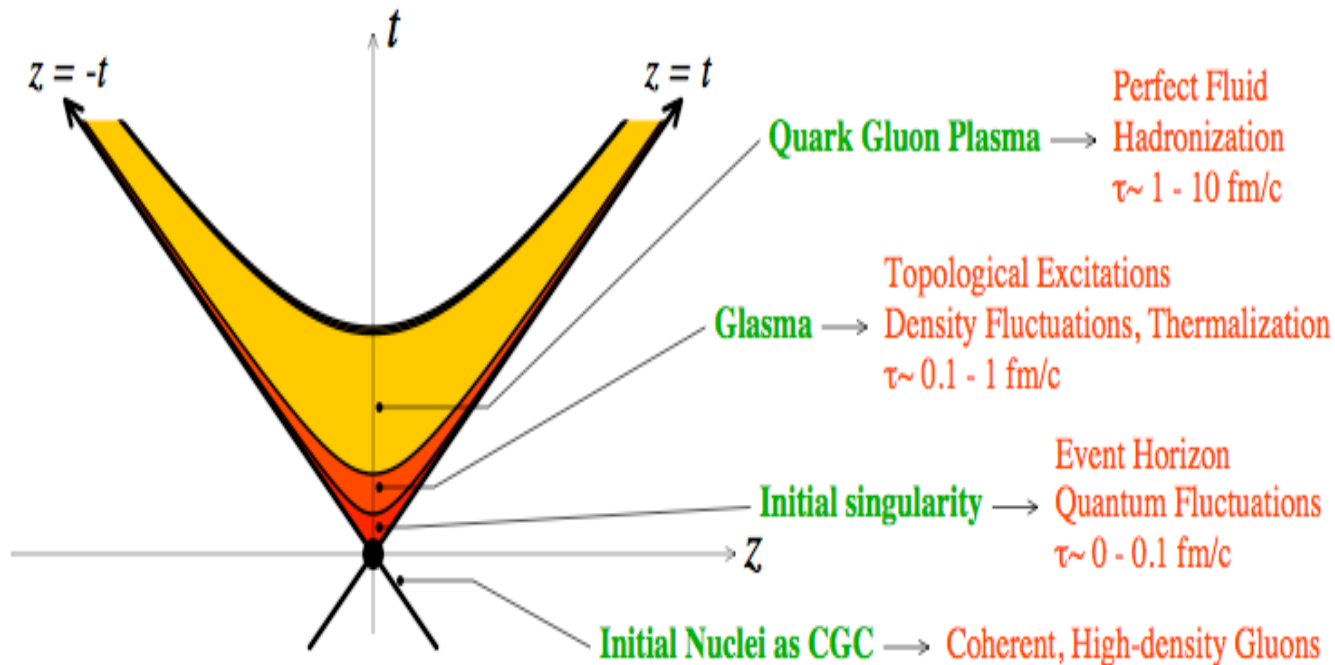
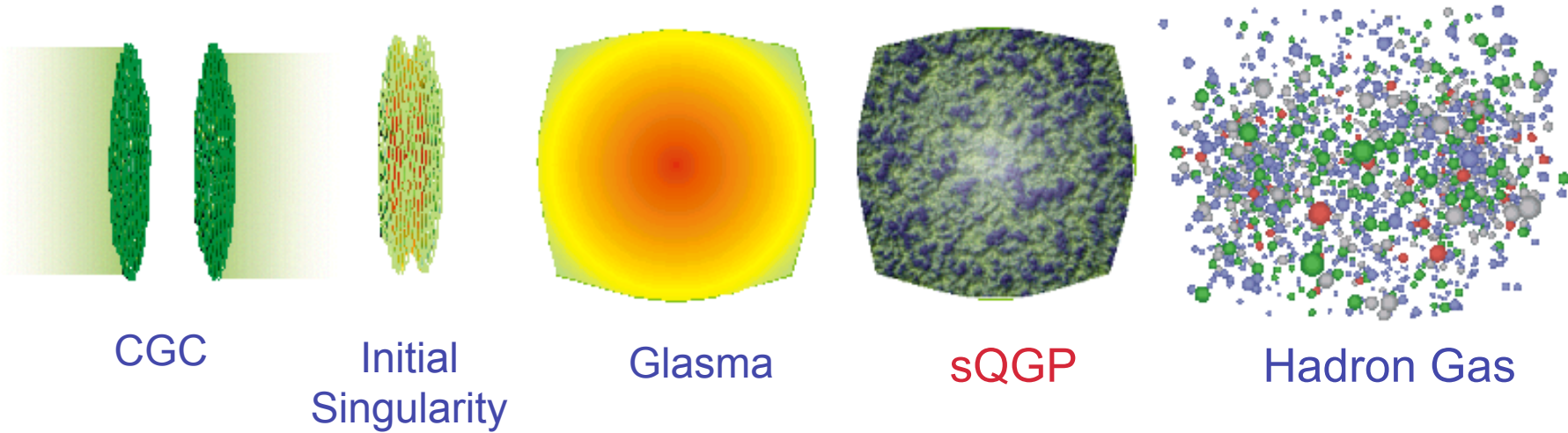


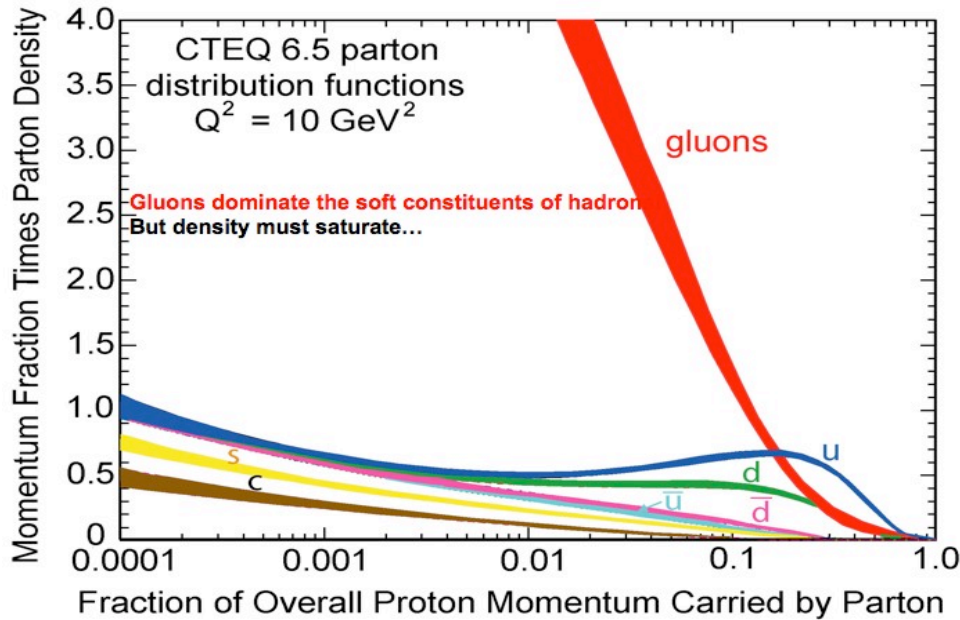
The Color Glass Condensate and Glasma

1

Art due to S. Bass



The Hadron Wavefunction at High Energy



Small x limit is high energy limit

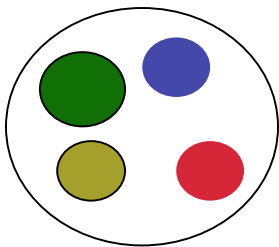
$$x = E_{\text{gluon}}/E_{\text{hadron}}$$

Where do all the gluons go?

Cross sections for hadrons rise very slowly with energy

$$\sigma_{\text{tot}} \sim \ln^2(E/\Lambda_{\text{QCD}})$$

$$\Lambda_{\text{QCD}} \sim 200 \text{ MeV}$$



Baryon:

3 quarks

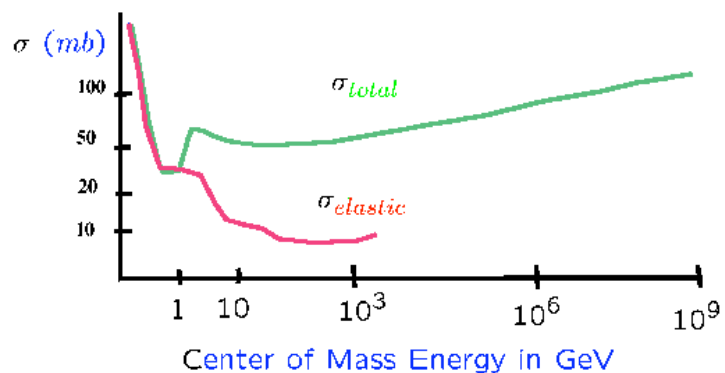
3 quarks 1 gluon

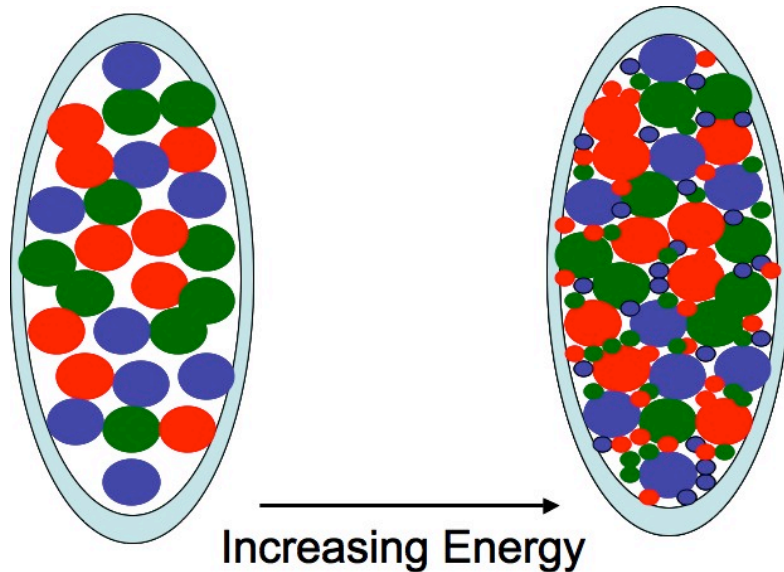
.....

3 quarks and lots of gluons



The total hadronic cross section:





But the gluon density rises much more rapidly!

The high energy limit is the high gluon density limit.

Surely the density must saturate for fixed sizes of gluons at high energy.

Color Glass Condensate

Color: Gluons

Condensate:

Glass:

The partons which make the CGC fields are moving fast \Rightarrow Lorentz time dilation \Rightarrow fields evolve slowly compared to natural times scales

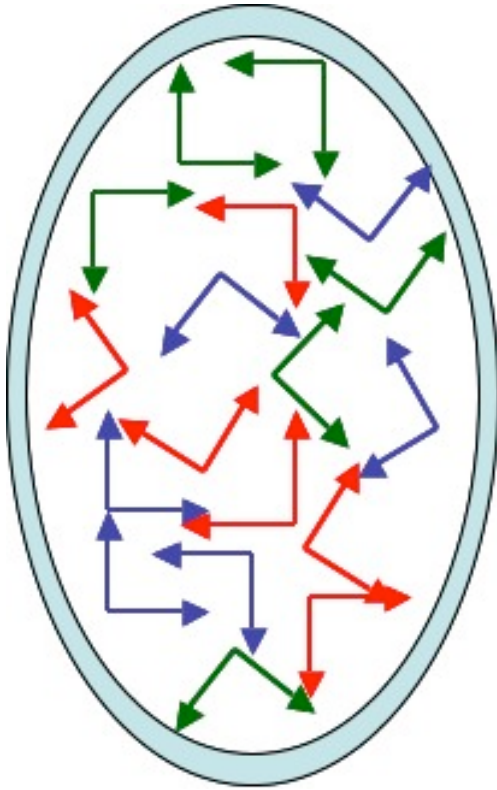
$$\frac{dN}{dyd^2p_T d^2x_T} = \rho \quad \text{Phase space density}$$

$$E = -\kappa\rho + \kappa'\alpha_S\rho^2 \Rightarrow \rho \sim 1/\alpha_S$$

Coupling weak because density is high

Fields are coherent and classical

What does a sheet of Colored Glass look like?



On the sheet $x^- = t - z$ is small

Independent of $x^+ = t + z$

F^{i+} is big F^{i-} is small

F^{ij} is 0(1)

so that the big fields are

$$\vec{E} \perp \vec{B} \perp \vec{z}$$

Lienard-Wiechart potentials

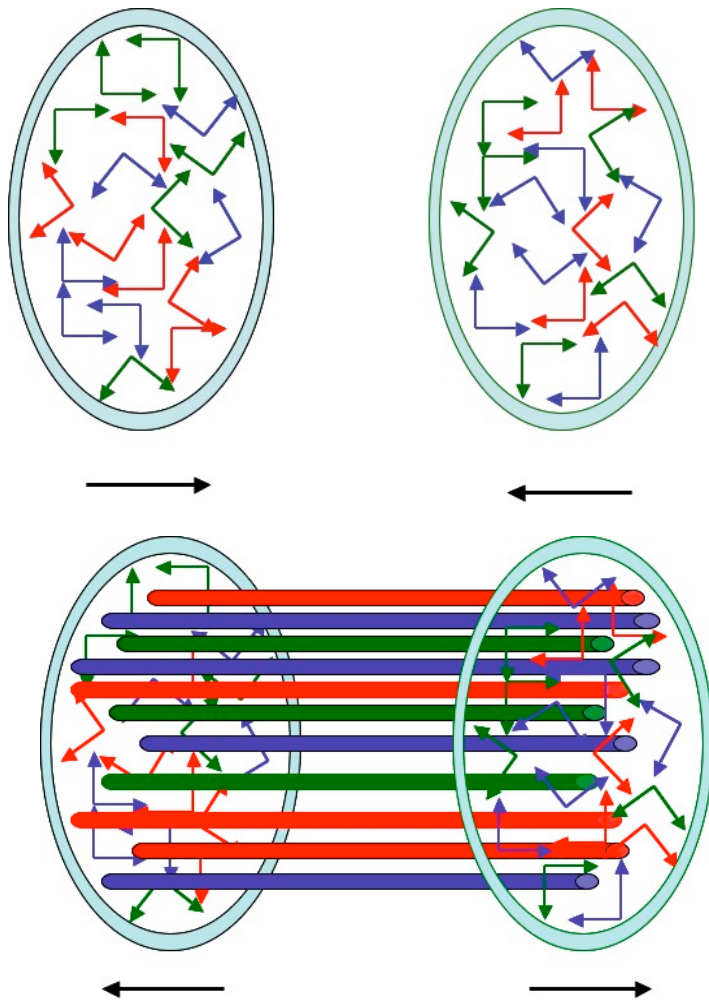
Random Color

Density of gluons per unit area

$$\frac{1}{\pi R^2} \frac{dN}{dy} \sim \frac{1}{\alpha_{strong}} Q_{sat}^2$$

CGC Gives Initial Conditions for QGP in Heavy Ion Collisions

5

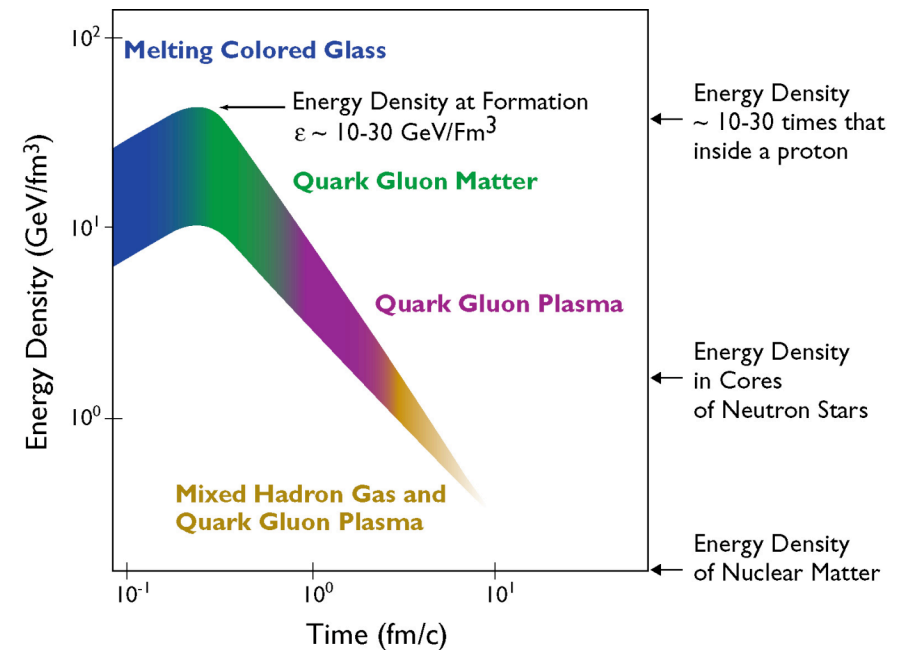


“Instantaneously” develop longitudinal color E and B fields

Two sheets of colored glass collide

Glass melts into gluons and thermalize

QGP is made which expands into a mixed phase of QGP and hadrons



Fields in longitudinal space:

$$F^{i+}$$

is a delta function on scales less than the
inverse longitudinal cutoff

A diagram showing two intersecting lines representing the longitudinal space axes. The line with a negative slope is labeled $x^- = 0$ at its upper end. The line with a positive slope is labeled $x^+ = 0$ at its upper end. Along the $x^- = 0$ line, the gauge field is given by $A^j = \frac{1}{i} U_1 \nabla^j U_1^\dagger$. Along the $x^+ = 0$ line, the gauge field is given by $A^j = \frac{1}{i} U_2 \nabla^j U_2^\dagger$.

Gluon distribution is at
scales larger than the
cutoff

$$G(k) \sim 1/p^+$$

$$G(k) = \langle a^\dagger(k) a(k) \rangle \sim \langle A(k) A(-k) \rangle$$

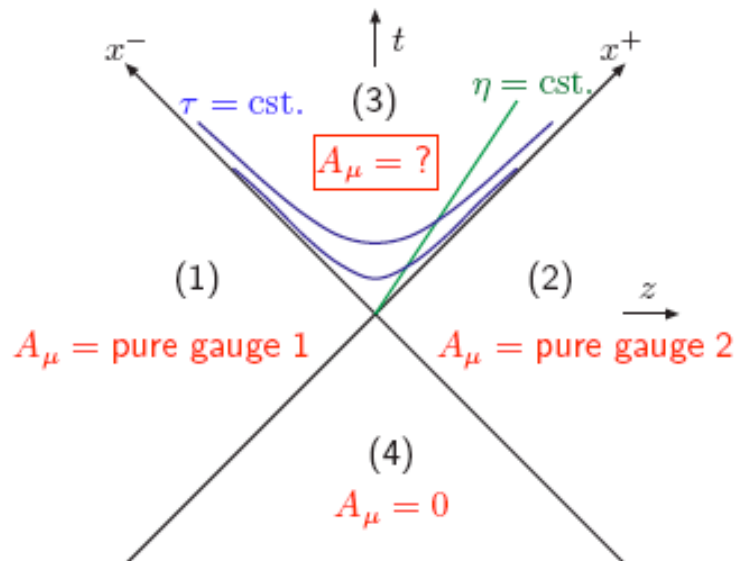
The Glasma:

7

Before the collision only transverse E and B
CGC fields

Color electric and magnetic monopoles

Almost instantaneous phase change
to longitudinal E and B



In forward light cone

$$A_1^i + A_2^i$$

generates correct sources on
the light cone

$$\nabla \cdot E = A \cdot E$$

$$\nabla \cdot B = A \cdot B$$

$$A_1 \cdot E_2$$

$$A_1 \cdot B_2$$

Equal strength for magnetic and
electric charge on average!

Glasma and Long Range Correlations

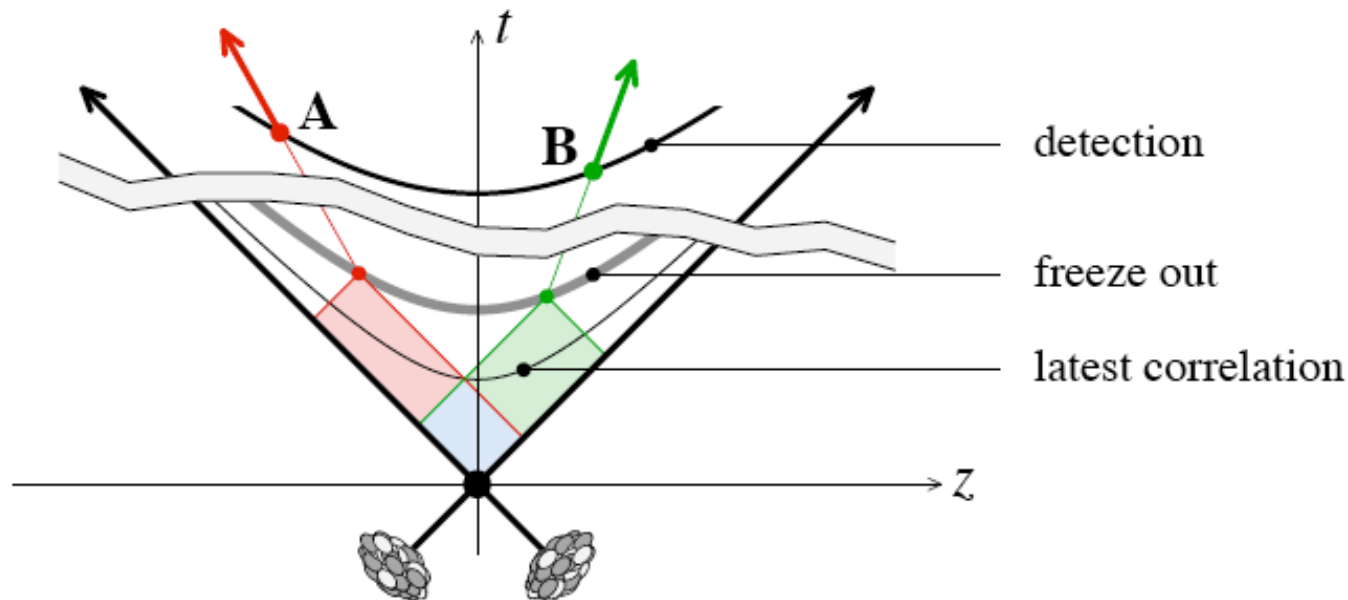
Glasma flux tubes formed at very early times

Long range correlation in rapidity

$$\tau_i = \tau_f e^{-\Delta y/2}$$

Impossible to set up long range correlations by final state interactions!

Long range correlation cannot be destroyed by final state interactions



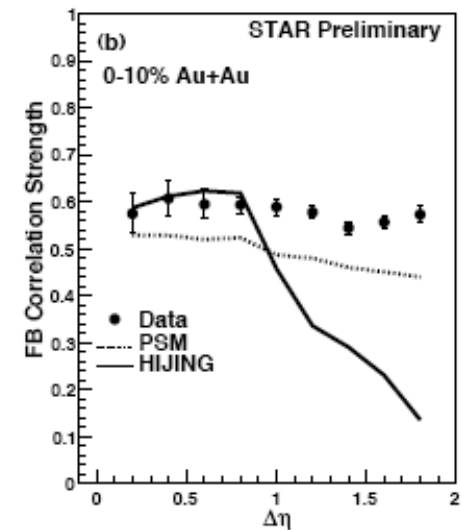
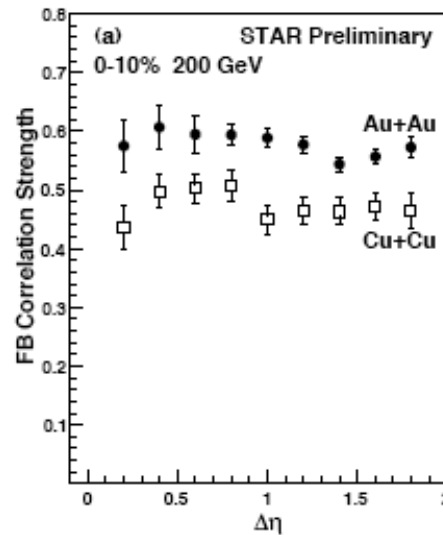
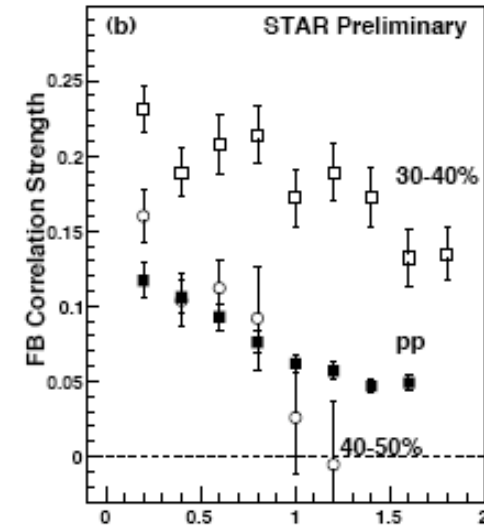
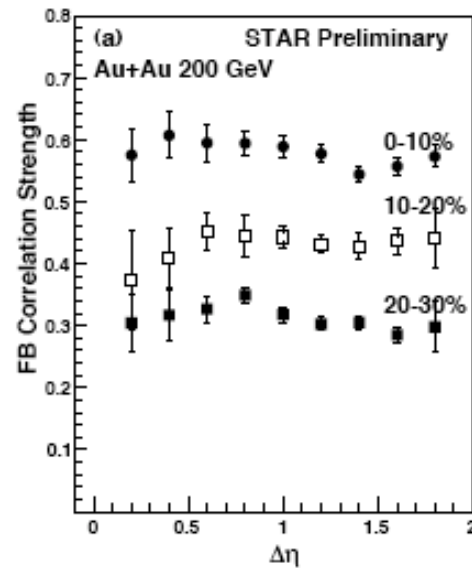
Forward-Backward Correlations:
$$\frac{\langle N_F N_B \rangle - \langle N_F \rangle \langle N_B \rangle}{\langle N_F^2 \rangle - \langle N_F \rangle^2}$$

Long range correlation present in proton-proton and proton-antiproton data.

Impact parameter effect

Correlation increases as function of centrality and of ion size

Hijing cannot explain the long range correlation



Intuitive Glasma Explanation:

Classical decay of a flux tube.

$$\frac{dN_{FT}}{dy} \sim R^2 / R_{FT}^2 \sim Q_{sat}^2 R^2$$

Multiplicity of each flux tube

$$1/\alpha_S$$

Classical contribution is
leading order and is strongly
correlated:

$$\langle N_F N_B \rangle \sim \langle N_F \rangle^2 \sim 1/\alpha_S^2$$

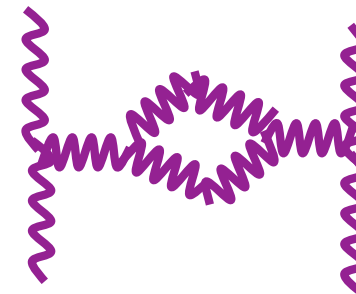
Short range correlation

$$O(1)$$

(Interference?)

\Rightarrow

$$\frac{dN_{gluons}}{dy} \sim \frac{1}{\alpha_S} Q_{sat}^2 R^2$$



LDM and Kovchegov,
Armesto, Pajares and LDM
Venugopalan and Gelis,
Fukushima and Hidaka

Forward-Backward Correlation

$$\frac{\langle N_F N_B \rangle - \langle N_F \rangle \langle N_B \rangle}{\langle N_F^2 \rangle - \langle N_F \rangle^2} \sim \frac{A/\alpha_S^2}{A/\alpha_S^2 + B} \sim \frac{1}{1 + \kappa \alpha_S^2}$$

Correlation strongest when centrality largest, since coupling is weakest for larger saturation momentum

Gluons are enhanced by inverse coupling squared because of Bose coherence in the Color Glass Condensate

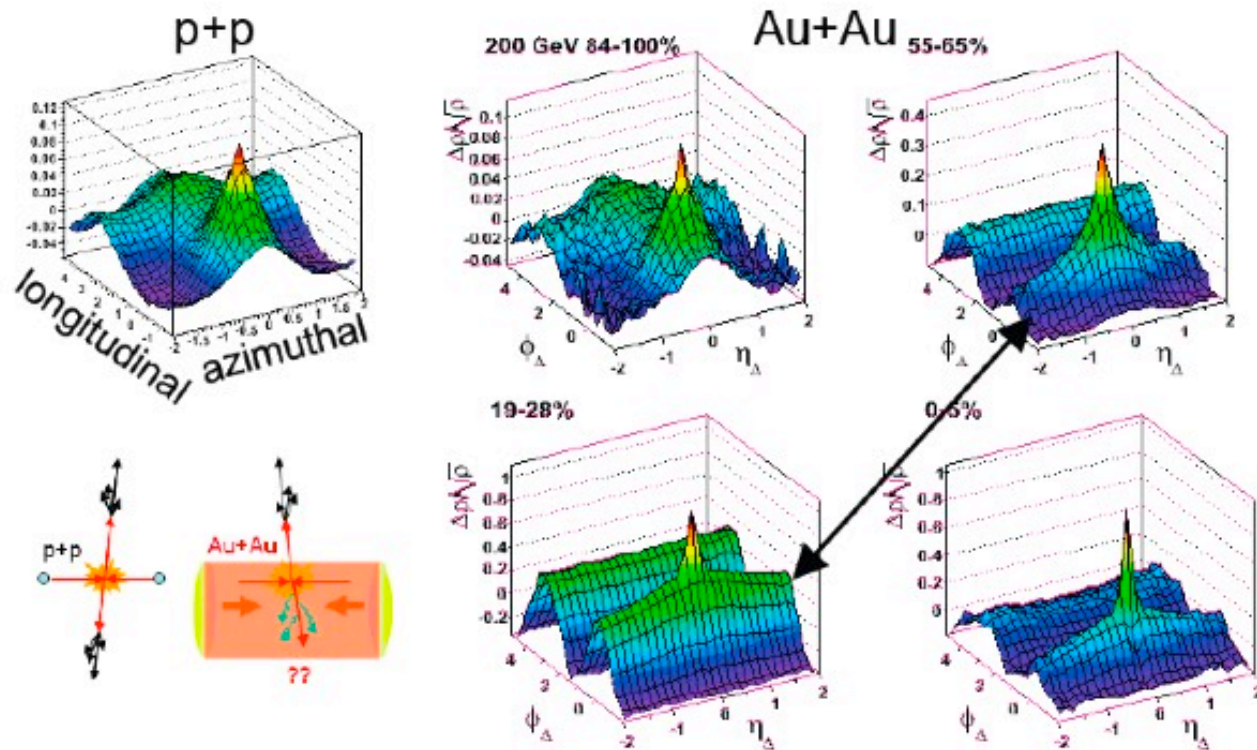
Baryons are Fermions and there is no enhancement. For central collisions there is little change relative to less central collisions

For baryons we therefore expect

$$\left\{ \frac{\langle N_F N_B \rangle - \langle N_F \rangle \langle N_B \rangle}{\langle N_F^2 \rangle - \langle N_F \rangle^2} \right\}_{baryons} \sim \frac{1}{1 + \kappa'}$$

The Ridge in Heavy Ion Collisions:

Correlations of all unique pairs of charged particles



M. Daugherty for STAR: QM2008

Is this evidence for Glasma flux tubes?

The Ridge

High p_T Ridge and Inclusive Ridge: Will analyze inclusive ridge

First compute the two particle spatial correlation:

$$r(x_1, x_2) = n_2(x_1, x_2) - n_1(x_1)n_1(x_2)$$

Convolute with hydrodynamic expansion via blast wave with parameters from Phenix

Scale resulting multiplicity distribution by total multiplicity to define $\frac{\Delta\rho}{\sqrt{\rho}}$

$$r \sim \frac{1}{\alpha_S^2} \frac{dN_{fluxtubes}}{dy} \quad \frac{dN}{dy} \sim \frac{1}{\alpha_S} \frac{dN_{fluxtubes}}{dy}$$

$$\frac{\Delta\rho}{\sqrt{\rho}} \sim \frac{1}{\alpha_S}$$

Ridge correlation is strong!

LM and Kovchegov,

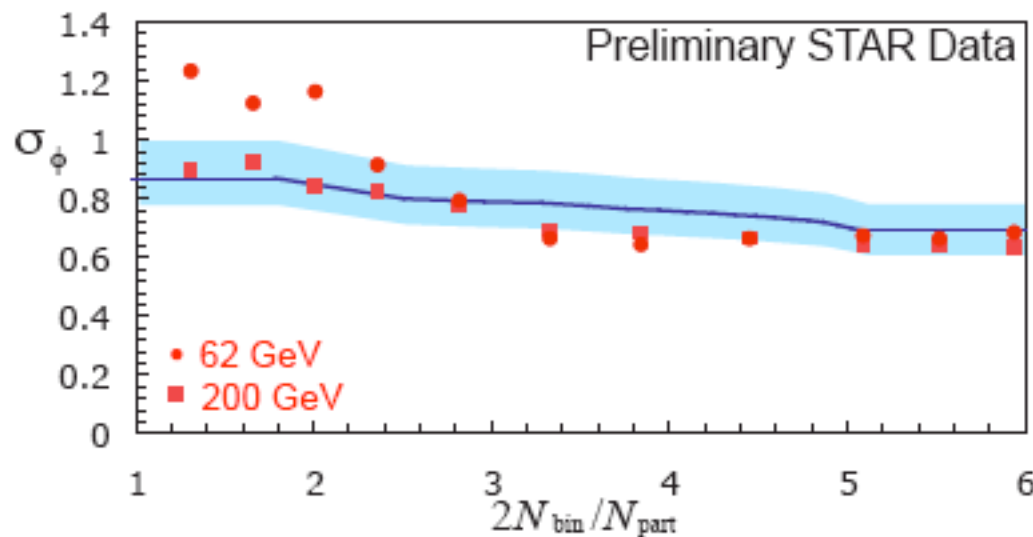
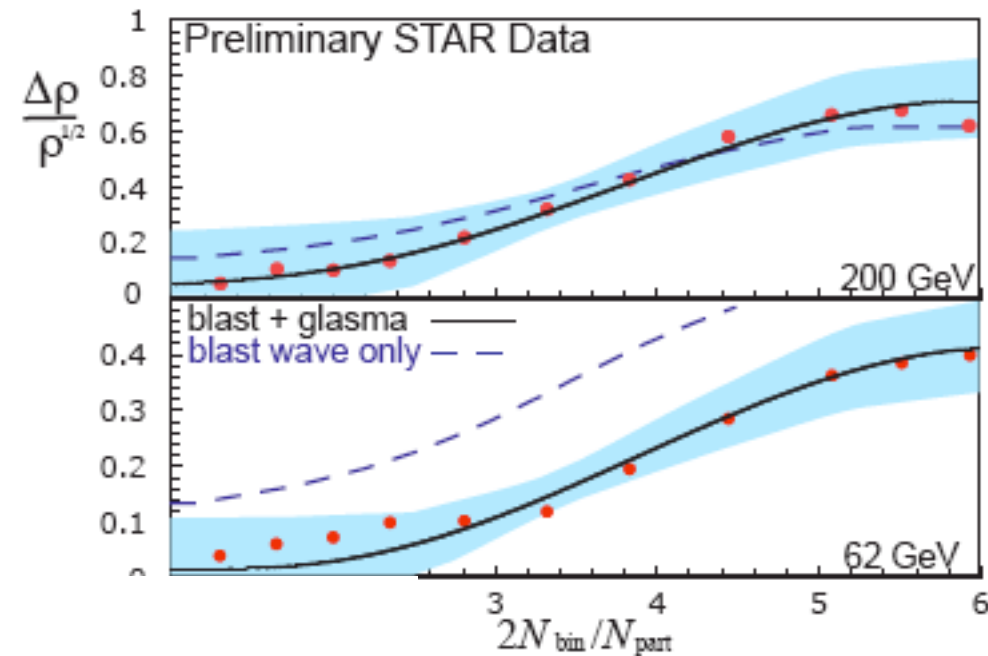
Dumitru, Gelis, LM and
Venugopalan

Gavin, LM and Moschelli

Assuming delta function in transverse size, flat in rapidity:

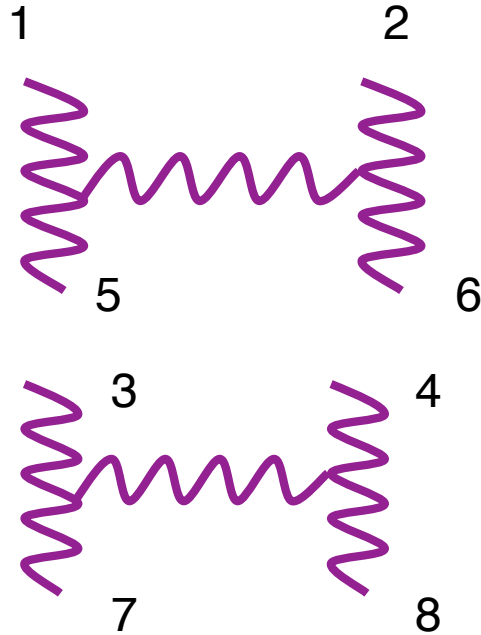
Naturally get right
magnitude.

Normalization fixed at
one energy and
centrality



Two particle correlation in
angle from flow effects

How is correlation generated ?



Naively disconnected,

But CGC averages over sources at
positions 1-8

Sources 1-4 can be contracted in pairs
as can 5-8

The contraction 1-2, 3-4, 5-6, 7-8
gives a disconnected diagram which
cancels the mixed event subtraction

There are however other contractions which do not cancel.

These correspond to interference diagrams associated with the
classical field, and are intrinsic to the classical field description

Questions:

How to describe the backwards ridge (associated with momentum conservation)?

Rapidity dependence of the ridge?

Experimentally is it truly long range?

What is the value of the un-determined numerical coefficient?

Can one compute the short range piece of the correlation to get a full description?

How to describe the high p_T ridge?

What can be done in LHC?

How is this related to forward-backward correlations seen in pp and $p\bar{p}$?

Could the effect be due to biasing arising from flow? (Akiba)