

## 1 Motivation

- Last decade: Overwhelming evidence that planar  $\mathcal{N} = 4$  super Yang-Mills theory is integrable [Minahan, Zarembo 2003] [Beisert, Kristjansen, Staudacher 2003] [Gromov, Kazakov, Vieira 2009]
  - Access non-perturbative regimes of a 4-dimensional quantum field theory
  - Recently: Remarkable closed formula for scattering amplitudes of the theory
  - Amplitudes are invariant under action of  $\infty$ -dimensional Yangian algebra
- ⇒ How to use standard methods of integrable systems to construct amplitudes?
- First step: We show that certain Yangian invariants can be obtained from a Bethe ansatz
  - Fruitful exchange of ideas between amplitudes and integrable systems

## 2 Basic Methods and Structures in Integrable Systems

### Quantum Inverse Scattering Method [Faddeev et al. 1980s]

- Powerful algebraic approach to integrable systems
- Main problem: Diagonalize the Hamiltonian of a spin chain
- Central object:  $R$ -matrix satisfying a Yang-Baxter equation

$$\begin{array}{c} \text{---} b \\ \diagup \quad \diagdown \\ c \quad a \end{array} = R_{ab}(z-y) R_{ac}(z) R_{bc}(y) = R_{bc}(y) R_{ac}(z) R_{ab}(z-y) = \begin{array}{c} \text{---} b \\ \diagdown \quad \diagup \\ c \quad a \end{array}$$

- Introduce monodromy and transfer matrix for a spin chain with  $N$  sites

$$M_a(z) = R_{aN}(z - v_N) \cdots R_{a1}(z - v_1) = \begin{array}{c} \text{---} a \\ | \quad | \quad | \\ N \quad \cdots \quad 2 \quad 1 \end{array}$$

$$T(z) = \text{tr}_a M_a(z) = M_{11}(z) + \dots + M_{mm}(z) = \begin{array}{c} \text{---} a \\ | \quad | \quad | \\ N \quad \cdots \quad 2 \quad 1 \end{array}$$

where  $R(z) = z1 + E_{\alpha\beta} \otimes J_{\beta\alpha}$  with  $\mathfrak{gl}_m$ -generators  $J_{\alpha\beta}$  and  $(E_{\alpha\beta})_{\mu\nu} = \delta_{\alpha\mu} \delta_{\beta\nu}$

⇒ Recover Hamiltonian as logarithmic derivative of transfer matrix

### Bethe Ansatz

- Method to diagonalize the transfer matrix
- Example: algebra  $\mathfrak{gl}_2$  and representation with  $\mathfrak{gl}_2$ -weights  $\Lambda = (\lambda, 0)$
- Bethe roots  $u_1, \dots, u_S$  are solutions of Bethe equations [Bethe 1931]

$$\prod_{i=1}^N \frac{u_k - v_i}{u_k - v_i + \lambda} = - \prod_{i=1}^S \frac{u_k - u_i + 1}{u_k - u_i - 1}$$

⇒ Bethe vector  $\Psi$  is an eigenvector of the transfer matrix given in terms of matrix elements of the monodromy acting on a vacuum state

$$\Psi = M_{12}(u_1) \cdots M_{12}(u_S) \Omega$$

### Yangian [Drinfeld 1985]

- Symmetry algebra of many integrable systems
- Yangian  $\mathcal{Y}(\mathfrak{gl}_m)$  is defined by the Yang-Baxter-like relation

$$R_{ab}(z-y) M_a(z) M_b(y) = M_b(y) M_a(z) R_{ab}(z-y)$$

where the  $\mathcal{Y}(\mathfrak{gl}_m)$ -generators are obtained from the expansion

$$M_{\alpha\beta}(z) = \delta_{\alpha\beta} 1 + M_{\alpha\beta}^{(1)} z^{-1} + M_{\alpha\beta}^{(2)} z^{-2} + \dots$$

⇒ Monodromy of a spin chain provides a realization of the Yangian

## 3 Details on the Bethe Ansatz for Yangian Invariants

### Characterization of Yangian Invariants

- Aim: Construct vectors  $\Psi$  annihilated by all generators of Yangian  $\mathcal{Y}(\mathfrak{sl}_m)$
- Invariants are characterized by condition in terms of  $\mathcal{Y}(\mathfrak{gl}_m)$ -generators

$$M_{\alpha\beta}(z) \Psi \propto \delta_{\alpha\beta} 1 \Psi$$

- Characterization implies directly that invariants are special eigenvectors of the transfer matrix

⇒ Construct invariants  $\Psi$  using Bethe ansatz

### Bethe Ansatz

- Example: algebra  $\mathfrak{gl}_2$  and a monodromy with even number of sites  $N = 2L$
- Need at least two different kinds of sites:  $L$  sites with representation  $\Lambda = (\lambda, 0)$  and  $L$  sites with dual representation  $\Lambda^* = (0, -\lambda)$
- Need special choice of inhomogeneities:  $v_{L+i} = v_i - \lambda - 1$

$$M_a(z) = \begin{array}{c} \Lambda^* \quad \cdots \quad \Lambda^* \quad \Lambda \quad \cdots \quad \Lambda \\ | \quad | \quad | \quad | \quad | \\ 2L \quad \cdots \quad L+1 \quad L \quad \cdots \quad 1 \end{array} a$$

⇒ Bethe equations can be solved explicitly!

- Bethe roots are  $u_{i+(j-1)L} = v_i - j$  where  $i = 1, \dots, L$  and  $j = 1, \dots, \lambda$

$$\begin{array}{c} \lambda \\ \cdot \quad \cdots \quad \cdot \\ u_{i+(\lambda-1)L} \quad u_{i+L} \quad u_i \quad v_i \end{array}$$

- Bethe vector  $\Psi$  is a Yangian invariant
- Construction can be generalized
- Arbitrary order of sites and odd number of sites
- Nested Bethe ansatz for algebra  $\mathfrak{gl}_m$

### Graphical Interpretation

- Understand Yangian invariant  $\Psi$  as “network” of  $R$ -matrices
- Pair of sites  $i, j$  with representations  $\Lambda, \Lambda^*$  and inhomogeneities  $v_i, v_j = v_i - \lambda - 1$  is connected by a line
- Intersection of two lines corresponds to an  $R$ -matrix satisfying a Yang-Baxter equation and “unitarity and crossing relations”

$$\begin{array}{c} \diagup \quad \diagdown \\ \diagdown \quad \diagup \end{array} = \begin{array}{c} \diagdown \quad \diagup \\ \diagup \quad \diagdown \end{array} \quad \text{---} \propto \text{---}$$

⇒ Properties of  $R$ -matrices imply immediately Yangian invariance of  $\Psi$

$$M_{\alpha\beta}(z) \Psi = \begin{array}{c} \text{---} \beta \\ | \quad | \quad | \\ N \quad \cdots \quad 1 \end{array} \propto \dots \propto \begin{array}{c} \text{---} \beta \\ | \quad | \quad | \\ N \quad \cdots \quad 1 \end{array} \alpha = \delta_{\alpha\beta} 1 \Psi$$

## 4 Main Results

- Invariants of Yangian  $\mathcal{Y}(\mathfrak{sl}_m)$  with certain irreducible  $\mathfrak{gl}_m$ -representations at each site of the monodromy can be constructed using Bethe ansatz
- Generalization of perimeter Bethe ansatz [Baxter 1987]: Invariants can be understood as partition functions of vertex models on random lattices
- Progress towards more explicit “coordinate expressions” of Bethe vectors for  $\mathfrak{gl}_3$  with mixed representations  $\Lambda$  and  $\Lambda^*$

## 5 Towards Super Yang Mills Amplitudes

### Scattering amplitudes

- Tree-level amplitudes for  $N$  particles in sector with  $K$  negative helicity gluons [Arkani-Hamed et al. 2009]

$$\int D^{K(N-K)} c \prod_{a=1}^K \prod_{\alpha=1}^8 \delta \left( \sum_{i=1}^N c_{ai} (W_\alpha)_i \right)$$

- $W_\alpha$  are supertwistor variables transforming in the fundamental representation of the superconformal algebra  $\mathfrak{psu}_{2,2|4}$
- Integration along contour in Grassmannian  $G(K, N) = \{K\text{-planes in } \mathbb{C}^N\}$
- Amplitudes are invariant under Yangian  $\mathcal{Y}(\mathfrak{psu}_{2,2|4})$  [Drummond et al. 2009]

### Outlook

- How to obtain amplitudes from a Bethe ansatz?
- ⇒ Are amplitudes “generating functions” for Yangian invariants with irreducible representations at each site?
- Excitingly Grassmannians play an important role in different areas of integrable systems
- $\tau$ -functions of KP-hierarchy [Sato 1981], partition function of vertex models [Foda et al. 2009], master  $T$ -operator [Alexandrov et al. 2011]
- ⇒ Relation to the Grassmannian appearing in scattering amplitudes?

## Benefits from Graduiertenkolleg

- Financial support to attend international conferences and schools
- Block courses with lectures on various topics in experimental and theoretical particle physics help to broaden the perspective

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