# **Instantons on Special Geometries**

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- Instantons in String Theory
- 2 Reduction to Matrix Equations
- **3 Special Example**
- **4** Summary

- Higher-dimensional Super-Yang-Mills theory appears in the low-energy limit of the heterotic superstring
- String compactification: Spacetime is decomposed as

$$M^{10} = M^{10-n} \times X^n$$

- Instanton equations on X<sup>n</sup> arise as conditions of supersymmetry preservation in heterotic superstring compactifications
- Instantons on special geometries can be lifted to solutions of heterotic supergravity (→ brane interpretation)

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- $\mathcal{A}$ : connection on a bundle over a manifold X with curvature  $\mathcal{F} = d\mathcal{A} + \mathcal{A} \wedge \mathcal{A}$ .
- Anti-self-duality equation in four dimensions (dim(X) = 4):

$$*\mathcal{F} = -\mathcal{F} \tag{1}$$

- In higher dimensions  $(\dim(X) = n)$ : Assume there exists a 4-form  $Q \in \Omega^4(X)$
- Instanton equation (generalized anti-self-duality equation):

$$*\mathcal{F} = -*Q \wedge \mathcal{F} \tag{2}$$

- Consider a coset space G/H, where G is a compact semisimple Lie group and H a closed Lie subgroup
- We work with the cone over this coset space

$$X := \mathcal{C}(G/H) = (\mathbb{R} \times G/H, g_{cone})$$

$$g_{cone} = e^{2\tau} \left( d\tau^2 + g_{G/H} \right)$$
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# **Special Geometries**

#### Why do we consider cones over coset spaces?

- A family of explicit higher-dimensional instanton solutions exists on Euclidean spaces and on certain conical manifolds.
- Cones over special holonomy manifolds are part of the manifolds on which certain branes are defined:

$$M = \mathbb{R}^{1,1} \times \mathbb{T}^{7-k} \times \mathcal{C}(G/H)$$

■ Instantons on cones over 5-dimensional Sasaki, 6-dimensional nearly Kähler, 7-dimensional 3-Sasakian and nearly parallel  $G_2$ -manifolds can be lifted to solutions of heterotic supergravity.

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■ Instanton equation  $*\mathcal{F} = -*Q \wedge \mathcal{F}$  in components:

$$\mathcal{F}_{0a} = -\frac{1}{2}Q_{0acd}\mathcal{F}_{cd} \qquad \qquad \mathcal{F}_{ab} = \frac{1}{2}(Q_{0eab}Q_{0ecd} - Q_{abcd})\mathcal{F}_{cd} \qquad (4)$$

Ansatz for the gauge potential

$$A = e^{i}I_{i} + e^{a}X_{a}(\tau) \tag{5}$$

 $X_a( au)$  matrix-valued function that satisfies the equivariance condition  $[I_i, X_a] = f_{ia}^b X_b$   $I_A$  generators of G  $e^A$  coframe on  $\mathcal{C}(G/H)$  structure constants of G

Curvature

$$\mathcal{F} = \frac{dX_a}{d\tau} e^0 \wedge e^a - \frac{1}{2} \left( f_{ab}^i l_i + f_{ab}^e X_e - [X_a, X_b] \right) e^a \wedge e^b \tag{6}$$

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  - $g \propto \delta_{AB} e^A e^B$  (induced by Cartan-Killing form)
  - $f = \frac{1}{2!} f_{abc} e^a \wedge e^b \wedge e^c$
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- The geometrical condition (8) is identically satisfied
- The differential equation (7) takes the form

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- Such branes are defined on manifolds of the form  $M = \mathbb{R}^{1,1} \times \mathbb{T}^{7-k} \times \mathcal{C}(G/H)$
- We have constructed instanton conditions on  $C(G/H) \subset M$  and solved them under special assumptions
- These solutions are building blocks for heterotic string vacua with explicitly non-trivial gauge field
- Outlook:
   Interpretation of instanton conditions in more general setups
   Lifting of the building blocks to heterotic supergravity solutions

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