

Flavour Physics at LHCb

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Overview

- Introduction
- CP-Violation Measurements
- B-Physics with LHCb
- Selected Decay Channels
- Summary

Why B-Physics ? – or – Where is all the Anti-Matter?

naive expectation: equal amounts of matter and anti-matter from the Big Bang, but...

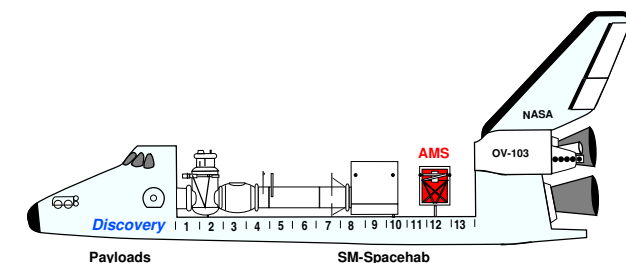
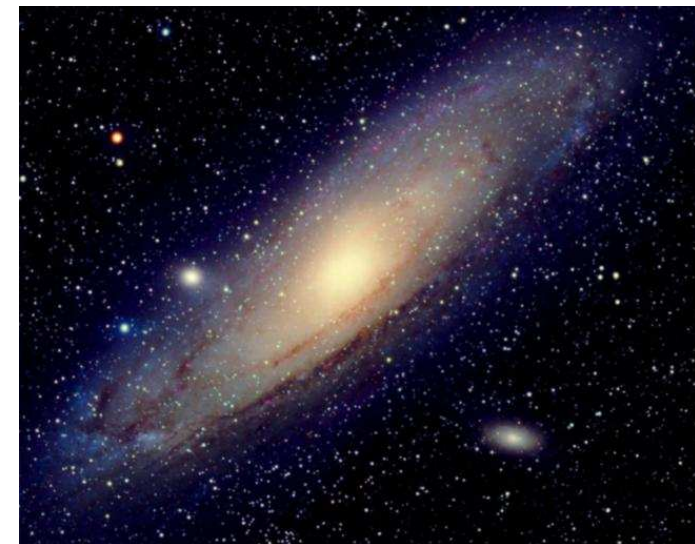
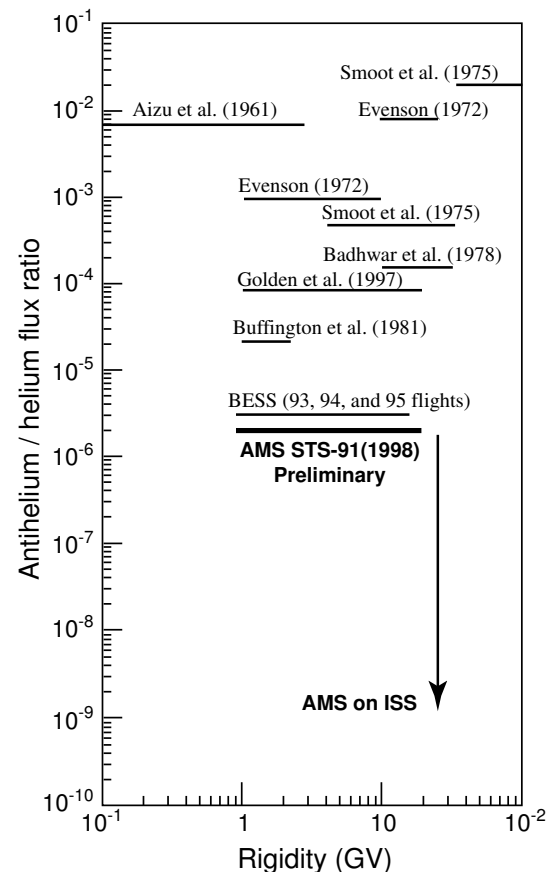
- for example, no annihilation radiation seen in γ -rays from space,
- or search primordial $\overline{\text{He}}$ so far without positive results...

look for:

- cosmic ray $\overline{\text{He}}$
- go into space
- AMS experiment

Result:

no $\overline{\text{He}}$ found



The necessary conditions to explain the matter dominance of the universe are:

1. C- and CP-violation
2. baryon-number violation
3. thermal non-equilibrium

Standard Model:

- C- and CP-violation exist in CKM-sector
- baryon-number violation via Spalurons and thermal non-equilibrium can be created in a 1st order phase transition during the early universe

Problem:

- SM-Higgs particle is too heavy to generate the 1st order phase transition and the CP-violation in the CKM-sector is too small to explain the matter dominance of the universe
 - extra sources of CP violation are needed
 - “New Physics” expected in CP-violation measurements

CP-Asymmetry A_{CP} for decays into a final state y :

$$A_{CP} = \frac{\Gamma(X \rightarrow y) - \Gamma(\bar{X} \rightarrow \bar{y})}{\Gamma(X \rightarrow y) + \Gamma(\bar{X} \rightarrow \bar{y})} \quad \text{with partial widths} \quad \Gamma(\cdot) = |a(\cdot)|^2$$

Consider mixing induced CP-violation in decays to a CP-eigenstate $y = y_{CP}$:

$$a(\textcolor{green}{X} \rightarrow y_{CP}) = a_m(\textcolor{green}{X} \rightarrow \textcolor{green}{X}) \cdot a_d(\textcolor{green}{X} \rightarrow y_{CP}) + a_m(\textcolor{green}{X} \rightarrow \textcolor{red}{\bar{X}}) \cdot a_d(\textcolor{red}{\bar{X}} \rightarrow y_{CP})$$

$$a(\textcolor{red}{\bar{X}} \rightarrow y_{CP}) = a_m(\textcolor{red}{\bar{X}} \rightarrow \textcolor{red}{\bar{X}}) \cdot a_d(\textcolor{red}{\bar{X}} \rightarrow y_{CP}) + a_m(\textcolor{red}{\bar{X}} \rightarrow \textcolor{green}{X}) \cdot a_d(\textcolor{green}{X} \rightarrow y_{CP})$$

The contributing amplitudes are:

$$\left. \begin{matrix} a_d(\textcolor{green}{X} \rightarrow y_{CP}) \\ a_d(\textcolor{red}{\bar{X}} \rightarrow y_{CP}) \end{matrix} \right\} = A e^{\pm i\omega} \quad \left. \begin{matrix} a_m(\textcolor{green}{X} \rightarrow \textcolor{green}{X}) \\ a_m(\textcolor{red}{\bar{X}} \rightarrow \textcolor{red}{\bar{X}}) \end{matrix} \right\} = \cos \frac{\Delta m t}{2} \quad \left. \begin{matrix} a_m(\textcolor{green}{X} \rightarrow \textcolor{red}{\bar{X}}) \\ a_m(\textcolor{red}{\bar{X}} \rightarrow \textcolor{green}{X}) \end{matrix} \right\} = i \sin \frac{\Delta m t}{2} e^{\pm i\phi}$$

with phase factors

■ phase from the decay: ω

■ mixing phase: ϕ



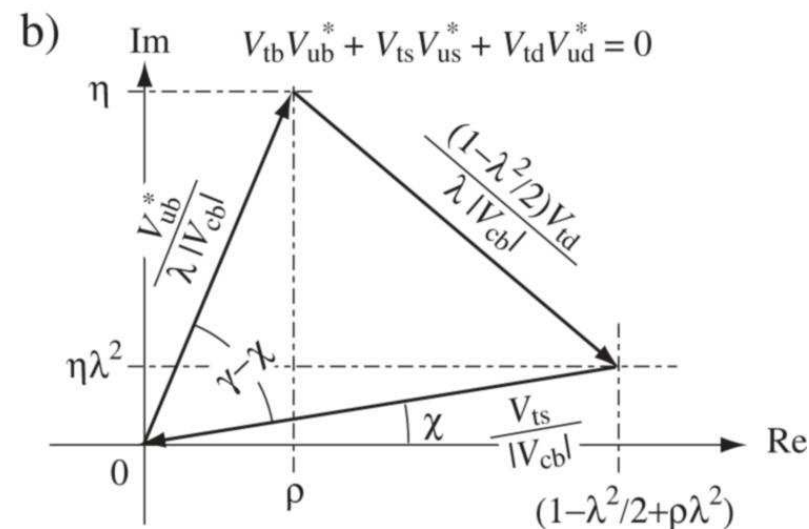
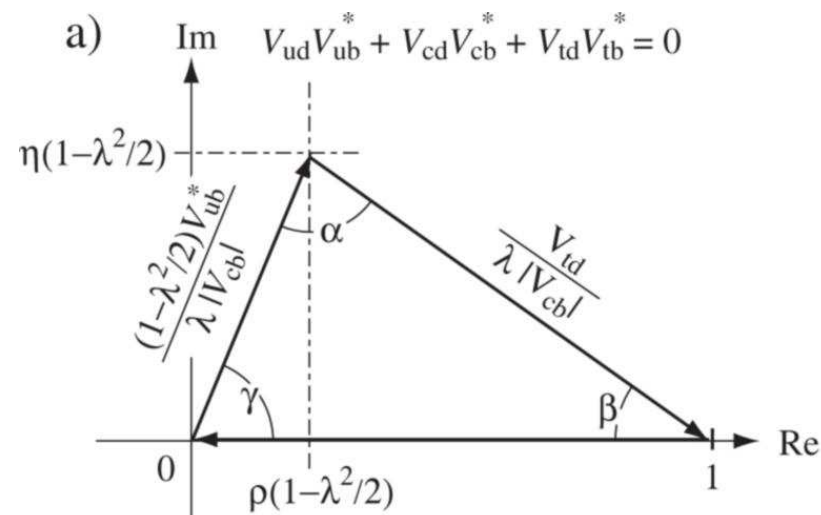
$$A_{CP} = -\sin(\Delta m t) \sin(\phi - 2\omega)$$

Standard Model: All phases arise from CKM Matrix elements

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

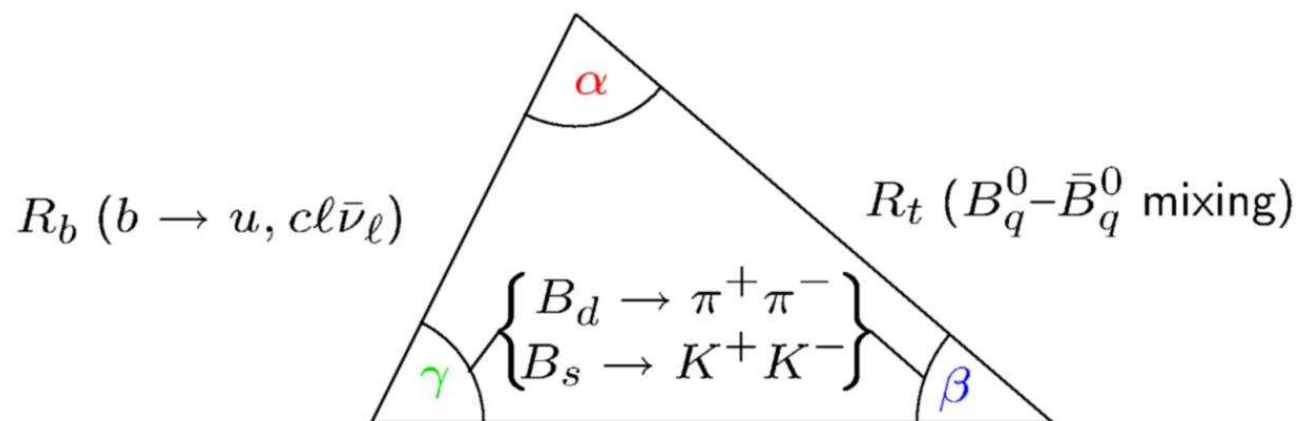
Wolfenstein-parametrization

- unitary matrix, i.e. $C_i \cdot C_j^* = R_i \cdot R_j^* = \delta_{ij}$, parametrized by 4 free parameters
- qualitative picture determined by $\lambda \approx 0.22$
- **unitarity triangles** $C_1 \cdot C_3^* = R_3 \cdot R_1^* = 0$
 - $\arg(V_{td}) = -\beta$ (B_d -mixing)
 - $\arg(V_{ub}) = -\gamma$ ($B_d \rightarrow \pi^+ \pi^-$ decays)
 - $\arg(V_{ts}) = \chi + \pi \approx \eta\lambda^2 + \pi$ (B_s -mixing)



note: $V \neq 1$ goes beyond Higgs search in addressing the origin of mass!

$B \rightarrow \pi\pi$ (isospin), $B \rightarrow \rho\pi$, $B \rightarrow \rho\rho$



$B \rightarrow \pi K$ (penguins)

$B_d \rightarrow \psi K_S$ ($B_s \rightarrow \psi\phi : \phi_s \approx 0$)

$B_u^\pm \rightarrow K^\pm D$
 $B_d \rightarrow K^{*0} D$
 $B_c^\pm \rightarrow D_s^\pm D$

only trees

$B_d \rightarrow \phi K_S$ (pure penguin)

$B_d \rightarrow D^{(*)\pm} \pi^\mp : \gamma + 2\beta$
 $B_s \rightarrow D_s^\pm K^\mp : \gamma + \phi_s$

only trees

from R. Fleischer hep-ph/0505018v1

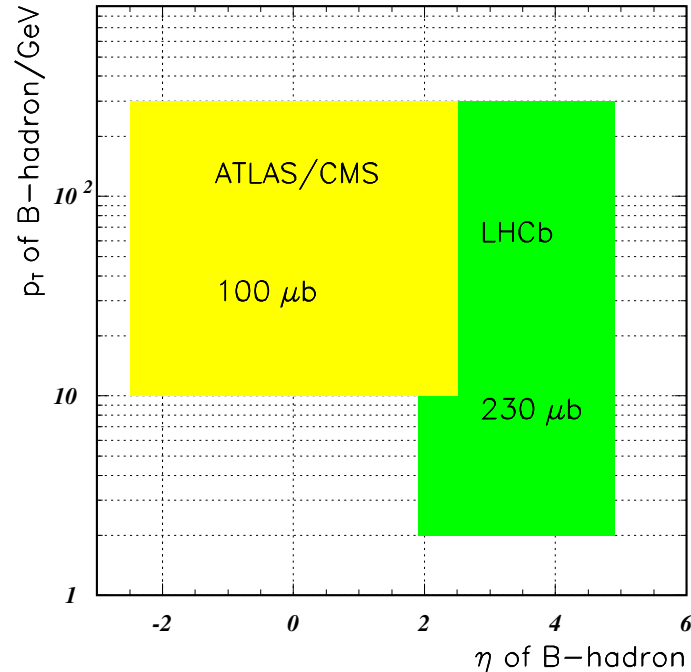
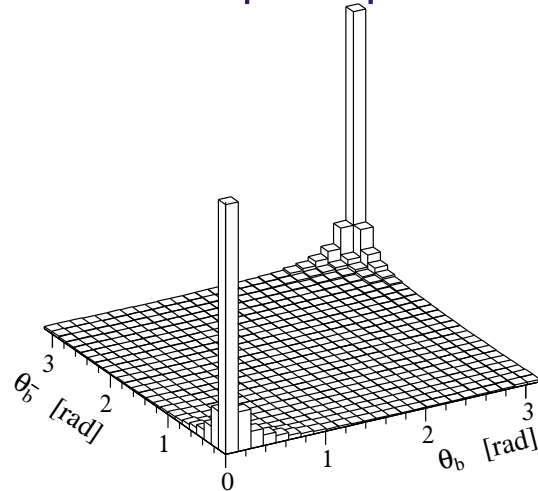
- compare measurements which are insensitive to new physics with measurements of the same quantity from a decay channel that is sensitive to new physics
 - β from $B_d \rightarrow J/\psi K_s$ - tree level dominates
 - β from $B_d \rightarrow \phi K_s$ - pure penguin
 - γ from $B_s \rightarrow D_s K$ - tree level dominates
 - γ from $B_d \rightarrow \pi K$ - with penguin contribution
- study observables which have a small expectation value in the standard model
 - CP-asymmetry in $B_s \rightarrow J/\psi \phi$
 - transitions sensitive to FCNC, for example
 - ✗ $B \rightarrow K^* \gamma$ i.e. $b \rightarrow s \gamma$
 - ✗ $B_d \rightarrow K^{0*} \mu^+ \mu^-$
 - ✗ rare decays $B_{d,s} \rightarrow \mu^+ \mu^-$
- compare UT from length and angle measurements

Note:

The discussion here focuses on mixing induced CP-violation. In addition also direct CP-violation can be measured, with usually different sensitivity to New physics...

$$A_{CP}(t) \sim A_{dir} \cos(\Delta mt) + A_{mix} \sin(\Delta mt)$$

Kinematics of b-quark production:



Challenges to the detector:

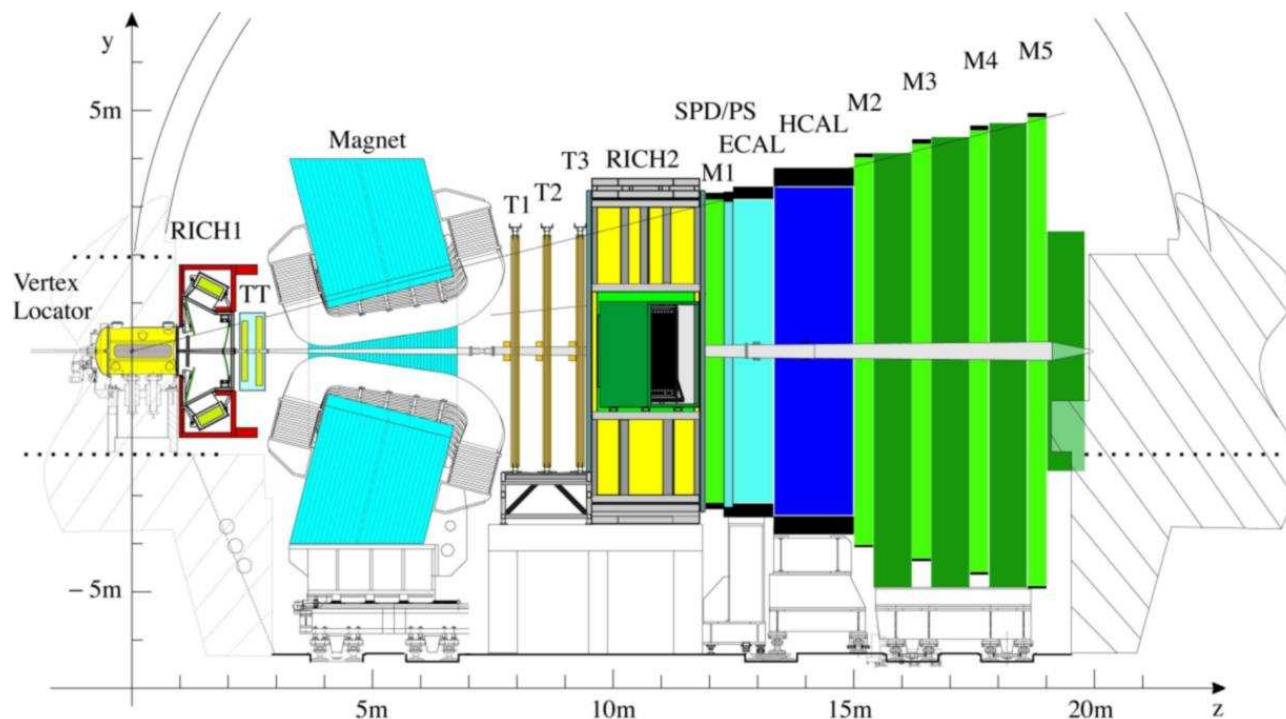
EFFICIENT TRIGGER
FOR NON-LEPTONS

DISTINGUISH π/K

$B \rightarrow D^*\pi, D^*3\pi$ $B \rightarrow \rho\pi$	$B \rightarrow DK^*$ $B \rightarrow K^*\gamma$	<p>RESOLVE x_s</p> $B_s \rightarrow J/\psi\phi$ $B_s \rightarrow J/\psi K_s$
$B_s \rightarrow D_s\pi$ $B_{d,s} \rightarrow D_{d,s}^+ D_{d,s}^-$	$B_s \rightarrow KK$ $B_s \rightarrow D_s K$ $B_s \rightarrow K^{*+} K^{*-}$	
$B \rightarrow \rho^+ \rho^-$ $B \rightarrow K^{*0} \bar{K}^{*0}$	$B \rightarrow \pi\pi$ $B \rightarrow K\pi$	

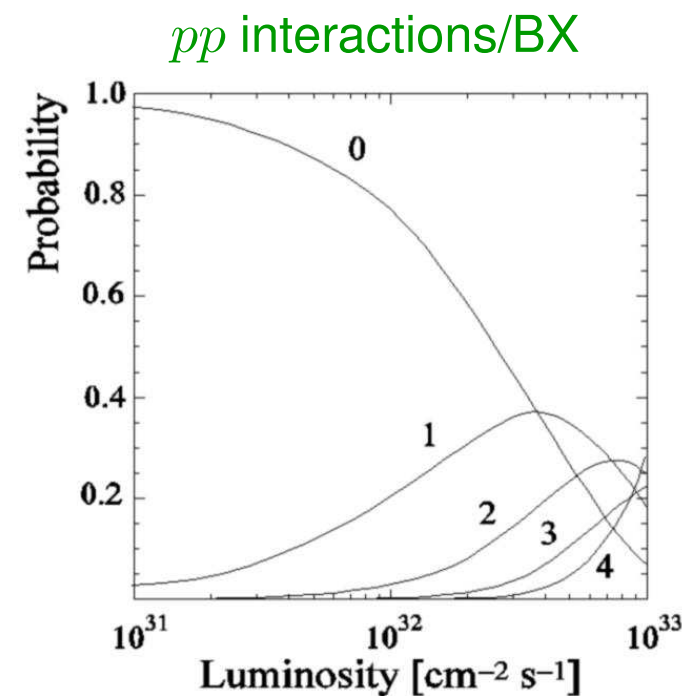
$B_d \rightarrow J/\psi K_s$ $B_d \rightarrow J/\psi \rho^0$ $B_{s,d} \rightarrow \mu^+ \mu^- \quad (\mathcal{O}(10^{-9})!)$
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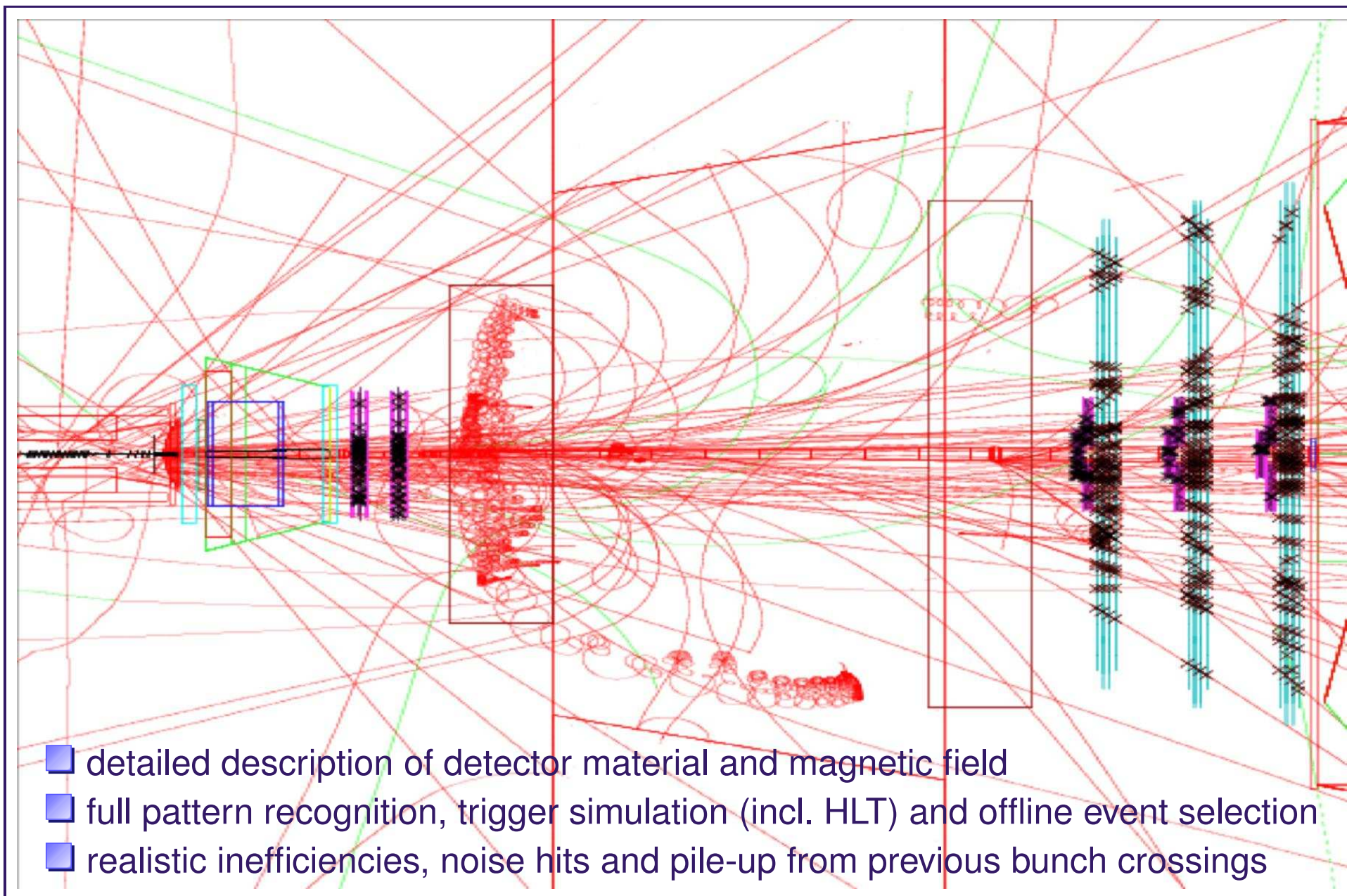
→ LHCb: optimized to meet these requirements



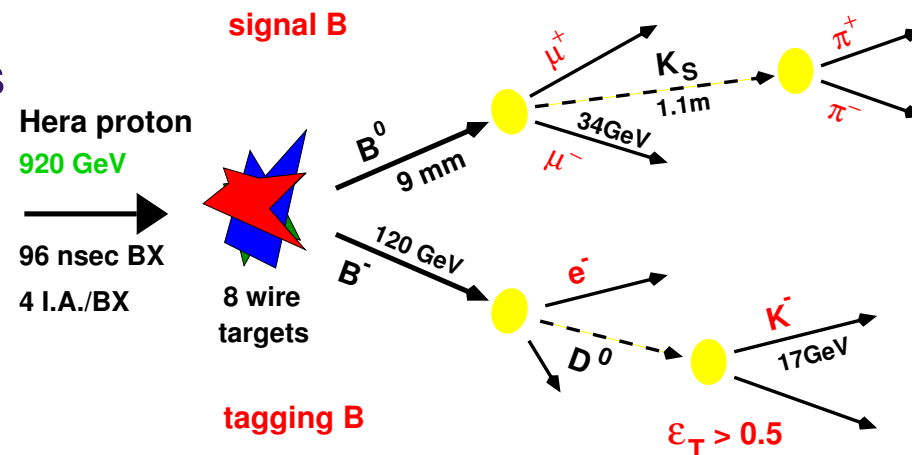
- forward spectrometer
- excellent particle-ID
- high vertex resolution
- flexible trigger
 - ✗ 1 MHz readout to HLT
 - ✗ 2 kHz to disk

- $\mathcal{L} = 2 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$, adjustable
- 10^5 b -events/second, $2 \cdot 10^{12}$ b -hadrons/year
- $B^0 : B^\pm : B_s : \Lambda_b \approx 0.8 : 0.8 : 0.2 : 0.2$
- lepton and hadron trigger
 - ✗ 200 Hz exclusive B-candidates
 - ✗ 600 Hz high mass Di-muons
 - ✗ 300 Hz D^* candidates
 - ✗ 900 Hz inclusive b -trigger





- obvious for charged B-mesons
- experimental challenge for neutral B-mesons
 - opposite side tagging
 - ✗ reconstruct flavour of “idler”-B
 - same-side tagging
 - ✗ self-tagging B^{**} -decays
 - ✗ fragmentation products



note: effective fraction of perfect tags

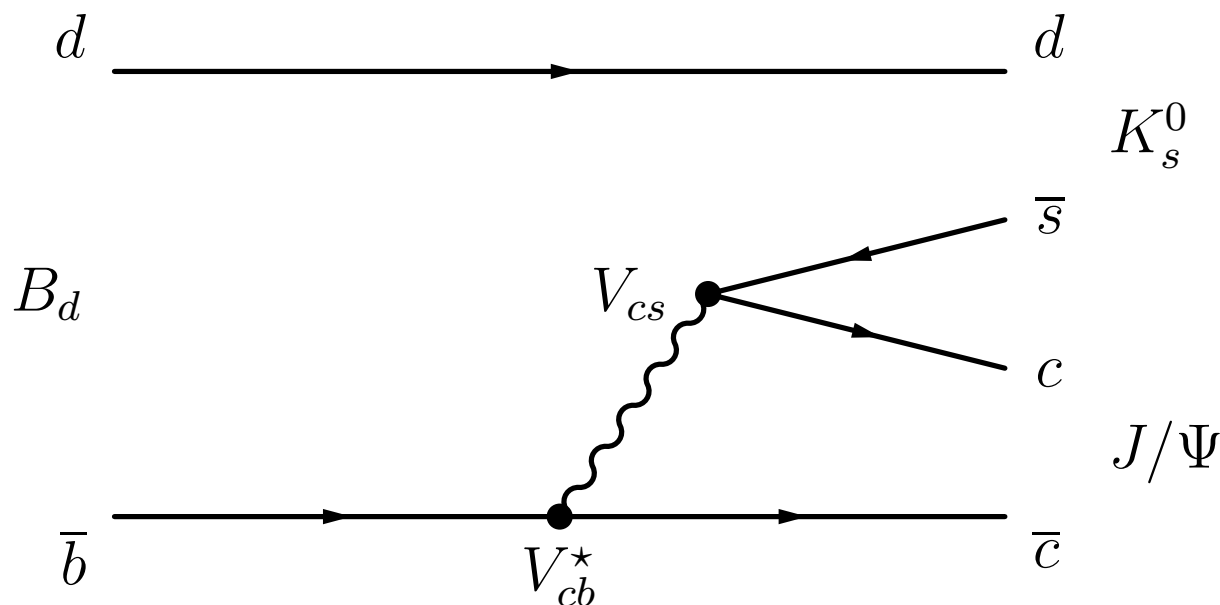
$$\varepsilon_{\text{eff}} = \varepsilon_{\text{tag}}(1 - 2w)^2$$

with

- ε_{tag} tagging efficiency
- w mistag fraction

Channel	ε_{tag} (%)	w (%)	ε_{eff} (%)
$B^0 \rightarrow \pi^+ \pi^-$	41.8 ± 0.7	34.9 ± 1.1	3.8 ± 0.5
$B^0 \rightarrow K^+ \pi^-$	43.2 ± 1.4	33.3 ± 2.1	4.8 ± 1.0
$B^0 \rightarrow J/\psi(\mu\mu)K_S^0$	45.1 ± 1.3	36.7 ± 1.9	3.2 ± 0.8
$B^0 \rightarrow J/\psi(\mu\mu)K^{*0}$	41.9 ± 0.5	34.3 ± 0.7	4.1 ± 0.3
$B_S^0 \rightarrow K^+ K^-$	49.8 ± 0.5	33.0 ± 0.8	5.8 ± 0.5
$B_S^0 \rightarrow \pi^+ K^-$	49.5 ± 1.8	30.4 ± 2.6	7.6 ± 1.7
$B_S^0 \rightarrow D_s^- \pi^+$	54.6 ± 1.2	30.0 ± 1.6	8.7 ± 1.2
$B_S^0 \rightarrow D_s^\mp K^\pm$	54.2 ± 0.6	33.4 ± 0.8	6.0 ± 0.5
$B_S^0 \rightarrow J/\psi(\mu\mu)\phi$	50.4 ± 0.3	33.4 ± 0.4	5.5 ± 0.3

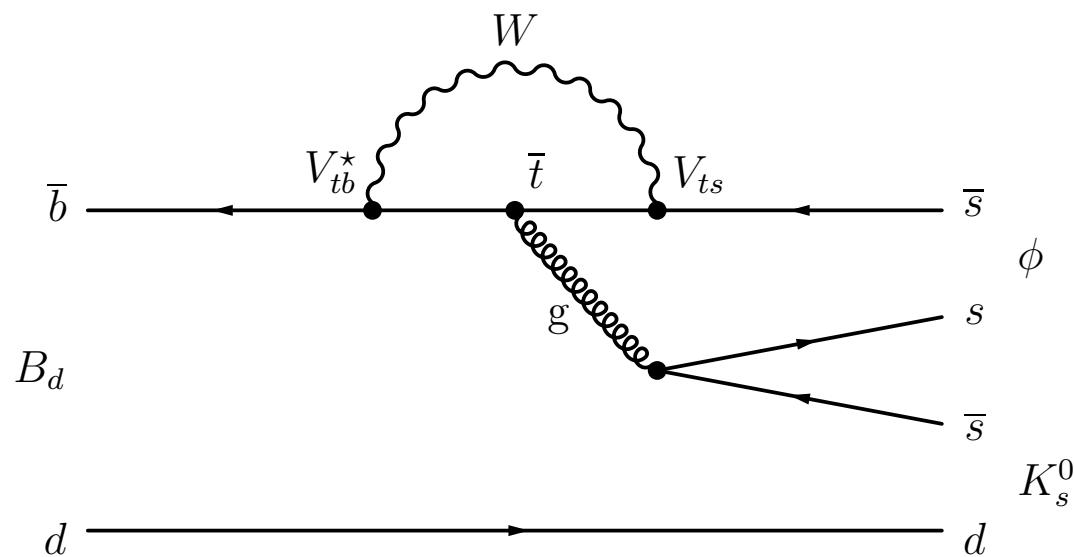
The “golden Decay” $B_d \rightarrow J/\Psi K_s$



- “classical” example for mixing induced CP-violation
- in Wolfenstein-parametrization no phase from decay
- in Standard-Model measures the B_d -mixing phase
- dominated by tree level contributions
- small or no New Physics contributions expected

→ Measurement of $\sin 2\beta$

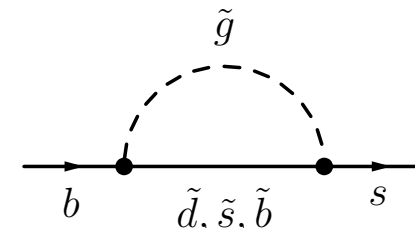
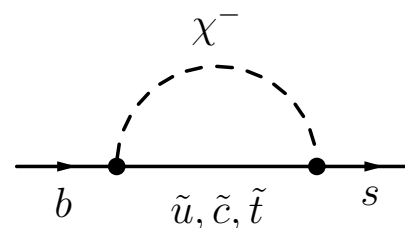
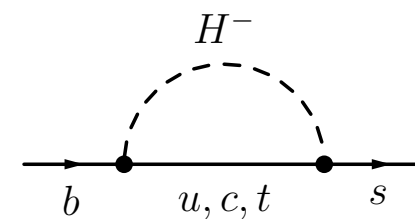
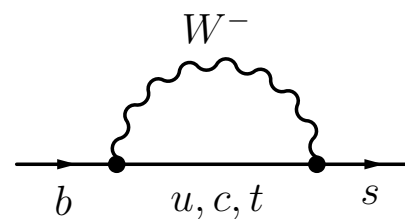
Standard Model: alternative way of measuring $\sin 2\beta$ in a pure penguin mode



the weak phase from decay is dominated by $\arg(V_{td})$ which is $\mathcal{O}(\lambda^2)$

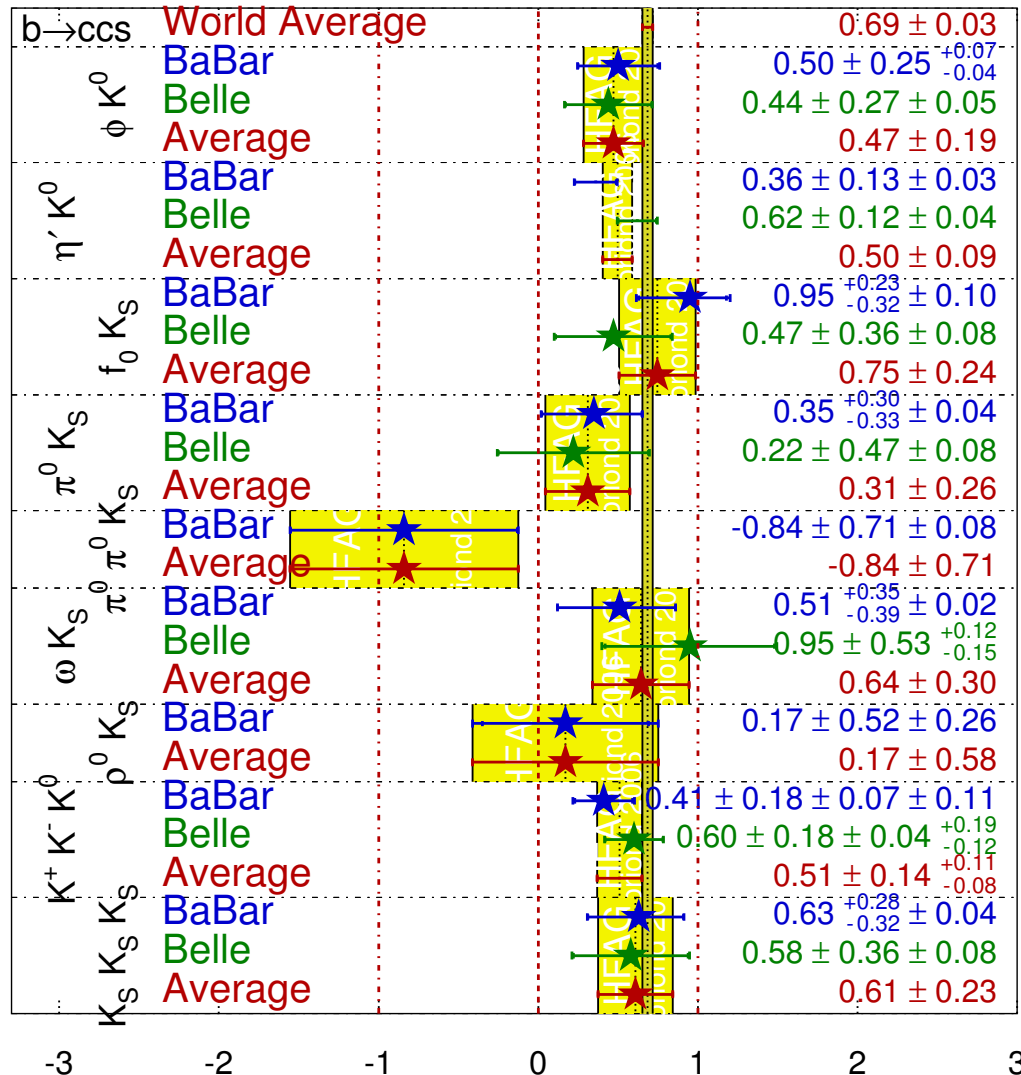
→ possible contributions from new physics:

- new heavy bosons with CKM-like couplings
- new contributions to FCNC transitions
 - ✗ would show up elsewhere as well
 - ✗ constrained by other limits



$\sin(2\beta^{\text{eff}})/\sin(2\phi_1^{\text{eff}})$

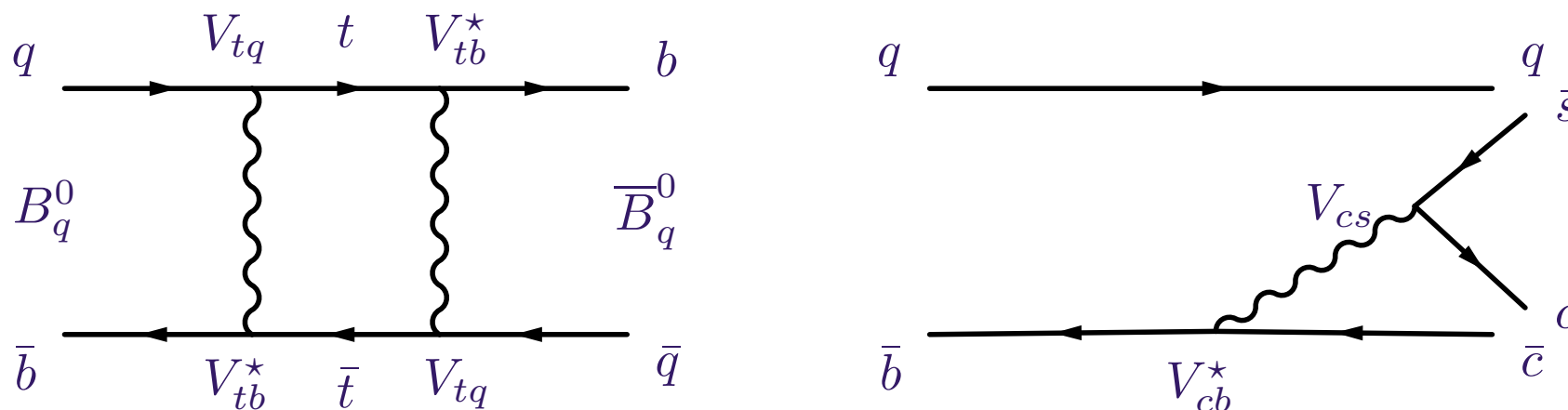
HFAAG
Moriond 2006
PRELIMINARY



- $n = 16$ measurements
- 14 below $b \rightarrow c\bar{c}s$
- consider the probability that from n measurements which fluctuate symmetrically around the $b \rightarrow c\bar{c}s$ average, at most $M = 2$ are above:

$$p = 2^{-n} \sum_{k=0}^M \binom{n}{k} \approx 0.0021$$

→ 3-sigma effect . . . stay tuned!



- similar to the “golden decay” $B_d \rightarrow J/\psi K_s$ with $d \rightarrow s$
- theoretically clean measurement, only very small penguin contribution
- only small CP-asymmetry expected in SM: $\phi_s = -\arg V_{ts}^2 \approx -2\eta\lambda^2 \approx -0.04$
- sensitive probe to NP, e.g.
 - gluinos and squarks in box diagram: $\sin \phi_s \sim 1$ (Ball et al., hep-ph/0311361)
 - up-type quark singlets: $\sin \phi_s \sim \lambda$ (Aguilar-Saavedra et al., hep-ph/0406151)
- mix of different CP eigenstates because two CP-odd vector mesons in final state
 - CP-even states from scalar 0 and longitudinal \parallel polarization
 - CP-odd states from transverse \perp polarization
- study also final states which are CP-eigenstates, e.g.
 - $B_s \rightarrow J/\psi \eta$, $B_s \rightarrow \eta_c \phi$, ...

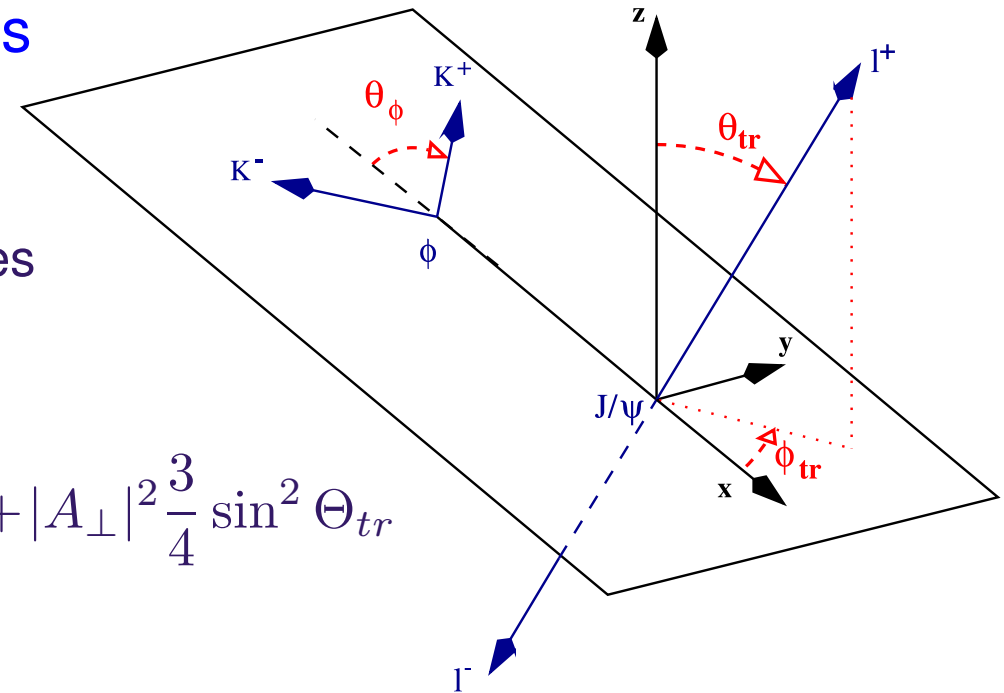
→ dealing with different CP-Eigenstates

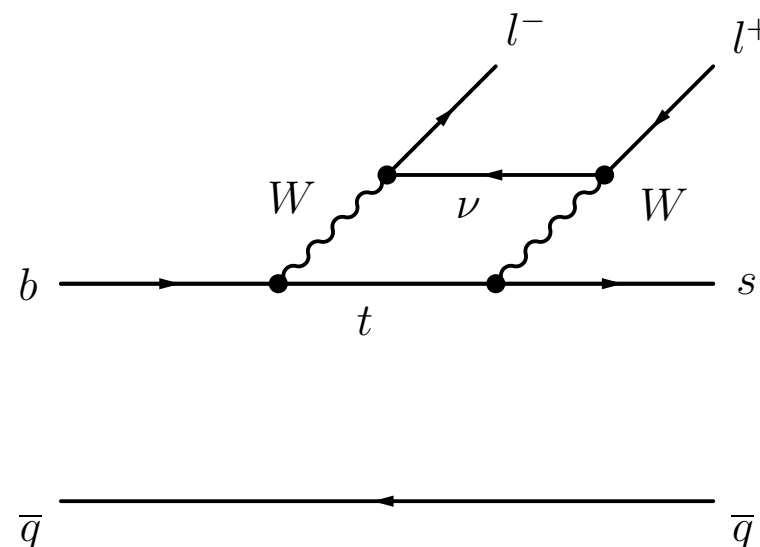
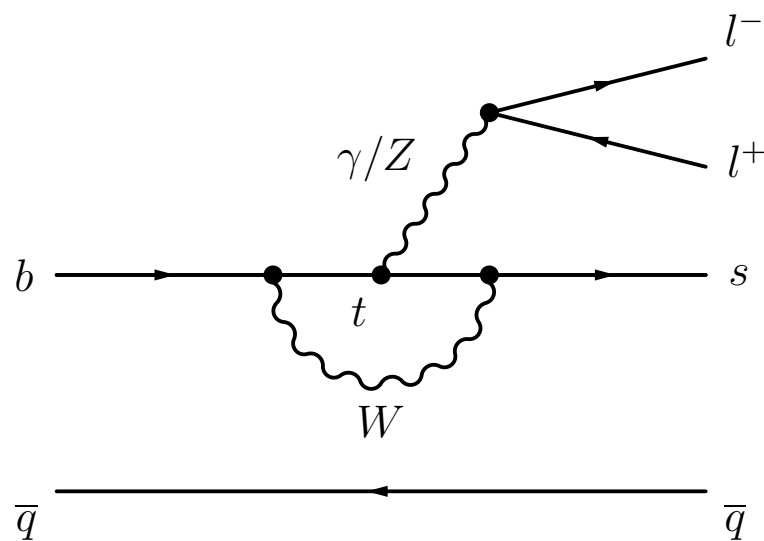
- standard measurement for CP-Eigenstates
- angular analysis for $B_s \rightarrow J/\Psi \phi$

$$\frac{d\Gamma}{d \cos \Theta_{tr}} \propto [|A_0|^2 + |A_{||}|^2] \frac{3}{8} (1 + \cos^2 \Theta_{tr}) + |A_{\perp}|^2 \frac{3}{4} \sin^2 \Theta_{tr}$$

→ Results (L. Fernandez/J. van Hunen)

Channel	CP-Eigenstate	$\sigma(\phi_s)[\text{rad}]$
$B_s \rightarrow J/\Psi \phi$	mixed	0.023
$B_s \rightarrow D_s D_s$	even	0.133
$B_s \rightarrow J/\Psi \eta(\pi^+ \pi^- \pi^0)$	even	0.129
$B_s \rightarrow J/\Psi \eta(\gamma \gamma)$	even	0.108
$B_s \rightarrow \eta_v \phi$	even	0.108
combined		0.021





■ many possible New Physics contributions

→ in $bs\gamma$ -vertex, which contributes also to $B \rightarrow K^* \gamma$ with $A_{CP}^{SM} \leq 0.01$

→ extra Z-bosons

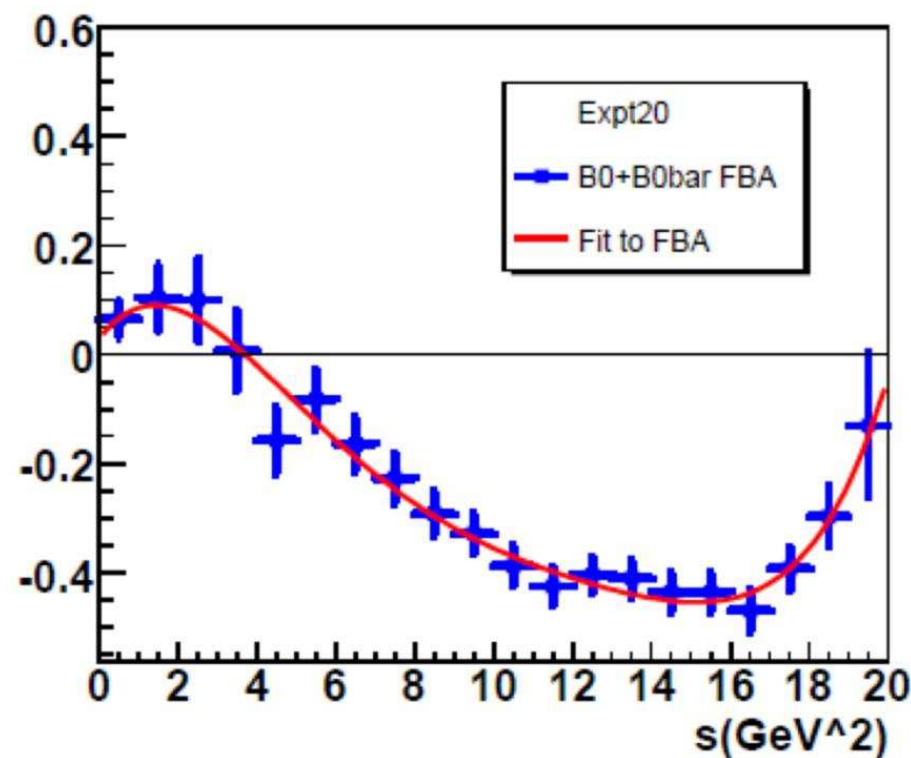
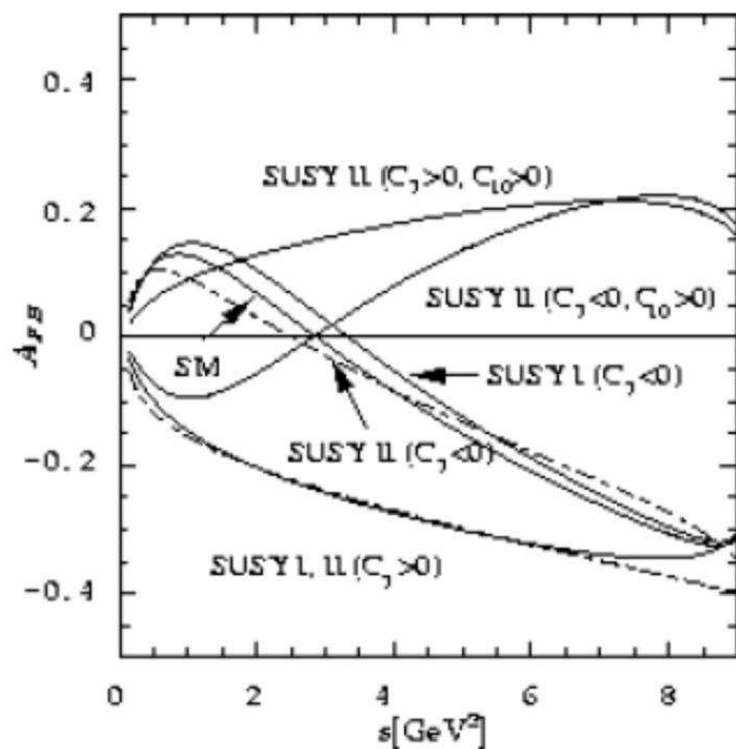
→ new particles in box diagrams

■ theoretically rather clean prediction for inclusive processes and ratios

■ still good accuracy for experimentally accessible exclusive final states

→ $B_d \rightarrow X \mu^+ \mu^-$ with $X = K^*, \rho, \phi$

→ study FB-asymmetry of the muons w.r.t. their common momentum ...



forward-backward asymmetry

- SM predicts zero around $s = 3 \text{ GeV}^2$
- very small hadronic uncertainties
- supersymmetry naturally expects no change of sign

LHCb sensitivity after 5 nominal years

- 22k events
- zero point $s_0 = 4.0 \pm 0.5 \text{ GeV}^2$
- constraint on C_7^{eff}

note: only continuum is shown – contributions from J/Ψ , Ψ' etc. are suppressed

- B-Physics at LHC is a excellent place to look for new physics
 - many ways to look for New Physics by overconstraining the SM
 - ✗ SM (tree level) dominated measurements
 - ✗ penguin dominated processes
 - ✗ box diagram dominated decays
 - comparison of results will reveal New Physics
 - sensitivity to new particles in quantum corrections up $\mathcal{O}(\text{TeV})$
- LHCb is optimized to exploit the B-Physics potential of LHC
 - dedicated to study of B-physics and rare decays
 - full physics potential from day-1 of LHC
 - precision measurements of couplings
 - “worst case” will be precision measurements of CKM matrix
 - ✗ historical precedent: Tycho Brahe . . .
- potential for charm physics not yet explored
- thinking has already started towards upgrades . . .

	Channel	Yield*	B_{bb}/S	Precision
γ	$B_s \rightarrow D_s K$	5.4k	<1	$\sigma(\gamma) \approx 14^\circ$
	$B_d \rightarrow \pi\pi$	26k	<0.7	
	$B_s \rightarrow KK$	37k	0.3	$\sigma(\gamma) \approx 6^\circ$
	$B_d \rightarrow D^0(K^-\pi^+)K^{*0}$	0.5k	<0.3	
	$B_d \rightarrow D^0(K^+\pi^-)K^{*0}$	2.4k	<2	$\sigma(\gamma) \approx 8^\circ$
	$B_d \rightarrow D_{CP}(K^+K^-)K^{*0}$	0.6k	<0.3	
	$B^- \rightarrow D^0(K^-\pi^+)K^-$	60k	0.5	$\sigma(\gamma) \approx 5^\circ$
	$B^- \rightarrow D^0(K^+\pi^-)K^-$	2k	0.5	
α	$B_d \rightarrow \pi^0\pi^-\pi^+$	14k	0.8	$\sigma(\alpha) \approx 10^\circ$
ϕ_s	$B_s \rightarrow J/\psi\Phi$	125k	0.3	$\sigma(\phi_s) \approx 2^\circ$
	$B_s \rightarrow J/\psi\eta$	12k	2-3	
	$B_s \rightarrow \eta_c\Phi$	3k	0.7	
Δm_s	$B_s \rightarrow D_s\pi$	80k	0.3	Δm_s up to 68 ps^{-1}
β	$B_d \rightarrow J/\psi K_S$	216k	0.8	$\sigma(\sin 2\beta) \approx 0.022$
rare decays	$B_d \rightarrow K^*\mu^+\mu^-$	4.4k	<2.6	$C_7^{\text{eff}}/C_9^{\text{eff}}$ with 13% error NP search $\sigma(A_{CP}^{\text{dir}}) \approx 0.01$
	$B_s \rightarrow \mu^+\mu^-$	17	<5.7	
	$B_d \rightarrow K^*\gamma$	35k	<0.7	