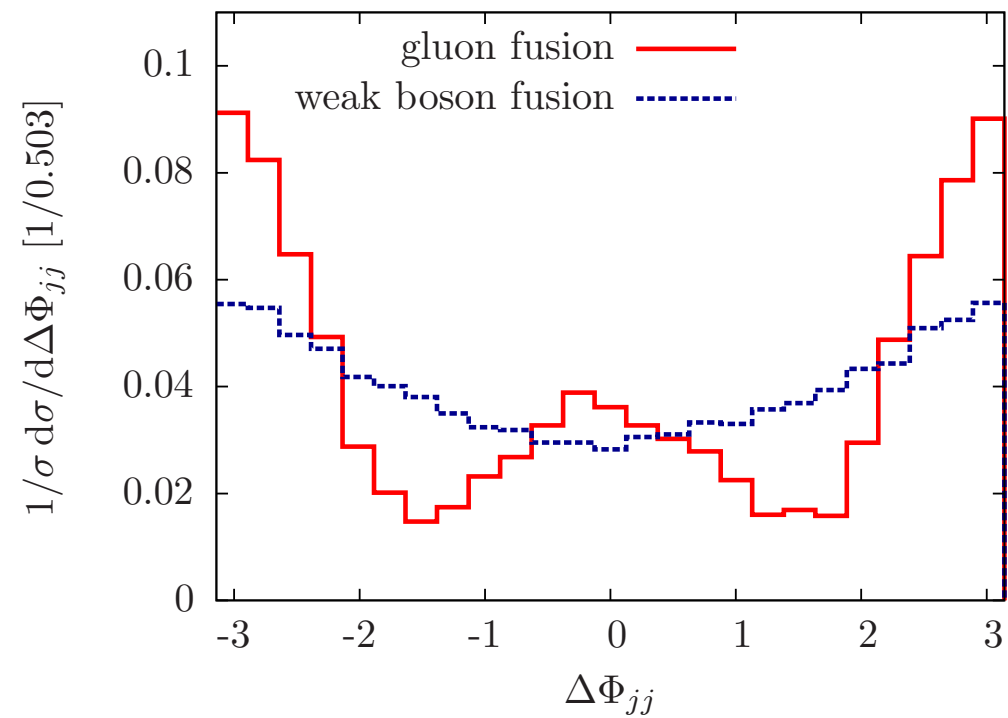
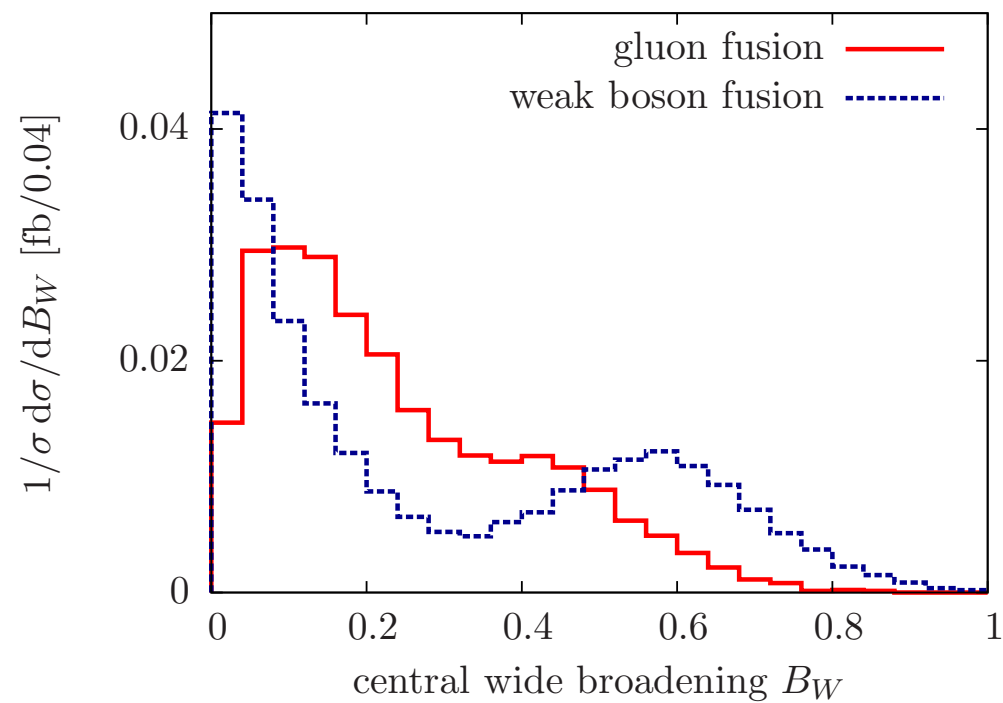
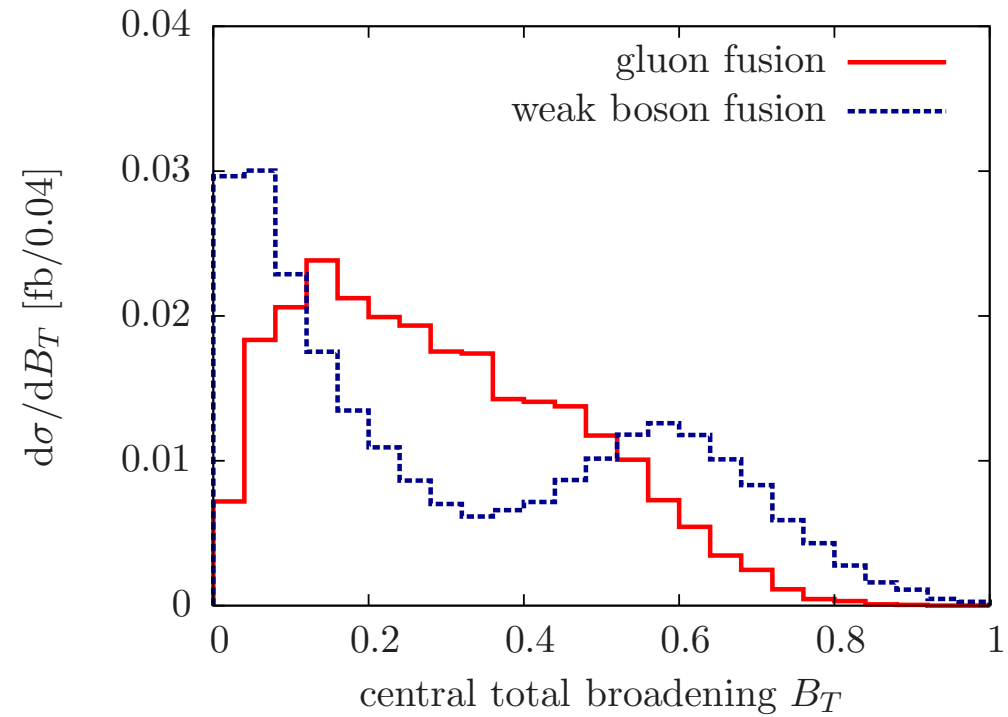
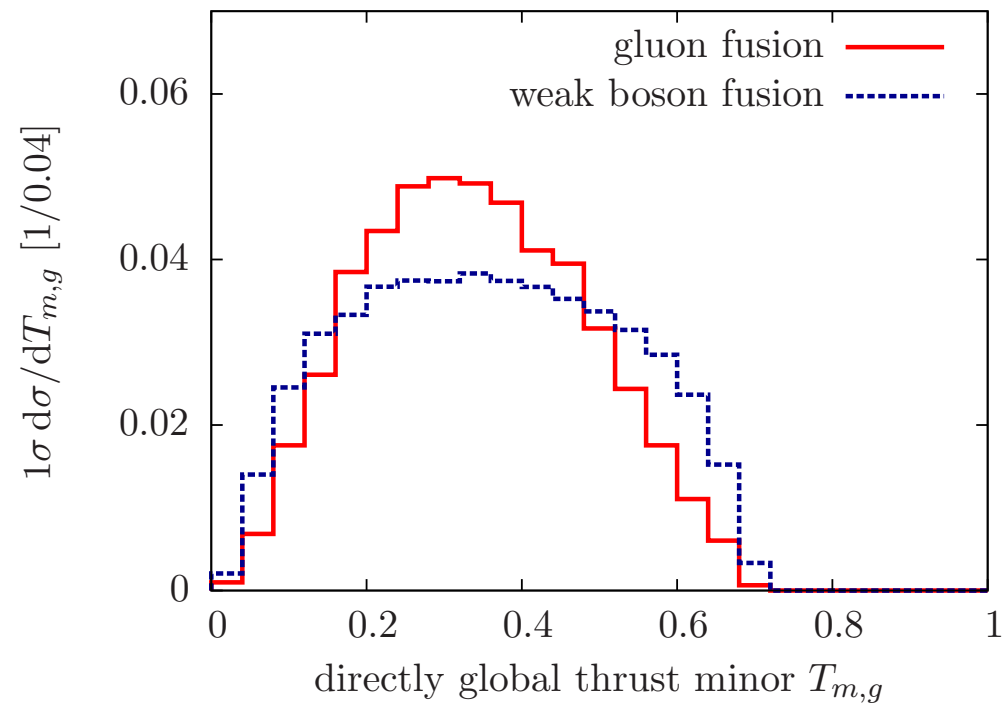
I. Use event shapes with constituents: (theorist's approach)

$$p_{T,i} \geq 1 \text{ GeV} \quad |\eta_i| \leq 4.5$$

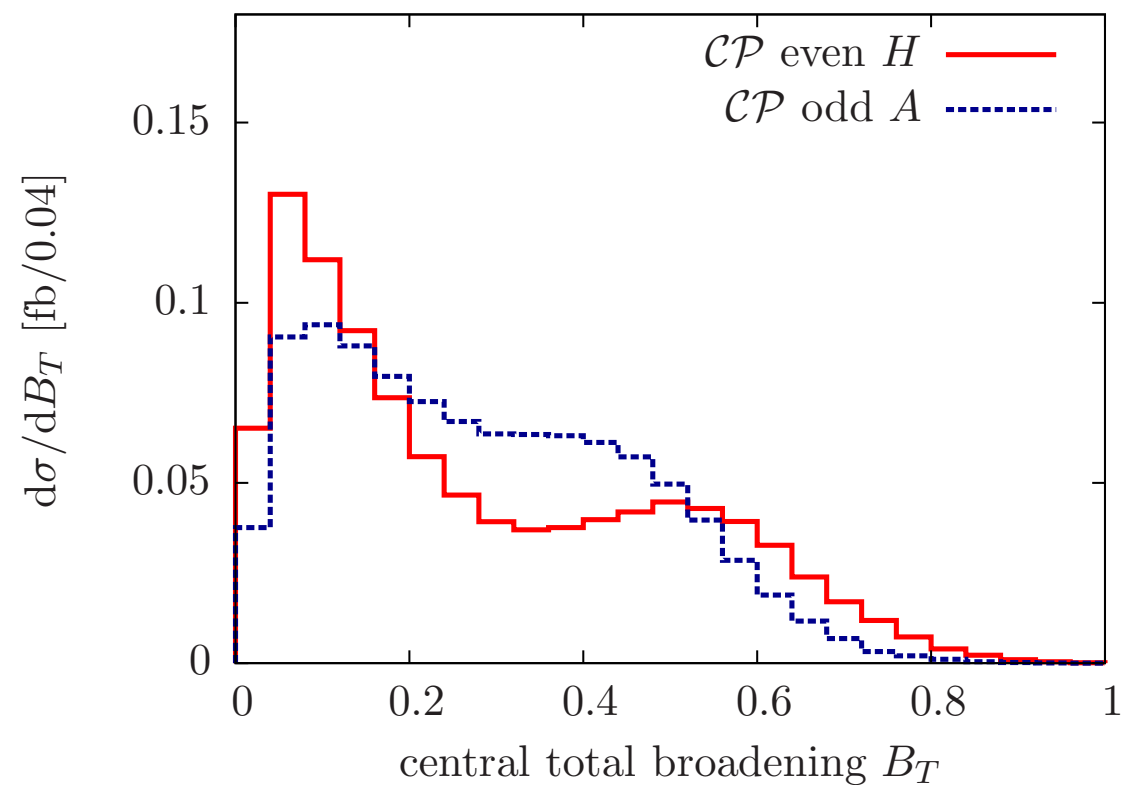
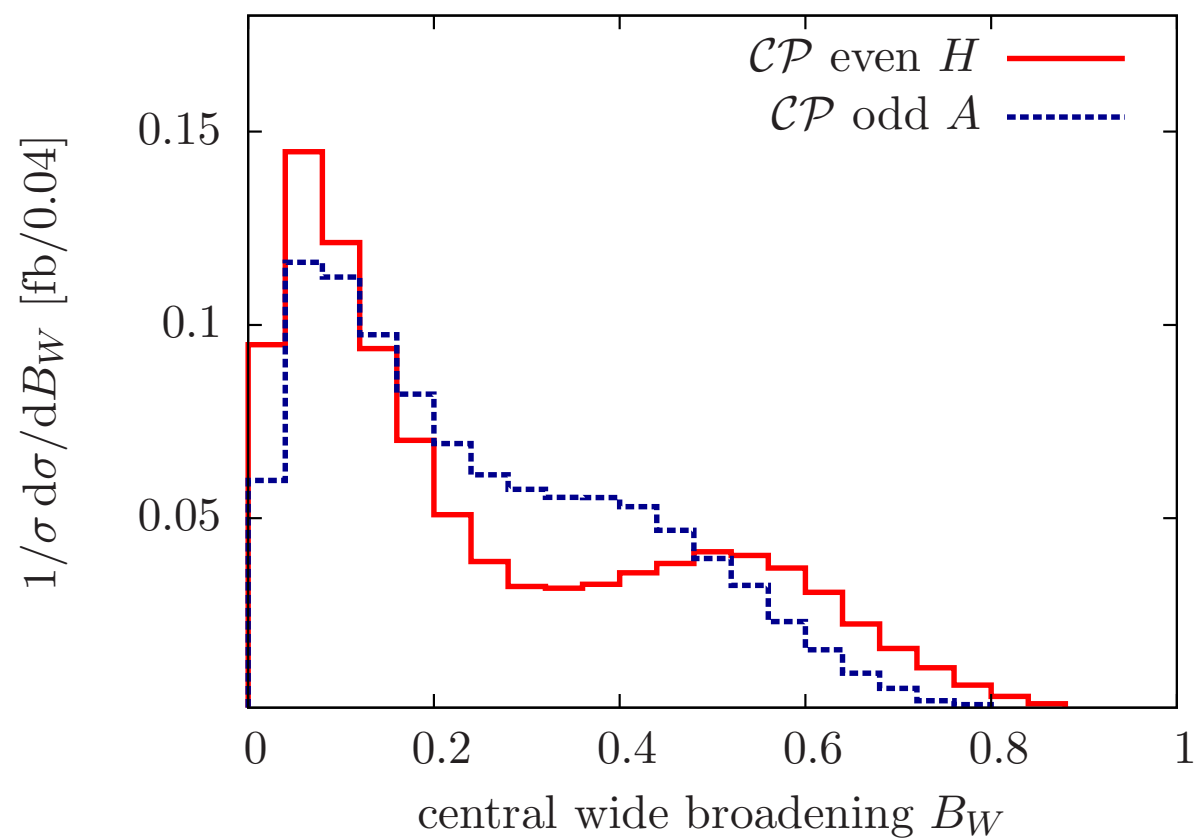
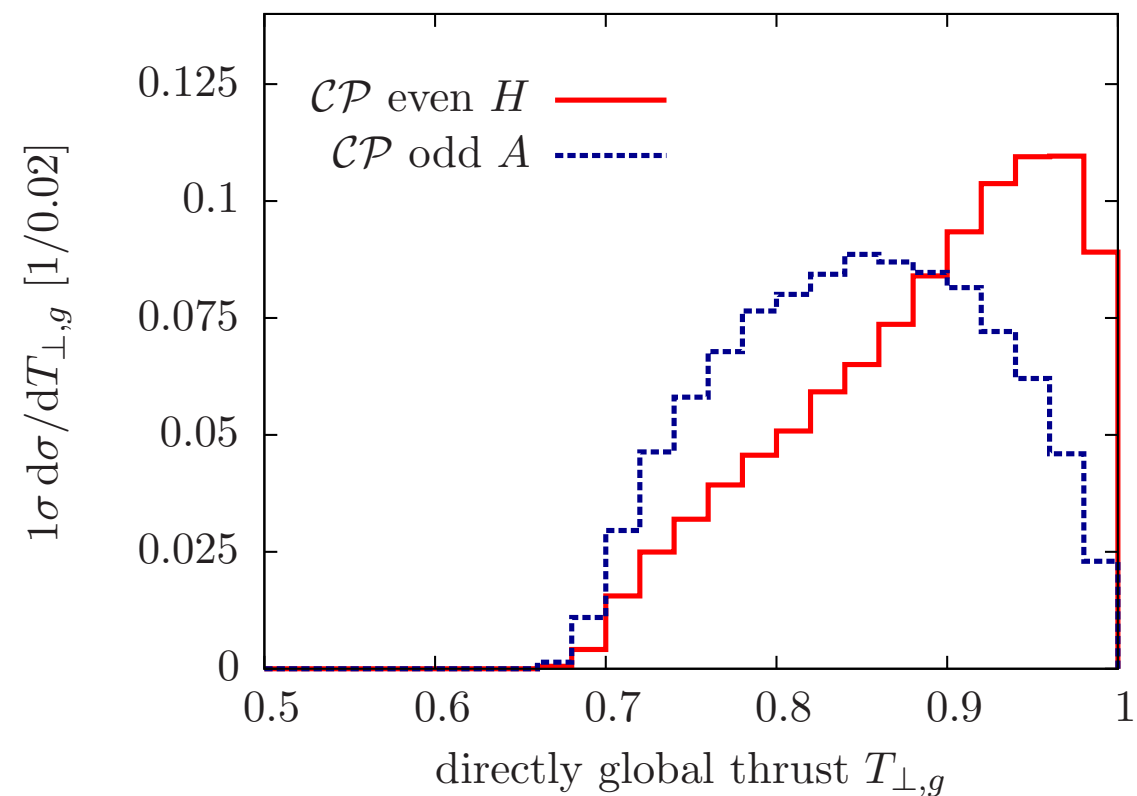
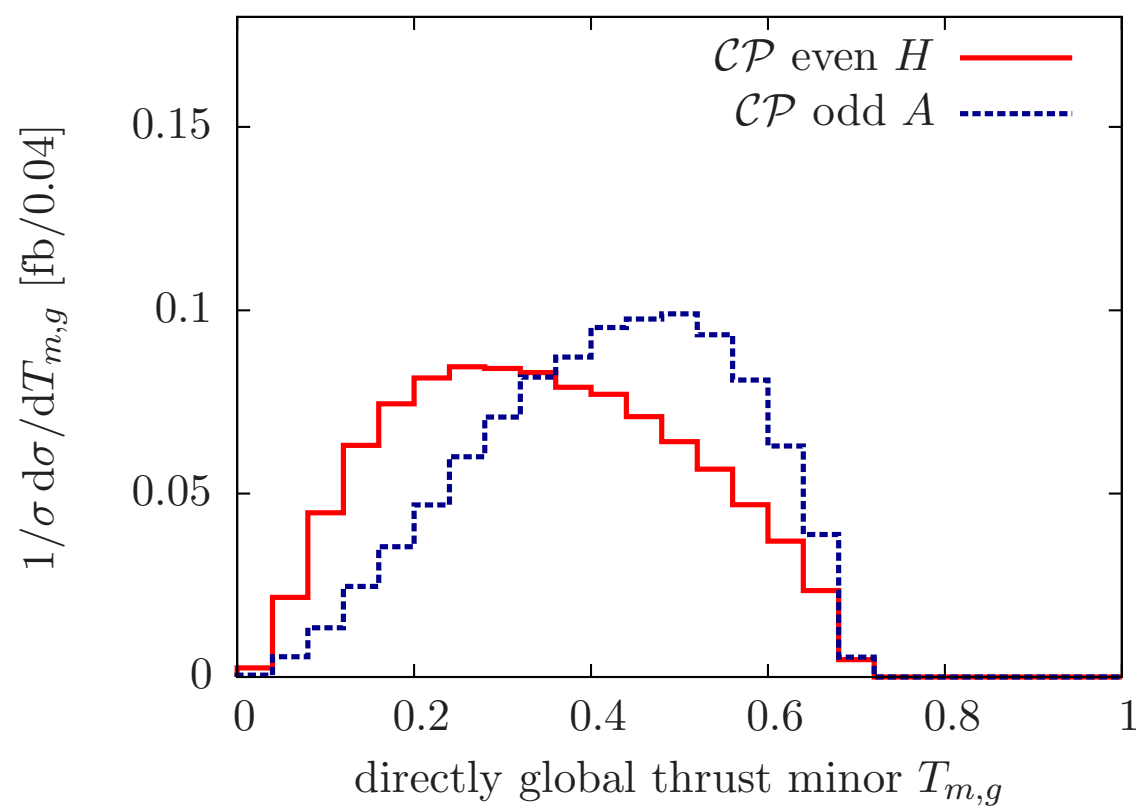
II. reduce pileup sensitivity, use $p_{T,j} \geq 40 \text{ GeV}$, if $2.5 \leq |y_j| \leq 4.5$, and
 $p_{T,j} \geq 10 \text{ GeV}$, if $|y_j| \leq 2.5$.

(compromise between theory and experiment)

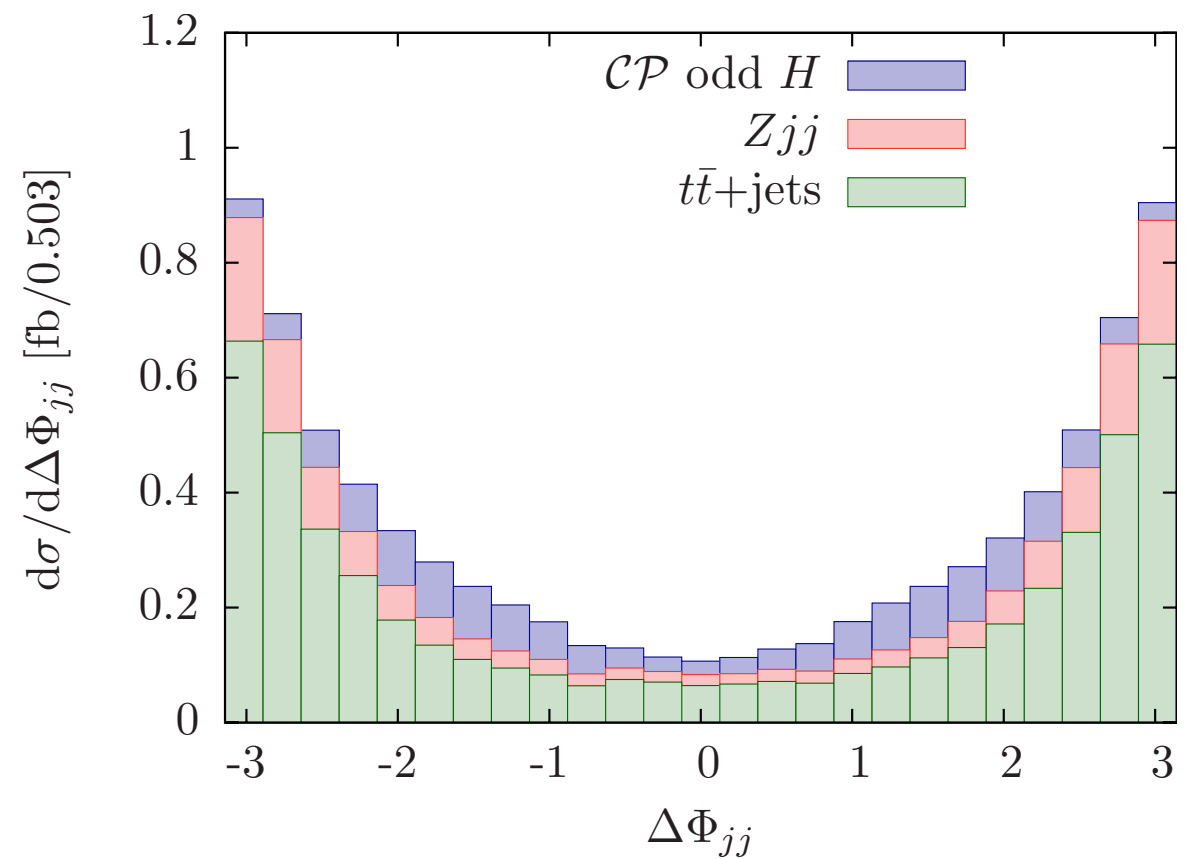
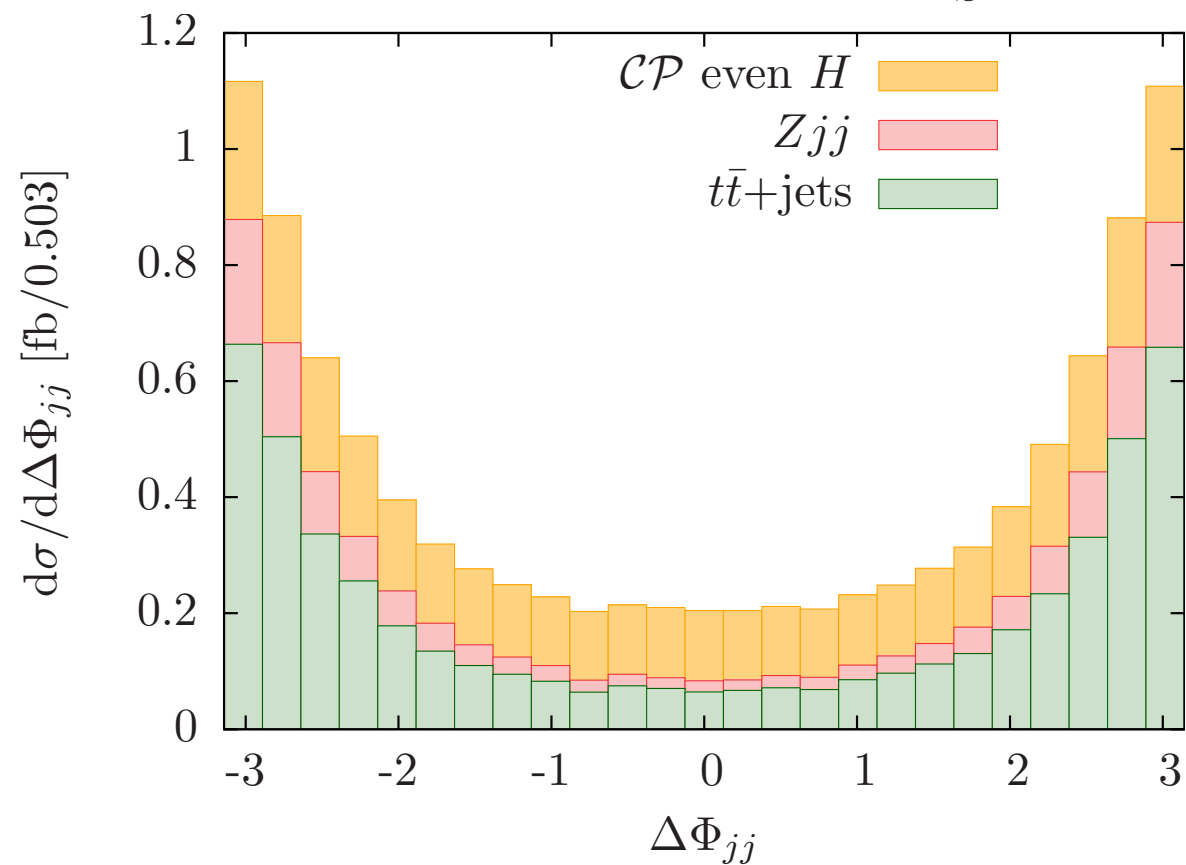
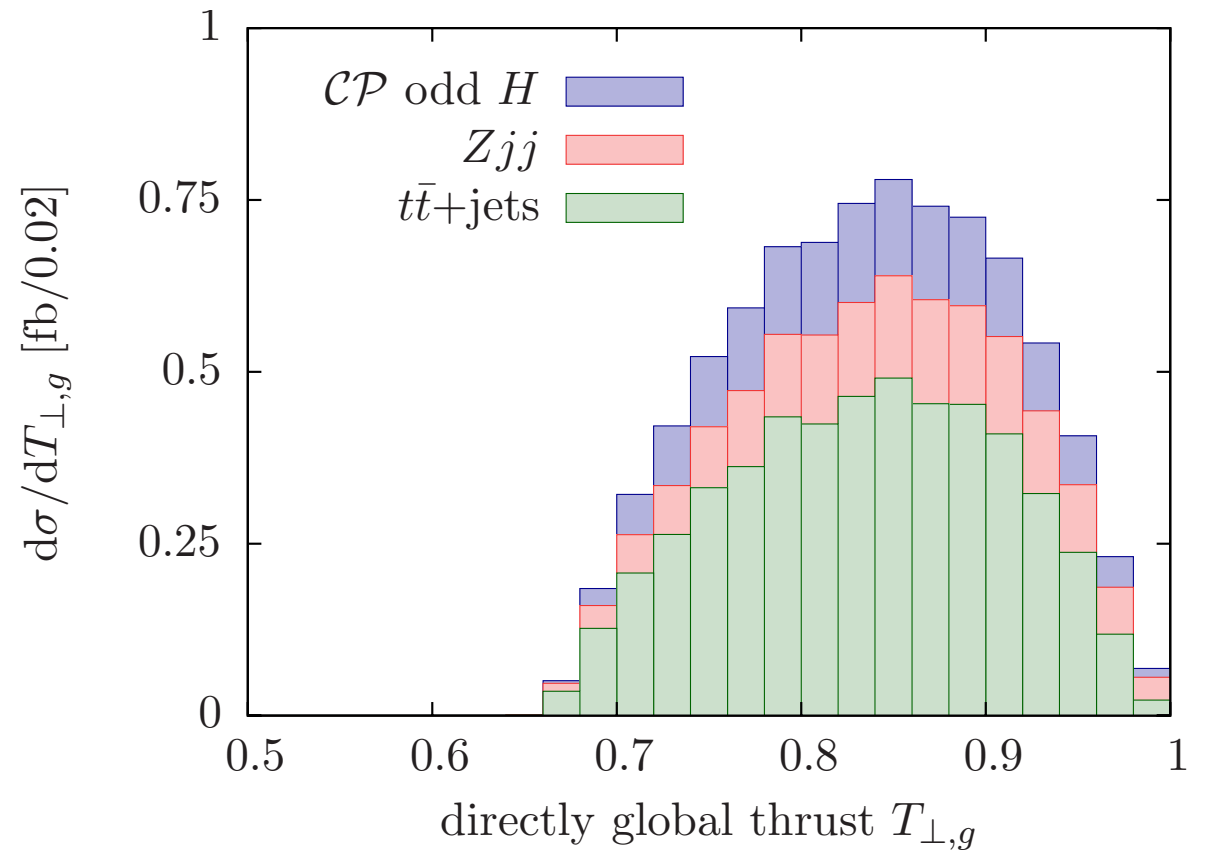
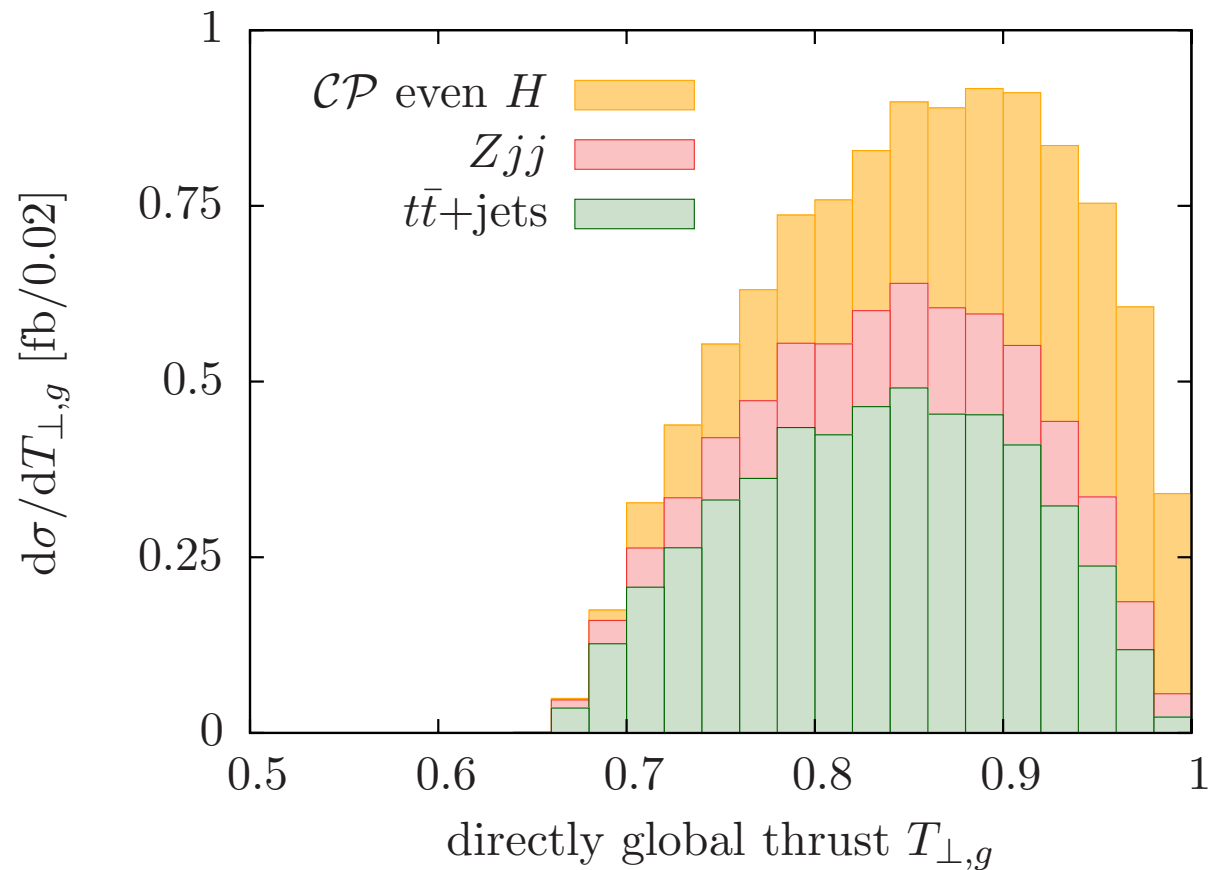
GF vs WBF



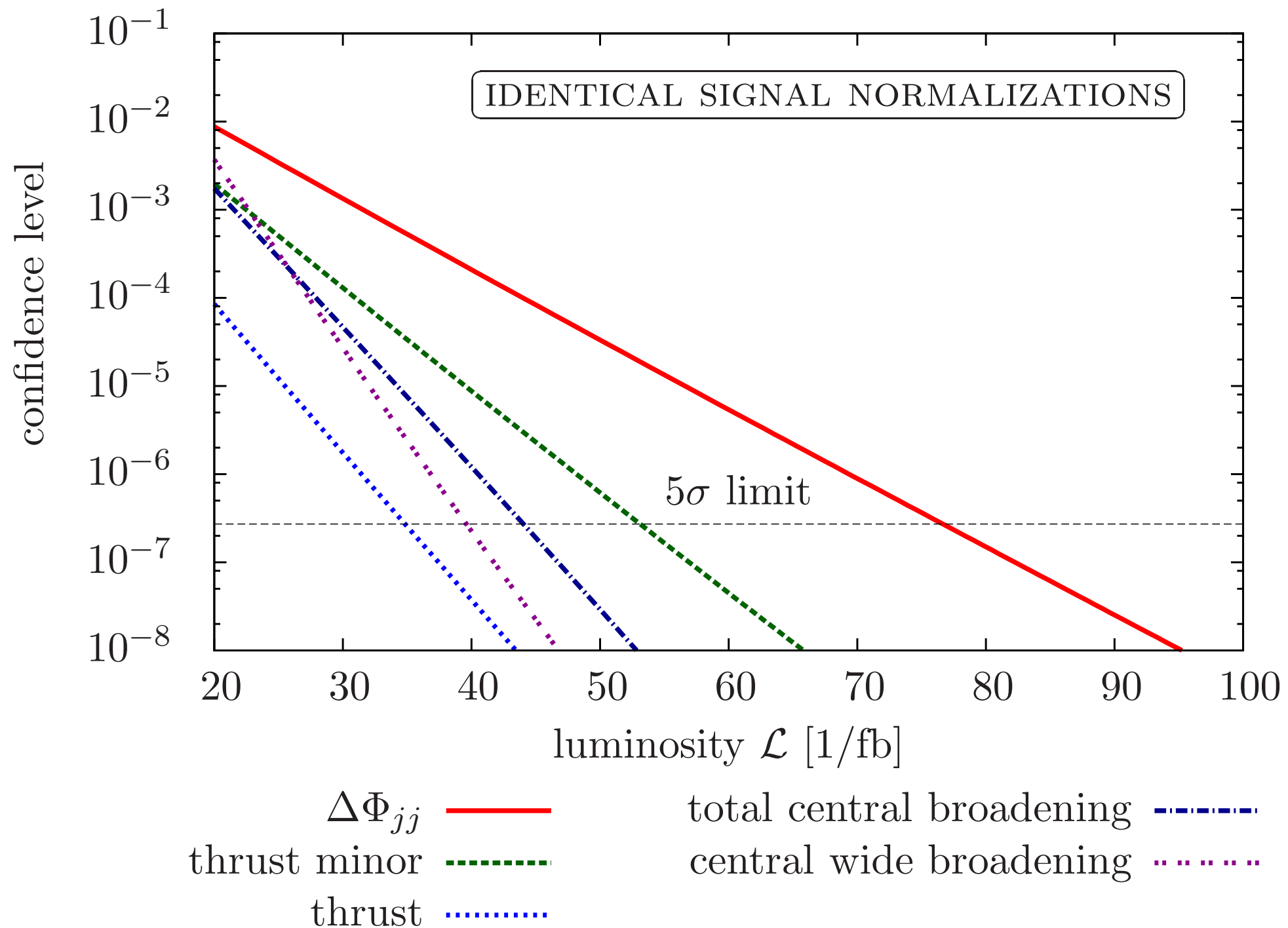
Distributions CP-odd vs CP-even



Distributions CP-odd vs CP-even



Sensitivity for discrimination between CP and CP-odd (normalized signal rates)



IV. Separation of production channels: GF vs WBF

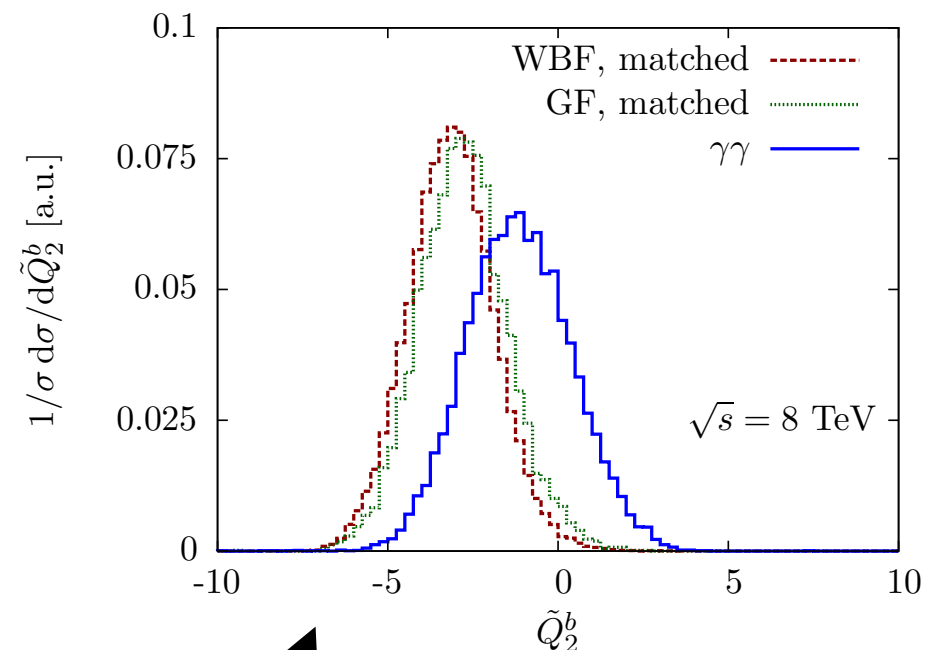
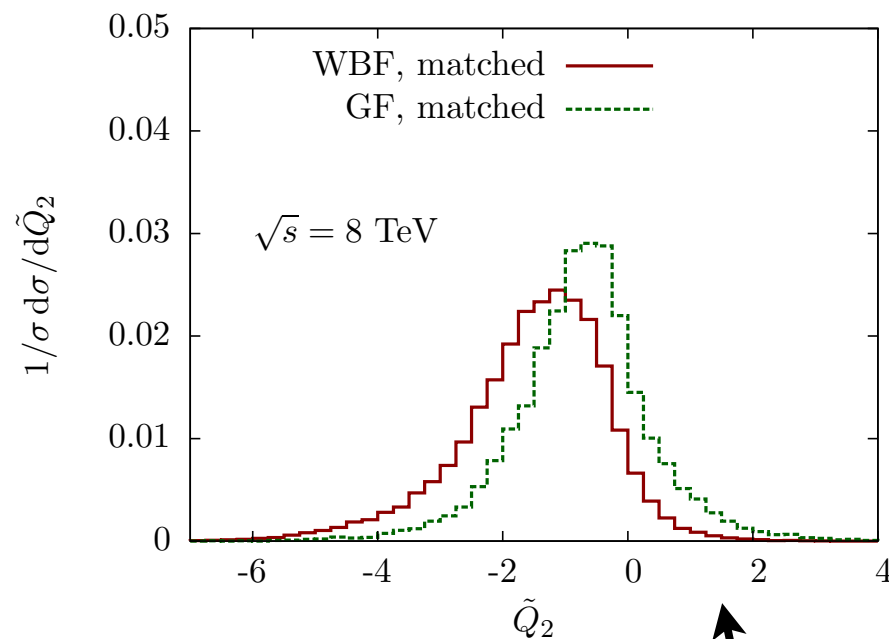
(“Polishing a shiny Higgs with matrix elements”)

[Andersen, Englert, MS
1211.3011]

Possible to construct discriminating observable based on matrix element information:

$$\tilde{Q}_n(p_1^\gamma, p_2^\gamma, \{p_n^j\}) = -\log \left[\frac{|\mathcal{M}^{\text{WBF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2}{|\mathcal{M}^{\text{GF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2} \right]$$

$$\tilde{Q}_n^b(p_1^\gamma, p_2^\gamma, \{p_n^j\}) = -\log \left[\frac{\{|\mathcal{M}^{\text{WBF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2 + |\mathcal{M}^{\text{GF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2\}}{|\mathcal{M}^{2\gamma}(pp \rightarrow \gamma\gamma j^n)|^2} \right]$$

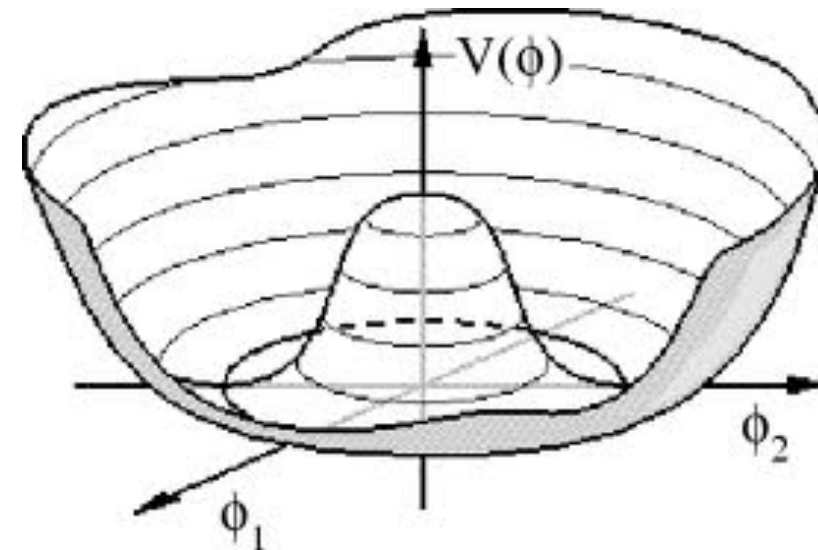


- Method separates GF vs WBF and reduces backgrounds
- Possible to loosen WBF-cuts in 8 TeV analysis to increase statistics and significance for CP measurements

V. Measuring the Higgs selfcoupling at the LHC

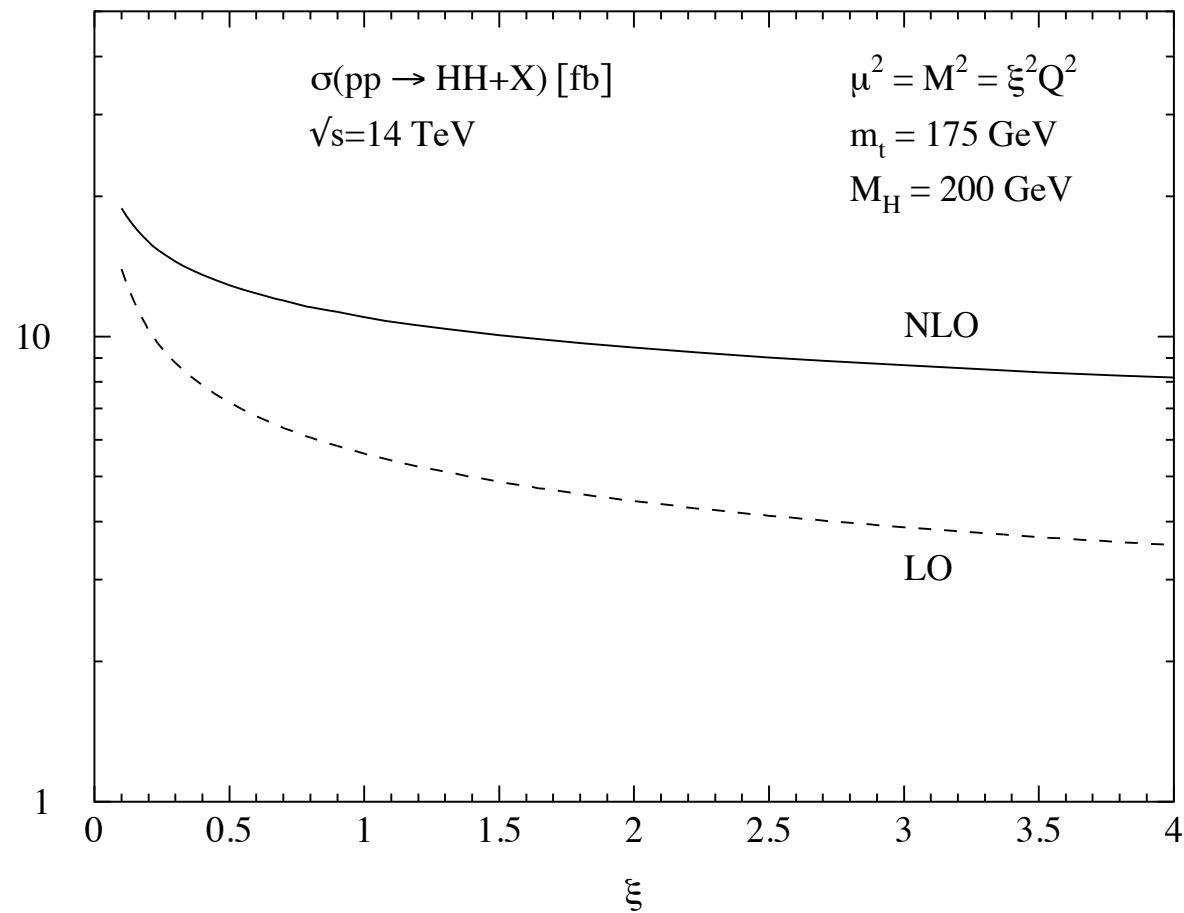
$$V(H^\dagger H) = \mu^2 H^\dagger H + \eta (H^\dagger H)^2$$
$$\supset \frac{1}{2} m_h^2 h^2 + \sqrt{\frac{\eta}{2}} m_h h^3 + \frac{\eta}{4} h^4$$

$$H = (0, v + h)^T / \sqrt{2} \quad \lambda_{\text{SM}} = \sqrt{\frac{\eta}{2}} m_h$$



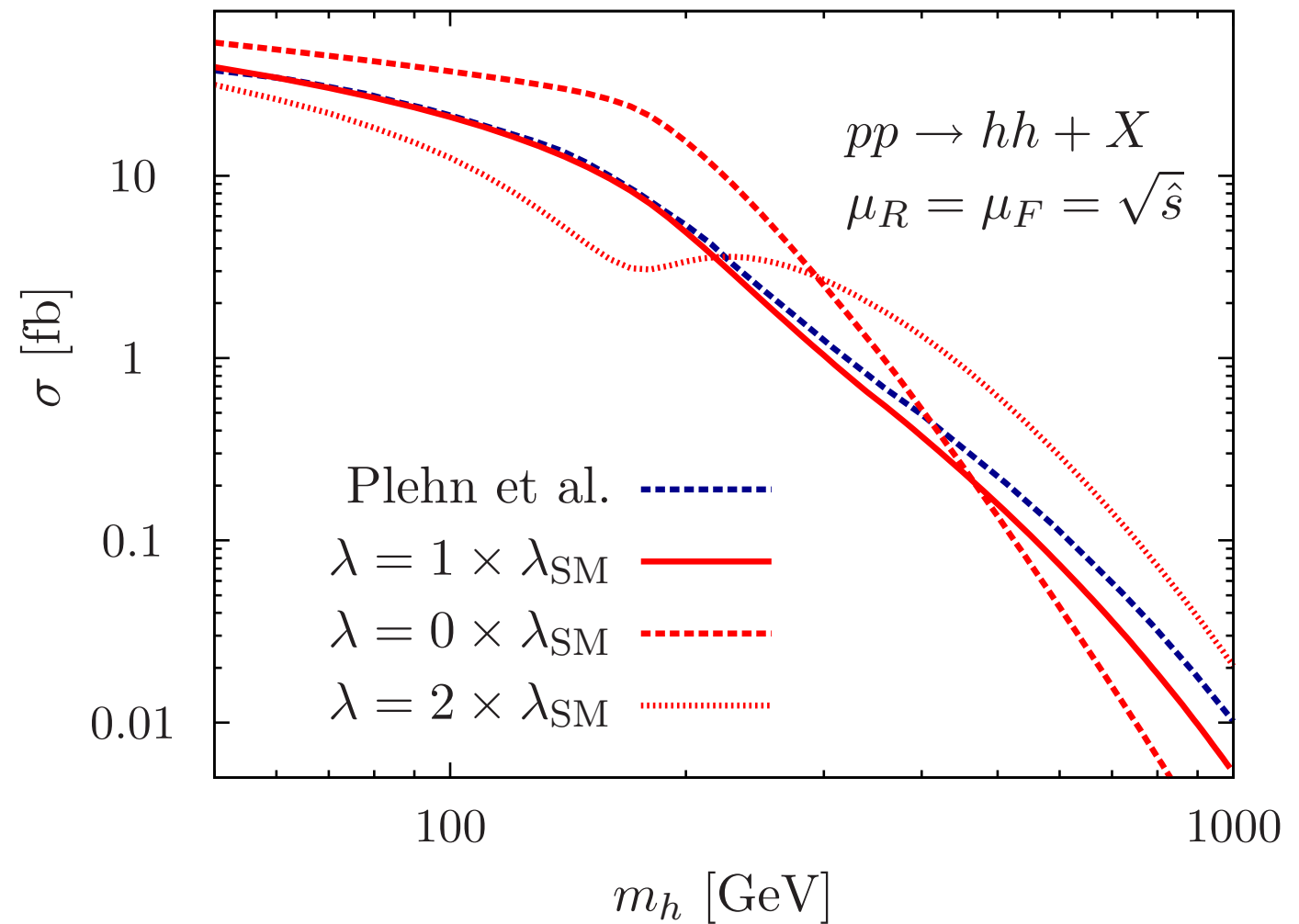
- For EWSB Higgs potential needed → measure selfcoupling
- Measuring Higgs potential 'holy grail' of Higgs physics
- λ_{HHHH} absolutely hopeless at LHC (and any of the others...)
- λ_{HHH} very difficult to measure at the LHC

Rate of di-Higgs production small!

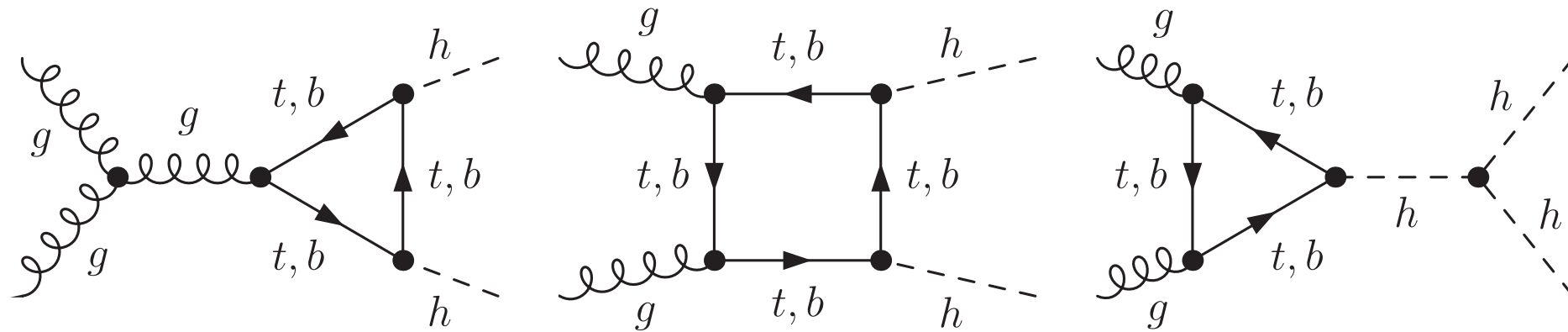


[Dawson, Dittmaier, Spira PRD 58 (1998)]

[Binoth, Karg, Kauer, Ruckl PRD 74 (2006)]



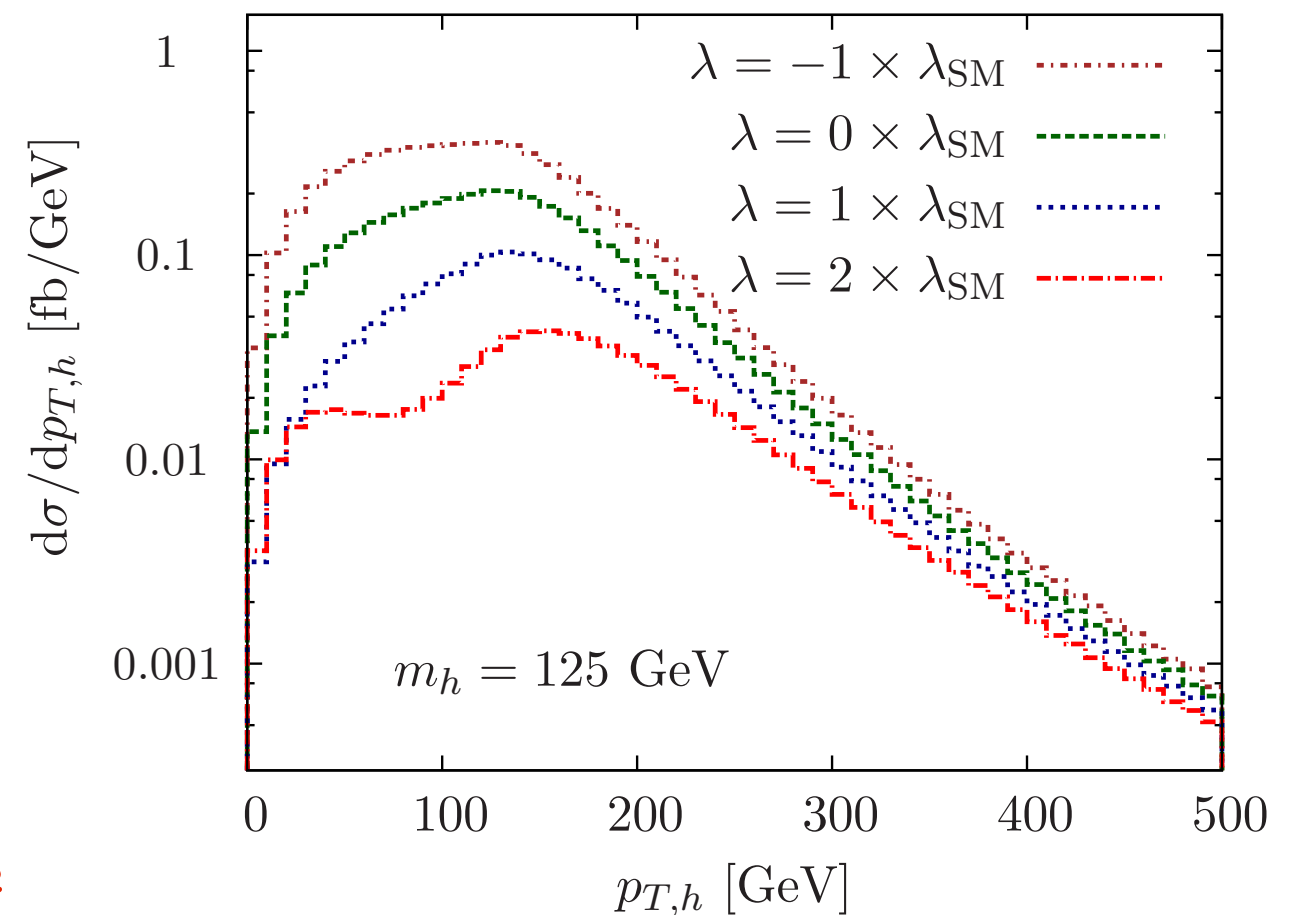
However, kinematic even more problematic



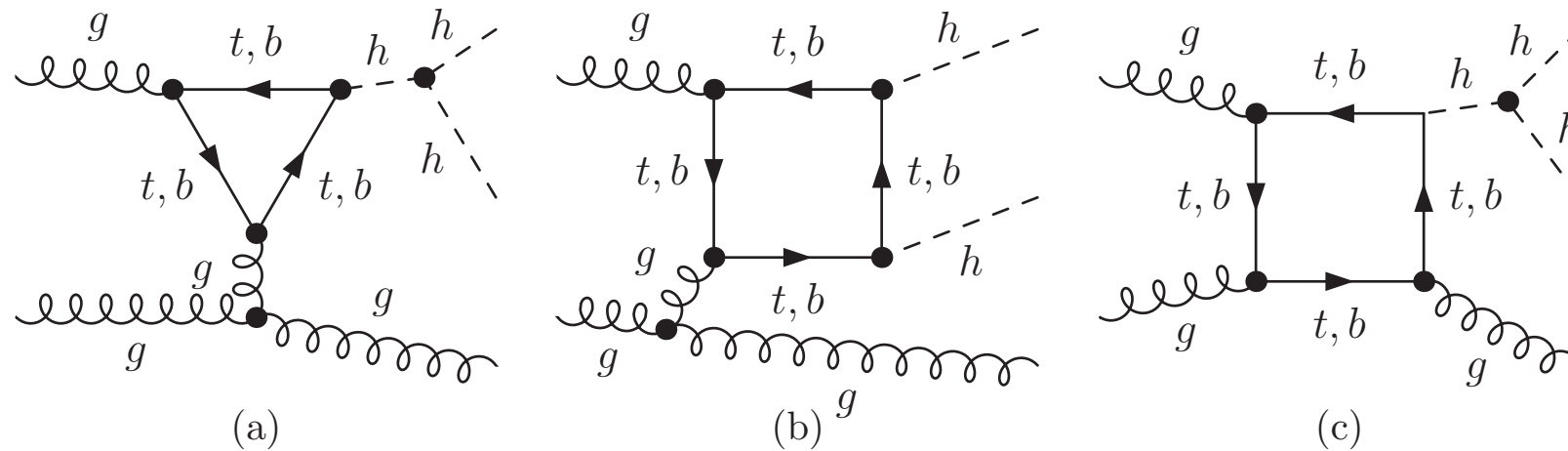
- Destructive interference in signal
- $1/s$ suppression for λ_{HHH} diagram
- But signal is naturally boosted



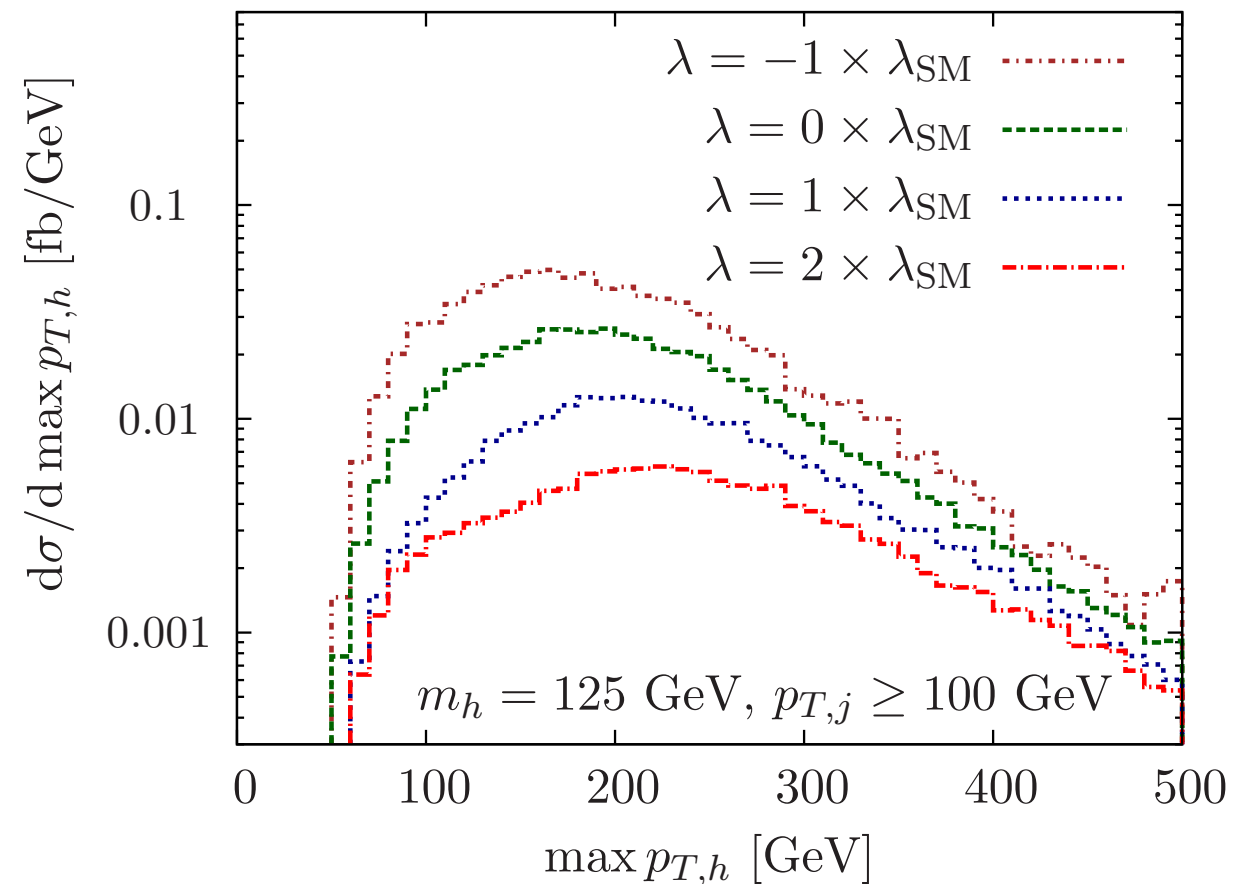
Need large BR due to small
signal but need to discriminate
from backgrounds



Novel idea: include a hard jet to make Higgs softer



- ameliorates the $1/s$ suppression
- rescues sensitivity on selfcoupling for boosted Higgs
- However, comes at the price of even smaller rate

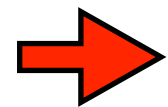


Several reconstruction approaches tried in

[Baur, Plehn, Rainwater PRD 69 (2004)]

[Dolan, Englert, MS JHEP 1210]

[Papaefstathiou, Yang, Zurita 1209.1489]



Most promising final states probably $\bar{b}b\gamma\gamma$, $\bar{b}b\tau^+\tau^-$, $4\tau(?)$

HH:

	$\xi = 0$	$\xi = 1$	$\xi = 2$	$b\bar{b}\tau\tau$	$b\bar{b}\tau\tau$ [ELW]	$b\bar{b}W^+W^-$	ratio to $\xi = 1$
cross section before cuts	59.48	28.34	13.36	67.48	8.73	873000	$3.2 \cdot 10^{-5}$
reconstructed Higgs from τ s	4.05	1.94	0.91	2.51	1.10	1507.99	$1.9 \cdot 10^{-3}$
fatjet cuts	2.27	1.09	0.65	1.29	0.84	223.21	$4.8 \cdot 10^{-3}$
kinematic Higgs reconstruction ($m_{b\bar{b}}$)	0.41	0.26	0.15	0.104	0.047	9.50	$2.3 \cdot 10^{-2}$
Higgs with double b -tag	0.148	0.095	0.053	0.028	0.020	0.15	0.48

HHj:

	$\xi = 0$	$\xi = 1$	$\xi = 2$	$b\bar{b}\tau^+\tau^-j$	$b\bar{b}\tau^+\tau^-j$ [ELW]	$t\bar{t}j$	ratio to $\xi = 1$
cross section before cuts	6.45	3.24	1.81	66.0	1.67	106.7	$1.9 \cdot 10^{-2}$
2 τ s	0.44	0.22	0.12	37.0	0.94	7.44	$4.8 \cdot 10^{-3}$
Higgs rec. from taus + fatjet cuts	0.29	0.16	0.10	2.00	0.150	0.947	$5.1 \cdot 10^{-2}$
kinematic Higgs rec.	0.07	0.04	0.02	0.042	0.018	0.093	0.26
2b + hh invariant mass + $p_{T,j}$ cut	0.010	0.006	0.004	<0.0001	0.0022	0.0014	1.54

Channels are being studied: [ATLAS-PHYS-PUB-2012-001]

Conclusions

- ▶ A new resonance has been found at the LHC, consistent with long predicted Higgs boson
- ▶ To confirm it is Higgs boson couplings and quantum numbers have to be measured
- ▶ Measuring hadronic decays (i.e. $H \rightarrow b\bar{b}$) and Higgs production in association with Jets is paramount for this program
- ▶ Jet substructure can help to perform this task
- ▶ Jet physics will be active field for long time to come

Di-Higgs production can discriminate between different BSM scenarios

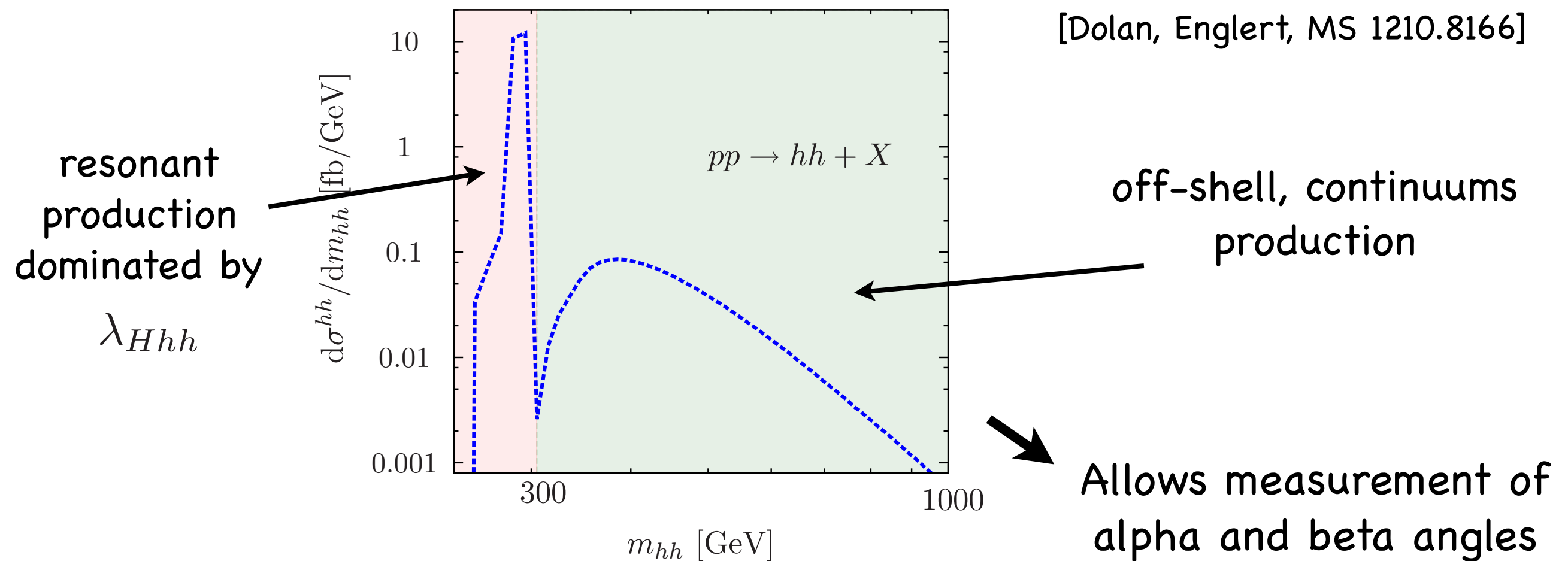
For example resonant di-Higgs production in the MSSM:

$$\lambda_{hhh} = 3 \cos 2\alpha \sin(\beta + \alpha)$$

$$\lambda_{Hhh} = 2 \sin 2\alpha \sin(\beta + \alpha) - \cos 2\alpha \cos(\beta + \alpha)$$

$$\lambda_{HHh} = -2 \sin 2\alpha \cos(\beta + \alpha) - \cos 2\alpha \sin(\beta + \alpha)$$

Assuming decoupling limit such that $M_H > 2 M_h$ and $\text{BR}(H \rightarrow hh) = 45\%$



Higgs mechanism a brief review

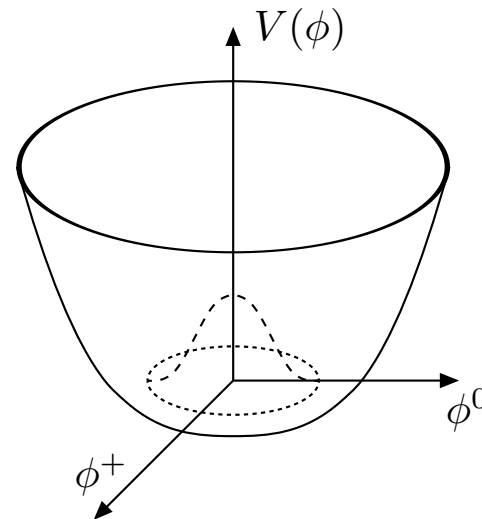
Purpose: explain existence of massive particles consistence with gauge invariance

Symmetry of the Lagrangian

$$SU(2)_L \times U(1)_Y$$

Higgs doublet

$$\Phi = \begin{pmatrix} H^+ \\ H^0 \end{pmatrix}$$



$$V(\Phi) = \lambda \left[\Phi^\dagger \Phi - \frac{v^2}{2} \right]^2$$

Symmetry of the vacuum

$$U(1)_{em}$$

vacuum expectation value

$$\langle \Phi \rangle = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$$

$$v = 246 \text{ GeV}$$

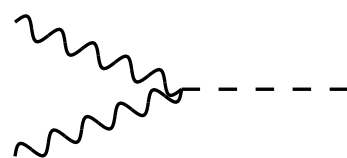
couplings to

Higgs mass: $m_H = \sqrt{2\lambda}v$

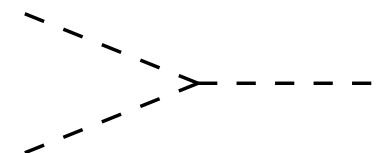
Higgs selfcouplings

gauge bosons

$$g_{VVH} = \frac{2m_V^2}{v}$$

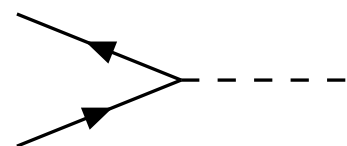


$$\lambda_{HHH} = -3 \frac{m_H^2}{v}$$



fermions

$$g_{f\bar{f}H} = \frac{m_f}{v}$$



$$\lambda_{HHHH} = -3 \frac{m_H^2}{v^2}$$

