Superconducting Radiofrequency Accelerating Structures

Lutz Lilje DESY -MPY-

> TAS2008 Part II

Largely taken from a lecture by Jacek Sekutowicz for SRF 2007, Beijing

Acknowledgement

- The main part is directly from Jacek's Beijing lecture on RF Cavity Design
 - For me this lecture came on rather short notice, therefore I took Jacek's lecture, which I think is excellent
- For the record you will be able to download the full lecture, beside the shortened version included here
 - e.g. omitted Numerical Simulation Methods

Outline of the lectures

Lecture 1

- Radiofrequency (RF) cavities
- Superconductivity basics
- RF superconductivity
- Limitations of superconducting RF cavities

Lecture 2

- RF Cavity Design
 - By J. Sekutowicz
- Mechanical Design
 - Lorentz-force Detuning

SRF 2007

Tutorial: Superconducting High & Cavities

Jacek Sekutowicz, DESY



Topics

- 1. Introduction and History (in brief)
- 2. RF Parameters
- 3. Criteria for Cavity Design
- 4. Multi-cell Structures and Weakly Coupled Structures
- 5. Tools for RF-design
- 6. LEC and Transient state
- 7. Performance tests
- 8. Mechanical Design
- 9. Final Remarks



Milestones that led to accelerators based on SRF

Supe

1908: He Liquefie

<u>1911:</u> House

1928-34: Discove V, Ti and



RF Acceleration

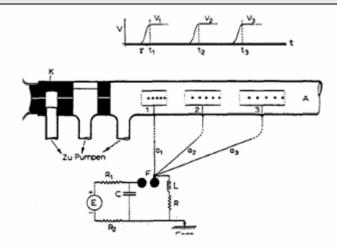


Fig. 1. Diagram of linear accelerator from Professor G. Ising's pioneer publication (1924) of the principle of multiple acceleration of ions.

invented Klystron.

1947: Luis Alvarez (USA)
Built first DTL (32 MeV protons).

1947: W. Hansen (USA)
Built first 6 MeV e-accelerator, Mark I (TW-structure).



Superconductivity

RF Acceleration

1961: Bill Fairbank (Stanford Univ.) presented the first proposal for a superconducting accelerator.

1964: Bill Fairbank, Alan Schwettman and Perry Wilson (Stanford University)
First acceleration of electrons with sc lead cavity.

1967: John Turneaure (Stanford University)
Epeak =70 MV/m and Q~10¹⁰ in 8.5 GHz cavity !!

1968-1981: Mike McAshan, Alan Schwettman, Todd Smith, John Turneaure and Perry Wilson (Stanford University)

Development and Construction of the Superconducting Accelerator SCA.



Dismantled Facilities

- 1. TRISTAN (32/49m)*
- 2. LEP (288/490m)
- 3. **HERA** (16/19m)

Operating Facilities

1. 3CA (4/28M)	1. SCA	(4/28m)
----------------	--------	---------

- 2. S-DALINAC (10/10m)
- 3. CESR (4/1.2 m)
- 4. CEBAF (320/160m)
- 5. KEK B-Factory (8/2.4m)
- 6. Taiwan LS (1+1/0.3m)
- 7. Canadian LS (1+1/0.3m)
- 8. **DIAMOND** (3/0.9m)
- 9. SOLEIL (4/1.7m)
- 10. TTF II (56/58m)
- 11. SNS (81/65m)
- 12. JLab-FEL (24/14m)
- 13. LHC (16/6m)
- 14. ELBE (6/6m)

*(Number of cavities/total active length)



Tomorrow Facilities

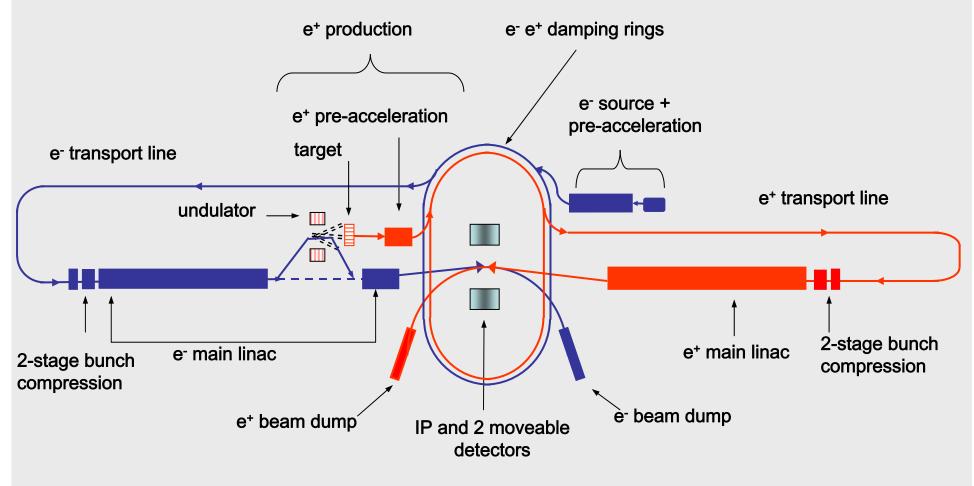
- 1. CEBAF-12GeV (400/216m)
- 2. SNS-upgrade (117/ 98m)
- 3. XFEL (800/832m)
- 4. ERL-Cornell (310/250m)
- 5. BESSY (144+7_{3-Harm}/152m)
- 6. 4GLS (~40/~42m)
- 7. RHIC-cooling (4/4 m)
- 8. Shanghai LS......
- 9. BEPC II (2/0.6m)

Day after Tomorrow Facilities

- 1. RIA (option 180/122m)
- 2. X-Ray MIT (option 176/184m)
- 3. LUX (~40/~50m)
- 4. FERMI Proton Linac (384/370m)
- 5. ERHIC.....
- 6. ELIC
- 7. ARC-EN-CIEL (48/50m)
- 8. ILC (~15764/~16395m)



ILC (~15764/~16395m)







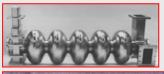
The "heart" of all mentioned facilities are sc standing wave (usually multi-cell) accelerating structures.



S-DALINAC 3 GHz



CESR/CEBAF 1.5 GHz





HEPL 1.3 GHz



TESLA/ILC 1.3 GHz



SNS ß=0.61,0.81, 0.805 GHz



HERA 0.5 GHz



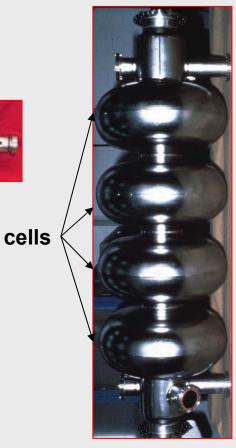
KEK-B 0.5 GHz



CESR 0.5 GHz



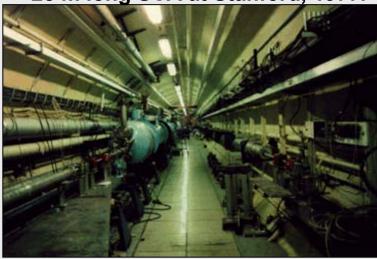
LEP 0.352 GHz





What is the progress in the 30 years and what do we need in the next 10 years?

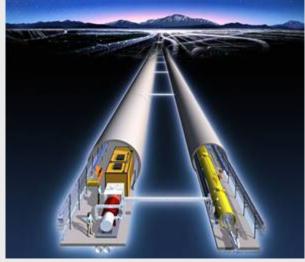
~ 28 m long SCA at Stanford, 1977.





E_{acc} ~2 (2.5) MV/m in cw (10% DF). 4 Structures 5.65m + capture + preaccelerator.

~ 21.6 km long ILC linac, 2015+





E_{acc} > 34 MV/m shown in several 9-cells in the cw test.

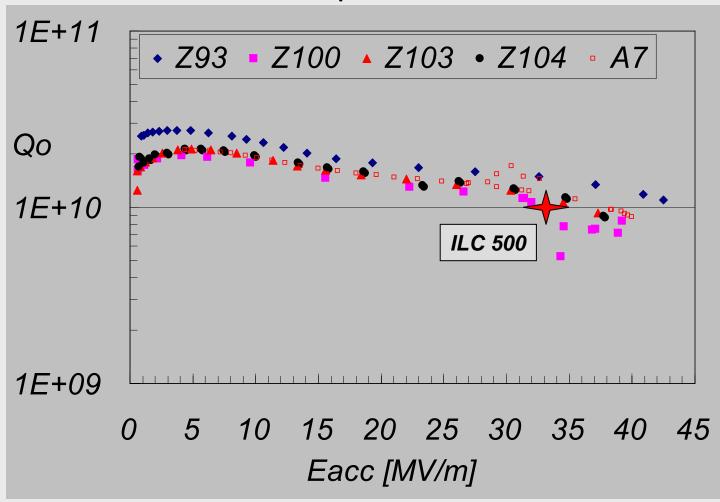


This gradient is required in all 15764 cavities.



Results at DESY and JLab (2007):

cw test result at 2K for 5 electropolished 9-cell TESLA cavities.





2.1 Cavities and their Eigenmodes

Cavity≡ an arbitrary volume, partially closed by the metal wall, capable to store the E-H energy



~ 3.95 GHz is the lowest frequency

First assumption:

1. Stored E-H fields are harmonic in time.

Maxwell equations for the harmonic, lossless case with no free charge in the volume

$$\nabla \times H = i\omega \varepsilon E$$

$$\nabla \times E = -i\omega \mu H$$

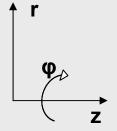
$$\nabla \cdot E = 0$$

$$\nabla \cdot H = 0$$



Second assumption (good approximation for the elliptical cavities):

2. The volume is cylindrically symmetric. We commonly use the (r, φ, z) coordinates.



z is conventional direction of the acceleration and symmetry axis

$$\begin{cases} \nabla_{c} \times H = i\omega\varepsilon E \\ \nabla_{c} \times E = -i\omega\mu H \\ \nabla_{c} \cdot E = 0 \\ \nabla_{c} \cdot H = 0 \end{cases}$$

$$\nabla_{\mathbf{c}} \times A = \vec{i}_{r} \left(\frac{1}{r} \frac{\partial A_{z}}{\partial \varphi} - \frac{\partial A_{\varphi}}{\partial z} \right) + \vec{i}_{\varphi} \left(\frac{\partial A_{r}}{\partial z} - \frac{\partial A_{z}}{\partial r} \right) + \vec{i}_{z} \left(\frac{1}{r} \frac{\partial (rA_{\varphi})}{\partial r} - \frac{1}{r} \frac{\partial A_{r}}{\partial \varphi} \right)$$

$$\nabla_{\mathbf{c}} \cdot A = \frac{1}{r} \frac{\partial (rA_{r})}{\partial r} + \frac{1}{r} \frac{\partial A_{\varphi}}{\partial \varphi} + \frac{\partial A_{z}}{\partial z}$$

$$\nabla_{\mathbf{c}} \cdot \mathbf{A} = \frac{1}{r} \frac{\partial (r \mathbf{A}_r)}{\partial r} + \frac{1}{r} \frac{\partial \mathbf{A}_{\varphi}}{\partial \varphi} + \frac{\partial \mathbf{A}_{\mathbf{z}}}{\partial \mathbf{z}}$$



Third assumption

3. For the acceleration are suitable field patterns with strong E along the beam trajectory. This ensures, by the proper phasing, maximal energy exchange between the cavity and beam.

TM0xx-like monopole modes have "very strong" E_z component on the symmetry axis. Fields of the monopole modes are independent on ϕ .

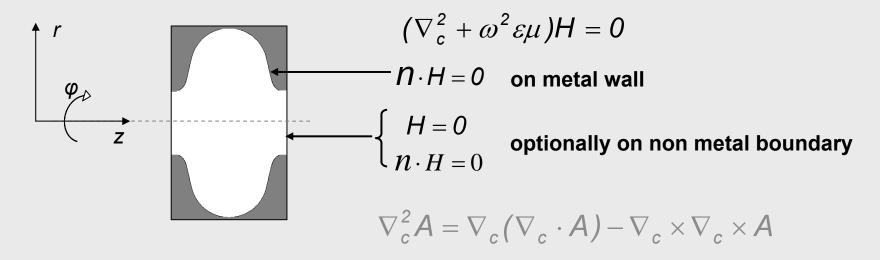
$$\frac{\partial \mathsf{E}}{\partial \varphi} = 0 \qquad \frac{\partial \mathsf{H}}{\partial \varphi} = 0$$

Non monopole (HOM) modes have component E_z = 0 on the symmetry axis. Their fields dependent on ϕ .



Maxwell equations + boundary conditions for E and H fields lead to the <u>Helmholtz</u> equation, which is an eigenvalue problem.

For H(r,z) field of a monopole mode the equation is:



There is infinity number of TM0xx solutions (modes) to the Helmholtz equation. All modes are determine by:

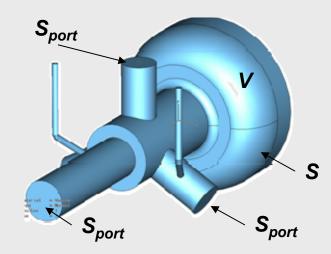
$$H_n(r,z)=[0, H_{\varphi,n}(r,z), 0],$$

$$E_n(r,z) = [E_{r,n}(r,z), 0, E_{z,n}(r,z)]$$

and frequency ω_n .



2.2 What are figures of merit for a cavity storing E-H energy?



 $W_n \equiv \text{stored energy of a mode } n : \{\omega_n, E_n, H_n\}$

$$W_n = 2\mu \int_V \frac{H_n^2}{4} dV = 2\varepsilon \int_V \frac{E_n^2}{4} dV$$

Quality Factors

The measure of the energy loss in the metal wall and due to the radiation via open ports:

Intrinsic $Q \equiv Qo$

$$Q_{0,n} = \frac{\omega_n \cdot W_n}{P_n} = \frac{\omega_n \cdot W_n}{\frac{R_{s,n}}{2} \int_{S} H_n^2 ds}$$

External
$$Q \equiv Q_{ext}$$

$$Q_{0,n} = \frac{\omega_n \cdot W_n}{P_n} = \frac{\omega_n \cdot W_n}{\frac{R_{s,n}}{2} \int_{S} H_n^2 ds} \qquad Q_{ext,n} = \frac{\omega_n \cdot W_n}{P_{rad,n}} = \frac{\omega_n \cdot W_n}{\frac{1}{2} \int_{S_{port}} E_n \times H_n ds}$$

Geometric Factor

The measure of the energy loss in the metal wall for the surface resistance $R_{s,n}=10$

$$G_{n} \equiv Q_{0,n} \cdot R_{s,n} = \frac{\omega_{n} \cdot W_{n} \cdot R_{s,n}}{P_{n}} = \frac{\omega_{n} \cdot W_{n}}{\frac{1}{2} \int_{S} H_{n}^{2} ds}$$

It is the ratio of the stored energy to the integral of $(H_n)^2$ on the metal surface.

2.3 What are figures of merit for the beam-cavity interaction?

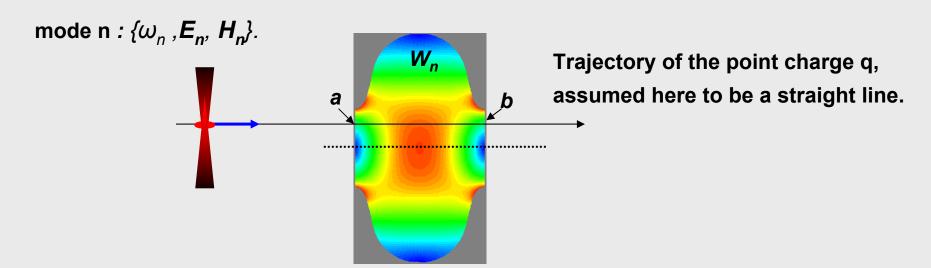
This interaction which is:

- **→** Acceleration
- Deceleration (ERL)
- HOMs excitation

can be described in the Frequency Domain (FD) or/and in Time Domain (TD).



 $(R/Q)_n$, a "measure" of the energy exchange between point charge and mode n (FD).



$$V_{n} = \sqrt{\left(\int_{z_{a}}^{z_{b}} E_{n,z} \sin(\frac{\omega_{n}}{\beta c}(z - z_{a}))dz\right)^{2} + \left(\int_{z_{a}}^{z_{b}} E_{n,z} \cos(\frac{\omega_{n}}{\beta c}(z - z_{a}))dz\right)^{2}}$$

$$(R/Q)_{n} \equiv \frac{V_{n}^{2}}{\omega_{n} W_{n}}$$



For the accelerating mode we often use the product of $G_{acc} \cdot (R/Q)_{acc}$, as a "measure" of the power P dissipated in the metal wall at the given accelerating voltage V_{acc} and the given surface resistance R_s .

$$\frac{P_{dissipated}}{V_{acc}^2} \equiv \frac{R_s}{G_{acc} \cdot (R/Q)_{acc}}$$

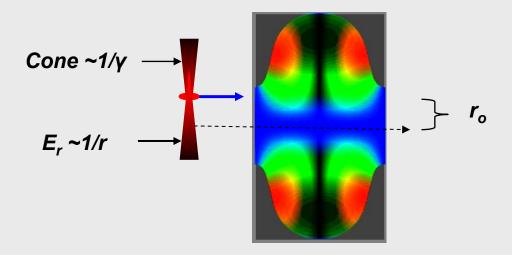
This is due to the surface quality;
Big improvement possible.



This is due to the geometry of cells; Moderate improvement possible.

Longitudinal and Transverse Loss Factors (TD)

Ultra relativistic point charge q passes empty cavity



- a. Density of the inducted charge on the wall depends on the distance to the beam trajectory.
- b. The non uniform charge density on the metal wall causes the current flow on the surface.



The amount of energy lost by charge q to the cavity is:

$$\Delta U_q = k_{\parallel} \cdot q^2$$
 for monopole modes (max. on axis)

$$\Delta U_q = k_{\perp} \cdot q^2$$
 for non monopole modes (off axis)

where k_{\parallel} and $k_{\perp}(r)$ are loss factors for the monopole and transverse modes respectively.

The induced E-H field (wake) is a superposition of cavity eigenmodes (monopoles and others) having the $E_n(r,\varphi,z)$ field along the trajectory.

Both description methods FD and TD are equivalent.

For individual mode n and point-like charge:

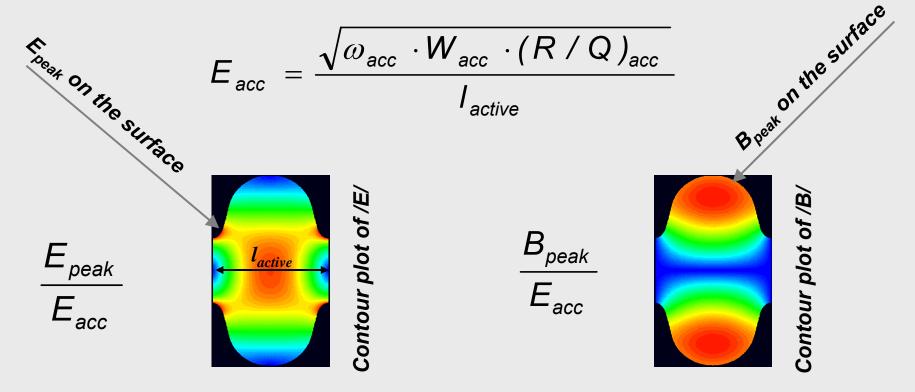
$$k_{||,n}^{p} = \frac{\omega_n \cdot (R/Q)_n}{4}$$
 Note the linac convention for (R/Q) definition.

Similar for other loss factors......



RF parameters of the accelerating mode having more practical background

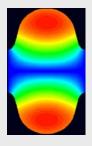
At stored energy W_{acc} the mean value of the accelerating gradient is:



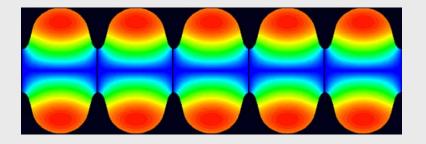
Ratio shows sensitivity of the shape to the field electron emission phenomenon.

Ratio shows limit in E_{acc} due to the breakdown of superconductivity (Nb ~190mT).

The last parameter, relevant for multi-cell accelerating structures, is the coupling k_{cc} between cells for the accelerating mode passband (Fundamental Mode passband).







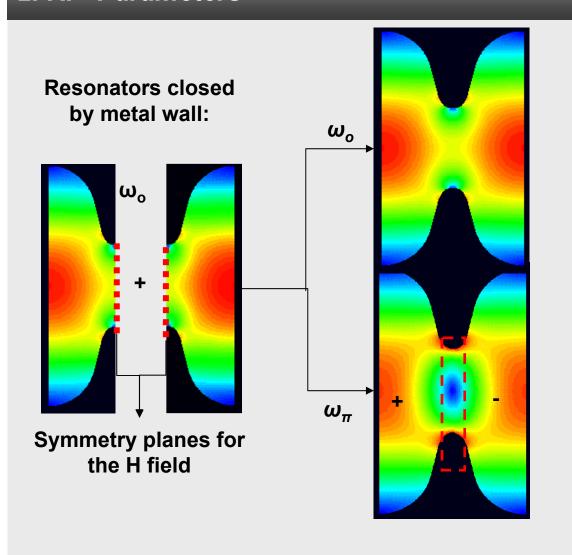
Single-cell structures are attractive from the RF-point of view:

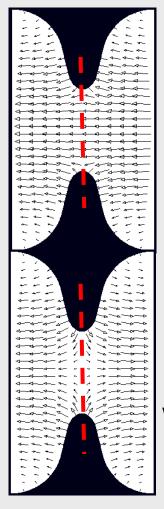
- Easier to manage HOM damping
- **→** No field flatness problem.
- → Input coupler transfers less power
- → Easy for cleaning and preparation
- → But it is expensive to base even a small linear accelerator on the single cell. We do it only for very high beam current machines.

Multi-cell structures are less expensive and offers higher real-estate gradient but:

- → Field flatness (stored energy) in cells becomes sensitive to frequency errors of individual cells
- **→** Other problems arise: HOM trapping...





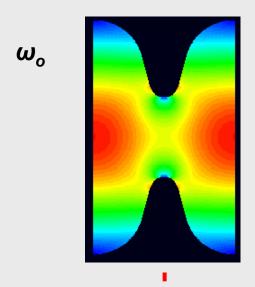


Symmetry plane for the H field

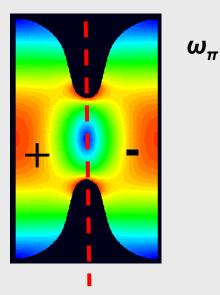
Symmetry plane for the E field which is an additional solution



The energy flux across the coupling region, refilling energy loss is proportional to the transverse components: H_{ω} and E_{r}





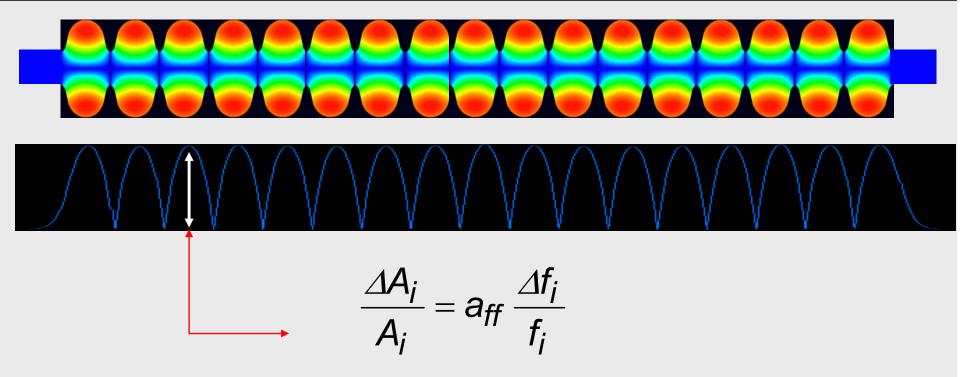


Small H_{ϕ} (due to the losses) + strong E_r at the symmetry plane

The normalized difference between these frequencies is a measure of the energy flow via the coupling region

$$k_{cc} = \frac{\omega_{\pi} - \omega_{0}}{\frac{\omega_{\pi} + \omega_{0}}{2}}$$





Field flatness factor for a structure made of N cells and coupling factor k_{cc}

$$a_{\rm ff} = \frac{N^2}{k_{cc}}$$

The above formulae estimate sensitivity of a multi-cell field profile to frequency errors of an individual cell for the accelerating mode (π -mode)



We will talk here about inner cells design because these cells "dominate" parameters of a multi-cell superconducting accelerating structure.

RF parameters summary:

FM :
$$(R/Q)$$
, G, E_{peak}/E_{acc} , B_{peak}/E_{acc} , k_{cc}

HOM:
$$k_{\perp}, k_{\parallel}$$

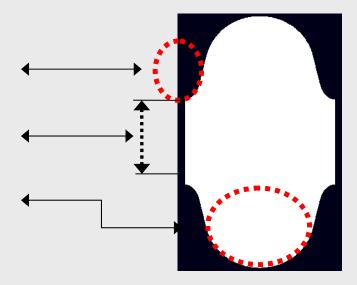
There are 7 parameters we want to optimize for an inner cell

Geometry:

iris ellipsis : half-axis h_r , h_z

iris radius : r_i

equator ellipsis: half-axis h_r , h_z



There is some kind of conflict <u>7 parameters</u> and only <u>5 variables</u> to "tune"

Criteria	RF-parameter	Improves when	Cavity examples
Operation at high gradient	E _{peak} / E _{acc} ↓ B _{peak} / E _{acc}	r _i Iris, Equator shape	TESLA, HG CEBAF-12 GeV
Low cryogenic losses	(R/Q) ·G	r _i ↓ Equator shape	LL CEBAF-12 GeV LL- ILC cavity
High I _{beam} ↔ Low HOM impedance	k_{\perp}, k_{\parallel}	r_i	B-Factory RHIC cooling

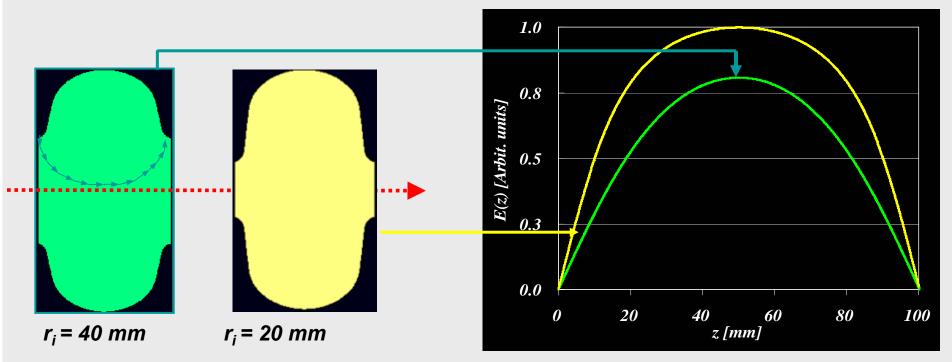
We see here that r_i is a very "powerful variable" to trim the RF-parameters of a cavity.



Why for a smaller aperture (r_i)

- (R/Q) is bigger
- E_{peak}/E_{acc} , B_{peak}/E_{acc} is lower?

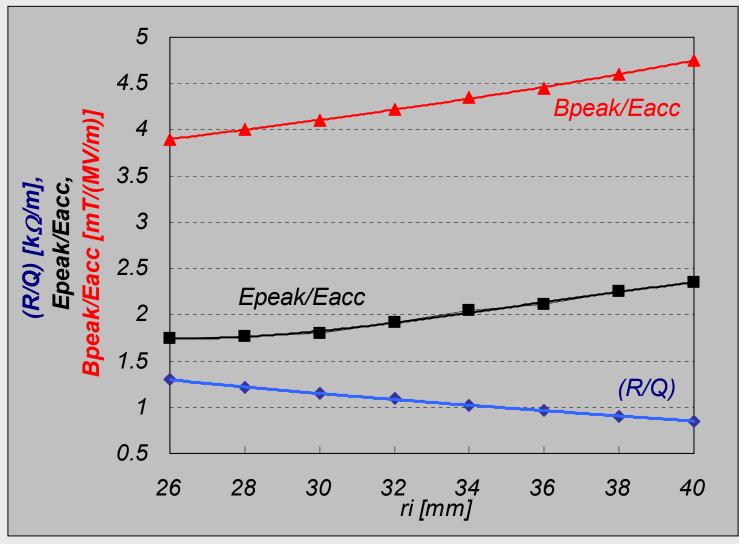
E_{acc} is higher at the same stored energy in the cell



 $E_z(z)$ for small and big iris radius



Example: f = 1.5 GHz

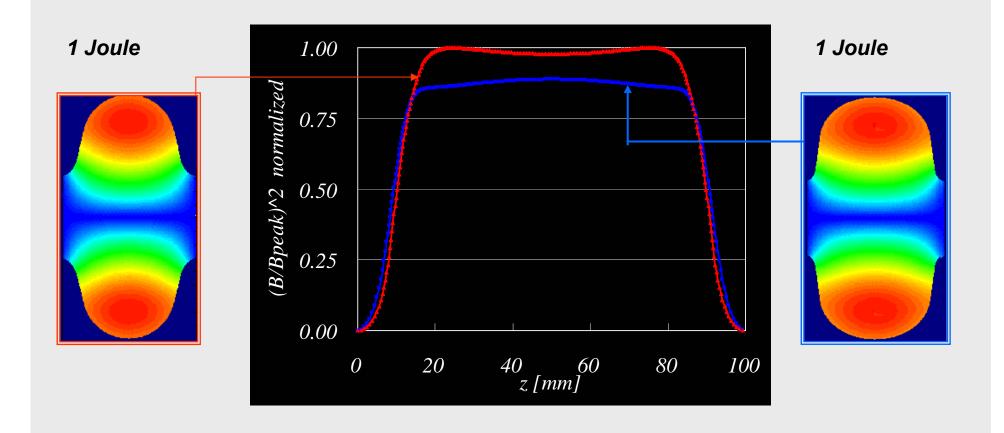


A. Mosnier, E. Haebel, SRF Workshop 1991



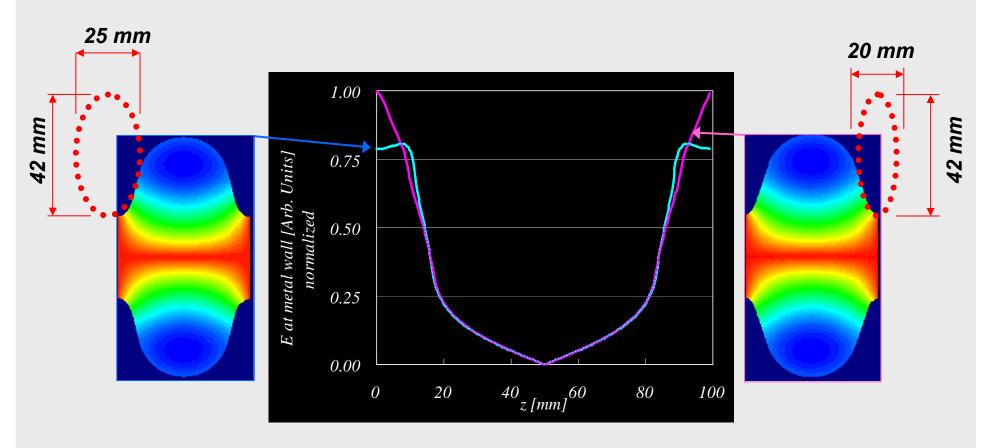
In addition to the iris radius r_i :

• B_{peak}/E_{acc} (and G) changes vs. the equator shape





Similarly: E_{peak}/E_{acc} changes vs. the iris shape

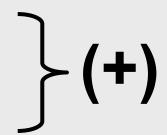


Both cells have the same: f_i (R/Q) and r_i



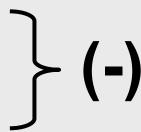
We know that a smaller aperture r_i makes FM:

- (R/Q) higher
- B_{peak}/E_{acc} , E_{peak}/E_{acc} lower

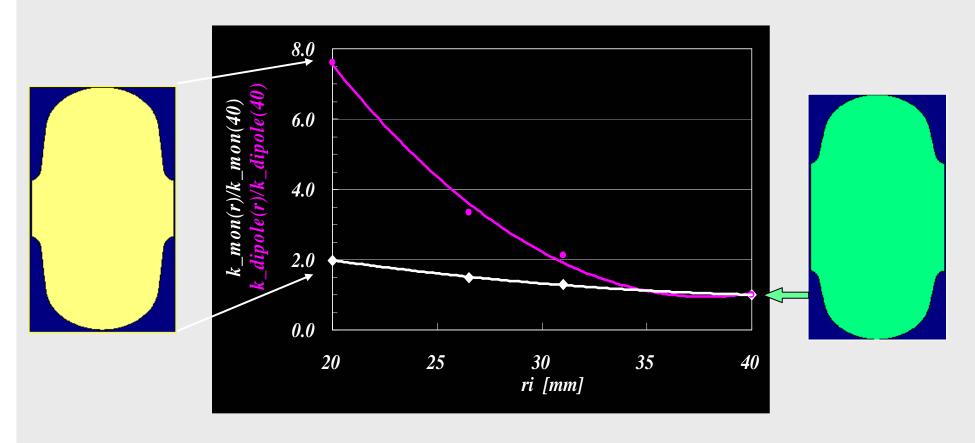


but unfortunately a smaller aperture r_i makes:

- HOMs impedances (k_{\perp} , k_{\parallel}) higher
- cell-to-cell coupling (k_{cc}) weaker



HOMs impedances $(k \perp, k \parallel)$



$$(R/Q) = 152 \Omega$$

$$B_{peak}/E_{acc} = 3.5 \ mT/(MV/m)$$

$$E_{peak}/E_{acc} = 1.9$$

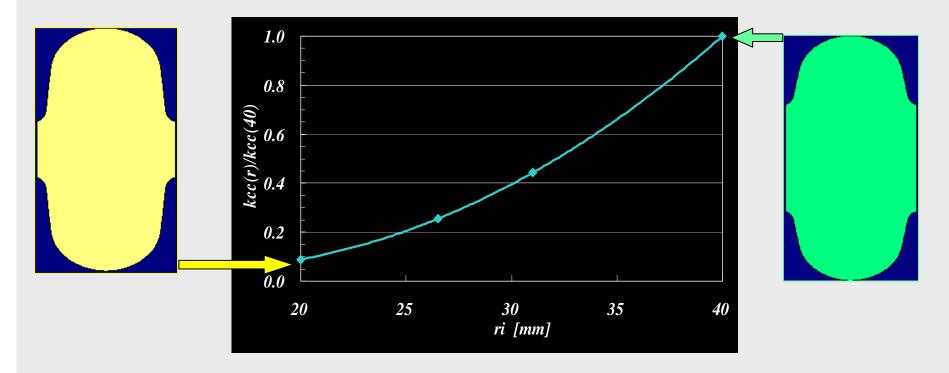
$$(R/Q) = 86 \Omega$$

$$B_{peak}/E_{acc} = 4.6 mT/(MV/m)$$

$$E_{peak}/E_{acc} = 3.2$$



Cell-to-cell coupling (kcc)



$$(R/Q) = 152 \Omega$$

$$B_{peak}/E_{acc} = 3.5 \ mT/(MV/m)$$

$$E_{peak}/E_{acc} = 1.9$$

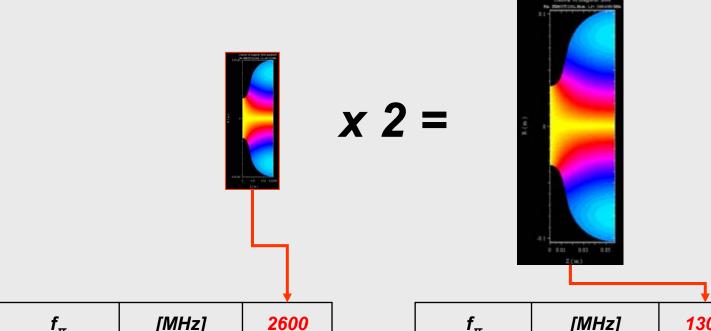
$$(R/Q) = 86 \Omega$$

$$B_{peak}/E_{acc} = 4.6 \text{ mT/(MV/m)}$$

$$E_{peak}/E_{acc} = 3.2$$



What about accelerating mode frequency of a superconducting cavity?



f_{π}	[MHz]	2600
R/Q	$[\Omega]$	57
r/q=(R/Q)/I	$[\Omega/m]$	2000
G	$[\Omega]$	271

f_{π}	[MHz]	1300
R/Q	$[\Omega]$	57
r/q=(R/Q)/I	[Ω/m]	1000
G	$[\Omega]$	271

 $r/q=(R/Q)/I \sim f$



From the formula, we learned before:

$$\frac{P_{\text{dissipated}}}{V_{\text{acc}}^2} \equiv \frac{R_{\text{s}}}{G_{\text{acc}} \cdot (R/Q)_{\text{acc}}}$$

one obtains:

$$P_{dissipated} = \frac{R_s \cdot V_{acc}^2}{G_{acc} \cdot (r/q)_{acc} \cdot I_{active}}$$

A higher frequency would be a good choice to minimize power dissipation in the metal wall when the length I_{active} and final energy V_{acc} are fixed.

Unfortunately this applies only to room temperature structures made of Cu, which $R_s \sim (f)^{1/2}$.

For superconductors like Nb:

$$R_s(f) = R_{res} + R_{BCS} = R_{res} + 0.0002 \cdot \frac{1}{T} \cdot (\frac{f[GHz]}{1.5})^2 \cdot exp(-\frac{17.67}{T})$$

and increase of $R_s \sim (f)^2$ for higher f must be compensated with lower temperature T.

This is why ILC (1.3GHz) will operate at 2K (1.8K), and HERA (0.5GHz) and LEP (0.352GHz) could operate at 4.2 K

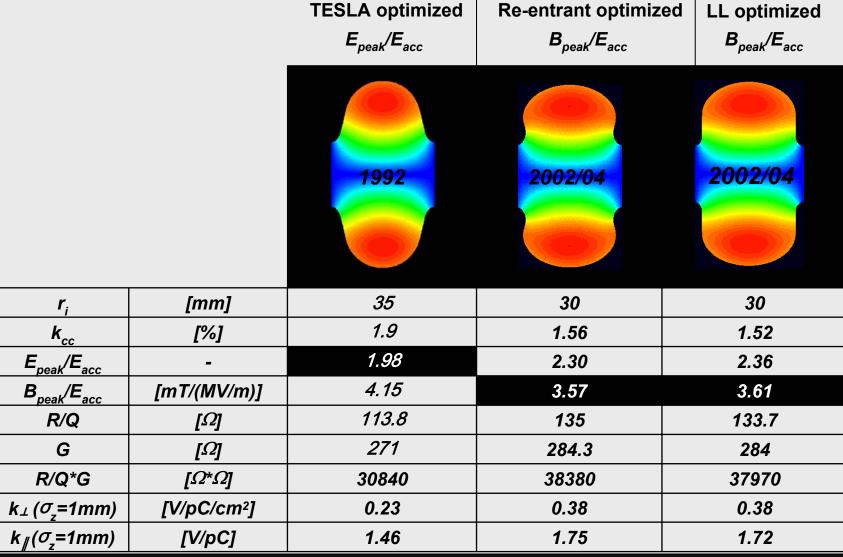


Examples of inner cells

		CEBAF Original Cornell ß=1	CEBAF -12 High Gradient ß=1	CEBAF -12 Low Loss ß=1	TESLA ß=1	SNS ß=0.61	SNS ß=0.81	RIA ß=0.47	RHIC Cooler ß=1
f _o	[MHz]	1448.3	1468.9	1475.1	1278.0	792.8	792.8	793.0	683.0
f_{π}	[MHz]	1497.0	1497.0	1497.0	1300.0	805.0	805.0	805.0	703.7
k _{cc}	[%]	3.29	1.89	1.49	1.9	1.52	1.52	1.52	2.94
E _{peak} /E _{acc}	-	2.56	1.96	2.17	1.98	2.66	2.14	3.28	1.98
B _{peak} /E _{acc}	[mT/(MV/m)]	4.56	4.15	3.74	4.15	5.44	4.58	6.51	5.78
R/Q	$[\Omega]$	96.5	112	128.8	113.8	49.2	83.8	28.5	80.2
G	$[\Omega]$	273.8	266	280	271	176	226	136	225
R/Q*G	$[\Omega^*\Omega]$	26421	29792	36064	30840	8659	18939	3876	18045
$k_{\perp} (\sigma_z = 1 mm)$	[V/pC/cm²]	0.22	0.32	0.53	0.23	0.13	0.11	0.15	0.02
$k_{/\!\!/}(\sigma_z=1mm)$	[V/pC]	1.36	1.53	1.71	1.46	1.25	1.27	1.19	0.85



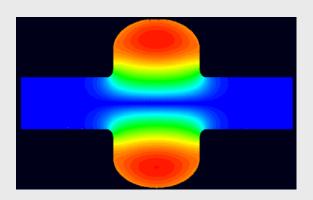
Evolution of inner cells proposed for the ILC collider:

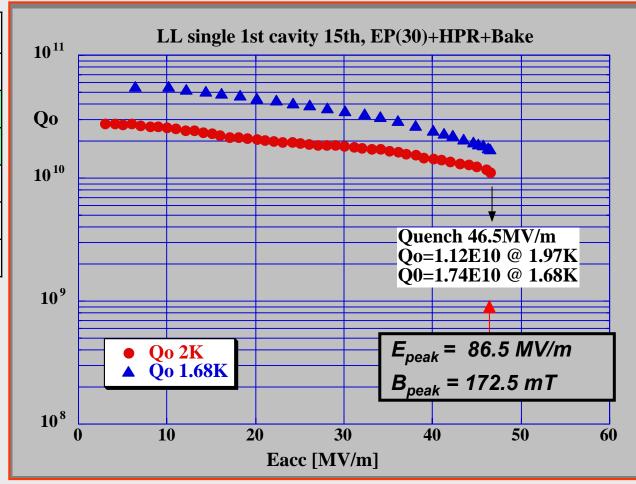




KEK test September 2005 !!!!!!!

		LL
f	[MHz]	1286.6
E _{peak} /E _{acc}	-	1.86
B_{peak}/E_{acc}	[mT/(MV/m)]	3.71
R/Q	$[\Omega]$	130.0
G	$[\Omega]$	279
Ø _{iris}	[mm]	61

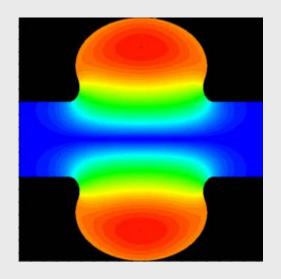


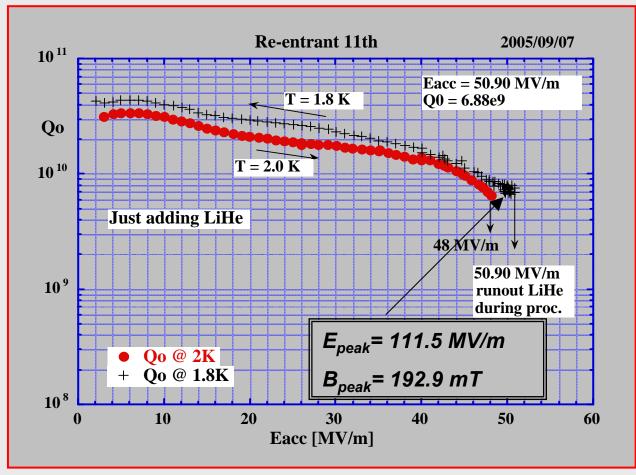




KEK tests September 2005

		RE
f	[MHz]	1278.6
E_{peak}/E_{acc}	-	2.19
B_{peak}/E_{acc}	[mT/(MV/m)]	3.79
R/Q	$[\Omega]$	126.0
G	$[\Omega]$	278
Ø _{iris}	[mm]	68



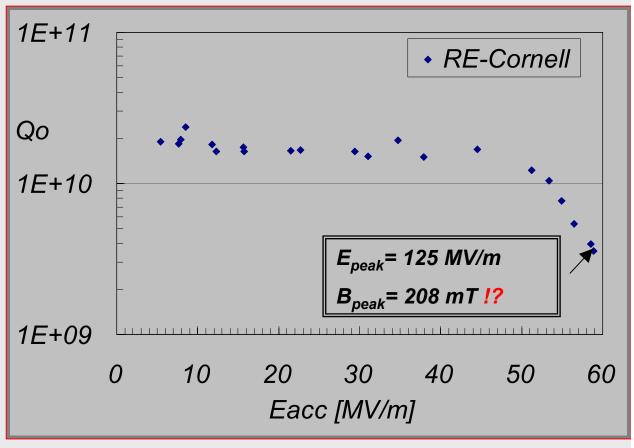




Cornell, test in March 2007 !!!!

		RE+Tubes
f	[MHz]	1300.3
E _{peak} /E _{acc}	-	2.11
B _{peak} /E _{acc}	[mT/(MV/m)]	3.53
R/Q	$[\Omega]$	126.0
G	$[\Omega]$	283.3
Ø _{iris}	[mm]	60







We re-call pros and cons for a multi-cell structure

- → Cost of accelerators are lower (less auxiliaries: LHe vessels, tuners, fundamental power couplers, control electronics)
- **→** Higher real-estate gradient (better fill factor)
- Field flatness vs. N
- → HOM trapping vs. N
- Power capability of fundamental power couplers vs. N
- Chemical treatment and final preparation become more complicated
- → The worst performing cell limits whole multi-cell structure



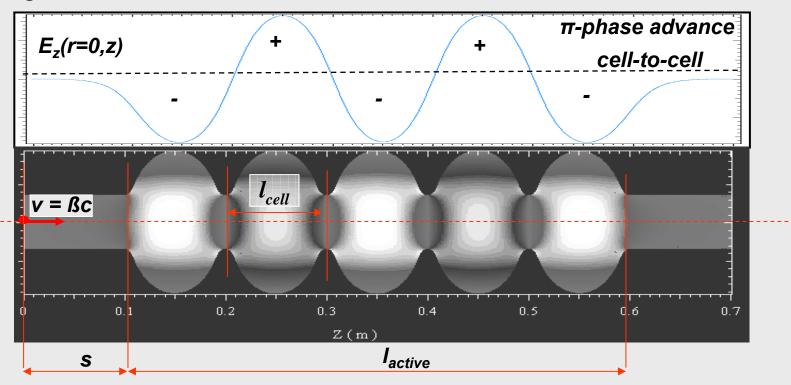
Break

We re-call pros and cons for a multi-cell structure

- → Cost of accelerators are lower (less auxiliaries: LHe vessels, tuners, fundamental power couplers, control electronics)
- **→** Higher real-estate gradient (better fill factor)
- Field flatness vs. N
- → HOM trapping vs. N
- Power capability of fundamental power couplers vs. N
- Chemical treatment and final preparation become more complicated
- → The worst performing cell limits whole multi-cell structure



Accelerating mode in a multi-cell structure



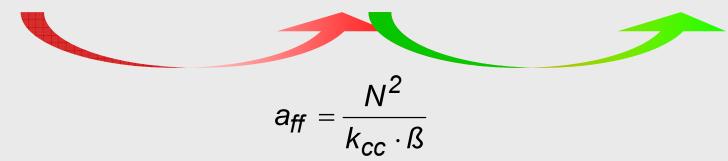
Synchronic acceleration and max of $(R/Q)_{acc}$ when:

- 1. $I_{active} = NI_{cell} = NcB/(2f)$ and
- 2. the injection takes place at an optimum phase φ_{opt} which ensures that particles arrive at the mid-plane of the first cell when E_{acc} reaches its maximum (+q passing to the right) or minimum (-q passing to the right).



Field flatness in a multi-cell structures

	Original Cornell	High Gradient	Low Loss	TESLA	SNS ß=0.61	SNS ß=0.81	RIA ß=0.47	RHIC
	N = 5	N =7	N =7	N=9	N=6	N=6	N=6	N=5
year	1982	2001	2002	1992	2000	2000	2003	2003
a _{ff}	1489	2592	3288	4091	3883	2924	5040	850



Many years of experience with: heat treatment, chemical treatment, handling and assembly allows one to preserve field profile, even in cavities with bigger N and weaker k_{cc}

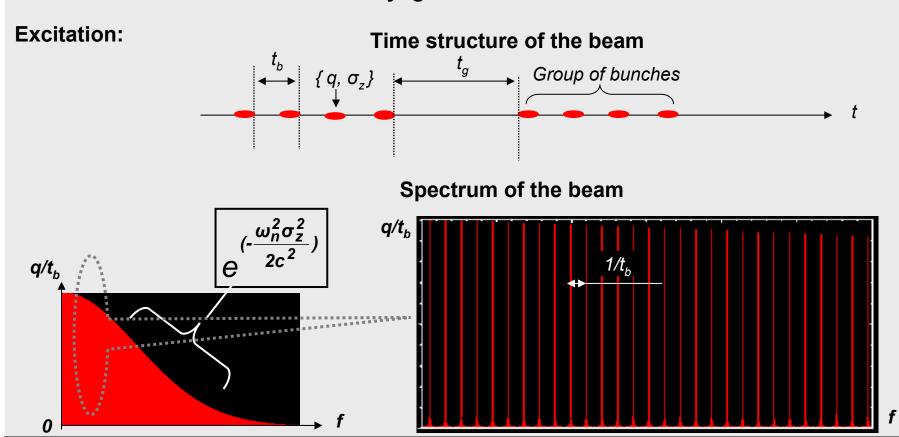
For the TESLA cavities: field flatness is better than 95 %



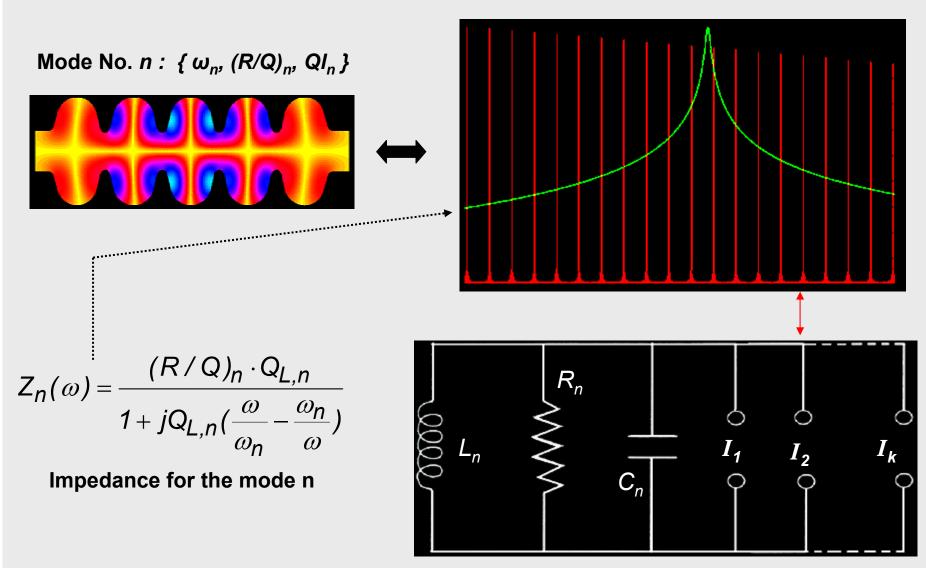
HOM trapping in a multi-cell structures

The excitation of HOMs by the accelerated beam causes:

- → Beam instabilities and/or dilution of emittance
- **→** Bunch-to-bunch energy modulation
- → Additional cryogenic loss









The power induced by "all" spectral lines (current sources) in mode No. n:

$$P_{n} = \frac{1}{2} \sum_{k} Z_{n}(\omega_{k}) \cdot I_{k}^{2}$$

$$Z_{n}(\omega) = \frac{(R/Q)_{n} \cdot Q_{L,n}}{1 + jQ_{L,n}(\frac{\omega}{\omega_{n}} - \frac{\omega_{n}}{\omega})} \quad \text{and} \quad \frac{1}{Q_{L,n}} = \frac{1}{Q_{0,n}} + \frac{1}{Q_{ext,n}} \leftarrow \text{Measure of the extracted power}$$

The HOM couplers, devices extracting the energy from the parasitic modes, are attached to cavities to mitigate these phenomenon.

The experience shows that, the HOM couplers and FM couplers can be attached to the beam tubes and must not be located at cells because this leads to the performance degradation.

Coaxial HOM coupler

Coaxial HOM coupler



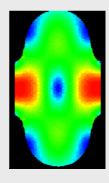


Waveguide HOM ports

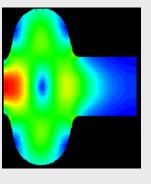
The HOM trapping mechanism is similar to the FM field profile unflatness mechanism:

- weak $k_{cc,HOM}$, cell-to-cell coupling for HOM
- → difference in the HOM frequency between the end-cell and inner-cell

f = 2385 MHz



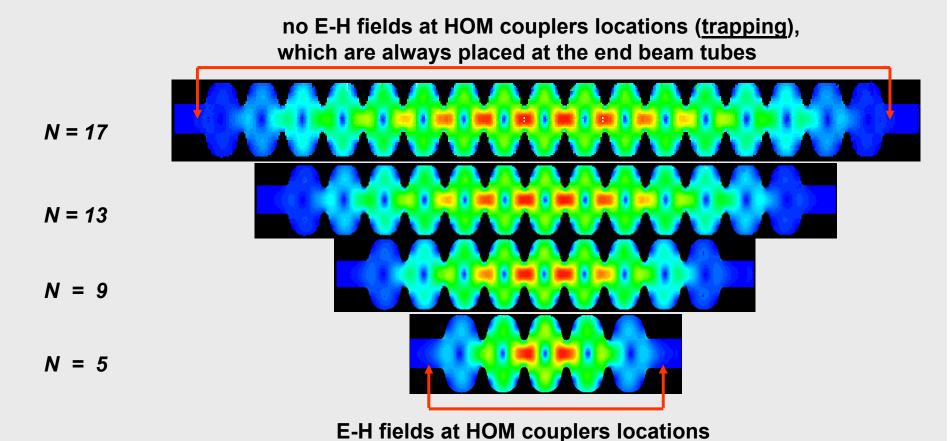
That is why they hardly resonate together



 $f = 2415 \, MHz$



Example: how N influences strength of the E-H fields at HOM couplers locations



Less cells in a structure helps always to reach low Qs of HOMs.



What additional to reducing N can we do to avoid the trapping?

Adjustment of end-cells

The geometry of end-cells differs from the geometry of inner cells due to the attached beam tubes, HOM- and input couplers.

Their function is multifold and their geometry must fulfill three requirements:

- field flatness and frequency of the accelerating mode
- → field strength of the accelerating mode at FPC location enabling operation with matched Qext
- → fields strength of dangerous HOMs ensuring their required damping by means of HOM couplers or/and beam line absorbers.

All three make design of the end-cells more difficult than inner cells.



1. Open irises of the inner cells and end-cells (bigger $k_{cc,HOM}$) and making shape of both very similar

Example: RHIC 5-cell cavity for the electron cooling:

Monopole mode k_c

$$f_{HOM} = 1394 MHz$$

 $f_{HOM} = 1403 MHz$

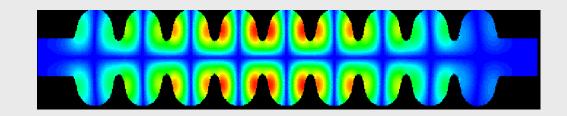
The method causes (relevant.



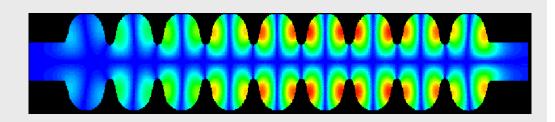
2. Tailor end-cells to equalize HOM frequencies of inner- and end-cells.

Example: TESLA 9-cell cavity, which has two different end-cells (asymmetric cavity)

The lowest mode in the passband $f_{HOM} = 2382 \text{ MHz}$

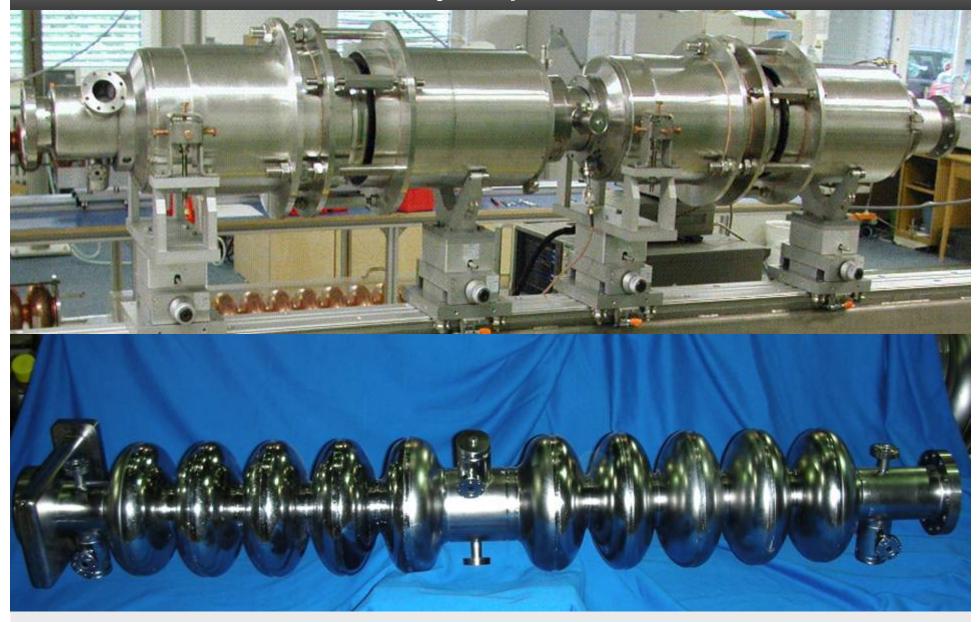


The highest mode in the passband $f_{HOM} = 2458 \text{ MHz}$

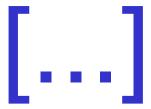


The method works for very few modes but keeps the (R/Q) value high of the fundamental mode.







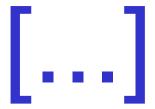


- Left out
 - Main coupler power estimation
 - Mode measurement for vertical test

List of multi-cell cavities ß=1 optimized for various criteria.

Criterion	Structure	Best parameter	Weakest parameter (point)	Comments
E _{acc}	HG: 1.5 GHz, N=7 TESLA: 1.3 GHz, N=9 ILC-LL: 1.3 GHz, N=9 ILC-RE 1.3 GHz, N=9	Epeak/Eacc= 1.96 Epeak/Eacc= 1.98 Bpeak/Eacc= 3.61 Bpeak/Eacc= 3.57	Real estate -Eacc Real estate -Eacc Real estate-Eacc, Epeak/Eacc Real estate-Eacc, Epeak/Eacc	Designed for I _{beam} < 10 mA, Pulse operation
Real estate E _{acc}	2x9 TESLA: 1.3 GHz, N= 18	Real estate-Eacc Epeak/Eacc= 2.0	Field flatness preservation Cleaning	New FPC design for 0.8 MW, Difficult to clean
P _{loss}	LL: 1.5 GHz, N= 7	Bpeak/Eacc= 3.7 (R/Q)·G	Not easy to clean, HOM damping	Designed for I _{beam} < 1 mA First LL-type cavity
Z _{HOM}	RHIC: 0.7 GHz, N= 5	Very low: k⊥ , k _∥ Epeak/Eacc= 1.98	Cryogenic losses	First multi-cell for I _{beam} ≈ 2 A





- Left out
 - Simulation methods
 - Vertical Testing

9. Final Remarks

- → Both RF- and Mechanical design are well understood
- → We have day by day better tools for designing of accelerating cavities
- → There is not a "golden" cavity suitable for all applications
- → Not all requirements can be fulfilled at once and cavities must be tailored to their applications.

References:

- 1. H. Padamsee, J. Knobloch, T. Hays, "RF-Superconductivity for Accelerators", Wiley Series in Beam Physics and Accelerator Technology, 1998.
- 2. Proceedings of all SRF Workshops
- 3. TESLA TDR, DESY-Report 2003.
- 4. J.S., "TESLA Superconducting Accelerating Structures" Institute of Physics, Journal of Measurement Science and Technology, 18 (2007) 2285-2292.



Mechanical Design: Detuning

- In many cases a deformation of a cavity leads to frequency change
- Vibrations = ,Microphonics'
 - Pumps, ground motion etc.
- Lorentz-force due to RF fields
 - High electromagnetic fields exert a mechanical force on the cavity walls
 - sometimes also called ,radiation pressure
 - For pulsed operation this needs to be compensated for efficient operation

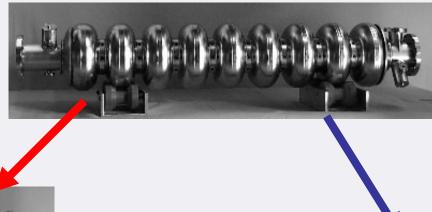


Basic Principle of Frequency Tuning

The frequency of a cavity is determined by the equator diameter: $f \sim 1/D$

The simplest way to change the diameter of a cavity is to change its length.

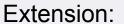
A length change of 3um results in a frequency change of 1kHz





Compression:

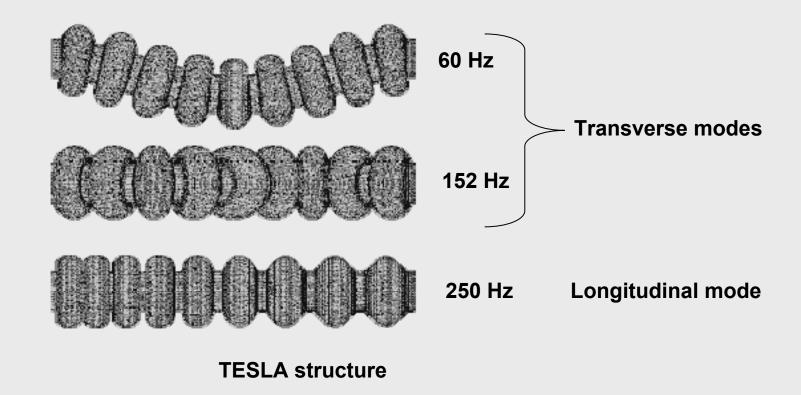
Frequency decreases



Frequency increases

Hananaka har

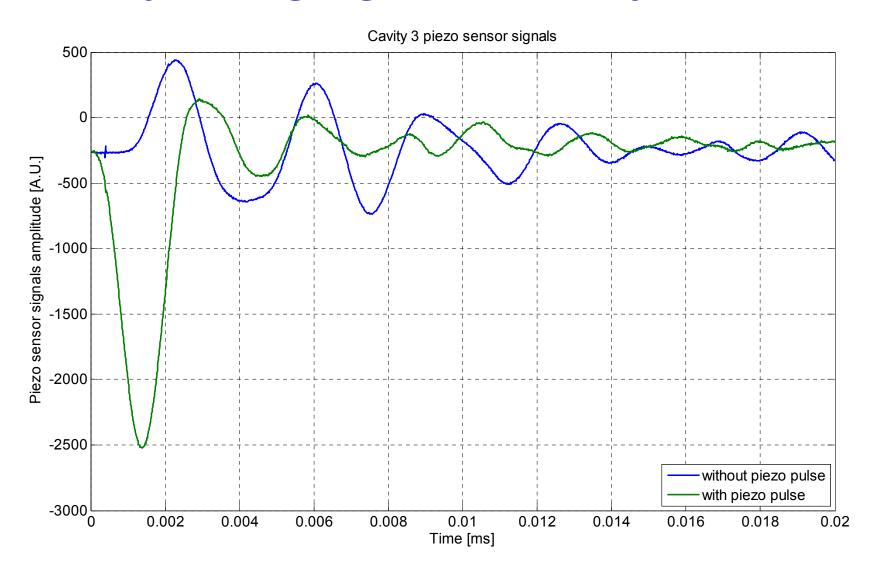
Mechanical Resonances of a multi-cell cavity



The mechanical resonances modulate frequency of the accelerating mode. Sources of their excitation: vacuum pumps, ground vibrations...



Cavity 'Ringing' excited by RF Pulse

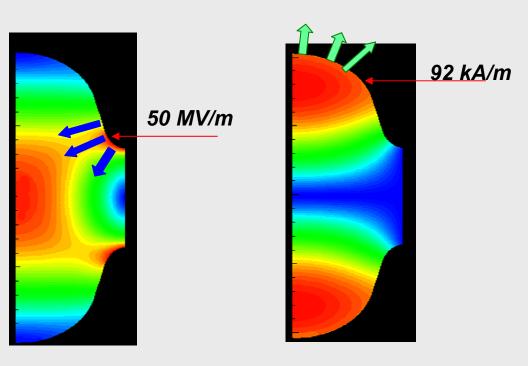


8. Mechanical Design

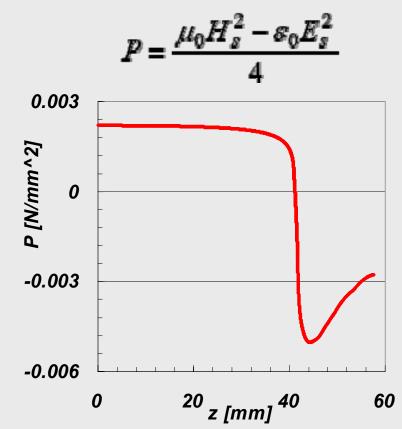
The mechanical design of a cavity follows its RF design:

- → Lorentz Force Detuning
- → Mechanical Resonances

Lorentz Force Detuning

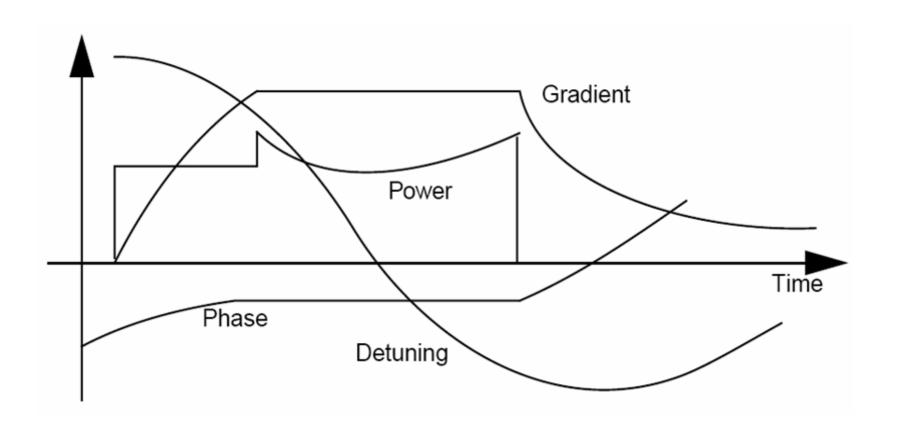




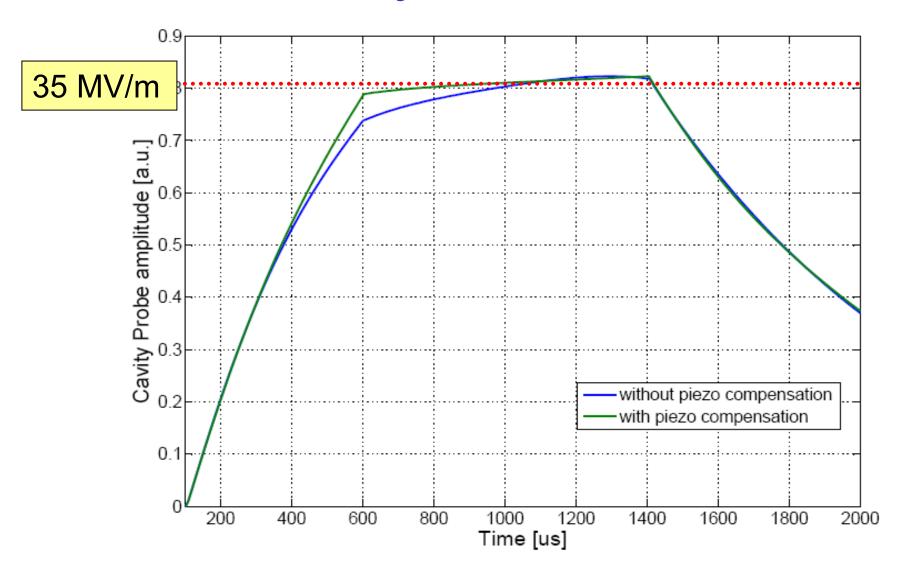




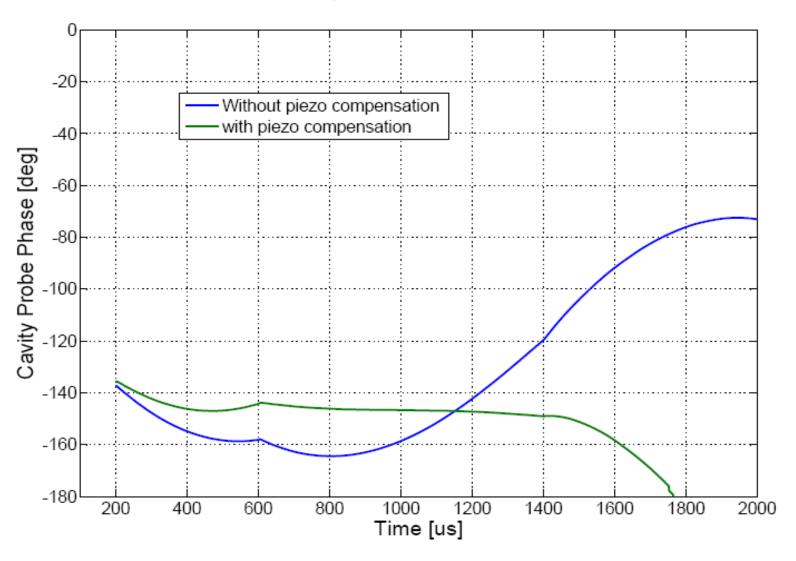
Lorentz-Force Detuning - Scheme



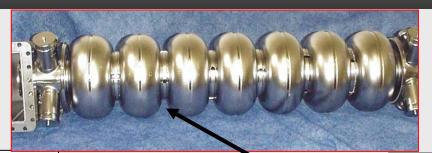
Cavity 3: Gradient

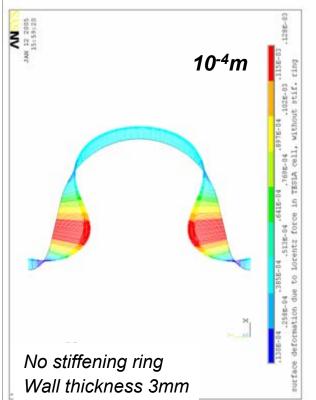


Cavity 3: Phase

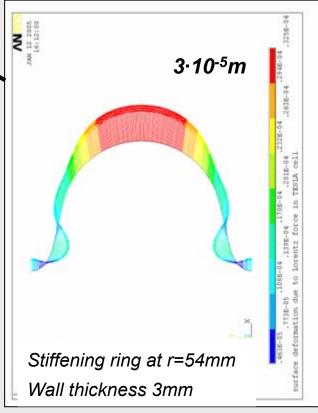


8. Mechanical Design





Surface deformation without and with stiffening ring (courtesy of I. Bonin, FERMI)



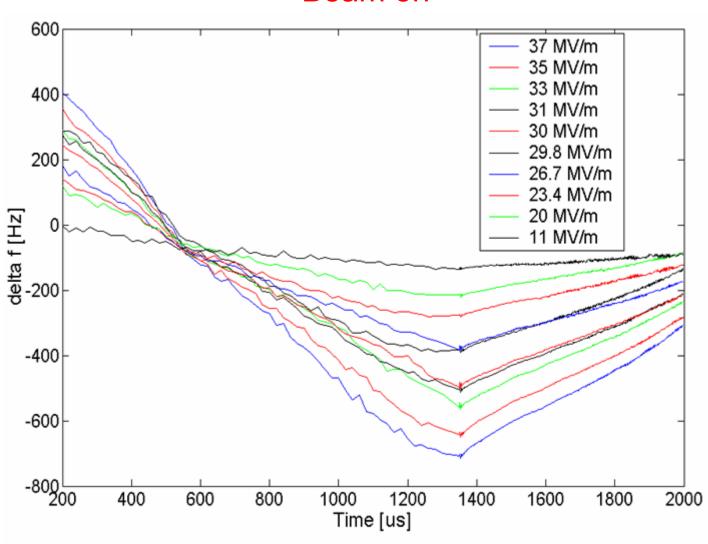
Essential for the operation of a pulsed accelerator $\Delta f = k_L (E_{acc})^2$

 $k_L = -1 \; Hz/(MV/m)^2$



Frequency detuning during RF pulse





Frequency detuning due Lorentz forces of the electromagnetic field in the cavities:

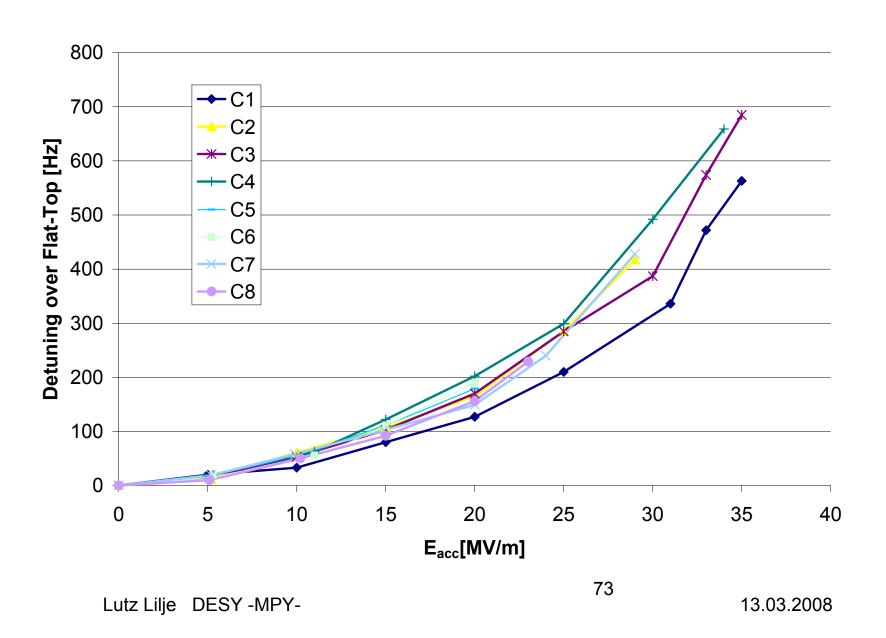
$$\Delta f = K \cdot E_{acc}^2$$

where $K \approx 1 \text{ Hz} / (MV/m)^2$

Remember: Cavity bandwidth with main coupler is ≈ 300 Hz

12

Lorentz Force Detunings in Module 6





Fast Tuner with Piezos

Maximal load:

- <5 kN

Operating temperature

- 4 -10 K in isolation vacuum
- Stroke (and capacity) reduced to about 10-20% of room temperature stroke

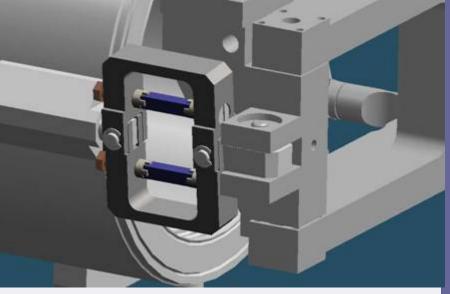
Preload force

0,8-1,2 kN/cm² for highest lifetime

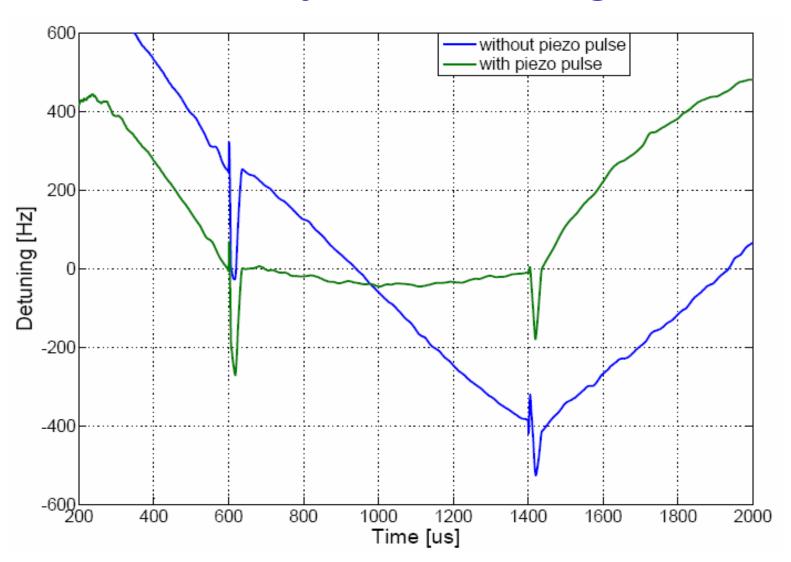
Sensor and actuator configuration

- One active element
- other element as sensor and for redundancy

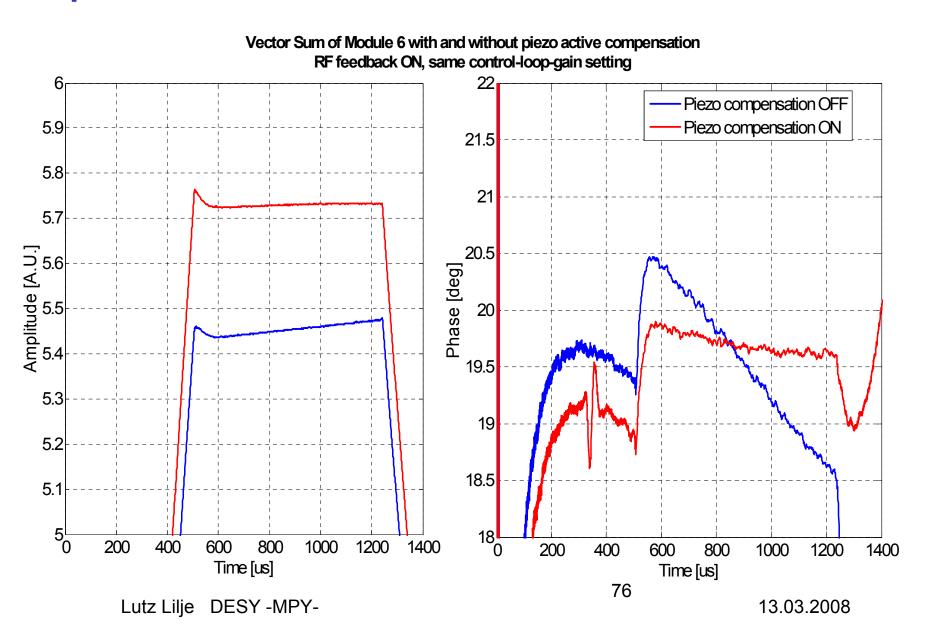




Cavity 3: Detuning

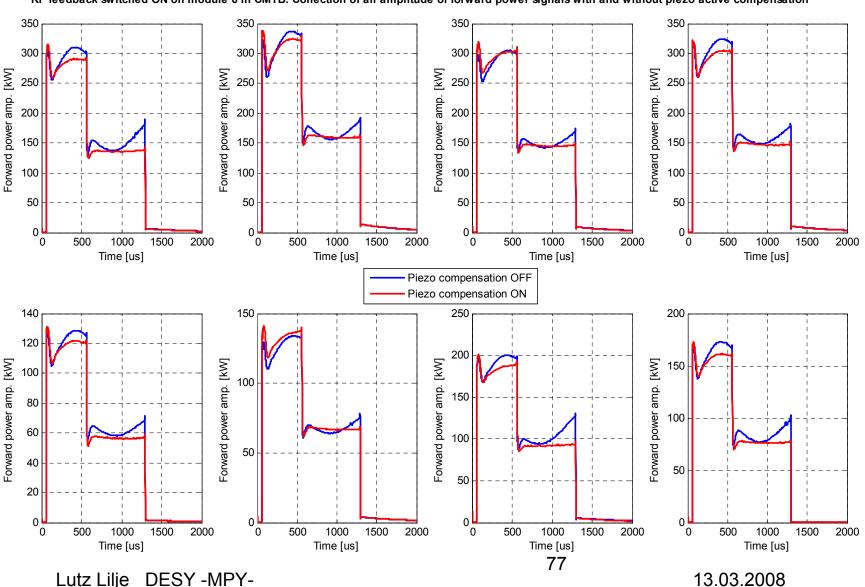


Operation of Full module – Vector-Sum



Operation of Full Module – Forward Power

RF feedback switched ON on module 6 in CMTB. Collection of all amplitude of forward power signals with and without piezo active compensation



Summary Mechanical Design

- Deformations of the cavity due to vibrations or operation need to be taken into account during design
 - E.g. stiffening rings

Thanks for your attention!