

The leading-order hadronic contribution to $(g-2)_{\mu}$ from $N_f=2+1+1$ twisted mass lattice QCD

Motivation

- ► anomalous magnetic moment of the muon, a_µ, serves as benchmark test of standard model of elementary particle physics (SM)
- ▶ key ingredient: leading-order hadronic vacuum-polarisation contribution a^{hvp}_µ
- -intrinsically nonperturbative \rightarrow lattice QCD computation highly desirable
- must be responsible for discrepancy since QED and weak contributions reliably computed in perturbation theoy
- largest source of uncertainty in SM computation



Method

- twisted mass lattice regularisation of QCD: alleviates discretisation effects by improving the continuum limit behaviour of *all* physical observables
- calculations based on configurations with four dynamical quark flavours generated by European Twisted Mass Collaboration (ETMC)
- computation of a^{hvp}_µ follows closely strategy of refs. [1,2] using an improved lattice definition of this quantity



[Feng, Jansen, Petschlies, Renner, Phys. Rev. Lett. **107**, 201

More Details

$$\begin{array}{l} & \lambda_{\mu}^{\text{hvp}} \text{ can be computed in Euclidean} \\ & \lambda_{\mu}^{\text{hvp}} = \alpha^2 \int_0^\infty \frac{dQ^2}{Q^2} w \left(\frac{Q^2}{m_{\mu}^2}\right) \Pi_{\text{R}}(Q^2) \\ & \mu \end{array}$$

$$\begin{array}{l} & \mu \\ & \mu \end{array}$$

$$\begin{array}{l} & \Pi_{\text{R}}(Q^2) = \Pi(Q^2) - \Pi(0) \text{ is renormalised hadronic vacuum polarisation} \end{array}$$

can be obtained from transverse part of hadronic vacuum polarisation tensor

$$\Pi_{\mu\nu}(Q) = \int d^4x e^{iQ \cdot (x-y)} \langle J_{\mu}^{\rm em}(x) J_{\nu}^{\rm em}(y) \rangle = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\mu\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) = (Q_{\mu}Q_{\mu\nu} - Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu}) \Pi(Q^2 + Q^2 \delta_{\mu\nu})$$

correlator of electromagnetic vector currents

[Hoecker, arXiv:1201.5093 [hep-ph], 2011]

 $J^{\rm em}_{\mu}(x) = \frac{2}{3}\overline{u}(x)\gamma_{\mu}u(x) - \frac{1}{3}\overline{d}(x)\gamma_{\mu}d(x) + \frac{2}{3}\overline{c}(x)\gamma_{\mu}c(x) - \frac{1}{3}\overline{s}(x)\gamma_{\mu}s(x)$

► on the lattice: point-split vector current \rightarrow conserved, satisfies lattice Ward-Takahashi identity $\partial_{\mu}^{*}J_{\mu}^{C}(x) = 0$

► to weaken quark mass dependence use redefinition

$$a_{\overline{\mu}}^{\text{hvp}} = \alpha^2 \int_0^\infty \frac{dQ^2}{Q^2} w \left(\frac{Q^2 H_{\text{phys}}^2}{H^2 m_{\mu}^2}\right) \Pi_{\text{R}}(Q^2)$$

► H is either H = 1 (standard definition) or any hadronic scale which can be computed at unphysical values of pion mass m_{PS}

For
$$m_{PS} o m_{\pi}$$
, $H o H_{
m phys}$, i.e. by definition $a_{\overline{\mu}}^{
m hvp} o a_{\mu}^{
m hvp}$

 $\blacktriangleright H = m_V$, mass of the $\rho\text{-meson}$, has proven to be especially beneficial

Preliminary Results



function

▶ only mass-degenerate up and down quarks in valence sector (left figure): value for light quark contribution to a^{hvp}_µ found for N_f = 2 in refs. [1,2] reproduced

very first lattice results worldwide for four-flavour contribution



- to a_{μ}^{hvp} (right figure)
- promising prospect for future experiments
- full agreement with result of dispersive analysis by Benayoun, David, DelBuono, Jegerlehner (Eur.Phys.J. C72 (2012) 1848) found
- systematic uncertainties remain to be quantified

Publications

[1] X. Feng, K. Jansen, M. Petschlies and D. B. Renner, "Two-flavor QCD correction to lepton magnetic moments at leading-order in the electromagnetic coupling" Phys. Rev. Lett. **107** (2011) 081802 [arXiv:1103.4818 [hep-lat]].

[2] D. B. Renner, X. Feng, K. Jansen and M. Petschlies, "Nonperturbative QCD corrections to electroweak observables", arXiv:1206.3113 [hep-lat].

[3] X. Feng, G. Hotzel, K. Jansen, M. Petschlies and D. B. Renner, "Leading-order hadronic contributions to a_{μ} and α_{QED} from $N_f = 2 + 1 + 1$ twisted mass fermions", PoS LATTICE2012 (2012) 174, [arXiV:1211.0828 [hep-lat]].

Selected Talks

"The leading-order hadronic contribution to the anomalous magnetic moment of the muon from twisted mass lattice QCD", seminar given at INFN Laboratori Nazionali di Frascati, Italy, January 2012

 \blacktriangleright "Leading-order hadronic contributions to a_{μ} and α_{QED} from $N_f = 2 + 1 + 1$ twisted mass fermions", Lattice 2012, Cairns, Australia, July 2012

Collaborations

Profit from the GK

- Xu Feng (KEK, Tsukuba, Japan), Dru B. Renner (Jefferson Laboratory, Newport News, USA), Marcus Petschlies (The Cyprus Institute, Nicosia, Cyprus)
- European Twisted Mass Collaboration (ETMC)
- ► SFB TR9 "Computational Particle Physics"

- \blacktriangleright Participation in Lattice 2012, the 30^{th} International Symposium on Lattice Field Theory
- Broader knowledge of the field due to discussions beyond the own group and informative lectures at the last block course

Contact Details and Further Information

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