

Landau gauge gluon and ghost propagators at non-zero temperature from lattice QCD



- Investigating the phase diagram of QCD at T > 0 on the lattice with the help of gluon and ghost propagators in Landau gauge.
- Providing ready-to use input data for Dyson-Schwinger or functional renormalization group eqs. for semi-analytic nonperturbative studies of the fermionic sector of QCD.
- Compare pure gauge theory, i.e. quenched QCD (\Rightarrow first order phase transition) with two-flavour QCD (\Rightarrow crossover phenomenon).
- Compute MOM-scheme renormalized gluon and ghost propagators in the non-perturbative range [0.4, 3.0] GeV across phase transition (crossover) at fixed pion mass values $m_{\pi} \simeq 320, 400, 470$ MeV and ∞ , resp.
- In the quenched case we study systematic effects: finite-size and Gribov copy effects, extrapolation to continuum limit.
- For full QCD provide proper fits of the momentum dependence for various temperature

Methods

- Lattice discretization of QCD with twisted mass formulation for fermionic action.
- Propagators computed from path integral represention with (hybrid) Monte Carlo method.
- Gauge potentials in terms of lattice link variables: $A_{\mu}(x + \hat{\mu}/2) = \frac{1}{2iag_0}(U_{x\mu} U_{x\mu}^{\dagger})|_{traceless}$
- Landau gauge by maximizing $F_U[g] = \frac{1}{3} \sum_x \sum_\mu \Re \operatorname{e} \operatorname{Tr} \left(g_x U_{x\mu} g_{x+\mu}^{\dagger} \right)$ w. r. to $g_x \in SU(3)$ employing simulated annealing + overrelaxation.
- Transverse and longitudinal gluon propagator in *q*-space (Matsubara frequency $\sim q_4 \equiv 0$): $D_T = \frac{1}{(d-2)8} \left\langle \sum_{i=1}^3 A_i^a(q) \ A_i^a(-q) - \frac{q_4^2}{\vec{q}^2} \ A_4^a(q) \ A_4^a(-q) \right\rangle.$ $D_L = \frac{1}{8} (1 + \frac{q_4^2}{\vec{q}^2}) \left\langle A_4^a(q) \ A_4^a(-q) \right\rangle$
- Ghost propagator G and its dressing function J from Faddeev-Popov operator $M = -\partial D$:



and pion mass values.

$G^{ab}(q) = a^2 \sum_{x,y} \langle e^{-2\pi i k \cdot (x-y)/L} [M^{-1}]_{xy}^{ab} \rangle = \delta^{ab} \ G(q) = \delta^{ab} \ \frac{J(q)}{q^2}.$

Some Details

Standard lattice plaquette action for pure $SU(3)\ {\rm gauge}\ {\rm theory}$

$$S_G = \beta \sum_{x} \sum_{\mu > \nu} \left[1 - \frac{1}{3} \operatorname{\mathfrak{Re}} \operatorname{Tr} \left(U_{x\mu} U_{x+\mu;\nu} U_{x+\nu;\mu}^{\dagger} U_{x\nu}^{\dagger} \right) \right], \ \beta = 6/g_0^2.$$

Symanzik-improved gauge action

$$S_G = \beta \sum_x [c_0 \sum_{\mu < \nu} (1 - \frac{1}{3} \operatorname{\mathfrak{Re}} \operatorname{Tr} U_{x\mu\nu}^{1 \times 1}) + c_1 \sum_{\mu \neq \nu} (1 - \frac{1}{3} \operatorname{\mathfrak{Re}} \operatorname{Tr} U_{x\mu\nu}^{1 \times 2})],$$

and twisted-mass fermionic action for full QCD

$$S_F = a^4 \sum_x \bar{\psi}(x) \left[(D[U] + m_0) + i \,\mu \, au_3 \,\gamma_5 \right] \psi(x).$$

Recent Results

Quenched QCD:



Quenched QCD: Systematic effects extrapolation of D_L and D_T to the continuum limit $(a \rightarrow 0)$:



Full two-flavour QCD:



Lattice momenta ((k, k, k, 0), k = 1, 2, 3), $\beta = 6.337$ and spatial size $N_{\sigma} = 48$ are fixed. Normalization applied with $J(q, T_{min})$, $T_{min} = 0.65 T_c$. Lowest panel corresponds to k = 1.

Strong (smoother) response for $D_L(q)$ $(D_T(q))$ at $T \simeq T_c$. Therefore, D_L good indicator for the deconfinement transition. J(q) very weakly temperature dependent at T_c .



Same as above but for D_T . l.h.s. $T = 0.86 T_c$; r.h.s. $T = 1.20 T_c$.

 $\blacksquare D_L$ and D_T nicely reach the continuum limit.

Finite volume and Gribov copy effects were seen to be negligible in the momentum range considered.

 $R_T(q,T) = D_T(q,T)/D_T(q,T_{\min})$ (left), $R_L(q,T) = D_L(q,T)/D_L(q,T_{\min})$ (middle) $R_G(q,T) = G(q,T)/G(q,T_{\min})$ (right) versus T for a few momentum values p (in [GeV]). Pion masses values (from top to bottom) are $m_{\pi} \simeq 316$, 398 and 469 MeV. Lattice size $32^3 \times 12$, spacing $a \le 0.09$ fm. Vertical bands indicate the chiral (T_{χ}) and deconfinement (T_{deconf}) pseudo-critical temperatures with their uncertainties.

- Quite smooth behaviour throughout the crossover region.
- D_L more sensitive than D_T , G again weakly T-dependent. D_L seems to behave smoother for decreasing m_{π} .

Publications

Talks and Posters

[1] R. Aouane, V.G. Bornyakov, E.-M. Ilgenfritz, V.K. Mitrjushkin, M. Müller-Preussker and A. Sternbeck, *Landau gauge gluon and ghost propagators at finite temperature from quenched lattice QCD*, Phys. Rev. D **85** (2012) 034501 [arXiv:1108.1735 [hep-lat]].

[2] R. Aouane, F. Burger, E.-M. Ilgenfritz, M. Müller-Preussker and A. Sternbeck, Landau gauge gluon and ghost propagators from lattice QCD with $N_f = 2$ twisted mass fermions at finite temperature, arXiv:1212.1102 [hep-lat], submitted to Phys. Rev. D.

Collaborations

Seminar talk: On the Landau Gauge Gluon Propagator in Lattice QCD at Finite Temperature, 49. Internationale Universitätswochen für Theoretische Physik, Schladming, 2011.

Poster: Landau Gauge Gluon and Ghost Propagators at Non-Zero Temperature on the Lattice, Workshop Quarks, Gluons, and Hadronic Matter under Extreme Conditions, St. Goar, 2011. tmfT Collaboration: HU Berlin, LNF/INFN Frascati, Univ. Bonn, Univ. Frankfurt/Main
 Collaboration with Univ. Regensburg, JINR Dubna, ITEP Moscow

Contact Details and Further Information

PhD student: Rafik Aouane, rafik@physik.hu-berlin.de **PhD advisors:** Prof. Michael Müller-Preussker, PD Dr.sc.nat. Harald Dorn

WWW: http://pha.physik.hu-berlin.de

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