

Introduction to Accelerator Physics

Part 2

Pedro Castro / Accelerator Physics Group (MPY)
Introduction to Accelerator Physics
DESY, 22nd July 2013

Differences between proton and electron accelerators

HERA (Hadron Electron Ring Accelerator) tunnel:

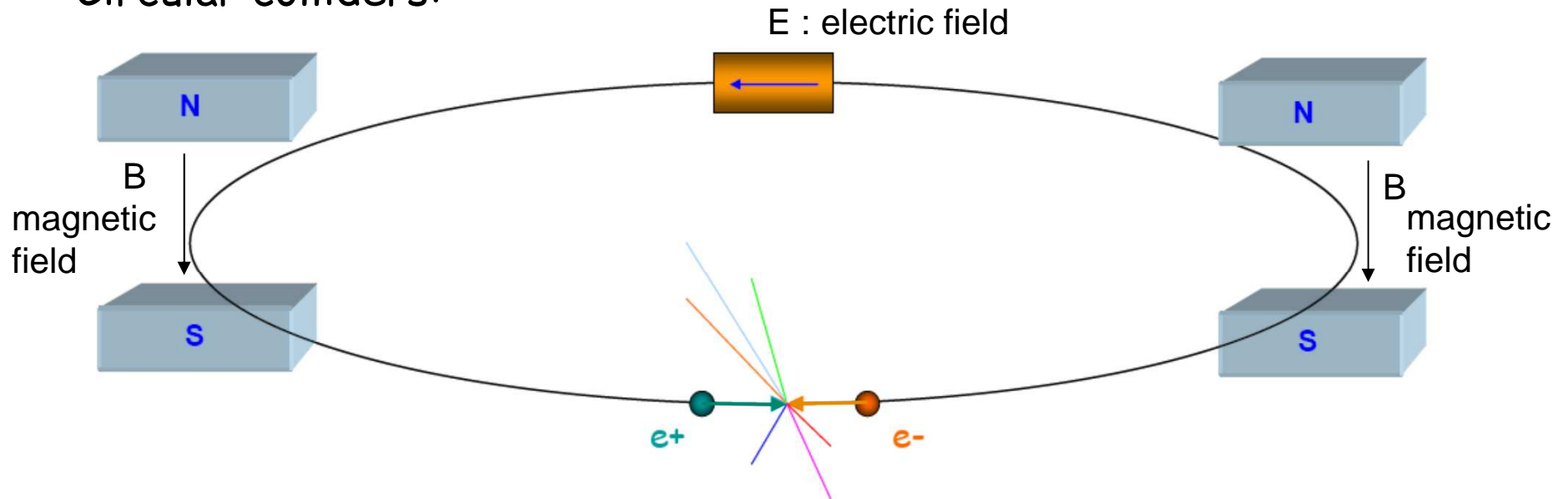


proton
accelerator
920 GeV

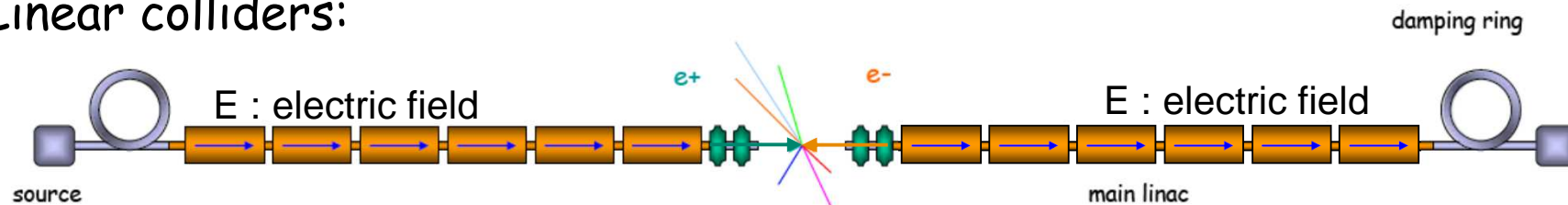
electron accelerator
27.5 GeV

Which collider is better?

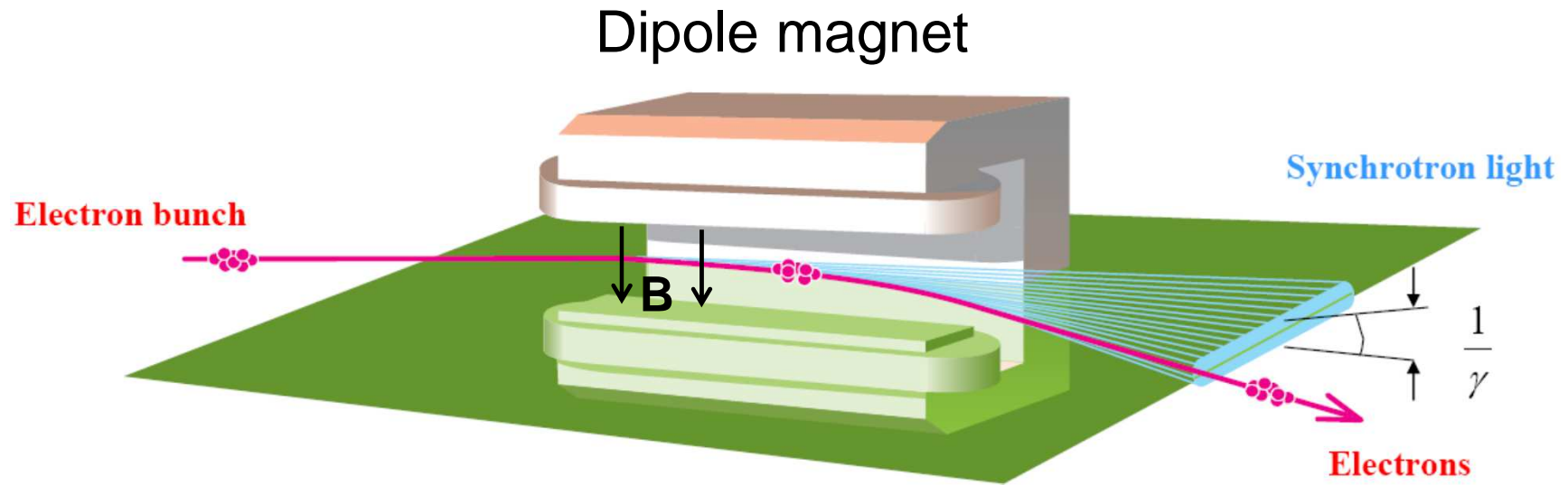
Circular colliders:



Linear colliders:



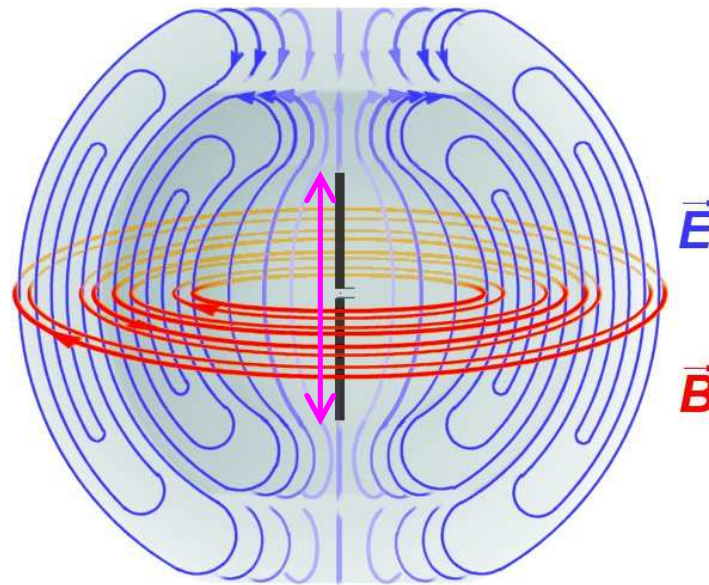
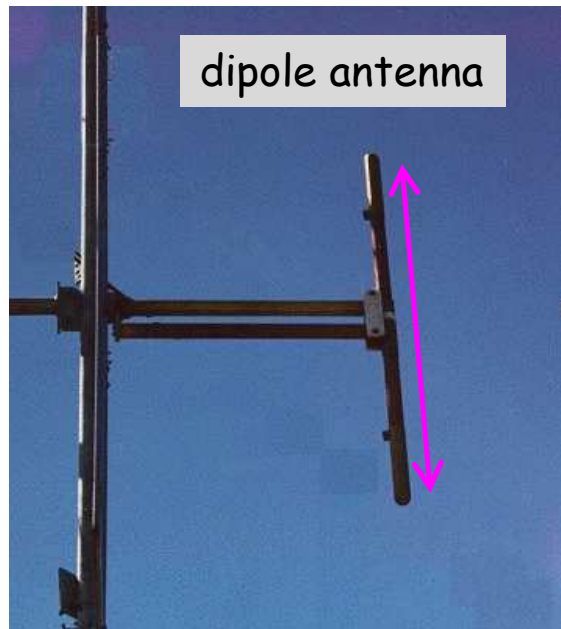
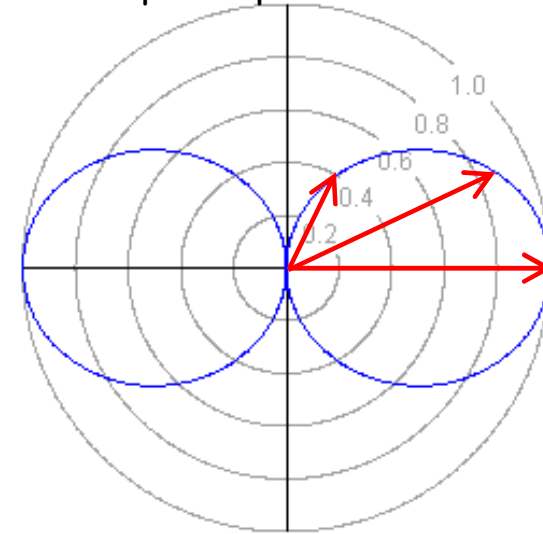
Synchrotron radiation



Radio antenna



radiation power pattern:



Radiation of a dipole antenna

local oscillator:

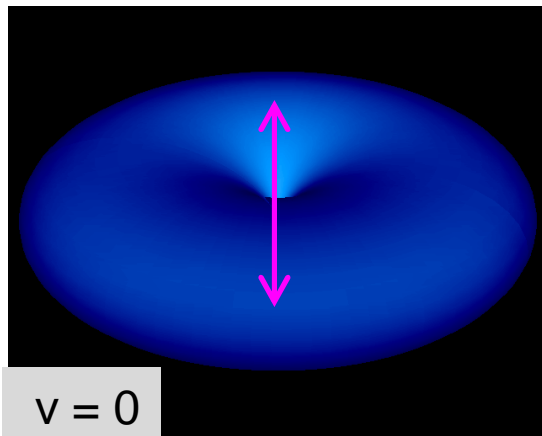
$$P = \frac{q^2 a^2}{12\pi\epsilon_0 c^3} \omega^4$$

(oscillation amplitude: $a < \lambda$)

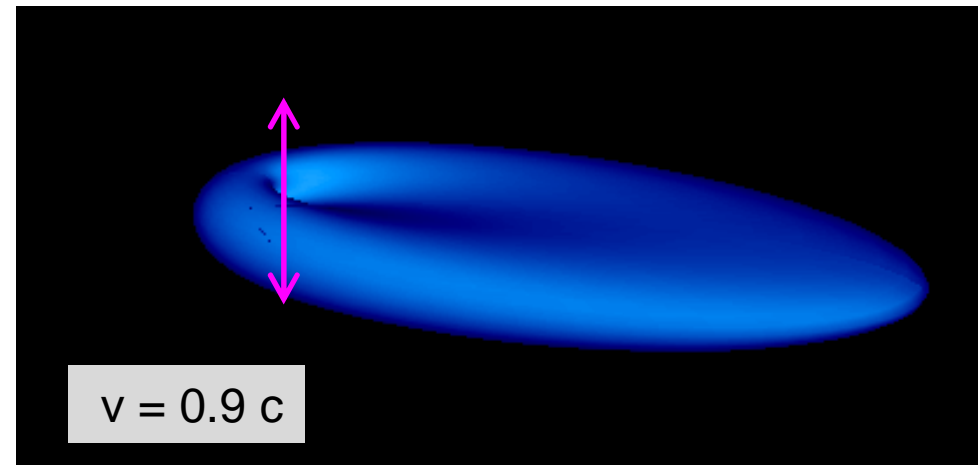
moving oscillator:

$$P = \frac{q^2 a^2}{12\pi\epsilon_0 c^3} \gamma^4 \omega^4$$
$$\gamma = \frac{E}{m_0 c^2}$$

Radiation of an oscillating dipole



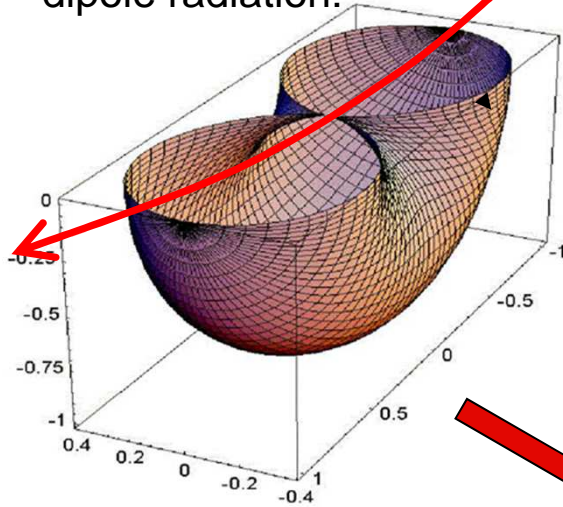
Radiation of a moving oscillating dipole



Radiation of a oscillating dipole under relativistic conditions

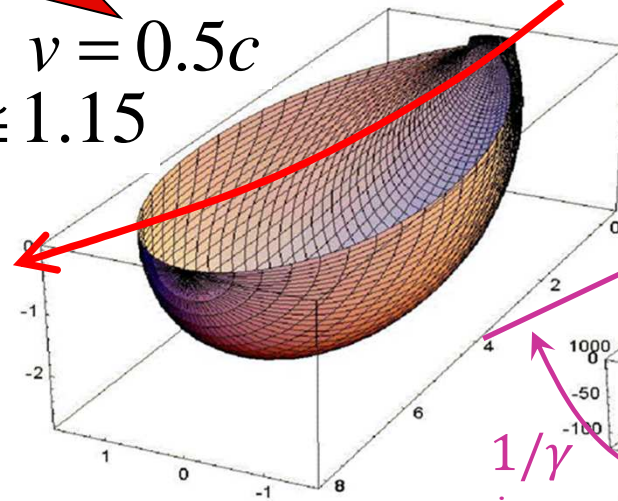
dipole radiation:

electron trajectory

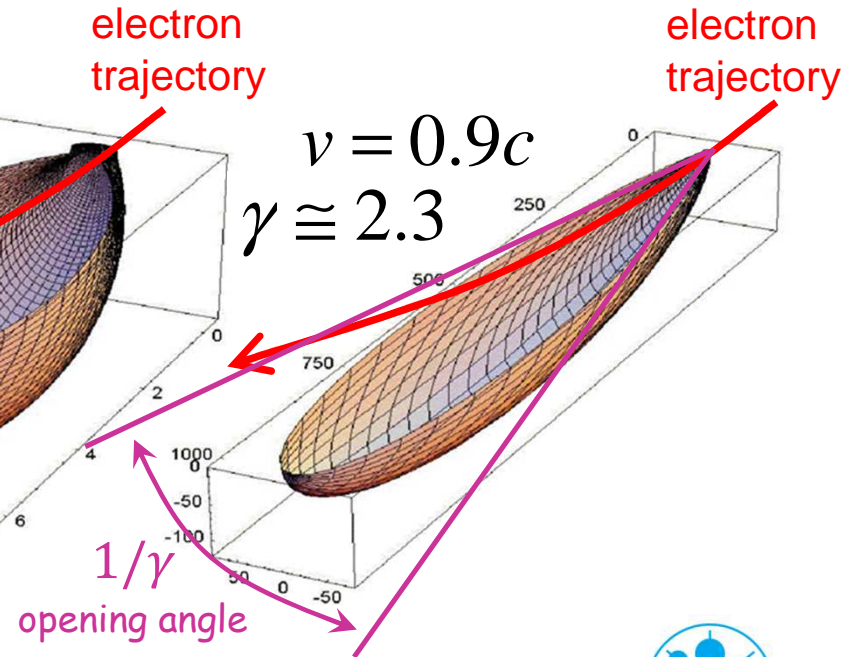


PETRA: $\gamma = 12000$

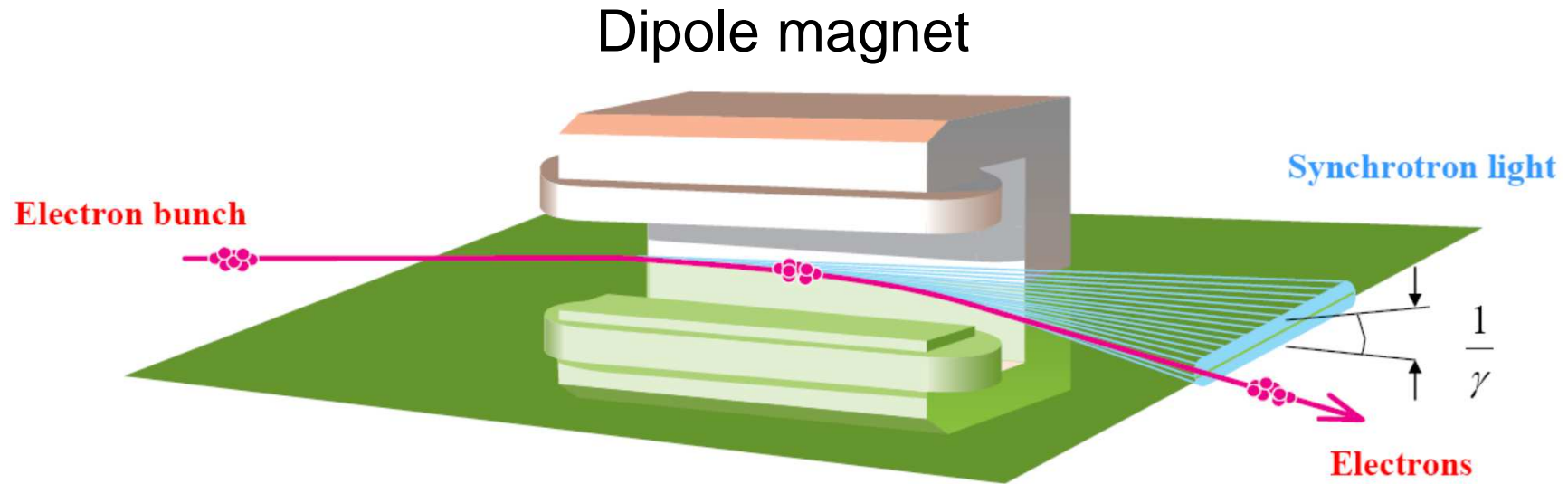
$v = 0.5c$
 $\gamma \cong 1.15$



$v = 0.9c$
 $\gamma \cong 2.3$



Synchrotron radiation



Power radiated by one electron in a dipole field B :

$$P = \frac{c q^2}{6\pi \epsilon_0} \frac{\gamma^4}{r^2}$$

$$\gamma = \frac{E}{m_0 c^2}$$

$$\frac{1}{r} = \frac{q B}{p}$$

vacuum permittivity

Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \quad \Rightarrow \quad \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

$$\gamma = 54000$$

$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

need acceleration = 87 MV per turn



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \quad \Rightarrow \quad \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

$$\gamma = 54000$$

$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

HERA proton ring:

$$r = 580 \text{ m}$$

$$E = 920 \text{ GeV}$$

$$\gamma = 980$$

$$\Delta E_{\text{turn}} \cong 10 \text{ eV (10}^{-9}\text{\%)}$$

← same →

need acceleration = 87 MV per turn



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \Rightarrow \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

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$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

HERA proton ring:

$$r = 580 \text{ m}$$

$$E = 920 \text{ GeV}$$

$$\gamma = 980$$

← same →

the limit is the max. dipole field = 5.5 Tesla

$$R = \frac{mv}{qB} = \frac{p}{qB} \Rightarrow p_{\text{max}} = RqB_{\text{max}}$$

$$\frac{1}{r} = \frac{qB}{p}$$

need acceleration = 87 MV per turn



Synchrotron radiation

Total energy loss after one full turn:

$$\Delta E_{\text{turn}} = \frac{q^2}{3\epsilon_0} \frac{\gamma^4}{r} \Rightarrow \Delta E_{\text{turn}} [\text{GeV}] = 6.032 \times 10^{-18} \frac{\gamma^4}{r[\text{m}]}$$

HERA electron ring:

$$r = 580 \text{ m}$$

$$E = 27.5 \text{ GeV}$$

$$\gamma = 54000$$

$$\Delta E_{\text{turn}} = 87 \text{ MeV (0.3\%)}$$

need acceleration = 87 MV per turn

LEP collider:

$$r = 2800 \text{ m}$$

$$E = 105 \text{ GeV}$$

$$\gamma = 205000$$

$$\Delta E_{\text{turn}} \cong 4 \text{ GeV (4\%)}$$

need 4 GV per turn !!



International Linear Collider (ILC)

Colliding beams with $E_{cm} = 500 \text{ GeV}$ (update to 1 TeV possible)

(not to scale)

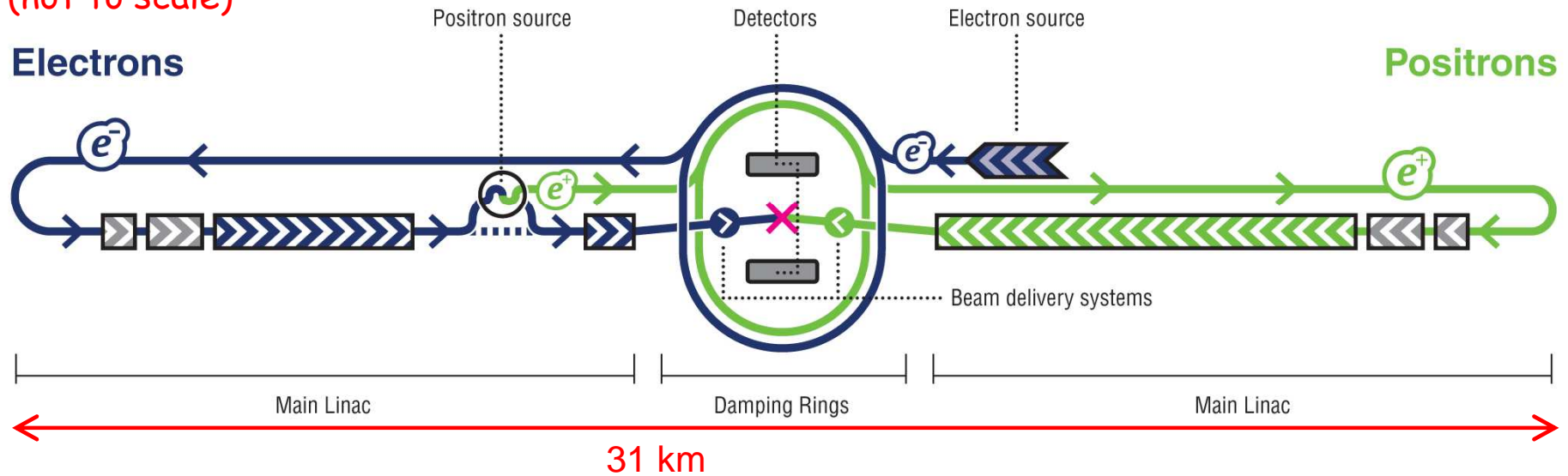
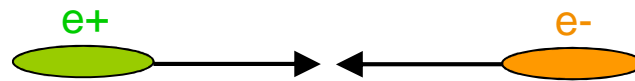
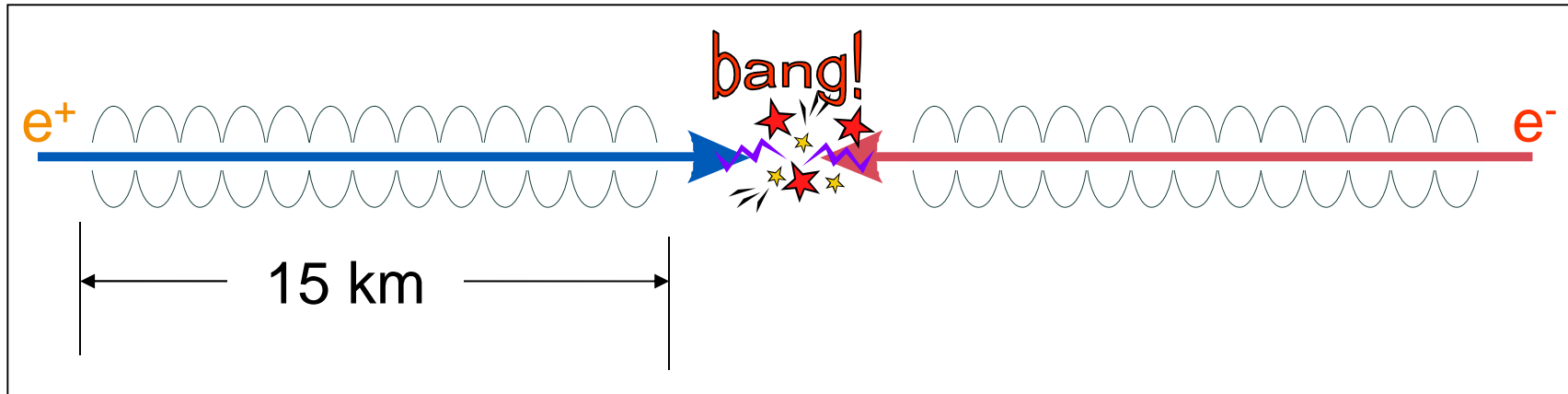


Figure of merit: Luminosity



production rate of a given event (for example, Z particle production):

$$R_Z = \frac{dN_Z}{dt} = \Sigma_Z \cdot L$$

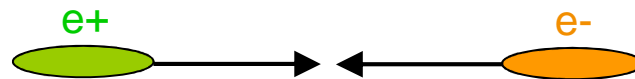
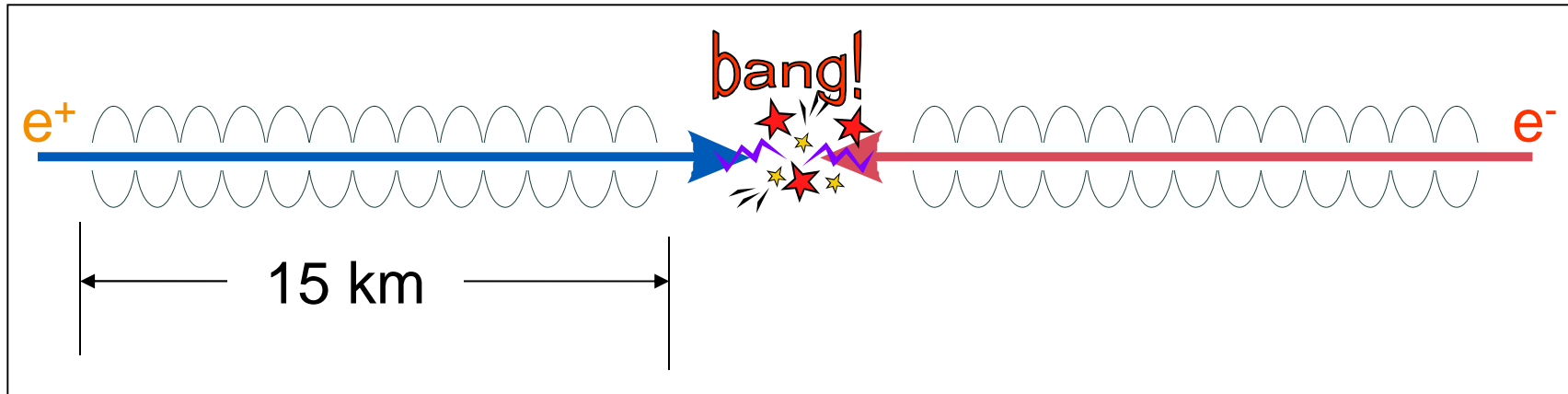
number of events

cross section of Z production

luminosity (independent of the event type)



Luminosity



number of colliding bunches per second

number of positrons per bunch

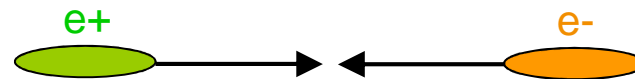
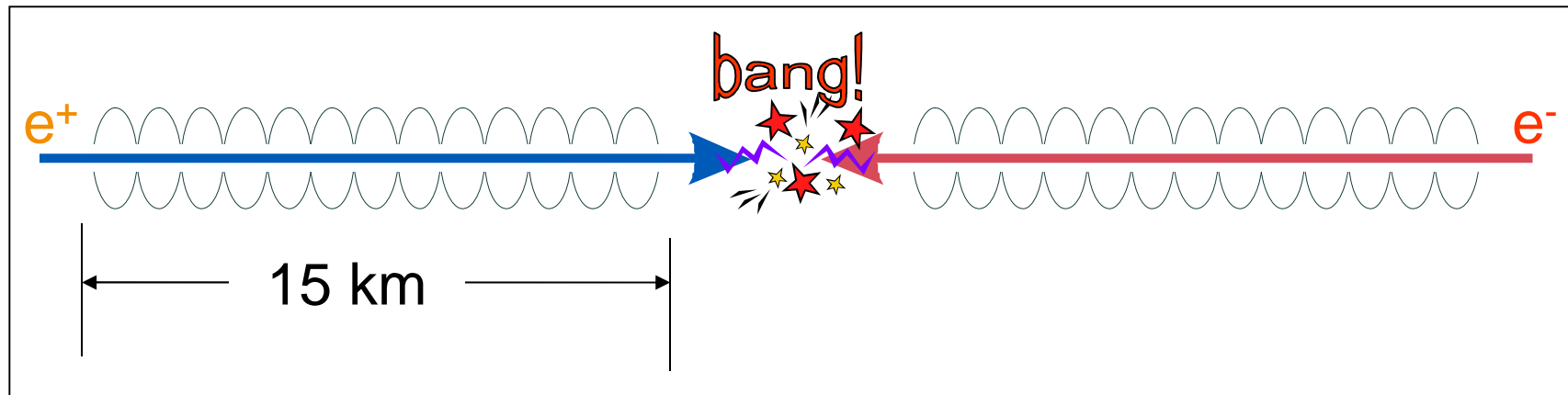
number of electrons per bunch

$$R_Z = \sum_z \cdot L = \sum_z \cdot \frac{f_b N_{e^+} N_{e^-}}{4\pi \sigma_x \sigma_y}$$

luminosity

transverse bunch sizes (at the collision point)

Luminosity



number of colliding bunches per second

number of positrons per bunch

number of electrons per bunch

$$R_Z = \sum_z \cdot L = \sum_z \cdot \frac{f_b N_{e^+} N_{e^-}}{4\pi \sigma_x^* \sigma_y^*} \cdot H_D$$

luminosity

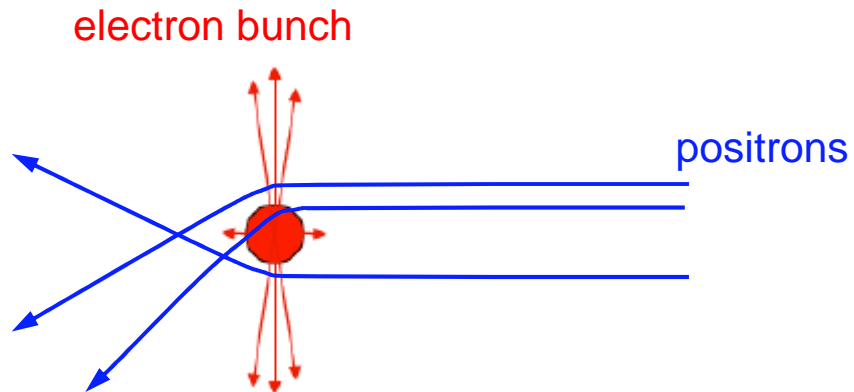
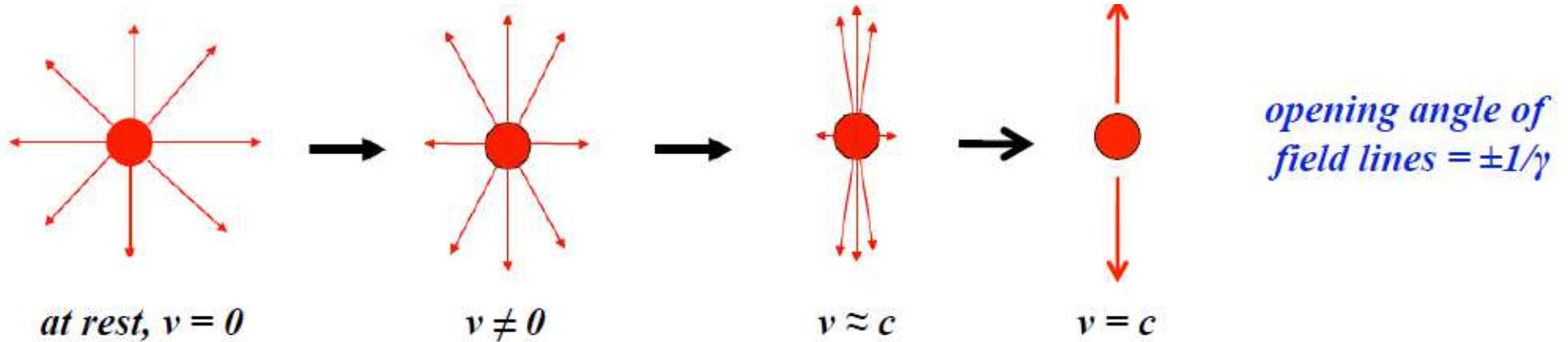
beam-beam enhancement factor

transverse bunch sizes (at the collision point *)

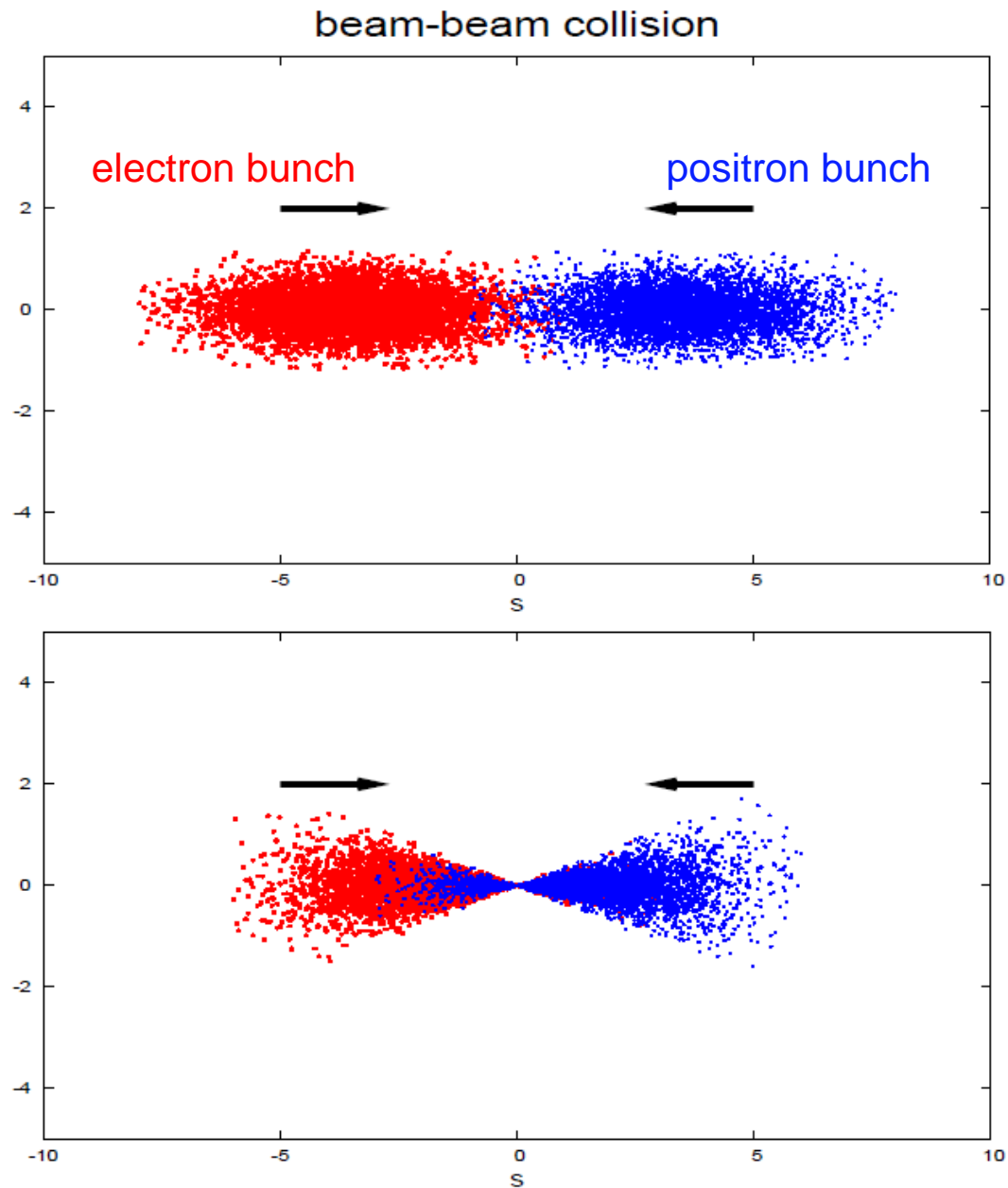


Luminosity enhancement factor H_D due to focusing of opposite beam

electric field of a charged particle (or bunch)



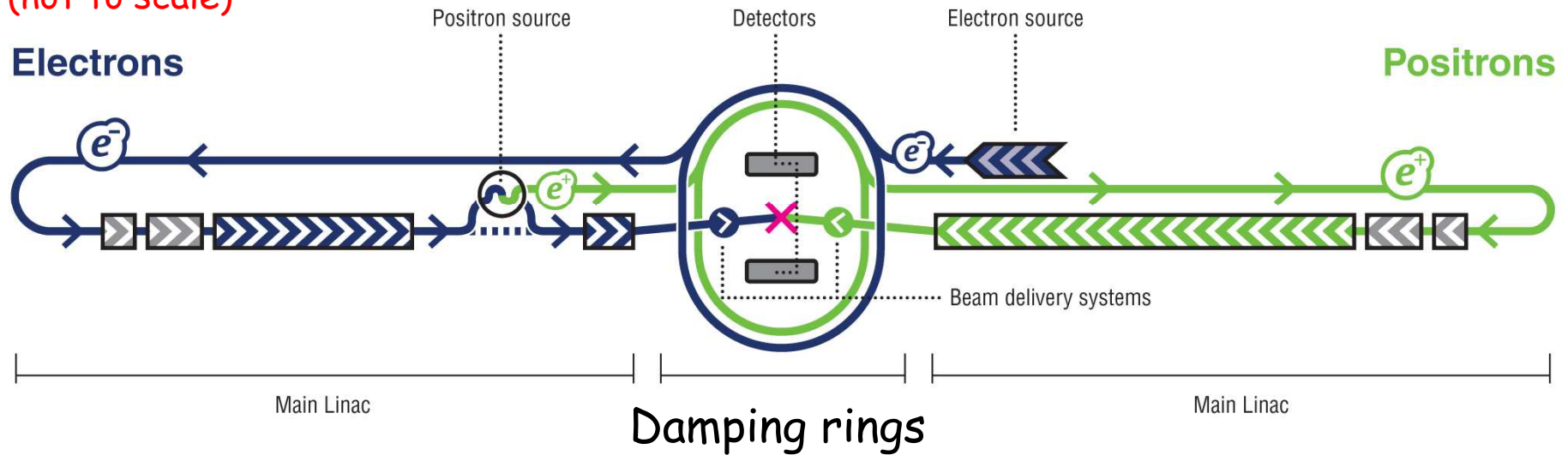
Luminosity enhancement factor H_D due to focusing of opposite beam



International Linear Collider (ILC)

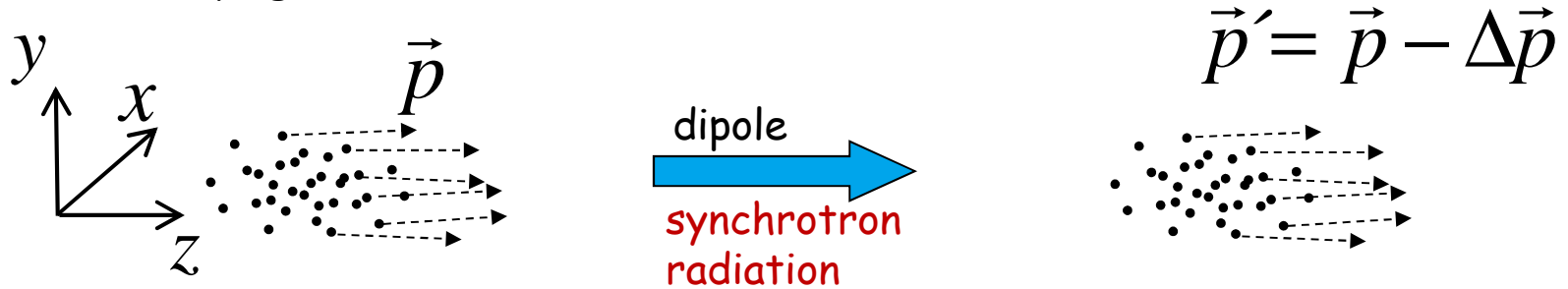
Colliding beams with $E_{cm} = 500 \text{ GeV}$ (update to 1 TeV possible)

(not to scale)

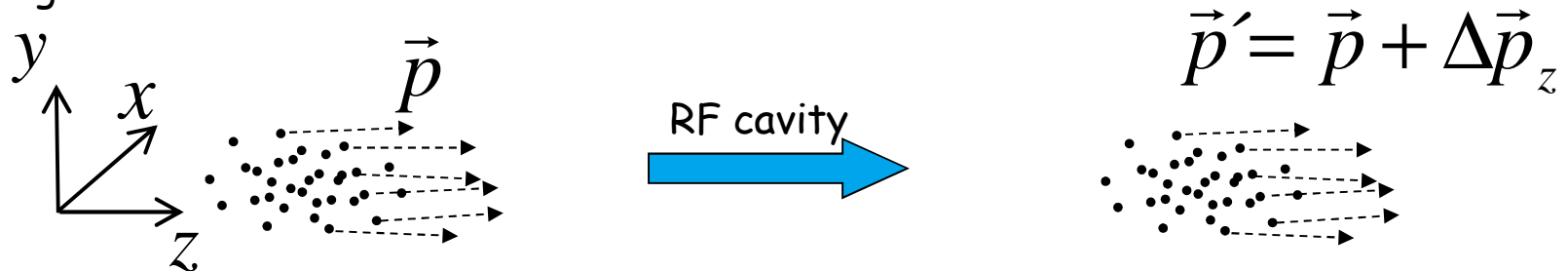


Damping rings

Radiation damping:

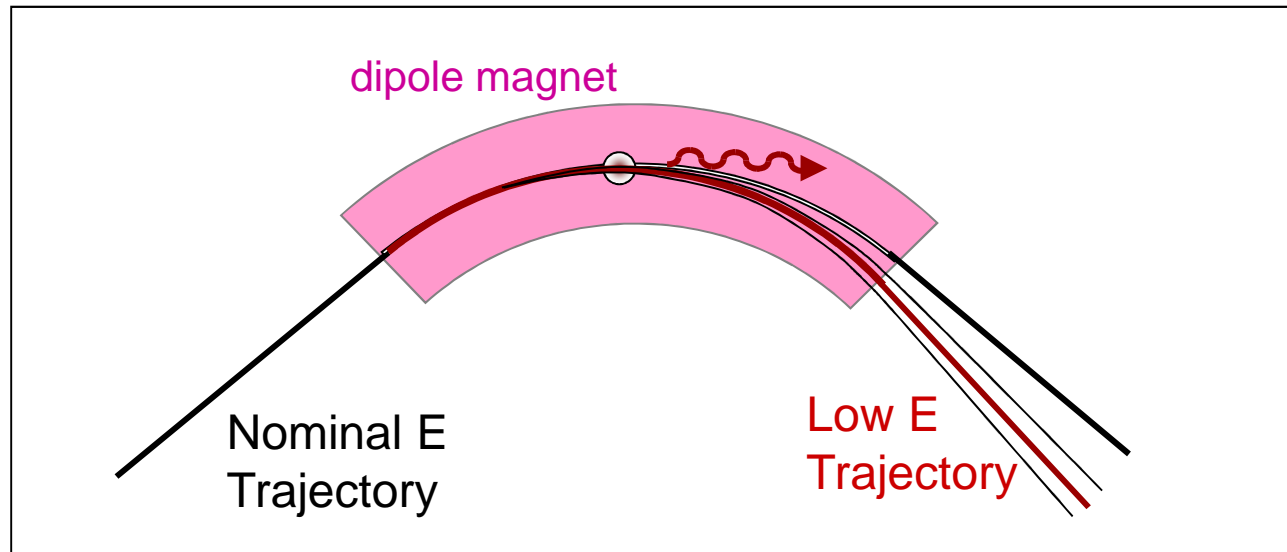


Longitudinal acceleration:



acceleration (only in z direction) $\frac{p_x}{|\vec{p}|}$ gets smaller

Quantum excitation

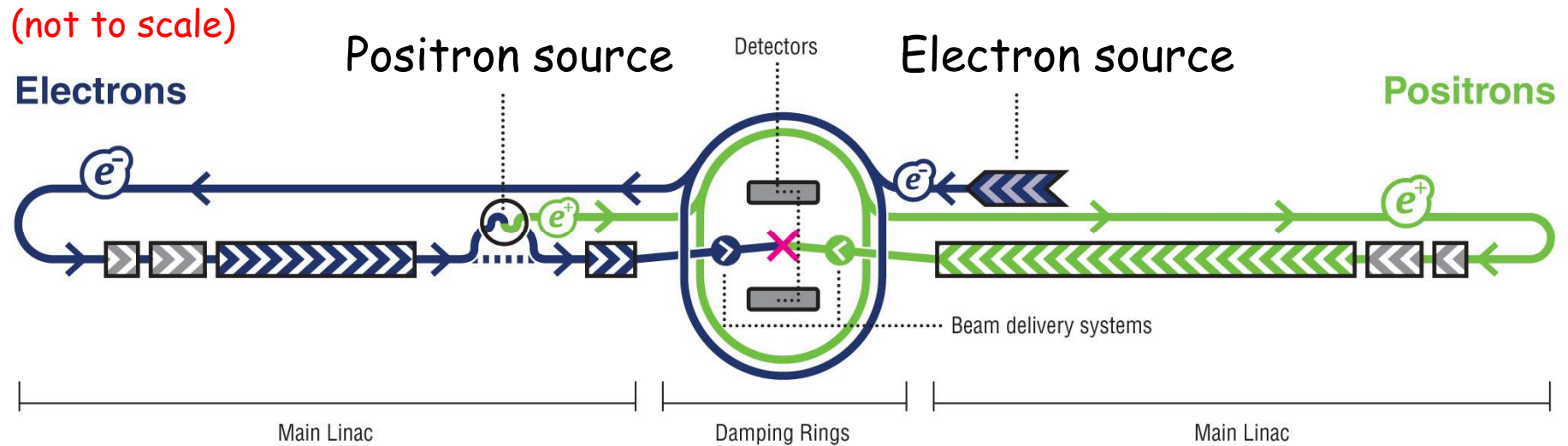


> Quantum excitation

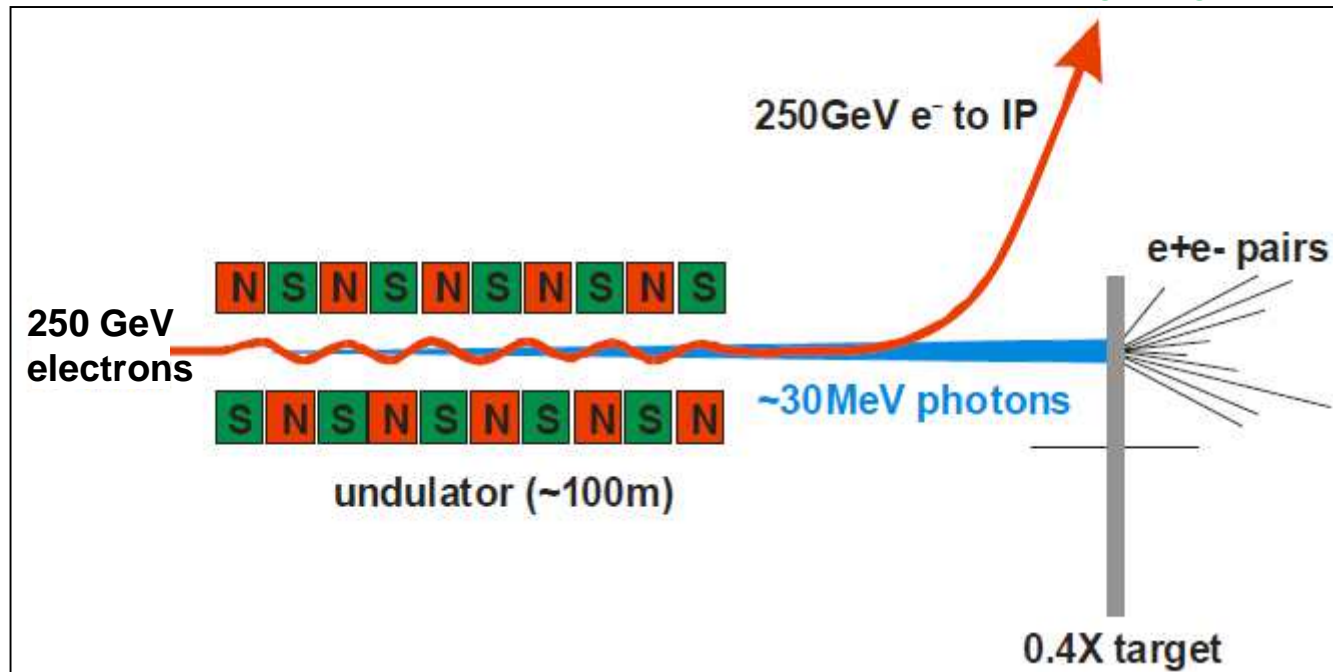
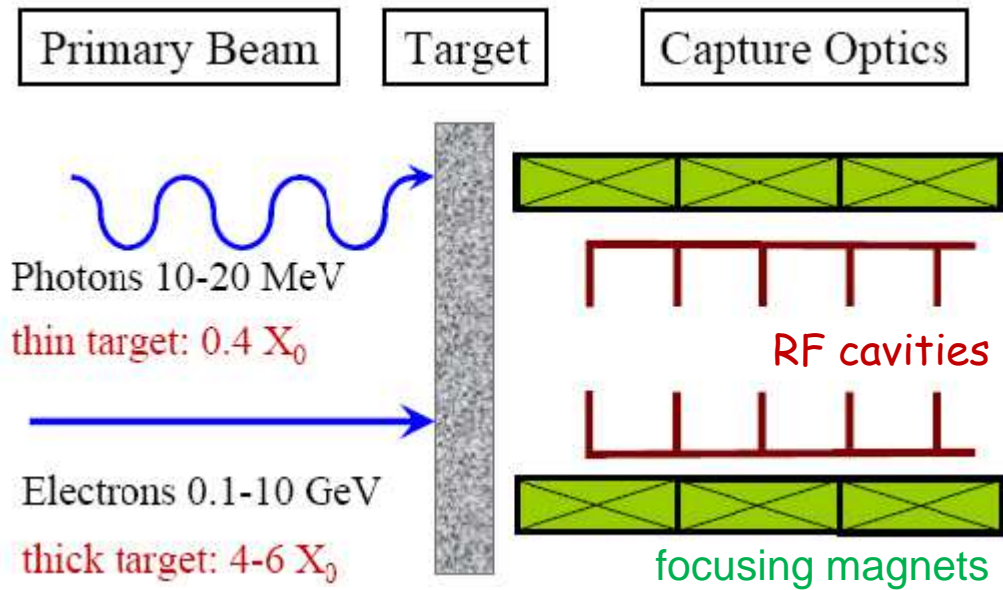
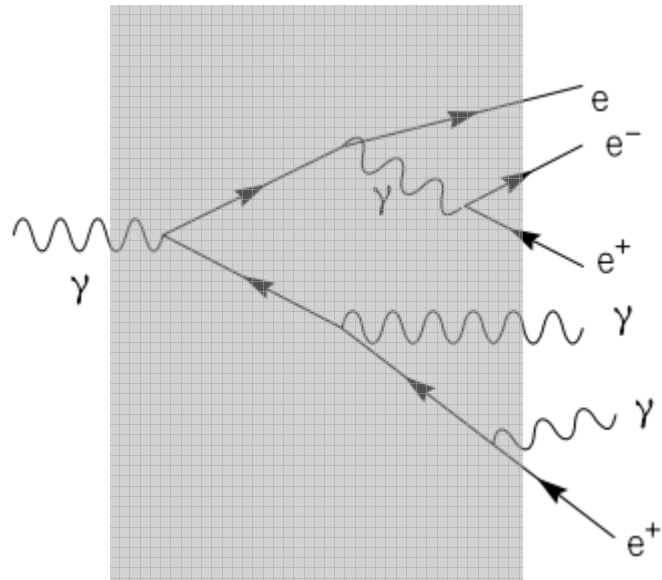
- Radiation is emitted in discrete quanta
- Number and energy distribution etc. of photons obey statistical laws
- → Increase beam size

International Linear Collider (ILC)

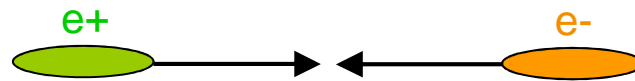
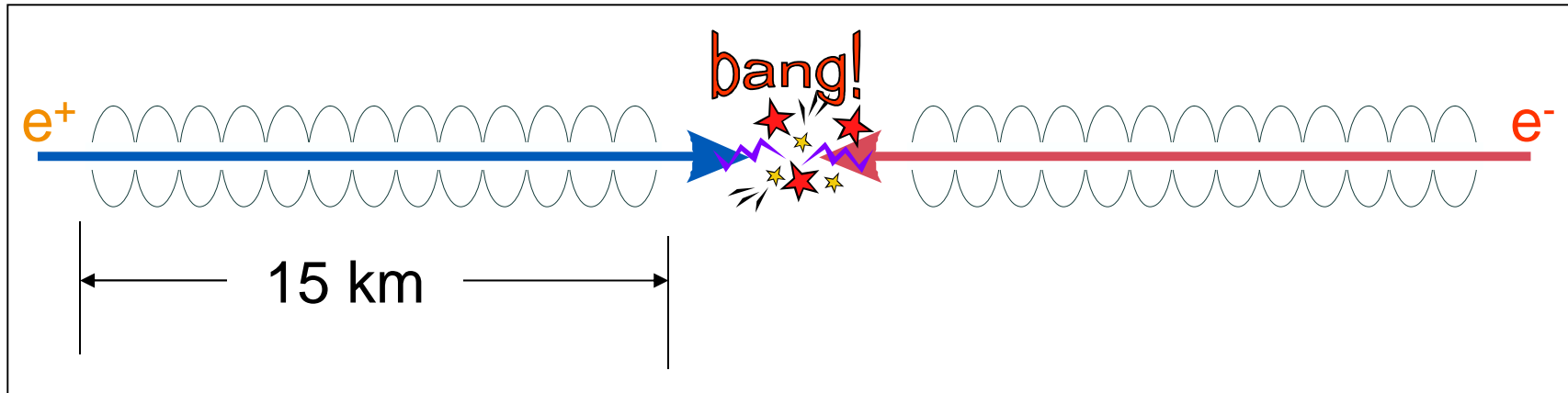
Colliding beams with $E_{cm} = 500 \text{ GeV}$ (update to 1 TeV possible)



Positron source



Luminosity



$$L = \frac{f_b N^2}{4\pi \sigma_x^* \sigma_y^*} \cdot H_D = \frac{\eta_{RF \rightarrow beam} P_{RF} N}{4\pi \sigma_x^* \sigma_y^* E_{cm}} \cdot H_D$$

beam power

RF power

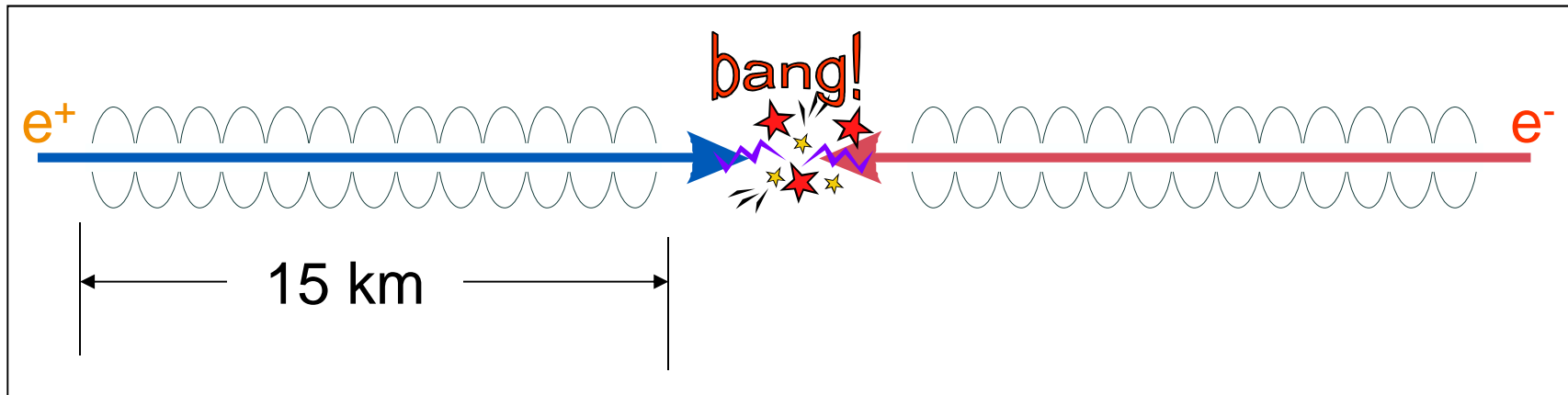
$$f_b N E_{cm} = P_{beam} = \eta_{RF \rightarrow beam} P_{RF}$$

conversion efficiency



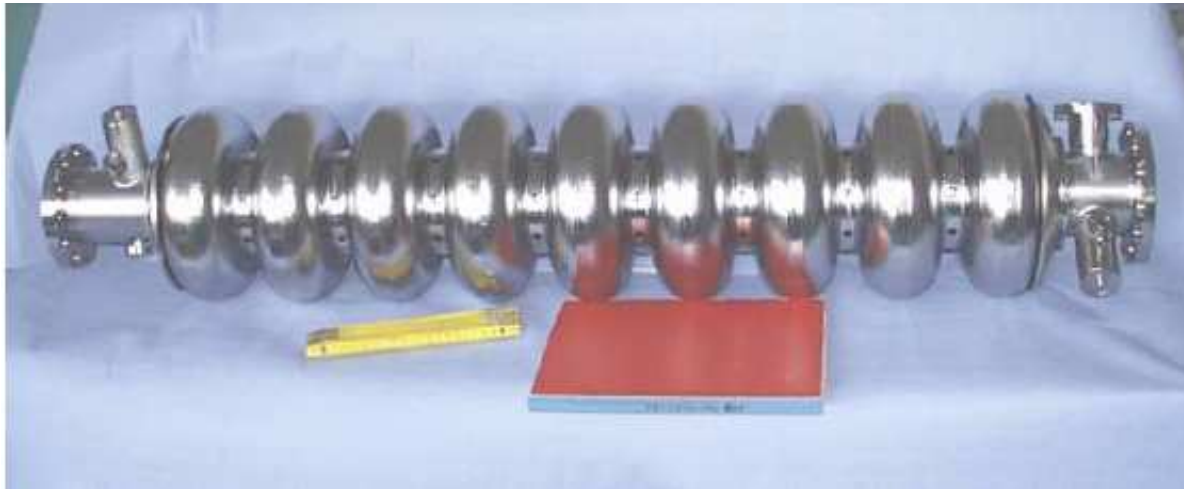
Project for a future e^-e^+ collider: ILC

The International Linear Collider

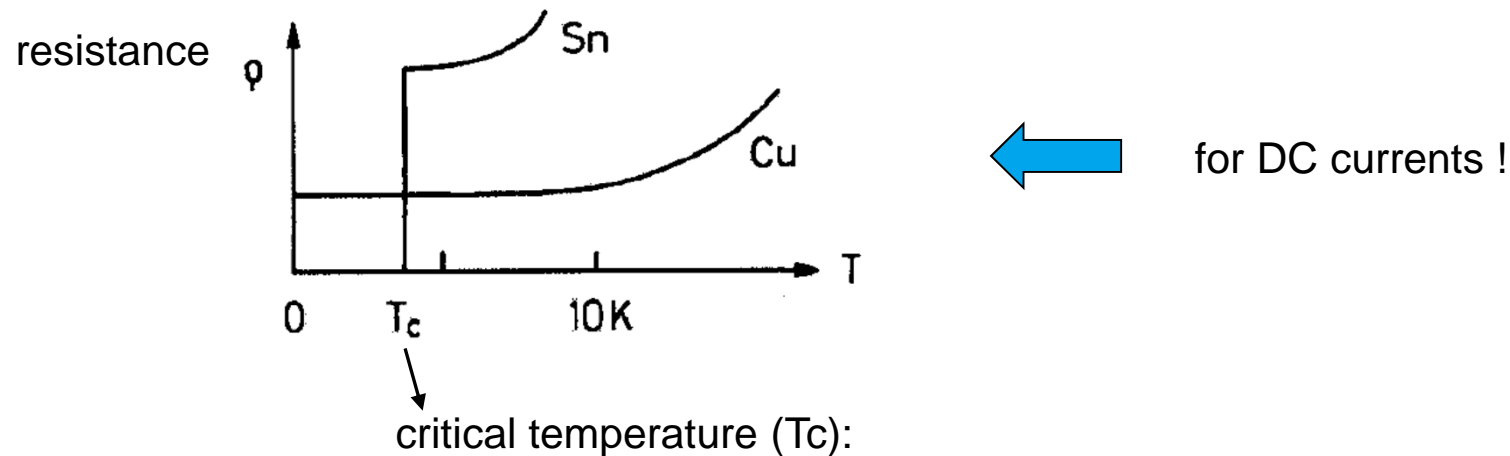


Colliding beams with $E_{CM} = 500 \text{ GeV}$

using superconducting cavities for acceleration:



Advantages of RF superconductivity



at radio-frequencies, there is a "microwave surface resistance"
which typically is 5 orders of magnitude lower than R of copper

Advantages of RF superconductivity

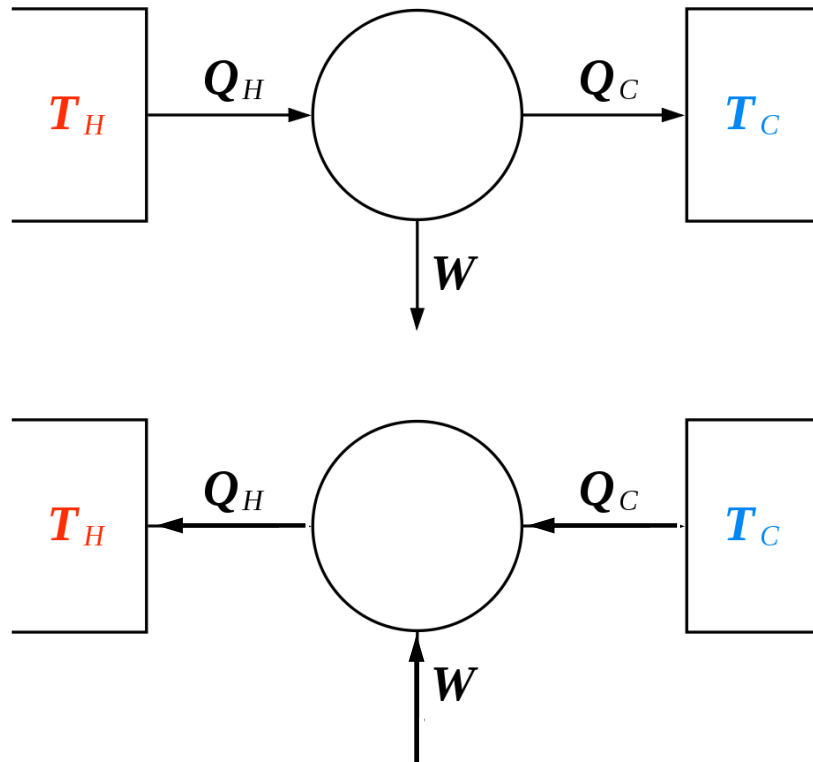
Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for $E = 1 \text{ MV/m}$	$0.4 + 1 \text{ W / m}$ at 2 K	60 kW / m	dissipated at the cavity walls



2nd law of Thermodynamics

“Heat cannot spontaneously flow from a colder location to a hotter location”



max. efficiency

$$\eta_c = \frac{T_H - T_C}{T_H}$$

Carnot efficiency:

$$\eta_c = \frac{T_C}{T_H - T_C}$$

most common applications

thermal power stations,
cars, ...

air conditioners,
refrigerators, ...



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for E = 1 MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
Carnot efficiency: $\eta_c = \frac{T}{300 - T} = 0.007$		x	cryogenics efficiency 20-30%
for E = 1 MV/m	1 kW / m	60 kW / m	
for E = 1 MV/m	1 kW / m	120 kW / m	including RF generation efficiency (50%)

reduction factor of >100 in (electrical) power



Advantages of RF superconductivity

Example: comparison of 500 MHz cavities:

	superconducting cavity	normal conducting cavity	
for E = 1 MV/m	0.4 + 1 W / m at 2 K	60 kW / m	dissipated at the cavity walls
Carnot efficiency: $\eta_c = \frac{T}{300 - T} = 0.007$ x			cryogenics efficiency 20-30%
for E = 1 MV/m	1 kW / m	60 kW / m	
for E = 1 MV/m and 20 mA beam	21 kW / m	80 kW / m	total RF power needed
for E = 1 MV/m	42 kW / m	160 kW / m	including RF generation efficiency (50%)

reduction factor of 4 in (electrical) power





Accelerator Research at DESY

SC Technology
(R&D on CW, cryo module test bench, ongoing)

SC Processes
(large series production & testing of sc cavities and modules, AMTF, ongoing)

Surface Technology
(cavity surfaces, CRISP, ongoing)

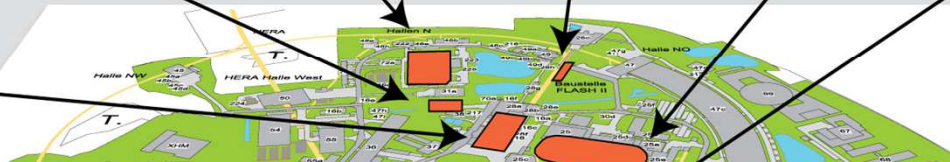
FEL Seeding
(sFLASH, FLASH2 seeding, ongoing)

LAOLA at FLASH2
(beam-driven plasma acceleration, 2016+)

SINBAD (ultra-short bunches, LAOLA, prototype table-top FEL, 2017+)

AXSIS (atto-second bunches, ICS, 2014+)

LAOLA at REGAE
(laser-driven plasma acceleration, 2013-2016)

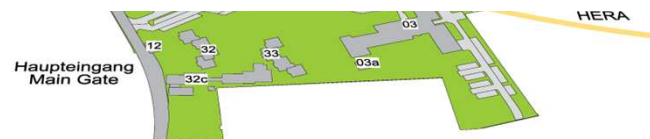


Laboratory for **L**aser- and beam-driven plasma **A**cceleration (LaoLA)
Relativistic **E**lectron **G**un for **A**tomic **E**xploration (REGAE)
Short **I**nnovative **B**unches and **A**ccelerators at **D**oris (SINBAD)
Attosecond **X**-ray **S**cience **I**maging and **S**pectroscopy (AXSIS)
Photo **I**njector **T**est Facility at DESY, Location **Z**euthen (PITZ)

Zeuthen Campus

Photo-Injector
(ongoing)

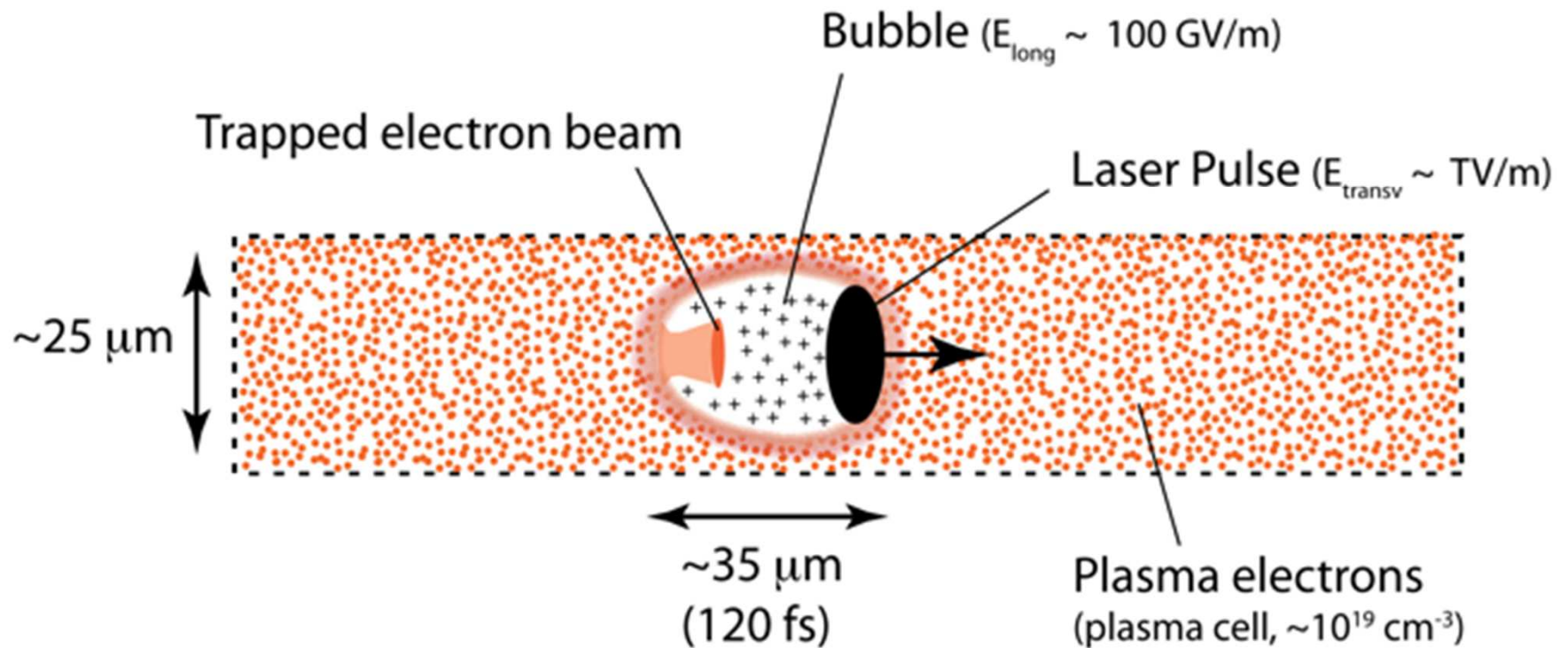
LAOLA at PITZ
(bunch modulation in plasma, 2013+)



Hamburg Campus



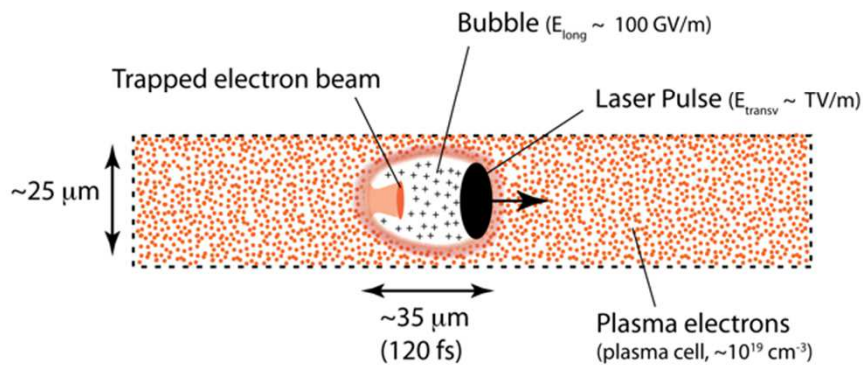
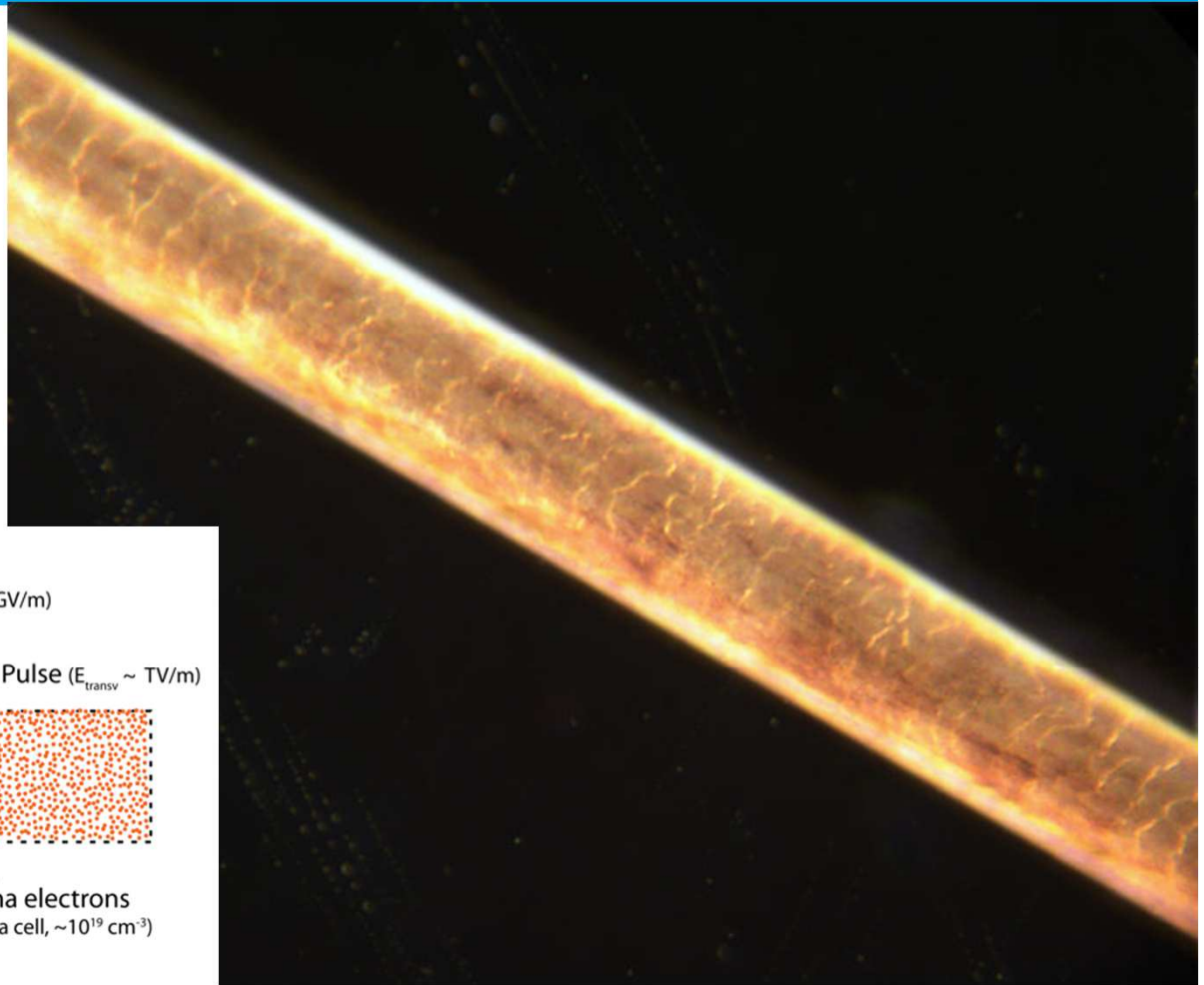
Laser-driven plasma acceleration



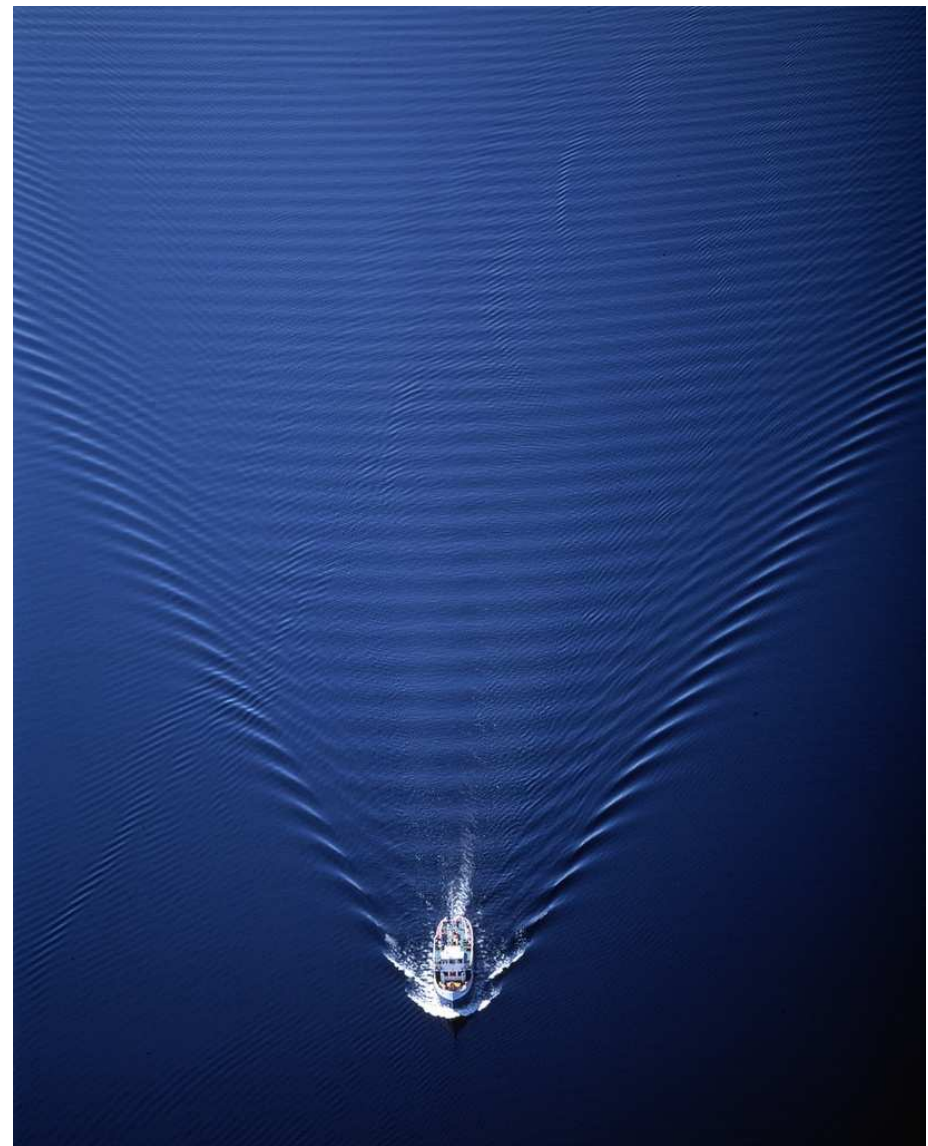
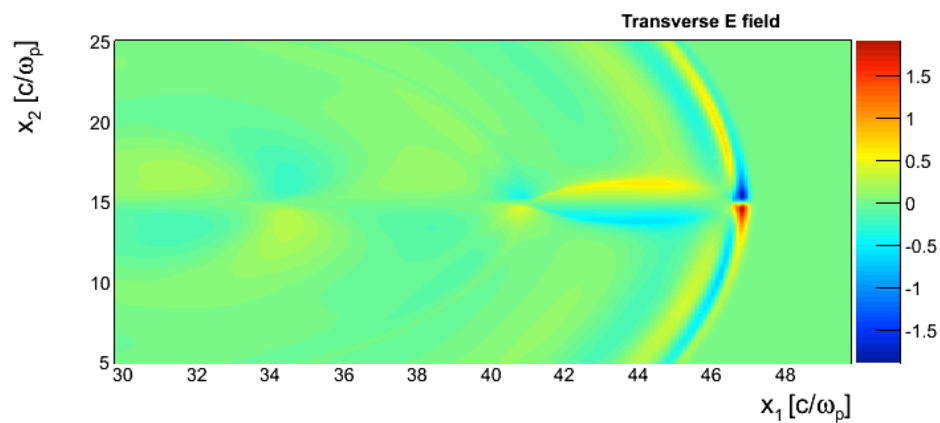
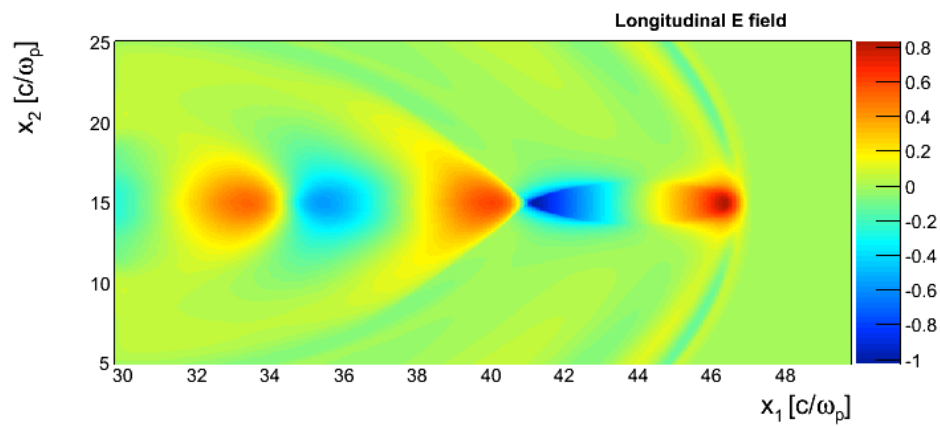
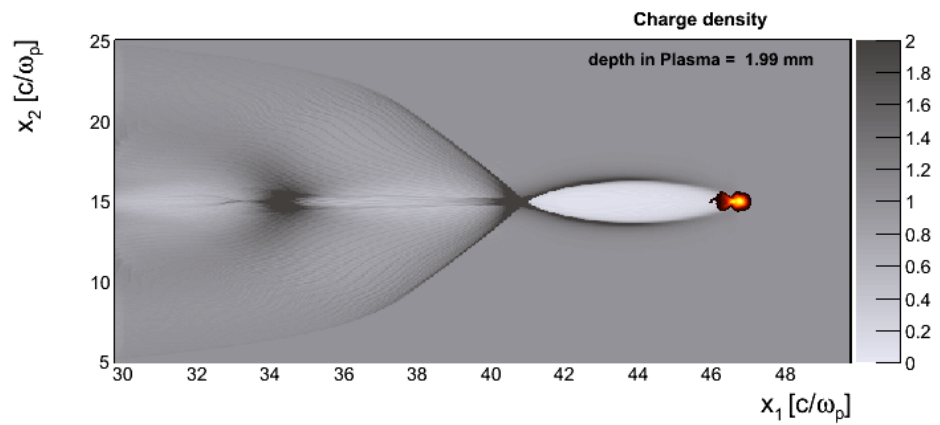
Acceleration (and focusing) within a hair width ...

E_{long} **100 GV/m**

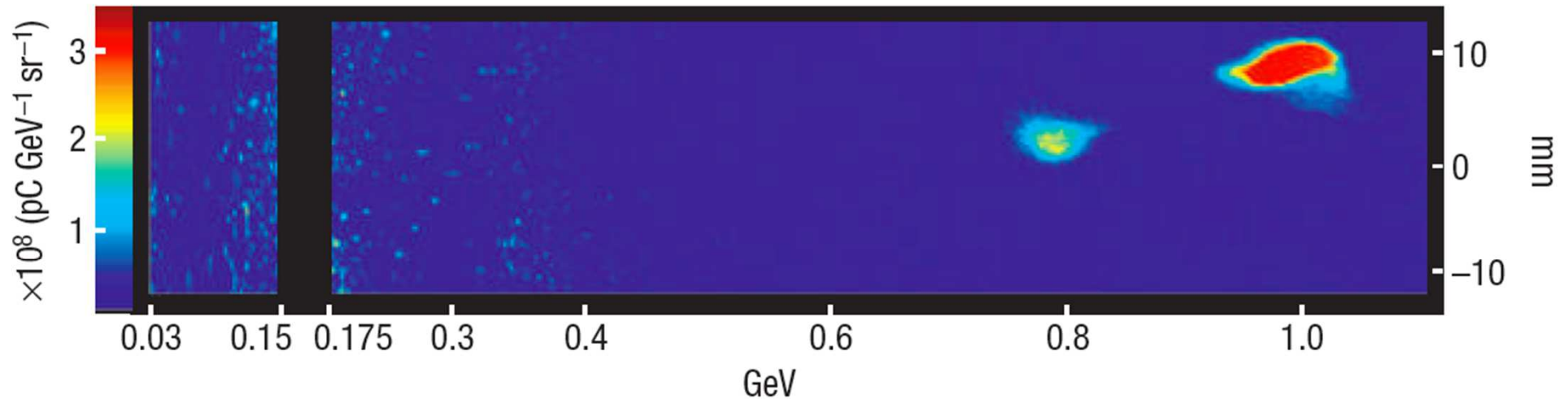
g_{transv} **300 MT/m**



Simulations



Breakthrough in the lab



LETTERS

GeV electron beams from a centimetre-scale accelerator

W. P. LEEMANS^{1*}, B. NAGLER¹, A. J. GONSALVES², Cs. TÓTH¹, K. NAKAMURA^{1,3}, C. G. R. GEDDES¹, E. ESAREY^{1*}, C. B. SCHROEDER¹ AND S. M. HOOKER²

¹Lawrence Berkeley National Laboratory, 1 Cyclotron Road, Berkeley, California 94720, USA

²University of Oxford, Clarendon Laboratory, Parks Road, Oxford OX1 3PU, UK

³Nuclear Professional School, University of Tokyo, 22-2 Shirane-shirakata, Tokai, Naka, Ibaraki 319-1188, Japan

*Also at: Physics Department, University of Nevada, Reno, Nevada 89557, USA

†e-mail: WPLEemans@lbl.gov

2-3 orders of magnitude higher acceleration fields than conventional RF cavities

smaller and cheaper accelerators are possible

nature physics | VOL 2 | OCTOBER 2006



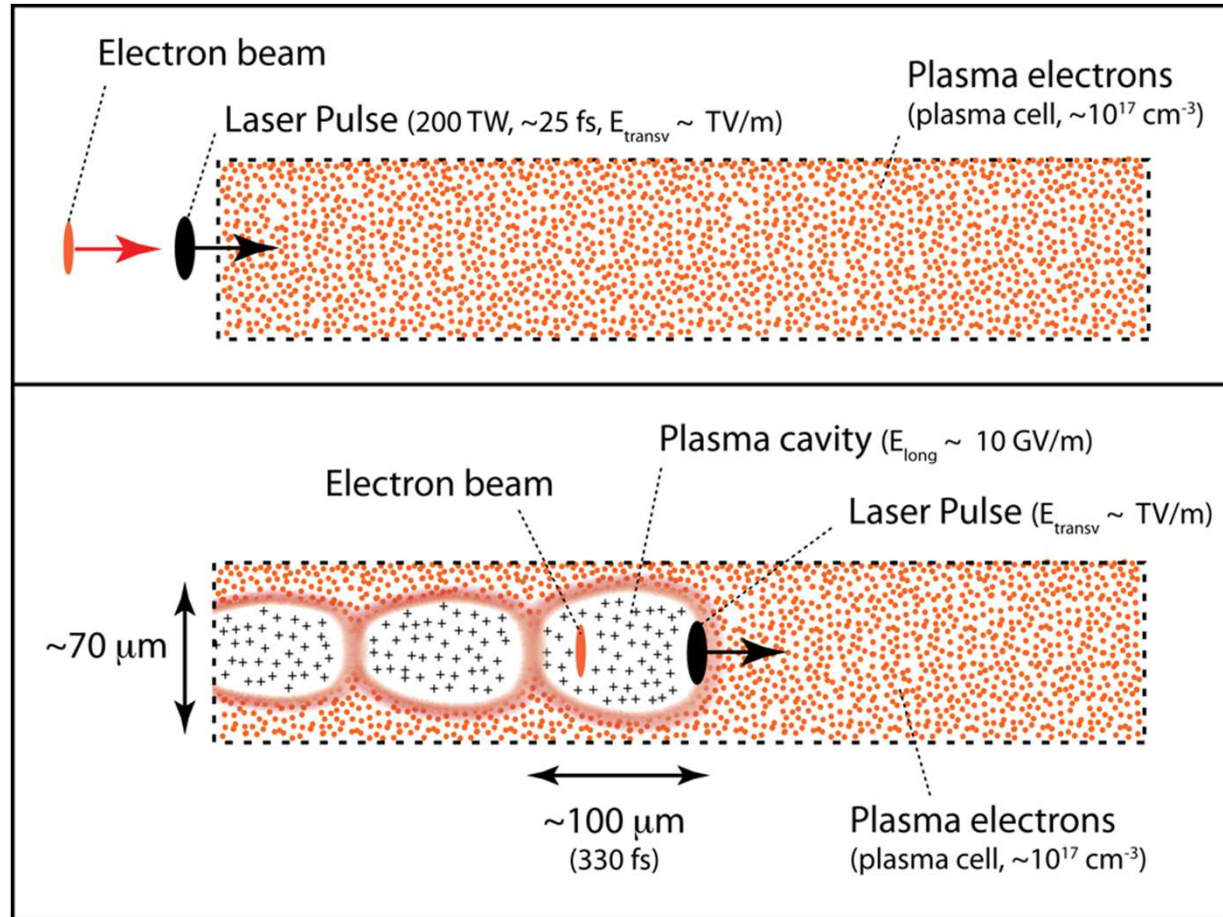
Challenges of plasma acceleration

Maximum energy limitation at 1 GeV or few GeV

Plasma self-injection is very unstable in energy, timing ...



Challenges of plasma acceleration



not yet proven !

Circular colliders (synchrotrons):

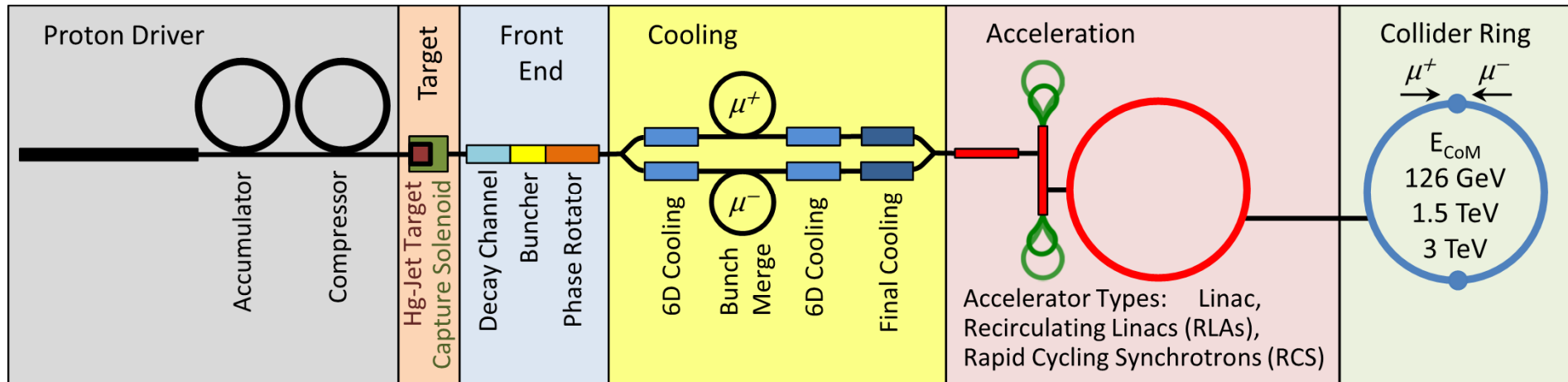
particle type	limitation
• proton synchrotrons 938 MeV/c ²	dipole magnet
• electron synchrotrons 0.511 MeV/c ²	synchrotron radiation
• muon synchrotrons 105.7 MeV/c ²	mean lifetime = 2.2 μs ?

Muon beams produced as tertiary beams:



Muon collider principle

Muon collider block diagram



Collider:

$$E_{cm} = 3 \text{ TeV}$$

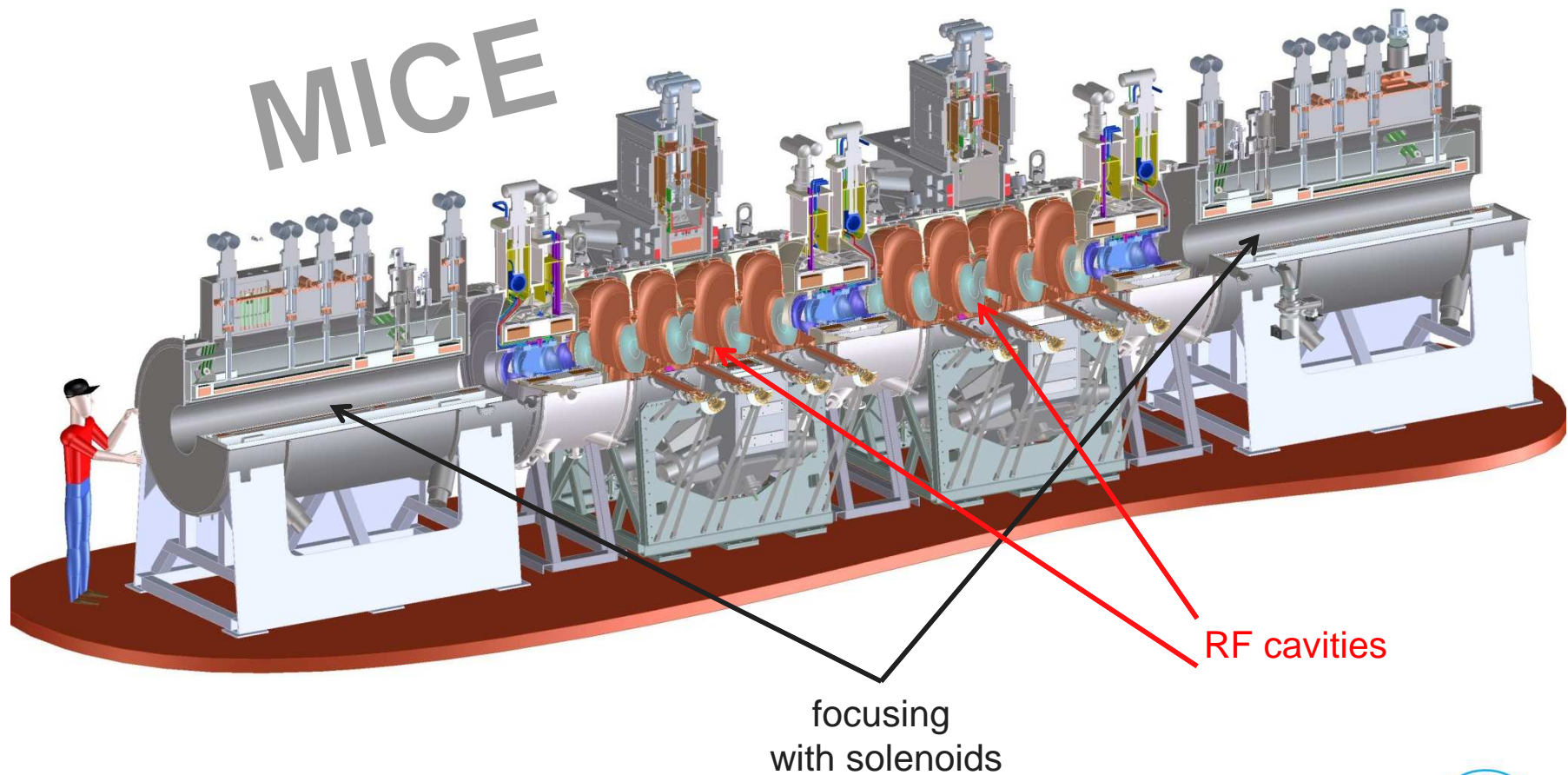
circumference = 4.5 km

$$L = 3 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$



MICE at FERMILAB (Chicago)

The Muon Ionization Cooling Experiment (MICE)
→ demonstrate the method and validate simulations



Summing-up

Basics of synchrotron radiation

particle type	limitation
• proton synchrotrons	dipole magnet
• electron synchrotrons	synchrotron radiation

International Linear Collider (ILC):

- luminosity eq.
- damping rings
- positron source
- power efficiency in superconducting cavities

Two very promising (and challenging) research areas in accel. physics:

- laser-driven plasma acceleration
- muon collider



Thank you for your attention

pedro.castro@desy.de

