Production amplitudes in N=4 SUSY and integrability

J. Bartels¹, $L.N.Lipatov^{1,3}$, $A.Sabio Vera^2$

Hamburg University¹, $CERN^2$, Petersburg Nuclear Physics Institute³

Content

- 1. Gluon reggeization
- 2. Effective action
- 3. BFKL Pomeron
- 4. BKP equation and integrability
- 5. Pomeron in N = 4 SUSY
- 6. BDS amplitudes at high energies
- 7. Absence of the Regge factorization
- 8. Mandelstam cuts
- 9. Open integrable Heisenberg spin chain

1 Gluon reggeization

Regge kinematics

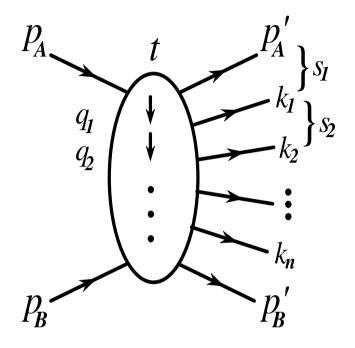
$$s = 4E^2 \gg -t = |q|^2 \approx E^2 \theta^2$$

Amplitude for the colored particle scattering

$$M_{AB}^{A'B'}(s,t)|_{LLA} = 2 g T_{A'A}^c \delta_{\lambda_{A'}\lambda_A} \frac{s^{1+\omega(t)}}{t} g T_{B'B}^c \delta_{\lambda_{B'}\lambda_B}$$

Gluon Regge trajectory in LLA

$$\omega(-|q|^2) = -\frac{\alpha_s N_c}{4\pi^2} \int d^2k \, \frac{|q|^2}{|k|^2 |q-k|^2} \approx -\frac{\alpha_s N_c}{2\pi} \, \ln \frac{|q^2|}{\lambda^2}$$



Multi-Regge amplitudes (F.,K.,L. (1975))

$$M_{2\to 1+n} \sim \frac{s_1^{\omega_1}}{|q_1|^2} g T_{c_2 c_1}^{d_1} C(q_2, q_1) \frac{s_2^{\omega_2}}{|q_2|^2} ... C(q_n, q_{n-1}) \frac{s_n^{\omega_n}}{|q_n|^2},$$

$$C(q_2, q_1) = \frac{q_2 q_1^*}{q_2^* - q_1^*}, \ \sigma_t = \sum_n \int d\Gamma_n |M_{2\to 1+n}|^2$$

2 Effective action approach

Gluon and Reggeized gluon fields

$$v_{\mu}(x) = -iT^{a}v_{\mu}^{a}(x), \ A_{\pm}(x) = -iT^{a}A_{\pm}^{a}(x)$$

Local gauge transformations

$$\delta v_{\mu}(x) = \frac{1}{g} [D_{\mu}, \chi(x)], \ \delta \psi(x) = -\chi(x) \ \psi(x), \ \delta A_{\pm}(x) = 0$$

Effective action for reggeized gluons (L., 1995)

$$S = \int d^4x \left(L_0 + L_{ind}^{GR} \right) , \ L_0 = i\bar{\psi}\hat{D}\psi + \frac{1}{2} Tr G_{\mu\nu}^2$$

$$L_{ind}^{GR} = -\frac{1}{g}\partial_{+} P \exp\left(-g\frac{1}{2}\int_{-\infty}^{x^{+}} v_{+}(x')d(x')^{+}\right) \partial_{\sigma}^{2} A_{-} + (+ \to -)$$

3 BFKL equation (1975)

Balitsky-Fadin-Kuraev-Lipatov equation

$$E \Psi(\vec{\rho}_1, \vec{\rho}_2) = H_{12} \Psi(\vec{\rho}_1, \vec{\rho}_2) , \ \sigma_t \sim s^{\Delta}, \ \Delta = -\frac{\alpha_s N_c}{2\pi} E$$

BFKL Hamiltonian

$$H_{12} = \ln |p_1 p_2|^2 + \frac{1}{p_1 p_2^*} (\ln |\rho_{12}|^2) p_1 p_2^*$$

$$+\frac{1}{p_1^*p_2}(\ln|\rho_{12}|^2)p_1^*p_2 - 4\psi(1), \ \rho_{12} = \rho_1 - \rho_2$$

Möbius invariance and conformal weights (L. (1986))

$$\rho_k \to \frac{a\rho_k + b}{c\rho_k + d},$$

$$m = \gamma + n/2, \ \widetilde{m} = \gamma - n/2, \ \gamma = 1/2 + i\nu$$

4 BKP equation (1980)

Bartels-Kwiecinski-Praszalowicz equation

$$E \Psi(\vec{\rho}_1, ..., \vec{\rho}_n) = H \Psi(\vec{\rho}_1, ..., \vec{\rho}_n) , H = \sum_{k < l} \frac{\vec{T}_k \vec{T}_l}{-N_c} H_{kl}$$

Holomorphic separability at large N_c (L. (1988))

$$H = h + h^*, \ h_{12} = \ln p_1 + \ln p_2 + \frac{1}{p_1} (\ln \rho_{12}) p_1 + \frac{1}{p_2} (\ln \rho_{12}) p_2 - 2\psi(1)$$

Holomorphic factorization of wave functions

$$\Psi(\vec{\rho}_1, \vec{\rho}_2, ..., \vec{\rho}_n) = \sum_{r,s} a_{r,s} \, \Psi_r(\rho_1, ..., \rho_n) \, \Psi_s(\rho_1^*, ..., \rho_n^*)$$

5 Integrability at $N_c \to \infty$

Monodromy and transfer matrices (L. (1993))

$$t(u) = L_1 L_2 ... L_n = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}, T(u) = A(u) + D(u),$$

$$L_k = \begin{pmatrix} u + \rho_k \, p_k & p_k \\ -\rho_k^2 \, p_k & u - \rho_k \, p_k \end{pmatrix}, \ [T(u), T(v)] = [T(u), h] = 0$$

Yang-Baxter equation (L. (1993))

$$t_{r_1'}^{s_1}(u) t_{r_2'}^{s_2}(v) t_{r_1 r_2}^{r_1' r_2'}(v - u) = t_{s_1' s_2'}^{s_1 s_2}(v - u) t_{r_2}^{s_2'}(v) t_{r_1}^{s_1'}(u), \ \hat{l} = u \hat{1} + i \hat{P}$$

Duality symmetry (L. (1999))

$$p_r \to \rho_{r+1,r} \to p_{r+1}$$

Heisenberg spin model (L. (1994); F., K.(1995))

6 Pomeron in N = 4 SUSY

BFKL kernel in two loops (F., L. (1998))

$$\omega = 4 \hat{a} \chi(n, \gamma) + 4 \hat{a}^2 \Delta(n, \gamma), \ \hat{a} = g^2 N_c / (16\pi^2),$$

Hermitian separability in N=4 SUSY (K.,L. (2000))

$$\Delta(n,\gamma) = \phi(M) + \phi(M^*) - \frac{\rho(M) + \rho(M^*)}{2\hat{a}/\omega}, M = \gamma + \frac{|n|}{2},$$

$$\rho(M) = \beta'(M) + \frac{1}{2}\zeta(2), \ \beta'(z) = \frac{1}{4} \left[\Psi'\left(\frac{z+1}{2}\right) - \Psi'\left(\frac{z}{2}\right) \right]$$

Maximal transcendentality (K.,L. (2002))

$$\phi(M) = 3\zeta(3) + \Psi''(M) - 2\Phi(M) + 2\beta'(M) \Big(\Psi(1) - \Psi(M)\Big),$$

$$\Phi(M) = \sum_{k=0}^{\infty} \frac{(-1)^k}{k+M} \left(\Psi'(k+1) - \frac{\Psi(k+1) - \Psi(1)}{k+M} \right)$$

7 Maximal helicity violation

BDS amplitudes in N=4 SUSY at $N_c\gg 1$ (2005)

$$A^{a_1,...,a_n} = \sum_{\{i_1,...,i_n\}} Tr T^{a_{i_1}} T^{a_{i_2}} ... T^{a_{i_n}} f(p_{i_1}, p_{i_2}, ..., p_{i_n}), f = f_B M_n$$

Invariant amplitudes

$$\ln M_n = \sum_{l=1}^{\infty} a^l \left(f^{(l)}(\epsilon) \left(-\frac{1}{2\epsilon^2} \sum_{i=1}^n \left(\frac{\mu^2}{-s_{i,i+1}} \right)^{\epsilon} + F_n^{(1)}(0) \right) + C^{(l)} \right) ,$$

$$a = \frac{\alpha N_c}{2\pi} \left(4\pi e^{-\gamma}\right)^{\epsilon}, \ C^{(1)} = 0, \ C^{(2)} = -\zeta_2^2/2, \ f^{(l)}(\epsilon) = \sum_{k=0}^{2} \epsilon^k f_k^{(1)}$$

Cusp anomalous dimension

$$f_0^{(l)} = \frac{1}{4} \gamma_K^{(l)}, \ f_1 = \beta(f) = -a\zeta_3/2 + a^2(2\zeta_5 + 5\zeta_2\zeta_3/3) + \dots$$

8 Steinman relations

Overlapping channels

$$(s_1, s_2) \ (2 \to 3); \ (s_1, s_2), \ (s_2, s_3), \ (s_{012}, s_2), \ (s_{123}, s) \ (2 \to 4)$$

Dispersion representation for $M_{2\rightarrow 3}$ in the Regge anzatz

$$M_{2\to 3} = c_1(-s)^{j(t_2)}(-s_1)^{j(t_1)-j(t_2)} + c_2(-s)^{j(t_1)}(-s_2)^{j(t_2)-j(t_1)}$$

Violation of the dispersion representation for $M_{2\rightarrow 4}$

$$M_{2\to 4} \neq d_1(-s)^{j_3}(-s_{012})^{j_2-j_3}(-s_1)^{j_1-j_2} + d_2(-s)^{j_1}(-s_{123})^{j_2-j_1}(-s_3)^{j_3-j_2}$$

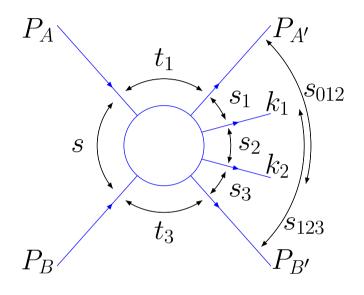
$$+d_3(-s)^{j_3}(-s_{012})^{j_1-j_3}(-s_2)^{j_2-j_1} + d_4(-s)^{j_1}(-s_{123})^{j_3-j_1}(-s_2)^{j_2-j_3}$$

$$+d_5(-s)^{j_2}(-s_1)^{j_1-j_2}(-s_3)^{j_3-j_2}, \ j_r = j(t_r)$$

Important relations

$$\Phi \equiv \frac{(-s)(-s_2)}{(-s_{012})(-s_{123})}, \ Li_2(1-\Phi)_{\Phi \to \exp(2\pi i)} = \ln(1-\Phi) \approx \ln\frac{(\vec{k}_1 + \vec{k}_2)^2}{s_2}$$

9 Regge factorization violation



$$M_{2\to 4}|_{s_2>0; s_1, s_3<0} = \exp\left[\frac{\gamma_K(a)}{4} i\pi \left(\ln\frac{t_1 t_2}{(\vec{k}_1 + \vec{k}_2)^2 \mu^2} - \frac{1}{\epsilon}\right)\right] \times \Gamma(t_1) \left(\frac{-s_1}{\mu^2}\right)^{\omega(t_1)} \Gamma(t_2, t_1) \left(\frac{-s_2}{\mu^2}\right)^{\omega(t_2)} \Gamma(t_3, t_2) \left(\frac{-s_3}{\mu^2}\right)^{\omega(t_3)} \Gamma(t_3)$$

10 Mandelstam cuts in j_2 -plane

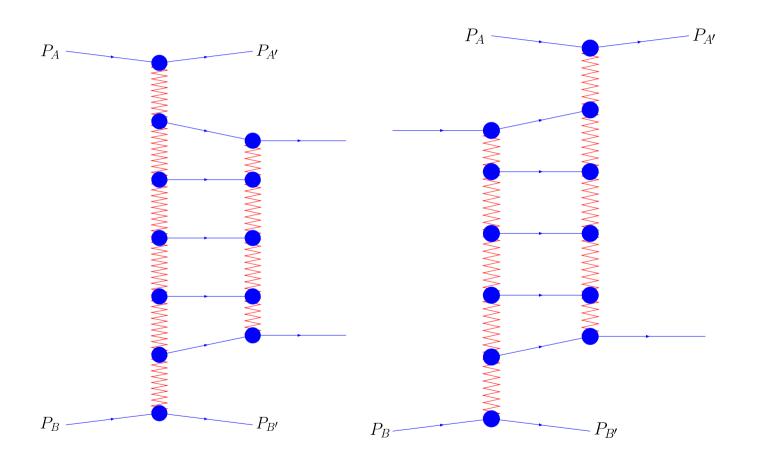


Figure 1: BFKL ladders in $M_{2\rightarrow 4}$ and $M_{3\rightarrow 3}$

11 BFKL equation in octet channels

Factorization of infrared divergencies in LLA

$$\lim_{\epsilon \to 0} M_{2 \to 4}^{LLA} = f_{2 \to 4}^{LLA} \lim_{\epsilon \to 0} M_{2 \to 4}^{BDS} ,$$

Renormalization of the intercept in the s_2 -channel

$$\Delta_2 = -a \left(E + \ln \frac{t_2}{\mu^2} - \frac{1}{\epsilon} \right)$$

BFKL hamiltonian for the partial wave f_{j_2}

$$H = \ln \frac{|p_1 p_2|^2}{|p_1 + p_2|^2} + \frac{1}{2} p_1 p_2^* \ln |\rho_{12}|^2 \frac{1}{p_1 p_2^*} + \frac{1}{2} p_1^* p_2 \ln |\rho_{12}|^2 \frac{1}{p_1^* p_2} + 2\gamma$$

Eigenfunctions and eigenvalues

$$\Psi_{n,\nu} = \left(\frac{p_1}{p_2}\right)^{i\nu+n/2} \left(\frac{p_1^*}{p_2^*}\right)^{i\nu-n/2}, \ E_{n,\nu} = 2Re\,\psi(i\nu + \frac{|n|}{2}) - 2\psi(1)$$

12 Multi-gluon states in octet channels

Holomorphic hamiltonian for n-gluon composite states

$$h = \ln(z_1^2 \partial_1) - 2\psi(1) + \ln \partial_{n-1} + \sum_{k=1}^{n-2} h_{k,k+1}, \ p_k = z_{k-1,k}, \ z_0 = 0, \ z_n = \infty$$

Pair hamiltonian of the spin chain

$$h_{k,k+1} = \ln(z_{k,k+1}^2 \partial_k) + \ln(z_{k-1,k}^2 \partial_k) - \ln(z_{k-1,k+1}^2) + 2\gamma$$

Monodromy matrix

$$t(u) = L_1(u)...L_{n-1}(u) = \begin{pmatrix} A(u) & B(u) \\ C(u) & D(u) \end{pmatrix}, L_k = \begin{pmatrix} u + z_k p_k & p_k \\ -z_k^2 p_k & u - z_k p_k \end{pmatrix}$$

Integrals of motion and Baxter equation for the open spin chain

$$[D(u), h] = 0, \ D(u)Q(u) = (u - i)^{n-1}Q(u - i)$$
14-1

13 Three-gluon composite state

Wave function in the coordinate representation

$$\Psi = z_2^{a_1 + a_2} (z_2^*)^{\widetilde{a}_1 + \widetilde{a}_2} \int \frac{d^2 y}{|y|^2} y^{-a_2} (y^*)^{\widetilde{a}_2} \left(\frac{y - 1}{y - z_2/z_1} \right)^{a_1} \left(\frac{y^* - 1}{y^* - z_2^*/z_1^*} \right)^{a_1}$$

Fourier transformation

$$\Psi = \int d^2p_1 d^2p_2 \exp(i\vec{p}_1\vec{z}_1) \exp(i\vec{p}_2\vec{z}_2) \Psi(\vec{p}_1, \vec{p}_2), \ E = E(a_1) + E(a_2)$$

Baxter-Sklyanin representation

$$\Psi^{t}(\vec{p}_{1}, \vec{p}_{2}) = P^{-a_{1}-a_{2}} (P^{*})^{-\tilde{a}_{1}-\tilde{a}_{2}} \int d^{2}u \, u \, \tilde{u} \, Q(u, \tilde{u}) \left(\frac{p_{1}}{p_{2}}\right)^{u} \left(\frac{p_{1}^{*}}{p_{2}^{*}}\right)^{u^{*}}$$

Baxter function

$$Q(u,\widetilde{u}) = \frac{\Gamma(-u)\Gamma(-\widetilde{u})}{\Gamma(1+u)\Gamma(1+\widetilde{u})} \frac{\Gamma(u-a_1)\Gamma(u-a_2)}{\Gamma(1-\widetilde{u}+\widetilde{a}_1)\Gamma(1-\widetilde{u}+\widetilde{a}_2)}, \int d^2u = \int d\nu \sum_{n=0}^{\infty} \frac{\Gamma(u-a_1)\Gamma(u-a_2)}{\Gamma(1-\widetilde{u}+\widetilde{u})\Gamma(1-\widetilde{u}+\widetilde{u}_2)}$$

14 Discussion

- 1. Reggeized gluons and BFKL equation.
- 2. Effective action for reggeized gluons.
- 3. Integrability of BKP equations at $N_c \to \infty$.
- 4. Remarkable properties of the BFKL kernel in N = 4 SUSY.
- 7. Breakdown of the Regge factorization for BDS amplitudes.
- 8. Mandelstam cuts in the planar diagrams.
- 9. Integrable open spin chain for gluon scattering amplitudes.