Calculating amplitudes in the multi-Regge regime of strongly coupled $\mathcal{N}=4$ SYM

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DESY Theory Workshop 2013

based on 1207.4204 and 1310.xxxx with J. Bartels and V. Schomerus



Particles, Strings, and the Early Universe Collaborative Research Center SFB 676



- planar $\mathcal{N} = 4$ integrable \leftrightarrow can compute observables for any coupling
- scattering amplitudes particularly interesting → techniques for less symmetric theories
- enormous progress on weak coupling side
- how do amplitudes behave at strong coupling?
 - interpolation to intermediate coupling
 - regime inaccessible (analytically) for QCD

Scattering Amplitudes via AdS/CFT

 $[{\sf Alday}/{\sf Maldacena}, {\sf A}/{\sf M}/{\sf Gaiotto}, {\sf A}/{\sf M}/{\sf Sever}/{\sf Vieira}]$



•
$$A \sim e^{-\frac{\sqrt{\lambda}}{2\pi} \operatorname{Area}}$$

• $k_i = x_{i-1} - x_i$
 \rightarrow polygon depends only on u

• $Y_{a,s}(\theta)$ generalised cross ratios \rightarrow for n gluons: 3n - 15 Y-functions

[Figure from 1002.2459]

$$\begin{split} \log Y_{a,s}\left(\theta\right) &= - \ m_s \ \cosh\theta \pm \ C_s \\ &+ \sum_{a',s'} \int d\theta' \ \mathcal{K}^{a,a'}_{s,s'} \ \left(\theta - \theta' + \ i\varphi_s - i\varphi_{s'} \ \right) \log\left(1 + Y_{a',s'}\left(\theta'\right)\right) \end{split}$$

 $=\frac{x_{ij}^2 x_{kl}^2}{x_{ik}^2 x_{il}^2}$

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$$\begin{split} \log Y_{a,s}\left(\theta\right) &= - \frac{m_{s}}{m_{s}} \cosh \theta \pm \frac{C_{s}}{C_{s}} \\ &+ \sum_{a',s'} \int d\theta' \frac{\mathcal{K}_{s,s'}^{a,a'}}{\mathcal{K}_{s,s'}^{a,a'}} \left(\theta - \theta' + \frac{i\varphi_{s} - i\varphi_{s'}}{i\varphi_{s}}\right) \log \left(1 + Y_{a',s'}\left(\theta'\right)\right) \end{split}$$

$$\begin{aligned} \mathsf{Area} &= \mathsf{A}_{\mathsf{div}}\left(x_{i}\right) + \mathsf{A}_{\mathsf{periods}}\left(m_{s},\varphi_{s}\right) + \Delta(u_{i}) \\ &+ \sum_{s} \int \frac{d\theta}{2\pi} \left|m_{s}\right| \cosh \theta \mathsf{log}\left[\left(1 + Y_{1,s}\right)\left(1 + Y_{3,s}\right)\left(1 + Y_{2,s}\right)^{\sqrt{2}}\right](\theta) \end{aligned}$$

- non-divergent piece: remainder function $e^{-\frac{\sqrt{\lambda}}{2\pi}R}$
- Y-system easy to solve numerically
- BUT: not solvable analytically for arbitrary kinematics!

Multi-Regge Limit



• for $2 \rightarrow n-2$ scattering: 3n-10 Mandelstam invariants

• Multi-Regge limit: rapidities of produced particles strongly ordered

ightarrow Regge limit in all subchannels $s \gg s_i \gg -t_i$

• $\mathcal{N} = 4$ dual conformal \rightarrow choose 3n - 15 cross ratios u_{as}

• kinematical analysis:

 $u_{1s} \to 1, u_{2s}, u_{3s} \to 0 \text{ with } \tilde{u}_{2s} = \frac{u_{2s}}{1 - u_{1s}} = \mathcal{O}(1), \quad \tilde{u}_{3s} = \frac{u_{3s}}{1 - u_{1s}} = \mathcal{O}(1)$

[1207.4204]

•
$$u_{as} = \left. \frac{Y_{2s}}{1+Y_{2s}} \right|_{\theta=i(k\pi/4-\varphi_s)}$$

- demand that cross ratios show behaviour predicted by MRL
- MRL corresponds to choice m_s large, $\varphi_s \to -(s-1)\frac{\pi}{4}$, $C_s \to \text{const.}$
- in this limit, integrals in Y-system can be neglected

$$\log Y_{a,s}\left(\theta\right)\cong -m_{s}\cosh\theta\pm C_{s}+\mathcal{O}(e^{-m})$$

• in this limit, R trivial

Multi-Regge regions

- 2^{n-4} regions, corresponding to the signs of $k_i^0 = E_i$
- different regions connected by analytic continuation in $s_i
 ightarrow s_i e^{ilpha}$



• for the above example: $s_{34} \rightarrow e^{i\pi}s_{34}$ $s_{56} \rightarrow e^{i\pi}s_{56}$, $s_{345} \rightarrow e^{i\pi}s_{345}$, $s_{456} \rightarrow e^{i\pi}s_{456}$ $\Rightarrow u_1 \rightarrow e^{-2\pi i}u_1$, $u_2 \rightarrow u_2$, $u_3 \rightarrow u_3$

Multi-Regge regions in the Y-system

- continuation in $u_{as} \sim$ continuation in m, C, φ
- (numerical) inversion of $u_{as} = \frac{Y_{2s}}{1+Y_{2s}}\Big|_{\theta=i(k\pi/4-\varphi_s)}$ to find paths for parameters
- solutions of $Y_{as}(\theta) = -1$ can cross real axis

$$\begin{split} \log Y_{a,s}\left(\theta\right) &= - \ m_{s} \cosh \theta \pm C_{s} \\ &+ \sum_{a',s'} \int d\theta' \mathcal{K}_{s,s'}^{a,a'} \left(\theta - \theta' + i\varphi_{s} - i\varphi_{s'}\right) \log \left(1 + Y_{a',s'}\left(\theta'\right)\right) \end{split}$$

• crossing leads to contributions to R

$$R' = \int \frac{d\theta}{2\pi} |m_s| \cosh \theta \log \left[(1 + Y_{a,s} \left(\theta \right)) \right] \pm i |m_s| \sinh \left(\theta_0 \right) + \dots$$

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$$\log Y_{a,s}(\theta) = -m_s \cosh \theta \pm C_s + \sum \log \mathcal{S}_{s,s'}^{a,a'}(\theta - \theta_0 + i\varphi_s - i\varphi_{s'}) \\ + \sum_{a',s'} \int d\theta' \mathcal{K}_{s,s'}^{a,a'}(\theta - \theta' + i\varphi_s - i\varphi_{s'}) \log (1 + Y_{a',s'}(\theta'))$$

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Example: 6-point case

[Bartels/Kotanski/Schomerus]



$$u_1 \rightarrow e^{-2\pi i} u_1, \ u_2 \rightarrow u_2, \ u_3 \rightarrow u_3$$

Solutions of $Y_3(\theta) = -1$ along continuation

• Remainder function has Regge behaviour:

$$e^{-rac{\sqrt{\lambda}}{2\pi}R_6}\sim \left((1-u_1)\sqrt{\widetilde{u}_2\widetilde{u}_3}
ight)^{rac{\sqrt{\lambda}}{2\pi}e_2}$$

•
$$e_2 = -\sqrt{2} + \frac{1}{2}\log(3 + 2\sqrt{2}) < 0$$

[1310.xxxx]



• strong coupling results consistent with weak coupling LLA identities!

- studied scattering amplitudes in strongly coupled $\mathcal{N}=4$ SYM
- showed that Y-system equations simplify in MRL
- calculated 6-point and 7-point remainder function for specific continuations →confirmed weak coupling LLA prediction

next steps:

- remaining 7-point continuations
- weak coupling: new contribution for 8-point case
- make contact with weak coupling results

