

An asymptotic safety scenario for gauged chiral Higgs-Yukawa models

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DESY September 26, 2013: Theory workshop

Motivations

Triviality of the Higgs sector:

Landau pole for the Higgs self-interaction

$$\beta_{\lambda}^{\text{1loop}} = a\lambda^2, \quad \lambda_{\Lambda} = \frac{\lambda_R}{1 - a\lambda_R \log \frac{\Lambda}{m_R}} \quad (a > 0)$$

$$\Lambda \rightarrow \infty \iff \lambda_R = 0 \text{ trivial QFT}$$

$$\Lambda < \infty \iff \lambda_R \neq 0 \text{ effective QFT}$$

confirmed by lattice for pure scalar theories

Hierarchy between Electroweak and Planck/GUT scales:

Quadratic running of the Higgs mass (vev^2 , $v = \langle \phi \rangle$)

$$\beta_{\kappa} \equiv \Lambda \partial_{\Lambda} \left(\frac{v^2}{2\Lambda^2} \right) = -2\kappa + \text{interactions}$$

@ Gaussian FP ($\kappa = \lambda = 0$) linearized RG flow with eigenvalues $\{-2, 0\}$

light Higgs masses \iff fine tuning in the UV

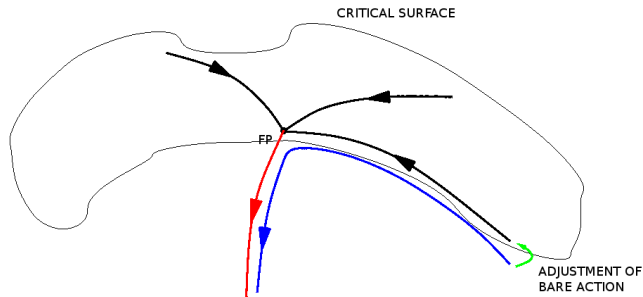
in SM: what happens **beyond perturbative expansions?**

Asymptotic safety

Building **fundamental** QFT's by means of RG FP G_{i*}

$$\Lambda \partial_\Lambda G_i(\Lambda)|_{G_*} = 0 \quad G_i(\Lambda) - G_{i*} = e^{-\theta_i \log(\Lambda/\Lambda_0)} + \dots$$

- r) $\theta_r > 0$ RELEVANT = attractive
(local coordinates of the UV critical surface)
- i) $\theta_i < 0$ IRRELEVANT = repulsive
(IR-dumped \rightarrow IR universality)
- m) $\theta_m = 0$ MARGINAL = nonlinear flow is needed



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Asymptotic freedom: GFP $G_{i*} = 0$ with interacting relevant or marginally relevant directions

Asymptotic safety (Weinberg '76): NGFP $G_{i*} \neq 0$ with a finite number of relevant or marginally relevant directions

- 1) no triviality
- 2) hierarchy? if small critical exponents \rightarrow slow UV running
- 3) more predictive power than perturbation theory?
- 4) limited accessible window for IR parameters

The model

Yukawa systems as toy models of the SM and of GUT's

Triggering AS in SSB regime

$$\beta_\kappa = -2\kappa + \text{bosons' interactions} - \text{fermions' interactions}$$

quasi-conformal behaviour of the vev \iff boson dominance

Chiral Yukawa: $N_L \times \psi_L, \quad 1 \times \psi_R, \quad 2N_L \times \phi$
global $U(1)_L \times U(1)_R, \quad SU(N_L)_L$

massless ϕ 's do not decouple

Higgs: gauging $SU(N_L)_L$ to kill the Goldstones

$$S = \int d^d x \left[\frac{1}{4} F_{\mu\nu}^i F^{i\mu\nu} + (D^\mu \phi)^\dagger (D_\mu \phi) + U(\phi^\dagger \phi) \right. \\ \left. + i(\bar{\psi}_L^a \not{D}^{ab} \psi_L^b + \bar{\psi}_R \not{\partial} \psi_R) + \bar{h} \bar{\psi}_R \phi^{\dagger a} \psi_L^a - \bar{h} \bar{\psi}_L^a \phi^a \psi_R \right]$$

No $U(1)_Y$ gauge sector here!

fRG method

1PI functional RG (Wetterich '91)

1) Regularization: IR mass-like regulator $R_k(p^2)$

2) Average effective action Γ_k

$$\begin{array}{ccc} k = 0 & \longleftarrow \longleftarrow \longleftarrow \longleftarrow \longleftarrow & k = \Lambda \\ \text{1PI effective action} & & \text{bare action} \end{array}$$

3) Flow equation $\partial_t \equiv k \partial_k$

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left[\left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_t R_k \right]$$

4) Truncation

$$\Gamma_k[\Phi] = \sum_i \bar{G}_i(k) \mathcal{O}_i[\Phi] \longrightarrow \partial_t \bar{G}_i(k) = \bar{\beta}_i(\bar{G}_j, k)$$

e.g. derivative expansion $\mathcal{L}_k = U_k(\Phi) + Z_\Phi \partial_\mu \Phi \partial^\mu \Phi + \dots$

5) Rescalings

$$G_i(t) \equiv \frac{k^{-D_i}}{Z_\phi^{n_\phi/2} \dots Z_W^{n_W/2}} \bar{G}_i(k) \longrightarrow \partial_t G_i = \beta_i(G_j)$$

Threshold effects & β 's

SSB: dynamically generated masses proportional to running couplings

$$\phi^4 \text{ 1-loop} \quad \mu_W^2 = \frac{g^2 \kappa}{2} \quad \mu_t^2 = h^2 \kappa \quad \mu_H^2 = 2\lambda \kappa \quad (N_L = 2)$$

$$\partial_t \lambda = \frac{1}{16\pi^2} \left\{ 12\lambda^2 + 4\lambda h^2 - 4h^4 - 9\lambda g^2 + \frac{9}{4}g^4 \right\}$$

ϕ^4 fRG @NLO-derivative expansion

$$\begin{aligned} \partial_t \lambda = \frac{1}{16\pi^2} \left\{ 12\lambda^2 \left(\frac{1}{4} + \frac{3}{4(1+\mu_H^2)^3} \right) + \frac{4\lambda h^2}{(1+\mu_t^2)^4} - \frac{4h^4}{(1+\mu_t^2)^3} \right. \\ \left. - 9\lambda g^2 \left(\frac{1}{2(1+\mu_W^2)} + \frac{1}{2(1+\mu_W^2)^2} \right) + \frac{9}{4} \frac{g^4}{(1+\mu_W^2)^3} + \dots \right\} \end{aligned}$$

Nonperturbative β 's for: λ κ h^2 g^2 and higher ϕ^{2n}
that is, for: μ_H^2 μ_W^2 μ_t^2 g^2 and higher ϕ^{2n}

Non-Gaussian fixed points

For $g^2 = 0$ also $\partial_t g^2 = 0$

Strategy: solve FP eqs. for matter sector at $g^2 \neq 0$ (and $\partial_t g^2 \neq 0$)

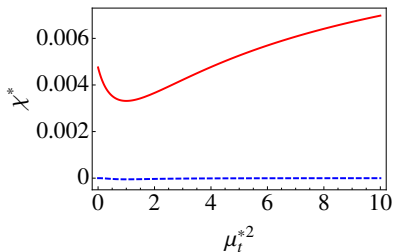
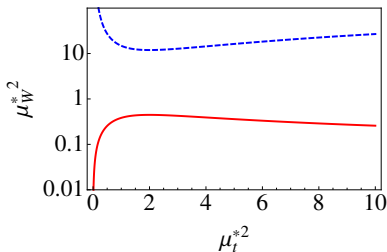
At least one solution exists

when $g^2 \rightarrow 0$, $\kappa_* \sim 1/g^2$ $h_*^2 \sim g^2$ $\lambda_* \sim g^4$

$\mu_{H_*}^2 \sim g^2$ $0 < \{\mu_{t_*}^2, \mu_{W_*}^2, \chi \equiv \frac{\mu_{H_*}^2}{g^2}\} < \infty$

$\partial_t \mu_t^2 = \frac{\mu_t^2}{\mu_W^2} (\partial_t \mu_W^2)$ **line of fixed points**

higher $\phi^{2n} \sim g^{2n}$ **flat potential**



UV FP-regime

Linearized RG flow close to FP's

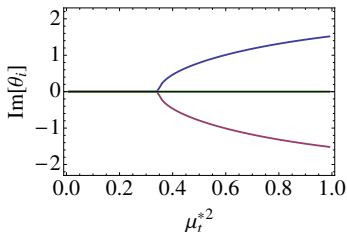
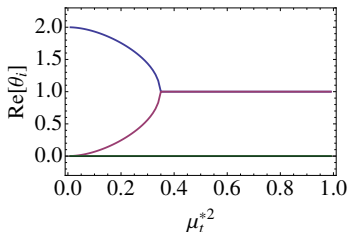
$$\partial_t G_i = B_i^j (G_j - G_j^*) + \dots, \quad B_i^j = \left. \frac{\partial \beta_i}{\partial G_j} \right|_*$$

Critical exponents $\theta_i =$ eigenvalues of $(-B_i^j)$

For $G_i \equiv \{g^2, \mu_W^2, \mu_t^2, \chi\}$

$\theta_3 = \theta_4 = 0$ (marginally relevant+exactly marginal)

$\theta_1 = \theta_2^* \neq 0$ (relevant)



partial smoothing of hierarchy
no improved predictive power

higher ϕ^{2n} have canonical exponents (irrelevant)

No triviality

@FP's all couplings vanish, the vev $\kappa \rightarrow \infty \iff$ finite μ 's

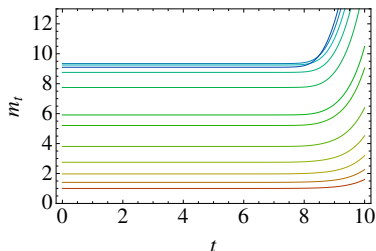
Interacting theory @NGFP's

1) @GFP massless gauge bosons and chiral fermions

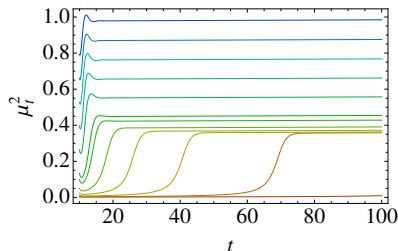
@NGFP dynamically generated masses

2) critical exponents different from GFP

Asymptotically safe trajectories connected to the Higgs phase



IR dimensionful freeze-out



UV dimensionless FP-regime

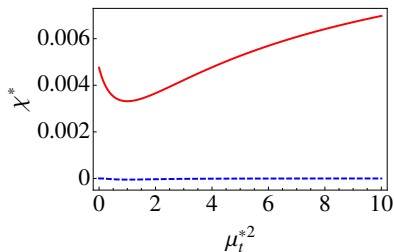
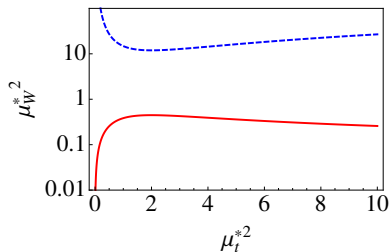
Flows to the electroweak scale

We neglect higher ϕ^{2n} & start the flow close to NGFP

@FP

$$\mu_{W*}^2 \sim \mu_{t*}^2$$

$$\chi_* \sim 10^{-2} \mu_{t*}^2$$



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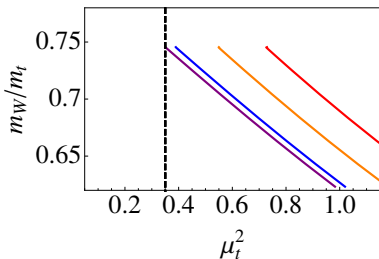
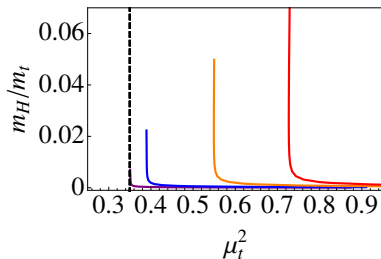
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For generic flows this is preserved in the IR.

For initial conditions close to a specific FP we can get $m_H/m_t \sim 0.7$



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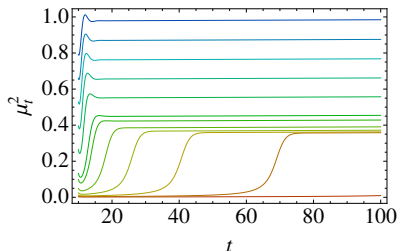
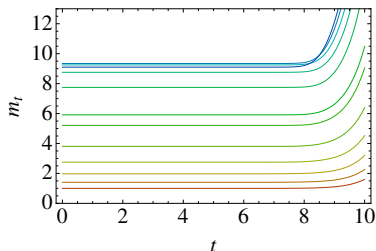
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These tuned trajectories exhibit a 'walking' regime



Conclusions

- 1) We studied the **nonperturbative RG** flow of **gauged chiral Higgs-Yukawa models**, interesting as toy models of the Higgs-top sector of the SM or of GUT's, by means of the 1PI fRG
- 2) We found **a line of NGFP's**, characterized by a flat scalar potential and finite dynamically generated masses, which seems solid against the relaxation of the required nonperturbative approximations
- 3) These FP's allow for the construction of **asymptotically safe theories** corresponding to RG trajectories interpolating between a standard Higgs phase and a UV FP-regime
- 4) The corresponding theories would show **no triviality problem**, would have a milder hierarchy problem, and would contain as many free parameters as in the conventional perturbative definition of the model
- 5) We have performed non-exhaustive scans of the **mass spectrum** emerging from the flow to the **Higgs** phase along these renormalizable trajectories