

The better NMSSM:

Naturalness and enhanced diphoton rates at Fermi and the LHC

Kai Schmidt-Hoberg



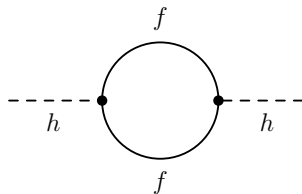
based on

1205.1509, 1208.1683, 1211.2835

DESY Theory Seminar

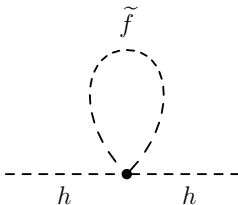
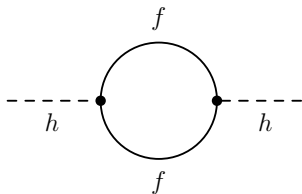
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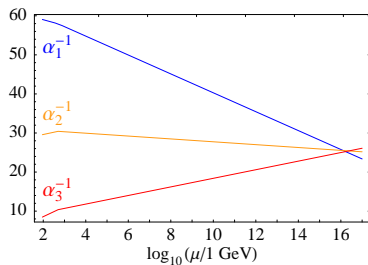
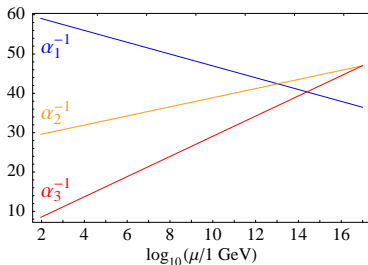
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Renewed interest in the NMSSM

MSSM superpotential

- The gauge-invariant superpotential terms of the MSSM include

$$\begin{aligned}\mathcal{W} = & \mu H_u H_d + \kappa_i L_i H_u \\ & + Y_e^{ij} H_d L_i E_j^c + Y_d^{ij} H_d Q_i D_j^c + Y_u^{ij} H_u Q_i U_j^c \\ & + \lambda_{ijk}^{(0)} L_i L_j E_k^c + \lambda_{ijk}^{(1)} L_i Q_j D_k^c + \lambda_{ijk}^{(2)} U_i^c D_j^c D_k^c \\ & + \kappa_{ij}^{(0)} H_u L_i H_u L_j + \kappa_{ijk\ell}^{(1)} Q_i Q_j Q_k L_\ell + \kappa_{ijk\ell}^{(2)} U_i^c U_j^c D_k^c E_\ell^c\end{aligned}$$

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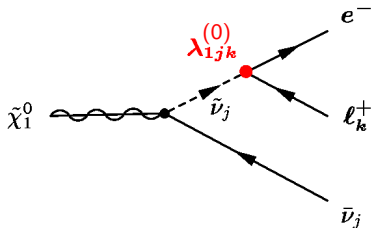
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- μ problem $\mu \sim M_{\text{EW}} \ll M_{\text{P}}$
- No stable dark matter candidate

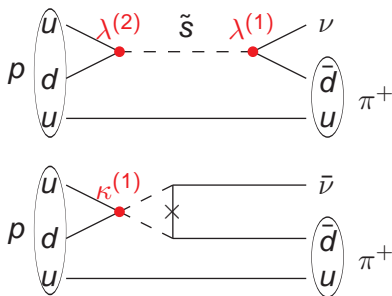


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- Proton decay



R parity

Standard symmetry in the MSSM:

\mathbb{Z}_2 **R parity** Farrar & Fayet; Dimopoulos, Raby, Wilczek

$$\begin{aligned}\mathcal{W} = & \mu H_u H_d + \cancel{\kappa_i} L_i H_u \\ & + Y_e^{ij} H_d L_i E_j^c + Y_d^{ij} H_d Q_i D_j^c + Y_u^{ij} H_u Q_i U_j^c \\ & + \cancel{\lambda_{ijk}^{(0)}} L_i L_j E_k^c + \cancel{\lambda_{ijk}^{(1)}} L_i Q_j D_k^c + \cancel{\lambda_{ijk}^{(2)}} U_i^c D_j^c D_k^c \\ & + \kappa_{ij}^{(0)} H_u L_i H_u L_j + \kappa_{ijkl}^{(1)} Q_i Q_j Q_k L_\ell + \kappa_{ijkl}^{(2)} U_i^c U_j^c D_k^c E_\ell^c\end{aligned}$$

- Experiment: $\kappa_{ijkl} < \frac{10^{-8}}{M_P}$
- By far not sufficient to make the proton stable!
- does not solve the μ problem...

NMSSM

- MSSM plus extra singlet S
- Standard symmetry for the NMSSM: global $\mathbb{Z}_3 + R$ parity

$$\mathcal{W}_{\text{NMSSM}} = \mathcal{W}_{\text{MSSM}}^{\mu=0} + \lambda S H_u H_d + \frac{1}{3} \kappa S^3$$

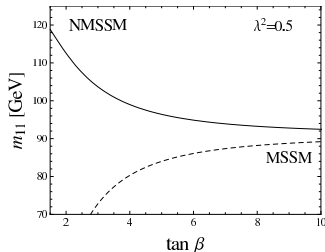
- Original motivation: Solve the μ problem: $\mu_{\text{eff}} = \lambda \langle S \rangle$

Nice feature:

Additional tree-level contribution to Higgs mass:

$$M_Z^2 \cos^2(2\beta) + \lambda^2 v^2 \sin^2(2\beta) + \text{radiative corrections}$$

- Upper bound only...
- perturbativity bound:
 $\lambda \lesssim 0.7$



Problems of the NMSSM

Problems:

- \mathbb{Z}_3 symmetry spontaneously broken once the Higgs fields acquire a vev
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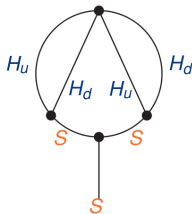
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- global symmetry not necessarily respected by gravity - Planck suppressed operators?

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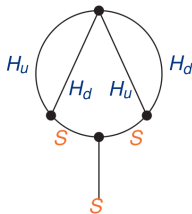


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Abel



The \mathbb{Z}_3 symmetric NMSSM is (basically) excluded!

Need another underlying symmetry!



arXiv:1009.0905

arXiv:1102.3595

N	\mathcal{W}	q_{10}	$q_{\overline{5}}$	q_{H_u}	q_{H_d}	q_S
4	2	1	1	0	0	2
8	2	1	5	0	4	6

The \mathbb{Z}_4^R case: MSSM part

$$\begin{aligned}
 \mathcal{W} = & \cancel{\kappa} \underbrace{H_u H_d}_0 + \cancel{\kappa_i} \underbrace{L_i H_u}_1 \\
 & + Y_e^{ij} \underbrace{H_d L_i E_j^c}_2 + Y_d^{ij} \underbrace{H_d Q_i D_j^c}_2 + Y_u^{ij} \underbrace{H_u Q_i U_j^c}_2 \\
 & + \cancel{\lambda_{ijk}^{(0)}} \underbrace{L_i L_j E_k^c}_3 + \cancel{\lambda_{ijk}^{(1)}} \underbrace{L_i Q_j D_k^c}_3 + \cancel{\lambda_{ijk}^{(2)}} \underbrace{U_i^c D_j^c D_k^c}_3 \\
 & + \kappa_{ij}^{(0)} \underbrace{H_u L_i H_u L_j}_2 + \cancel{\kappa_{ijk\ell}^{(1)}} \underbrace{Q_i Q_j Q_k L_\ell}_0 + \cancel{\kappa_{ijk\ell}^{(2)}} \underbrace{U_i^c U_j^c D_k^c E_\ell^c}_0
 \end{aligned}$$

The \mathbb{Z}_4^R case: MSSM part

$$\begin{aligned}
 \mathcal{W} = & \frac{1}{M_{\text{P}}^2} \underbrace{\langle \mathcal{W} \rangle H_u H_d}_2 + \cancel{\kappa_i} \underbrace{L_i H_u}_1 \\
 & + Y_e^{ij} \underbrace{H_d L_i E_j^c}_2 + Y_d^{ij} \underbrace{H_d Q_i D_j^c}_2 + Y_u^{ij} \underbrace{H_u Q_i U_j^c}_2 \\
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 \end{aligned}$$

- need $\langle \mathcal{W} \rangle \sim m_{3/2} M_{\text{P}}^2$ to cancel cosmological constant
- $\mu \sim \langle \mathcal{W} \rangle / M_{\text{P}}^2 \sim m_{3/2}$
- proton decay operators suppressed by $\langle \mathcal{W} \rangle / M_{\text{P}}^4 \sim 10^{-15} / M_{\text{P}}$
- matter parity exact

Singlet part

- Allowed superpotential terms for $\mathbb{Z}_{4/8}^R$

$$\mathcal{W}_{\mathbb{Z}_{4/8}^R} = \mathcal{W}_{\text{MSSM}}^{\mu=0} + \lambda S H_u H_d + \kappa S^3 + M^2 S$$

\mathbb{Z}_4^R : reintroduces hierarchy problem if $M \gtrsim M_{EW}$

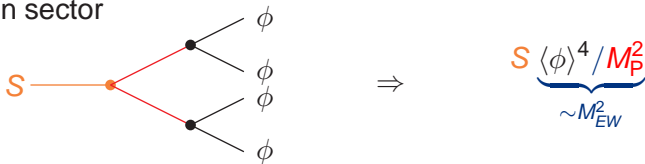
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\mathbb{Z}_4^R : reintroduces hierarchy problem if $M \gtrsim M_{EW}$

- assume only trilinear couplings in \mathcal{W} and intermediate VEVs $\langle \phi \rangle$ in hidden sector



- However: linear term can also be generated radiatively
- divergent tadpoles arise from even (odd) terms in the super (Kähler) potential Abel

\Rightarrow Hierarchy problem not reintroduced

What about the domain wall problem?

$$\langle \mathcal{W} \rangle \sim m_{3/2} M_{\text{P}}^2 \sim 10^{13} \text{ GeV}$$

\Rightarrow corresponding domain walls inflated away

$$\begin{aligned} \Delta \mathcal{W}_{\mathbb{Z}_4^R} &= \langle \mathcal{W} \rangle + \langle \mathcal{W} \rangle^2 \mathcal{S} + \langle \mathcal{W} \rangle \mathcal{S}^2 + \langle \mathcal{W} \rangle H_u H_d \\ &\sim m_{3/2} M_{\text{P}}^2 + m_{3/2}^2 \mathcal{S} + m_{3/2} \mathcal{S}^2 + m_{3/2} H_u H_d \end{aligned}$$

$$\mathcal{W}_{\text{GNMSSM}} = \mathcal{W}_{\text{NMSSM}} + \mu H_u H_d + \frac{1}{2} \mu_s \mathcal{S}^2$$

- Phenomenological implications?

The GNMSSM

- Upper bound on the lightest Higgs mass same as in NMSSM
- Difference? Look at naturalness
- Correct electroweak symmetry breaking requires (MSSM case):

$$m_Z^2 = -2(m_{h_u}^2 + |\mu|^2) + \mathcal{O}(1/\tan\beta^2)$$

- Cancellations of unrelated parameters required
- Standard definition of fine-tuning Ellis et al; Barbieri, Giudice

$$\Delta \equiv \max \text{Abs}[\Delta_p], \quad \Delta_p \equiv \frac{\partial \ln v^2}{\partial \ln p} = \frac{p}{v^2} \frac{\partial v^2}{\partial p}$$

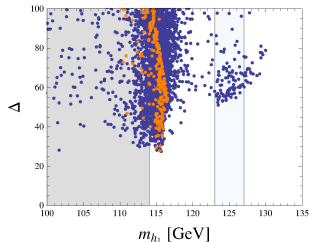
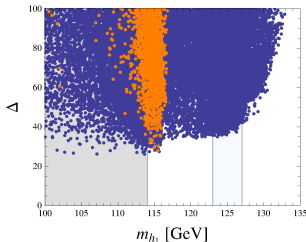
- p is ‘fundamental’ parameter of the theory

Fine-Tuning in the GNMSSM

- Outcome depends on what is the 'fundamental' theory, i.e. what is assumed for SUSY breaking
- Simplest case: 'CGNMSSM' with $m_{1/2}, m_0, A_0, \tan \beta, \lambda, \kappa, v_S, \mu_S$.
- Full 'state of the art' evaluation with *SPHENO*

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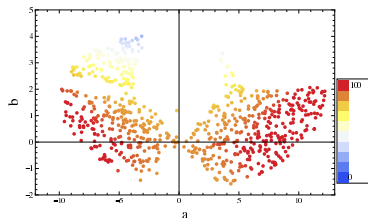
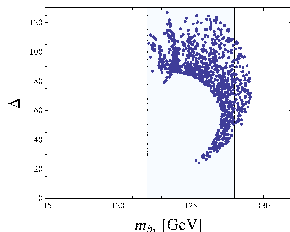
- Pheno: MSSM like with heavier Higgs, singlets rather heavy.
- stau coannihilation region for relic abundance

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- Other option - allow for non-universal gaugino masses with
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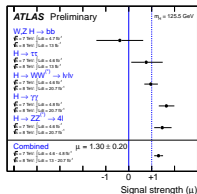
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- more freedom in neutralino sector: DM and compressed spectra
- Low fine-tuning requires small $\tan \beta$ and large λ !
- What are the implications for the Higgs sector?

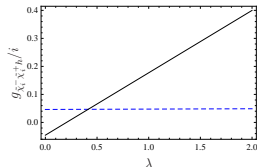
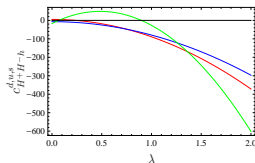
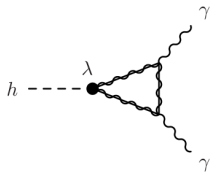
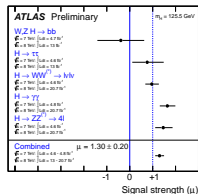
Enhanced diphoton rate

- Two easy options to enhance the di-photon rate
 - suppress partial width into bb to increase BR
 - light** charged particles in the loop to increase partial width into $\gamma\gamma$
- only option in MSSM: 'light stau scenario' Carena et al
- needs very large $\mu \cdot \tan \beta \Rightarrow$ unnatural, danger of CCB vacuum
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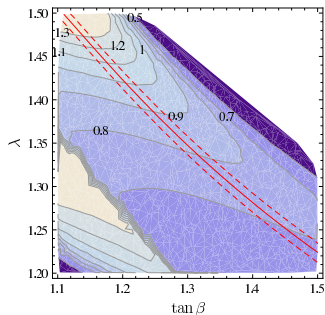
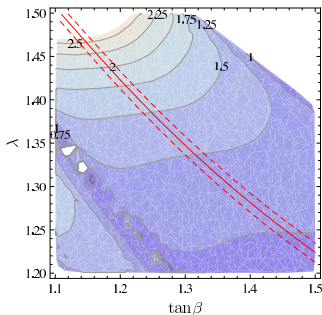
- Enhancement via charginos requires (small) singlet component

Full Analysis

- Potentially very large effects for large λ
- Use low energy variables to see what is possible
- take into account the Higgs mass, production cross section, ...

Full Analysis

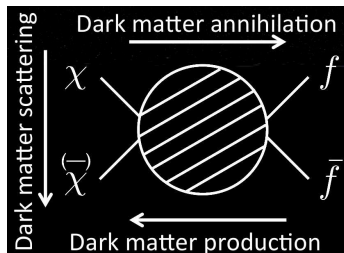
- Potentially very large effects for large λ
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- Enhancement of the diphoton signal by a factor 2 with SM like WW, the correct Higgs mass etc. possible
- main enhancement due to light charginos (higgsinos)...

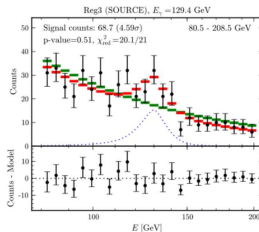
Dark matter indirect detection

- Indirect detection experiments look for the products of DM annihilation in regions of high DM density (e.g. the galactic center)
- Photons particularly interesting - preserve spatial information
- Photon line would be smoking gun for DM



Fermi LAT

- DM annihilation seen by Fermi?



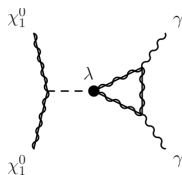
Weniger

Significance of the signal: around 4σ

Inferred dark matter properties:

- $m_\chi \sim 130 \text{ GeV}$ (135 GeV)
- large annihilation cross section into photons, $\langle\sigma v\rangle \sim 10^{-27} \text{ cm}^3/\text{s}$
- Any particle physics models?

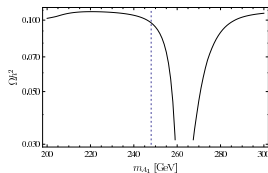
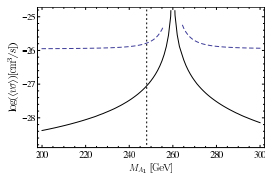
- GNMSSM: charginos show up in similar diagrams as before...



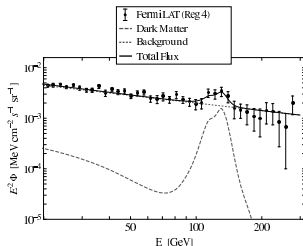
- Picture: mainly singlino neutralino, mainly singlet CP odd Higgs
- NOT possible in the usual NMSSM! Das, Ellwanger
- GNMSSM: possible to realise 130 GeV DM with correct $\langle\sigma v\rangle_{\gamma\gamma}$
 - with correct relic abundance
 - consistent with continuum photon flux
 - consistent with direct detection constraints
 - consistent with electroweak observables
 - 125 GeV Higgs with slightly enhanced diphoton rate

How fine-tuned?

- how much on resonance exactly?



- No cancellations required for singlet like CP odd scalar
- Observational picture...



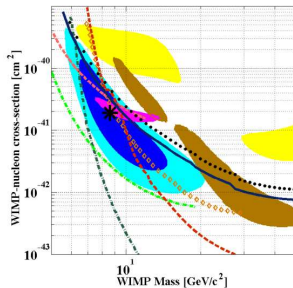
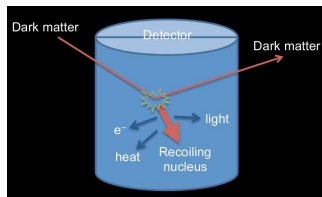
SUSY Summary

- 125 GeV Higgs makes MSSM singlet extensions very interesting
- \mathbb{Z}_3 -symmetric NMSSM (basically) excluded
- Interesting NMSSM like structure with additional mass terms based on discrete R symmetry
- Reasonable fine-tuning for measured Higgs mass, correct relic density,...
- Small fine-tuning corresponds to large λ
- Large λ allows for significant enhancement of $h \rightarrow \gamma\gamma$ while keeping $h \rightarrow WW$ SM like
- Also allows for a simultaneous explanation of the Fermi line

Still work to be done - join in!

Dark matter in CDMS-Si?

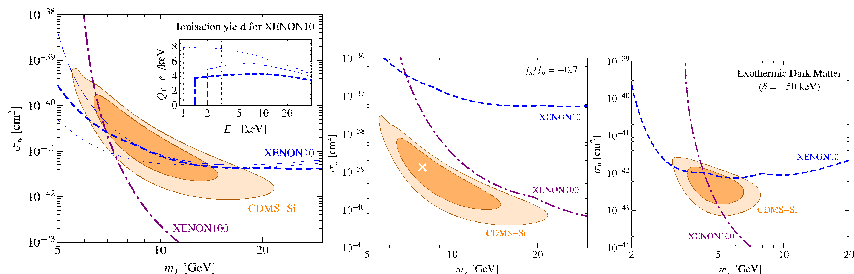
- Recently at the direct detection front...



- CDMS-Si has observed three events with a background estimate of < 0.62 events 1304.4279
- Interpretation in terms of dark matter preferred with about 3σ
- In tension with bounds from XENON10/100

Dark matter in CDMS-Si?

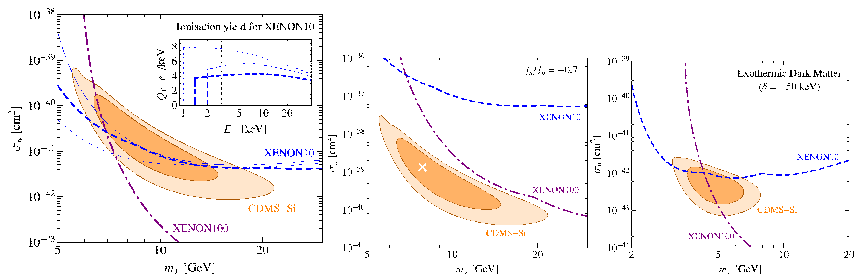
- We find significantly weaker bounds from XENON10 1304.6066
- In contact with the XENON collaboration...



- Take home message: There is a dark matter signal which is not ruled out even with standard assumptions

Dark matter in CDMS-Si?

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Thank You!