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Four-point functions of R-currents in AdS/CFT-correspondence

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- 2 R-currents in $\mathcal{N} = 4$ SYM in the Regge limit
- 3 R-currents in supergravity







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AdS/CFT-correspondence

- AdS/CFT-correspondence is a relation between a conformal field theory in d dimensions and string theory in d+1-dimensional Anti de Sitter space
- the case of interest is d=4: then the CFT is $\mathcal{N} = 4$ SYM theory with gauge group SU(N)and string theory is Type IIB superstring theory in an $AdS_5 \times S_5$ background
- in the low energy limit of string theory supergravity is a good approximation
- $\mathcal{N} = 4$ SYM lives on the 4-dim boundary of AdS_5

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Maldacena, 1997; Witten, 1998
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Duality

- the coupling constant of SYM is g_{YM} and of string theory g_s
- they are related by $g_{YM}^2 = g_s$
- for the supergravity limit one needs $g_s \rightarrow 0$, $N \rightarrow \infty$ and $\lambda = g_s N$ fixed and large
- in a SU(N) gauge theory the limit corresponds to the 't Hooft limit, $N \to \infty$, where $\lambda = g_{eff} = g_{YM}^2 N$ is fixed
- hence the AdS/CFT-correspondence is a duality because it relates the strong coupling regime of CFT to the weakly coupled string theory and vice versa

- an operator \mathcal{O} with conformal dimension Δ in $\mathcal{N} = 4$ SYM, the theory at the boundary, corresponds to a field ϕ of mass *m* in supergravity theory in $AdS_5 \times S_5$
 - the value of ϕ at the boundary is ϕ_0
 - meaning in the gauge theory: field φ₀ is a source for the operator O, the field couples to O via ∫_{S^d} φ₀O

Partition function $Z_{\mathcal{O}}[\phi_0]$

$$Z_{\mathcal{O}}[\phi_0] = \int \mathcal{D}[\mathsf{SYM fields}] \exp(-S_{\mathcal{N}=4\,SYM} + \int d^4 x \mathcal{O}(x)\phi_0(x))$$
$$\Rightarrow \langle \mathcal{O}(x_1) \dots \mathcal{O}(x_n) \rangle = \frac{\delta^n}{\delta \phi_0(x_1) \dots \delta \phi_0(x_n)} Z_{\mathcal{O}}[\phi_0]|_{\phi_0=0}$$

Witten prescription

$$= \int \mathcal{D}[SYM \, fields] \exp(-S_{\mathcal{N}=4SYM} + \int d^4 x \mathcal{O}(x)\phi_0(x))$$

= $Z_{class}[\phi_0]_{AdS} = \exp(-S_{sugra}[\phi[\phi_0]])$

n-point functions in AdS/CFT

$$\langle \mathcal{O}(\mathbf{x}_1) \dots \mathcal{O}(\mathbf{x}_n) \rangle = \frac{\delta^n}{\delta \phi_0(\mathbf{x}_1) \dots \delta \phi_0(\mathbf{x}_n)} \exp(-S_{\text{sugra}}[\phi[\phi_0]])|_{\phi_0=0}$$

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Witten diagrams

- computing n-point functions means evaluating Witten diagrams
- three-point function as one example:



• the boundary of AdS₅ is represented by a circle, wavy lines are gauge boson propagators

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• two kind of propagators: bulk-to-boundary propagators, between an interior point and a boundary point, and bulk-to-bulk propagators, between two interior points



2 R-currents in $\mathcal{N} = 4$ SYM in the Regge limit

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Summary

$\mathcal{N} = 4 \text{ SYM}$

- non-Abelian gauge theory in four dimensions with $\mathcal{N}=4$ supersymmetry
- contains a vector multiplet in the adjoint representation of the gauge group $SU(N_C)$
- theory offers a $SU_R(4)$ global symmetry called R-symmetry \Rightarrow R-currents
- the fields are
 - 1 vector field A_{μ} , scalar of $SU_R(4)$
 - 4 chiral spinors λ_l in the fundamental representation of $SU_R(4)$
 - 6 real scalars X_M in the vector representation of $SU_R(4)$

Scattering of R-currents (1)

 \bullet interested in four-point correlation functions in $\mathcal{N}=4$ SYM

Momentum space four-point function

$$i(2\pi)^{4}\delta(p_{A}+p_{B}-p_{A'}-p_{B'})A_{R}(s,t) = \int \prod_{i} d^{4}x_{i}e^{-ip_{A}\cdot x_{A}-ip_{B}\cdot x_{B}+ip_{A'}\cdot x_{A'}+ip_{B'}\cdot x_{B'}} \times \langle J_{R}^{\mu A}(x_{A})J_{R}^{\nu B}(x_{B})J_{R}^{\mu' A'}(x_{A'})J_{R}^{\nu' B'}(x_{B'})\rangle$$

- corresponds to scattering of R-currents $J_R^{\mu A}(x)$
- $\bullet\,$ R-current scattering is the analog of $\gamma^*\gamma^*\text{-scattering}$ in QCD

$\gamma^*\gamma^*$ -scattering in QCD

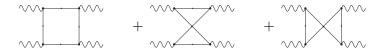
- elastic scattering of two virtual photons at high energies
- each photon splits into a quark-antiquark pair
- the decay of the photons is mediated by the electromagnetic current j^{μ}
- leads to the computation of the four-point function of j^{μ}



- the computation is performed in the Regge limit: $s \rightarrow \infty$ and t < 0, fixed
- in this limit: conjectured duality of the BFKL pomeron and the graviton, Kotikov, Lipatov et al., 2005 integrability of BKP-states



• in the Regge limit one-loop diagrams are subleading $\sim \log^2 s$



• three-loop diagrams give the leading contribution to the amplitude $\sim \alpha_s^2 s$



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Impact factors

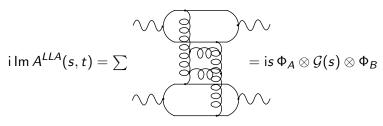
Amplitude for three-loop diagrams

$$A(s,t) = is \int \frac{d^2\mathbf{k}}{(2\pi)^2} \frac{\Phi_A^{\lambda_A \lambda_{A'} a a'}(\mathbf{k}, \mathbf{q} - \mathbf{k}) \Phi_B^{\lambda_B \lambda_{B'} a a'}(\mathbf{k}, \mathbf{q} - \mathbf{k})}{\mathbf{k}^2 (\mathbf{q} - \mathbf{k})^2}$$

- the leading term of the amplitude is purely imaginary
- the amplitude A(s, t) factorizes into two impact factors
- the impact factors are

LLA resummation

• Leading-Logarithmic Approximation is summing up radiative corrections to A(s, t)



- $\bullet~\mathcal{G}(s)$ is the Green's function of the BFKL equation
- the four-point function satisfies Regge factorization in LLA
- the two exchanged gluons become reggeized

• the BFKL Pomeron is the bound state of the reggeized gluons Balitsky, Fadin, Kuraev, Lipatov

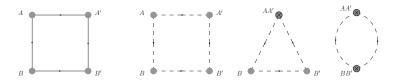
Scattering of R-currents (2)

• scattering of R-currents means evaluation of the amplitude

$$A_R(s,t) \sim \langle J_R^{\mu A}(x_A) J_R^{
u B}(x_B) J_R^{\mu' A'}(x_{A'}) J_R^{
u' B'}(x_{B'}) \rangle$$

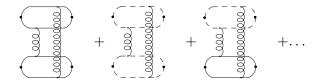
- \bullet fermions and scalars provide contributions to the amplitude in $\mathcal{N}=4~\text{SYM}$
- **QUESTIONS** for R-current scattering:
 - Which diagrams are the leading ones?
 - Does the amplitude $A_R(s, t)$ factorize? If so:
 - What are the impact factors?

One-loop diagrams



- contributions from fermions and scalars to the correlation function, altogether 3 fermionic and 12 scalar diagrams
- ullet one-loop diagrams are again subleading at high energy $\sim \log^2\!s$
- one-loop diagrams are UV finite because of supersymmetry, the divergencies of the fermionic and the scalar part cancel each other

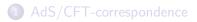
Impact factors



- at high energy in $\mathcal{N}=4$ SYM three-loop diagrams are dominating $\sim s$
- the amplitude factorizes, one gets separately fermion impact factors Φ_F and scalar impact factors Φ_S
- \bullet additional scalar diagrams like $\left(\frac{1}{2} \frac{1}{2} \right)^2$ are suppressed $\sim 1/s^2$
- the scalar impact factor is

Results

- there are four impact factors $\Phi = \Phi_F + \Phi_S$ to compute: $\Phi^{LL'}, \Phi^{TT'}, \Phi^{LT'}$ and $\Phi^{TL'}$
- the results for $\Phi^{LL'}$ and $\Phi^{TT'}$ have got the same structure as the impact factors of $\gamma^*\gamma^*\text{-scattering}$
- a new feature is the vanishing of the helicity changing impact factors: $\Phi^{LT'} = \Phi^{TL'} = 0$
- the simplification is due to supersymmetry: the fermionic and scalar parts cancel each other



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Partition function for J_R

- SU(4) R-symmetry currents J^{μA}_R of SYM are gauge invariant operators of type O
- they correspond to SU(4) gauge fields A^{μA} on the supergravity side with the value A^{μA}₀ at the boundary
- one gets the partition function

$$Z_{J_{R}^{A}}[A_{0}] = \int \mathcal{D}[\text{SYM fields}]\exp(-S_{\mathcal{N}=4} \text{SYM} + \int d^{4}x J_{R}^{A} A_{0}^{A})$$
$$= \exp(-S_{sugra}[A^{A}[A_{0}^{A}]])$$

Witten diagrams



- Witten diagrams for the four-point function of R-currents
- there are only tree diagrams in the supergravity limit
- the currents are inserted at points of the boundary of AdS_5
- the interaction takes place in the interior
- one has boundary-to-bulk propagators for the gauge boson and bulk-to-bulk propagators for the gauge field and the graviton



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- we computed the four-point function of R-currents in ${\cal N}=4$ SYM in the Regge limit, computation of impact factors
- next steps:

evaluation of corresponding Witten diagrams,



generalization to, e.g. six-point functions of R-currents (BKP Green's function, integrability)