Orbifolds and topological defects

Daniel Plencner

LMU Munich

September 26th, 2013

based on arXiv:1307.3141 with I. Brunner and N. Carqueville

Daniel Plencner (LMU Munich)

Motivation

- orbifolds with symmetry group G can be described via defects A_G
- generalization: allow any defect A with appropriate algebraic structure
- \bullet standard results on orbifolds recovered in simpler/more conceptual way \rightsquigarrow carries over to the generalized setting
- case study: $\mathcal{N} = 2$ Landau-Ginzburg models:
 - description of many $\mathcal{N}=2$ CFTs, stringy regime of CY compactifications
 - explicit description of defects via matrix factorizations
 - ► compute arbitrary topological correlators in the generalized orbifold theory: bulk/boundary correlators (~→ eff. superpotentials, D-brane charges), defect actions on bulk fields, ...

Defects in 2D field theories

- A **defect** X is a 1D interface between two theories \mathcal{T}_1 and \mathcal{T}_2 together with a gluing condition on fields
- defects can form **junctions** \rightsquigarrow junction fields: $Y_1 \searrow \cdots \nearrow Y_r$

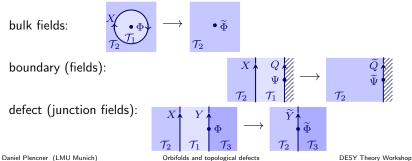
- usually want to preserve some symmetry (e.g. conformal, SUSY, ...)
- X is a **topological defect** if all correlators inv. under deformations of X. use this to map objects of \mathcal{T}_1 to objects of \mathcal{T}_2 (and vice versa):

X

 \mathcal{T}_1

3 / 17

 \mathcal{T}_2



Defect description of orbifolds

• ordinary orbifolds: theory \mathcal{T} with finite symmetry group G \rightsquigarrow orbifold theory \mathcal{T}/G :

orbifold projection

 $P_{\text{orb}}|\phi_a\rangle \equiv \frac{1}{G}\sum_{h\in G} h|\phi_a\rangle \stackrel{!}{=} |\phi_a\rangle$

twisted fields

$$\phi_g(e^{2\pi i}z) = g\phi_g(z)$$

• defect perspective:

 $\phi^{(1)} = g\phi^{(2)}$ $I_{h} \overset{I_{g}}{\bigoplus} \overset{I_{g}}{\bigoplus} \overset{I_{g}}{=} \phi_{g} \overset{I_{g}}{\bigoplus} \forall h \in G$ $\phi^{(1)} = \phi^{(2)}$

• Note: I_q is always topological since $q \in G$ is a symmetry of \mathcal{T}

Daniel Plencner (LMU Munich)

Generalized orbifolds

Assemble I_g to $A_G = \bigoplus_{g \in G} I_g$ twisted fields – elements in Hom (I, A_G) orbifold projector ϕ A_G $P_{orb}(\phi) = A_G \phi$ I

Generalized orbifolds – allow any (symmetric) separable Frobenius algebra *A*: [Fröhlich, Fuchs, Runkel, Schweigert '09], [Carqueville, Runkel '12]

$$A \otimes A \to A, \quad \downarrow : I \to A, \quad \bigwedge = \bigwedge, \quad \bigwedge = \left| = \right|$$

 \rightsquigarrow associative OPE, unique vacuum, non-degenerate bulk pairings, $P_{\sf orb}, \ldots$

(symmetry):

 \rightsquigarrow spectral flow operator

Daniel Plencner (LMU Munich)

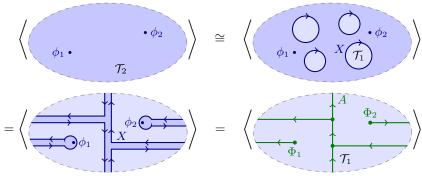
Ă |=| Ă

Generalized orbifolds

• important class $A = X^{\dagger} \otimes X$, where $X : \mathcal{T}_1 \to \mathcal{T}_2$ is a top. defect with

invertible dim(X), i.e. $X \xrightarrow{X} \bullet c1$, $c \in \mathbb{C} \setminus \{0\}$

• \mathcal{T}_1 and \mathcal{T}_2 are related via generalized orbifold:



- "twisted fields" in $\mathcal{T}_1 \leftrightarrow$ bulk fields in \mathcal{T}_2
- similarly with D-branes (defects) twisted branes in $\mathcal{T}_1 \leftrightarrow$ branes in $\mathcal{T}_2 \leftrightarrow$ equivalence of D-brane categories

Daniel Plencner (LMU Munich)

Landau-Ginzburg models

• $\mathcal{N}=2~$ LG models are 2D QFTs with action:

$$S = \int d^2 z \, d^4 \theta \, K(X_i, \bar{X}_i) + \frac{1}{2} \left(\int d^2 z \, d^2 \theta \, W(X_i) |_{\bar{\theta}^{\pm} = 0} + c.c. \right)$$

 $W(X_i)$ – superpotential

- LG not conformal, but for W quasi-homogeneous \rightsquigarrow flow to an IR fixed pt.
- **CFT in IR** characterized solely by W. Can extract information about the CFT from properties of W, e.g. chiral primary fields $\leftrightarrow \frac{\mathbb{C}[X_i]}{(\partial W)}$
- Many N = 2 CFTs described as IR fixed pts. of LG orbifolds, in particular, **CY compactifications** in stringy regime of Kähler moduli space.
- from here on: topologically B-twisted LG models

Defects in LG models

- Defects between LG models with superpotentials $W(x_1, ..., x_n)$, $V(z_1, ..., z_m)$ described by matrix factorizations of V - W. [Brunner, Roggenkamp '07] [Khovanov, Rozansky '04] [Kapustin, Li '02]
- A matrix factorization of a polynomial $p(u_i)$ is a pair (X, d_X) , where X is a \mathbb{Z}_2 -graded $\mathbb{C}[u_i]$ -module and d_X is an odd operator s.t.

$$d_X^2 = p(u_i) \cdot 1_X$$

• Junction fields Φ from X to Y given by maps in the **cohomology** Hom(X, Y) of the operator

$$D_{XY}\Phi = d_Y\Phi - (-1)^{|\Phi|}\Phi \, d_X$$

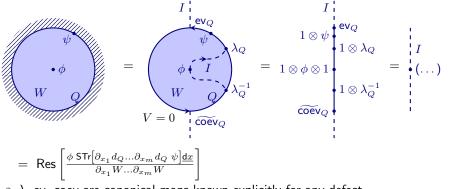
• Product (chiral ring) – composition of maps, defect fusion – \otimes of m.f.





Defects in LG models

- one can compute arbitrary topological correlators using defects and m.f.
- first express everything in terms of defects, then evaluate using m.f. by composing horizontally (⊗) and vertically (○)
- e.g. disk correlator:



- $\lambda,$ ev, coev are canonical maps known explicitly for any defect [Carqueville, Murfet '12]
- natural language for top. defects bicategories with adjoints

Daniel Plencner (LMU Munich)

LG orbifolds - bulk sector

study twisted **RR g.s.** and **(c,c)-fields** unprojected RR g.s. in g-th sector $|\phi_g\rangle = \prod_{\Theta_i^g \in \mathbb{Z}} (X_i)^{l_i} |0\rangle_{RR}^g$

defect/m.f. perspective:

compute cohomology $\operatorname{Hom}(I,I_g)$

$$\begin{split} \phi_g = \prod_{\Theta_i^g \in \mathbb{Z}} (X_i)^{l_i} \prod_{\Theta_i^g \notin \mathbb{Z}} \omega_i^g \\ [\text{Brunner, Roggenkamp '07}] \end{split}$$

[Intriligator, Vafa '90]

action of
$$h \in G$$

 $hX_ih^{-1} \equiv h_i^j X_i = e^{2\pi i \Theta_i^h} X_i$
 $h|0\rangle_{RR}^g = \det(h) e^{2\pi i \sum_{\Theta_i^h \in \mathbb{Z}} \Theta_i^g} |0\rangle_{RR}^g$

evaluate the diagram

$$P_{\rm orb}^{RR}(\phi) = \left(\begin{array}{c} & \\ \phi \\ \end{array} \right)^{A_G}$$

 \rightsquigarrow reproduces the spectrum and phases above

(c,c)-fields similarly - unprojected spectrum iso to RR g.s. via spectral flow,

but different representation of G – reproduced by $P_{\rm orb}^{(c,c)}(\phi) = \begin{pmatrix} A_G \\ \phi \end{pmatrix}$

[Brunner, Carqueville, DP '13]

Generalized LG orbifolds - bulk sector

- generalization: allow any defect A with separable Frobenius algebra structure
- the RR g.s. and (c,c) fields are then images of the projectors

$$P_{\mathsf{orb}}^{RR}(\phi) = \left(\begin{array}{c} \downarrow A \\ \phi \downarrow \end{array} \right) = \left(\begin{array}{c} \downarrow A \\ \phi \downarrow \end{array} \right) = \left(\begin{array}{c} \downarrow A \\ \phi \downarrow \end{array} \right)$$

$$P_{\mathsf{orb}}^{(c,c)}(\phi) = \left(\begin{array}{c} \downarrow A \\ \phi \downarrow \end{array} \right)$$

$$A = \left(\begin{array}{c} \downarrow A \\ \phi \downarrow \end{array} \right)$$
Nakayama automorphism $\Rightarrow \gamma_A = 1_A$: {RR g.s.} \cong {(c,c)} (spectral flow)

• (c,c) fields endowed with a commutative product $\phi_1 \cdot \phi_2 = \phi_1 + \phi_2$

• topological bulk pairing:
$$\langle \phi_1, \phi_2 \rangle = \left\langle \begin{array}{c} & & \\ &$$

computes pairing on RR ground states

consistency check: nondegenerate in the generalized orbifold theory

Daniel Plencner (LMU Munich)

 γ

LG orbifolds - boundary sector

 a boundary Q in LG orbifold is described by a G-equivariant matrix factorization, i.e. there is a representation γ of G on Q, s.t.

$$\gamma \, d_Q(gx_i) \, \gamma^{-1} = d_Q(x_i)$$

• for generalized orbifolds, boundaries are A-modules

$$X : X \otimes A \longrightarrow X \quad \text{s.t.} \quad = \quad , \quad = \quad ,$$

boundary fields module maps

$$\begin{array}{c} Y \\ \Phi \\ X \end{array} = \begin{array}{c} Y \\ \Phi \\ X \end{array}$$

• for $A = A_G$ this reproduces G-equivariant matrix factorizations [Carqueville, Runkel '12]

LG orbifolds - topological disk correlators

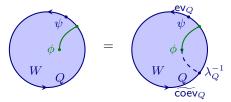
• for ordinary LG orbifolds proposal by [Walcher '04]

$$\left\langle \phi_g \right\rangle_Q = \operatorname{Res}\left[rac{\phi_g^{\mathsf{inv.}} \operatorname{STr}[\gamma \, \partial_1 d_{\bar{Q}} \dots \partial_r d_{\bar{Q}}]}{\partial_1 \overline{W} \dots \partial_r \overline{W}}
ight]$$

with $\phi_g^{\text{inv.}}$ the polynomial part of ϕ_g , and \overline{W} , $d_{\bar{Q}}$ are W, d_Q with non-invariant variables set to zero.

checks: Cardy condition, comparison with known D-brane charges

• defect approach:



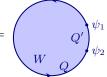
• reproduces above proposal for $A = A_G$

Generalized LG orbifolds - boundary sector

• topological disk correlators:

$$\left\langle \phi\psi\right\rangle _{Q} \ = \left(\begin{array}{c} \psi \\ \phi \\ \psi \\ W \\ Q \end{array} \right)$$

$$\left\langle \psi_1, \psi_2 \right\rangle_{Q,Q'} \; = \;$$



• consistency checks: boundary paring nondegenerate:

boundary chiral sector paired with Ramond sector

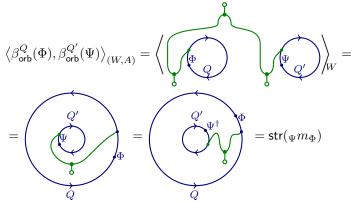
$$\left\langle -,-\right\rangle _{Q,Q'}:\operatorname{Hom}(Q,Q')\times\operatorname{Hom}(Q',{}_{\gamma_A}Q[n])\longrightarrow \mathbb{C}$$

nontrivial Serre functor $S_A = {}_{\gamma_A}(-)$ twisting the A-action:

where again γ_A is the Nakayama automorphism $\gamma_A = \bigcap_{A} \bigcap_{A}$

Cardy condition

- boundary-bulk map $\beta^Q_{\rm orb}(\Psi) := \bigvee_{\mathbf{v}} \bigvee_{\mathbf{v}} Q$
- Theorem: The Cardy condition holds for (generalized) LG orbifolds, i.e. $\langle \beta^Q_{orb}(\Phi), \beta^{Q'}_{orb}(\Psi) \rangle_{(W,A)} = \operatorname{str}(_{\Psi}m_{\Phi})$ Proof:



using only the general (algebraic) properties of A and Q,Q^\prime

Daniel Plencner (LMU Munich)

Orbifold equivalences between LG models

• Recall: if there is a top. defect $X: V \to W$ with

X $\bullet 1$ W $\bullet c1$ W

s.t. $c \equiv \dim(X) \in \mathbb{C} \setminus \{0\}$, one can describe W as a generalized orbifold of V with $A = X^{\dagger} \otimes X$

• in LG models amounts to finding matrix factorisation X of W(x) - V(z) s.t.

$$\dim(X) = \operatorname{Res}\left[\frac{\operatorname{STr}\left[\prod_{i} \partial_{x_{i}} d_{X} \prod_{j} \partial_{z_{j}} d_{X}\right] \underline{dz}}{\partial_{z_{1}} V, \dots, \partial_{z_{m}} V}\right] \neq 0$$

- constructed explicitly between A-↔D-type and A-↔E-type singularities [Carqueville, Runkel '12], [Carqueville]
- Task: classify defects with invertible $\dim(X) \rightsquigarrow$ new equivalences between LG models and their D-brane categories $(\operatorname{mod}(X^{\dagger} \otimes X, V) \cong \operatorname{mf}(W))$ beyond the rational case

Summary & Outlook

- describe orbifolds via defects ~> generalized orbifolds
- one can recover standard results on LG orbifolds in this approach
- in particular, rigorously derive expressions for all topological correlators (e.g. RR-charges of D-branes, eff. superpotentials) and a new, simpler proof of Cardy condition
- all this can be carried over to the generalized setting
- several consistency checks: nondegeneracy of bulk and boundary pairings, Cardy condition

Outlook:

- understand discrete torsion and its effect on D-branes/defects from the generalized perspective [Brunner, Carqueville, DP in progress]
- find **new equivalences** between theories via a generalized orbifold construction (e.g. between different **CY** compactifications)

Daniel Plencner (LMU Munich)