

# Strong phase transition in two-Higgs-doublet models

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PLANCK 2013  
Bonn – May 21, 2013

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- Not satisfied in SM unless  $m_h \lesssim m_W$ , which is not the case.
- CP violation from CKM matrix too small.

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- Correct BAU can be obtained for simplified cases and for particular combinations of parameters.
- But what happens in the general case?

# 2HDM

- Avoid FCNC  $\Rightarrow \mathbb{Z}_2$  symmetry:  $\Phi_1 \rightarrow -\Phi_1$ ,  $\Phi_2 \rightarrow \Phi_2$ .

	$u_R$	$d_R$	$e_R$
Type I	+	+	+
Type II	+	-	-
Type X	+	+	-
Type Y	+	-	+

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Type X	+	+	-
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- For PT, only top-quark needs to be considered.  
Then models differ only in phenomenological constraints on their parameter space. These come mainly from  $B$ -physics, so

$$\begin{aligned} \text{Type I} &\sim \text{Type X}, \\ \text{Type II} &\sim \text{Type Y}. \end{aligned}$$

- For simplicity, consider CP conserving case only.

$$\begin{aligned}
 V_{tree}(\Phi_1, \Phi_2) = & -\mu_1^2 \Phi_1^\dagger \Phi_1 - \mu_2^2 \Phi_2^\dagger \Phi_2 - \frac{\mu^2}{2} (\Phi_1^\dagger \Phi_2 + H.c.) + \\
 & + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \\
 & + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} \left[ (\Phi_1^\dagger \Phi_2)^2 + H.c. \right].
 \end{aligned}$$

- EW minimum:  $\langle \Phi_1 \rangle = \begin{pmatrix} 0 \\ v \cos \beta \end{pmatrix}$ ,  $\langle \Phi_2 \rangle = \begin{pmatrix} 0 \\ v \sin \beta \end{pmatrix}$ .
- Physical parameters:
  - $v \approx 174$  GeV and  $M \equiv \frac{\mu}{\sqrt{\sin(2\beta)}}$ .
  - Masses:  $m_{h^0}, m_{H^0}, m_{A^0}, m_{H^\pm}$ .
  - $\beta$  is the mixing angle between  $(G^+, H^+)$  and between  $(G^0, A^0)$ .
  - Likewise,  $\alpha$  is the mixing angle between  $(h^0, H^0)$ .  
It is here defined such that  $\alpha = \beta \iff h^0 = h_{SM}$ .

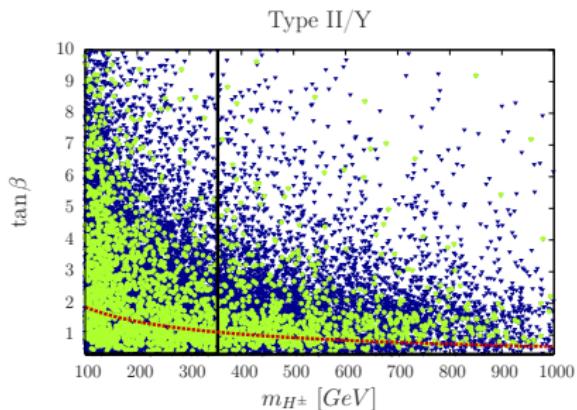
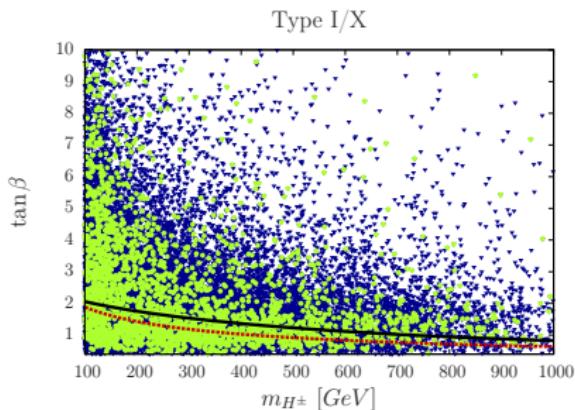
$$V = V_{tree} + V_{CW} + V_{CT} + V_T.$$

- Fix  $m_{h^0} = 125$  GeV

$$\begin{aligned} 0.4 &\leq \tan\beta \leq 10, \\ -\frac{\pi}{2} &< \alpha \leq \frac{\pi}{2}, \\ 0 \text{ GeV} &\leq \mu \leq 1 \text{ TeV}, \\ 100 \text{ GeV} &\leq m_{A^0}, m_{H^\pm} \leq 1 \text{ TeV}, \\ 150 \text{ GeV} &\leq m_{H^0} \leq 1 \text{ TeV}. \end{aligned}$$

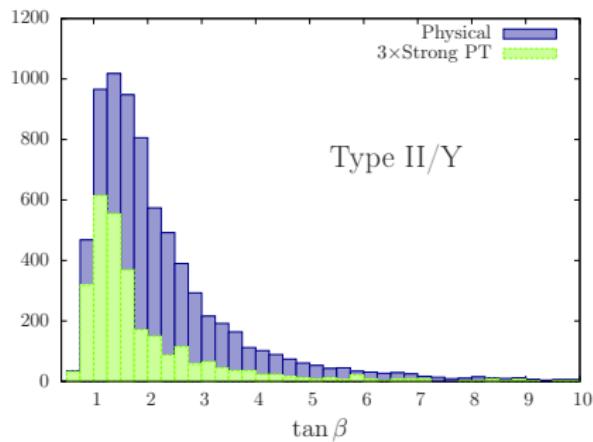
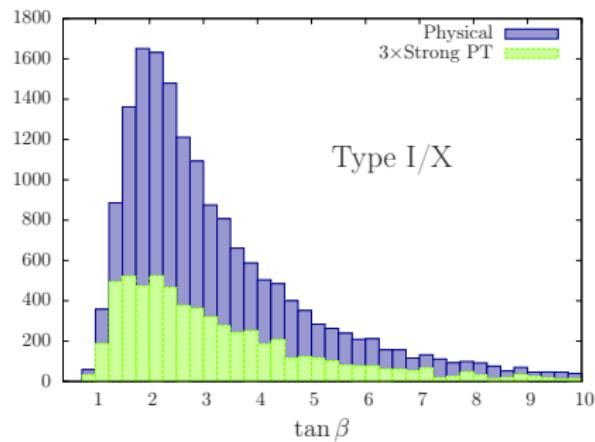
- Constraints:

- EW precision:  $\rho - 1 \approx 0$ ;
- $\lambda_i < 4\pi$ ;
- metastability;
- $B^0 - \bar{B}^0$  (red/dashed) and  $\bar{B} \rightarrow X_s \gamma$  (black/full).



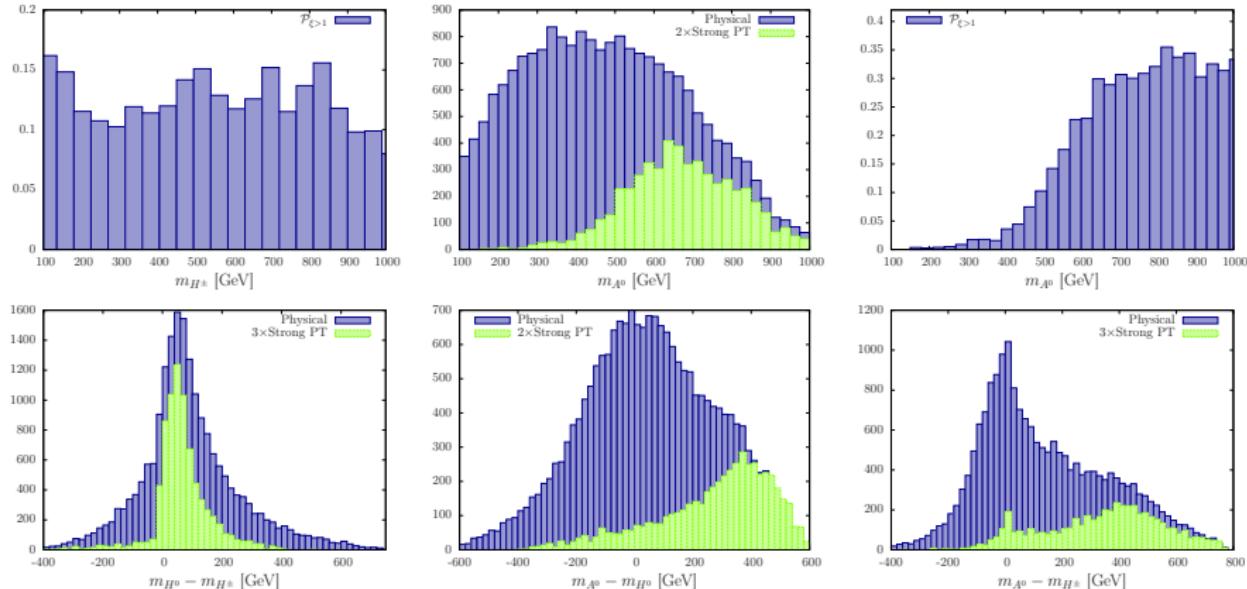
Type II/Y:  $m_{H^\pm} \geq 360$  GeV [Hermann et al., arXiv:1208.2788].

# Results: $\tan \beta$



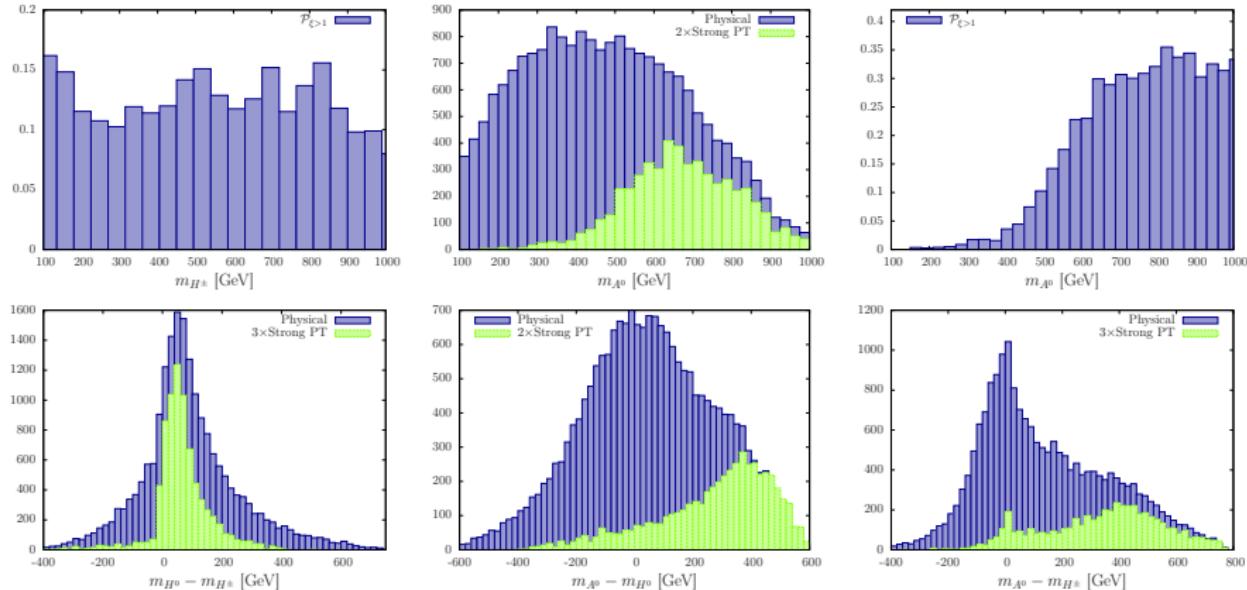
Preference for  $\tan \beta \lesssim 3$  is excellent for baryogenesis, since  
 $n_B \sim (\tan \beta)^{-2}$ .

# Results: Masses



$m_{H^\pm}$  hardly influences the phase transition.

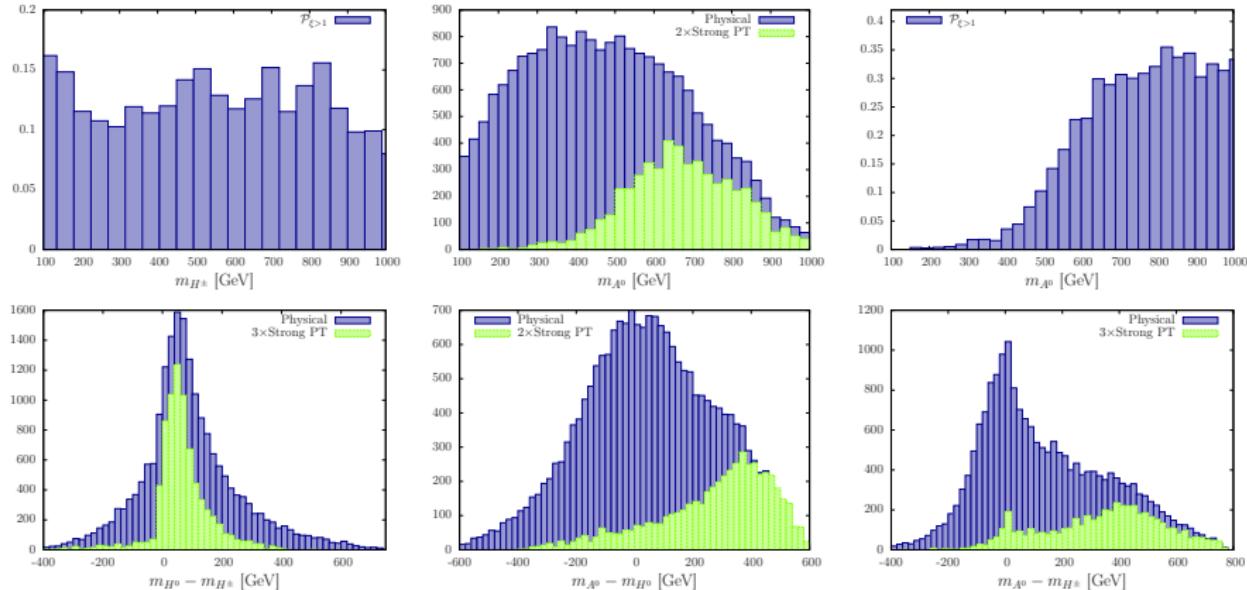
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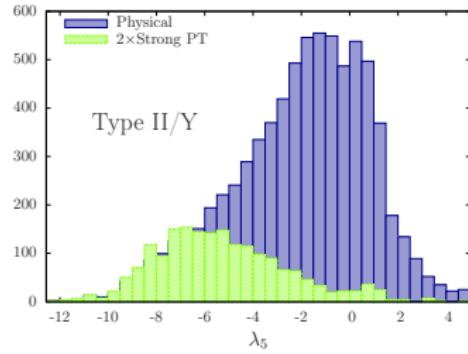
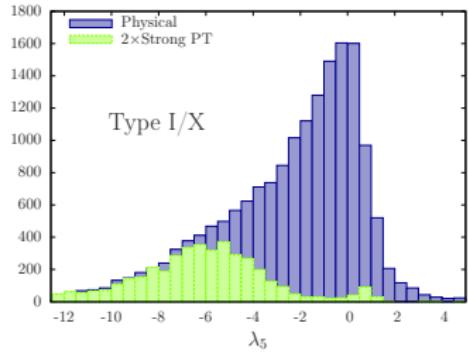
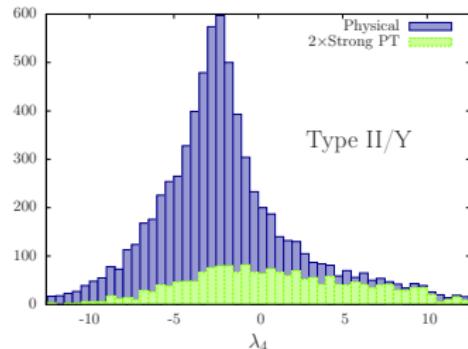
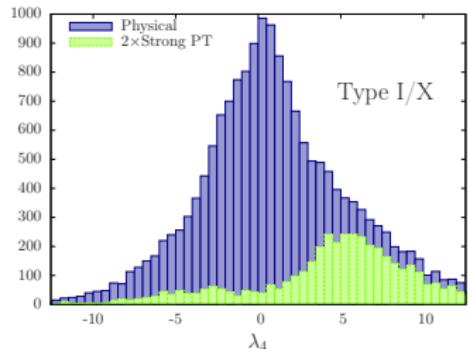
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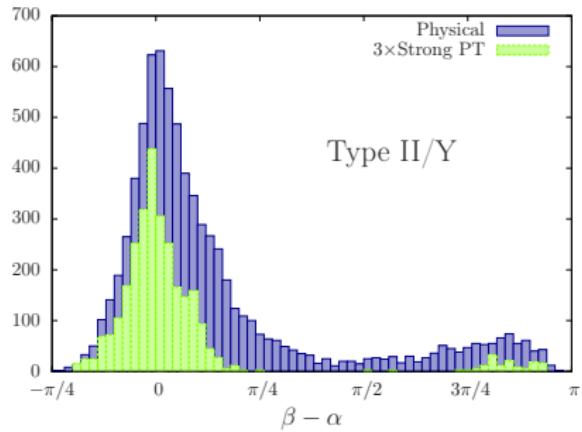
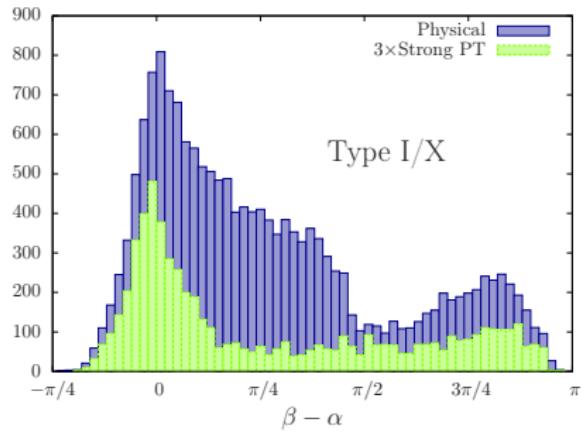
Strong PTs also prefer hierarchy  $m_{A^0} > m_{H^0} \gtrsim m_{H^\pm}$ .

# Results: Couplings

$$\lambda_4 = \frac{1}{2v^2} (M^2 + m_{A^0}^2 - 2m_{H^\pm}^2), \quad \lambda_5 = \frac{1}{2v^2} (M^2 - m_{A^0}^2).$$

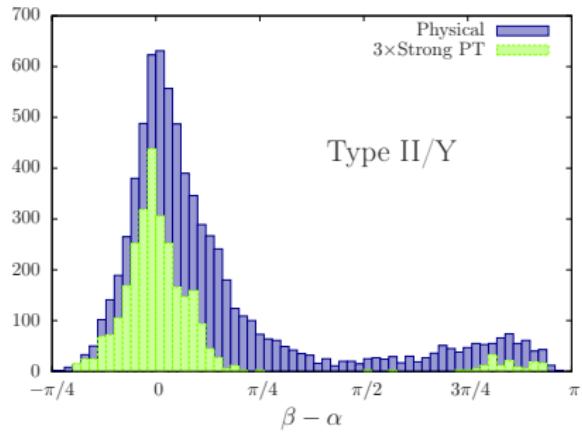
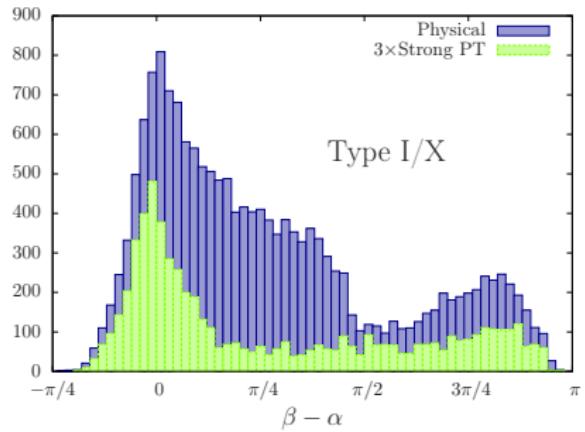


# Results: $\beta - \alpha$



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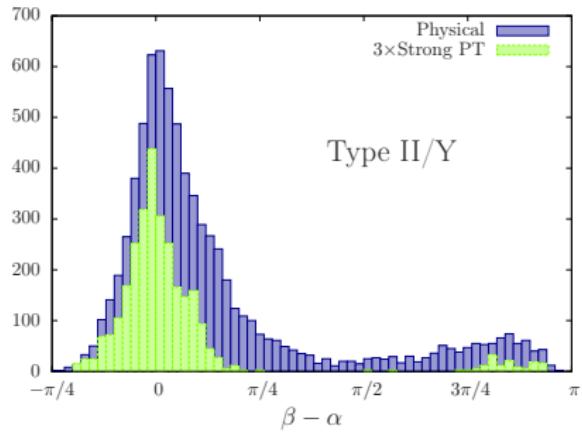
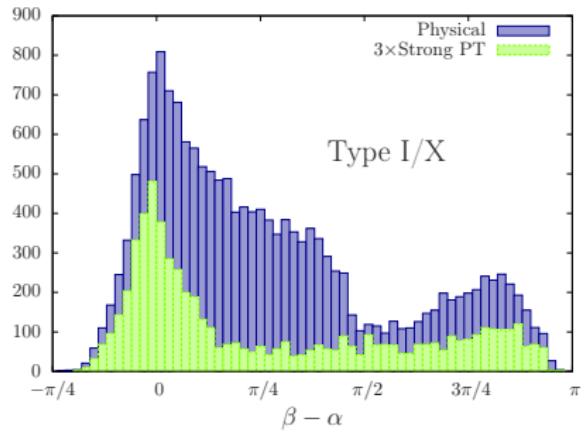
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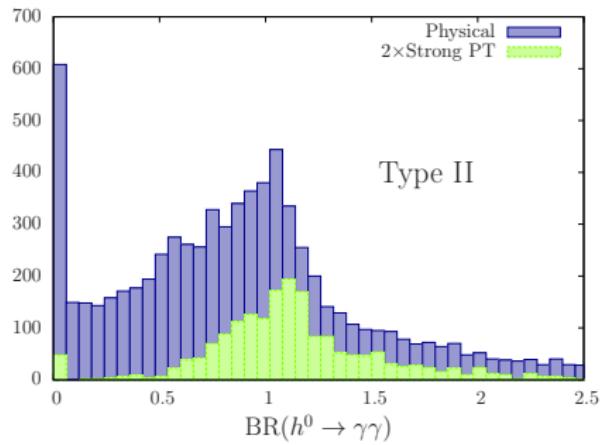
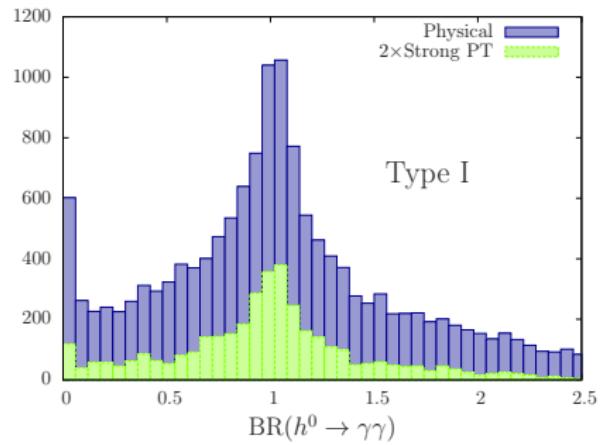


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Moreover,  $\alpha \lesssim \beta$  is slightly preferred.

# Results: $h^0 \rightarrow \gamma\gamma$



Coupling of  $h^0$  to  $b$  and  $\tau$ :

$$\text{Type I: } \frac{\sin \alpha}{\sin \beta}, \quad \text{Type II: } \frac{\cos \alpha}{\cos \beta}.$$

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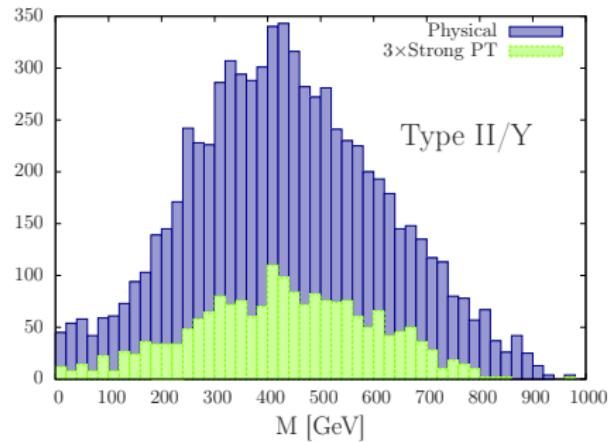
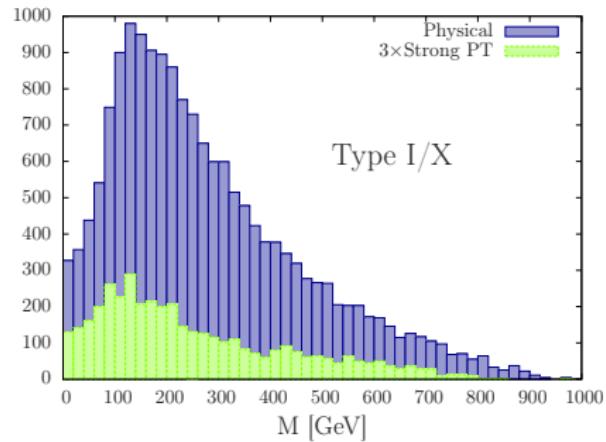
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Thank you!

# Appendix – Scale of new physics

$$M \equiv \frac{\mu}{\sqrt{\sin(2\beta)}}$$



Type II/Y:  $m_{H^\pm} \geq 360$  GeV.

# Appendix

- Surviving points after each step of tests:

	Total	EW precision	$\lambda_i < 4\pi$	Metastability	Strong PT
Absolute	$6.3 \cdot 10^6$	$1.2 \cdot 10^6$	$1.4 \cdot 10^5$	$2.6 \cdot 10^4$	$3.7 \cdot 10^3$
Relative	100%	19.1%	2.3%	0.41%	0.059%

- Physical fields:

$$\begin{aligned} G^+ &= \cos \beta \varphi_1^+ + \sin \beta \varphi_2^+ && (\text{charged Goldstone}), \\ H^+ &= -\sin \beta \varphi_1^+ + \cos \beta \varphi_2^+ && (\text{charged scalar}), \\ G^0 &= \cos \beta \eta_1 + \sin \beta \eta_2 && (\text{neutral Goldstone}), \\ A^0 &= -\sin \beta \eta_1 + \cos \beta \eta_2 && (\text{pseudo-scalar}), \\ h^0 &= \cos \alpha h_1 + \sin \alpha h_2 && (\text{lightest scalar}), \\ H^0 &= -\sin \alpha h_1 + \cos \alpha h_2 && (\text{heaviest scalar}). \end{aligned}$$

where

$$\Phi_i = \begin{pmatrix} \varphi_i^+ \\ h_i + i\eta_i \end{pmatrix}.$$