

Hidden SUSY from Gauge Coupling Unification

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in collaboration with S. Krippendorf, H.P. Nilles, M. Ratz

arXiv:1306.0574 (2013)

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- BUT: heavy sparticles reintroduce fine-tuning

$$\Delta M_{H_u}^2|_{\text{stop}} \propto -y_t^2 m_t^2 \log\left(\frac{\Lambda}{\text{TeV}}\right)$$

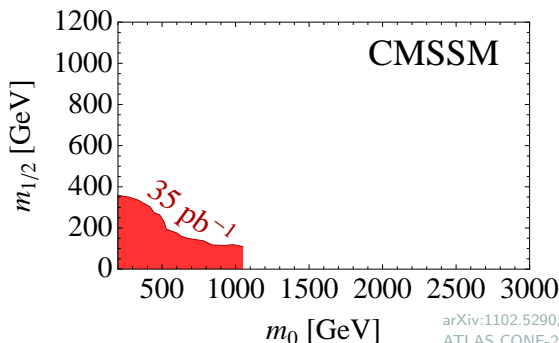
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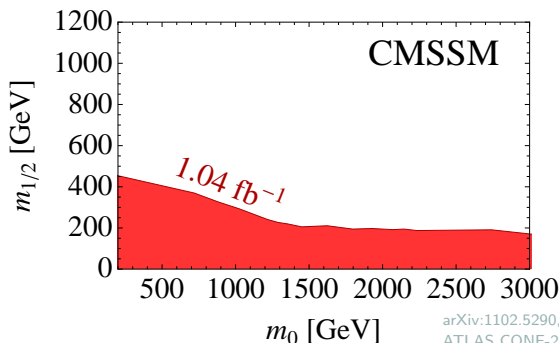
arXiv:1102.5290, arXiv:1109.6572
ATLAS CONF-2013-047

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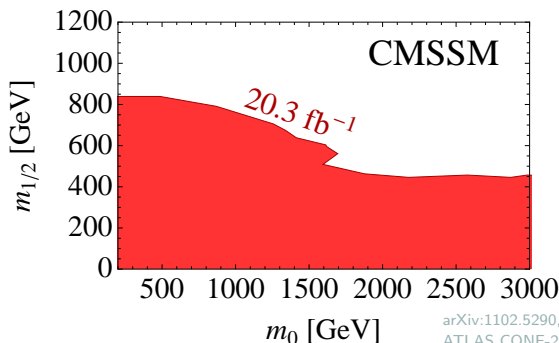
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Dark Matter

- Conserved R-parity \hookrightarrow dark matter candidate
- thermal production

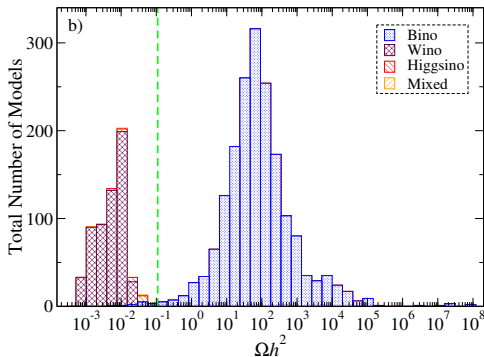
$$\Omega h^2 \simeq 0.1 \frac{\text{pb}}{\langle \sigma v \rangle}$$

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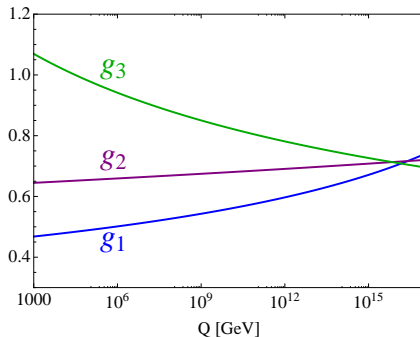
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Baer et al., JHEP 1010 (2010)

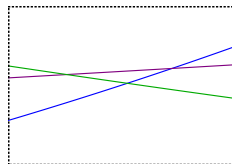
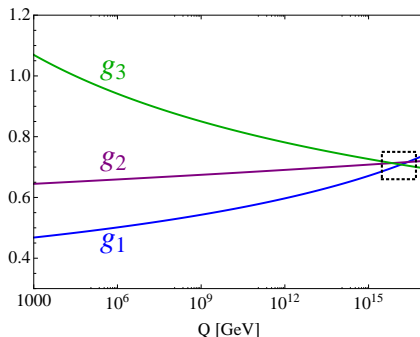
Gauge coupling unification

- gauge coupling unification strong motivation for Supersymmetry
- running gauge couplings in the MSSM



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- BUT: g_i miss by a few per cent in minimal models

- 1 Analytic discussion of gauge coupling unification
- 2 Precision gauge unification in realistic models
- 3 Phenomenology: Implications for LHC, dark matter

Gauge couplings at the high scale

$$\frac{1}{g_i^2(M_{\text{GUT}})} = \frac{1}{g_i^2(M_Z)} - \frac{b_i}{8\pi^2} \ln\left(\frac{M_{\text{GUT}}}{M_Z}\right) + \frac{1}{g_{i,\text{Thr}}^2} + \dots$$

Diagram illustrating the renormalization group equation for gauge couplings at the high scale. The equation is shown with terms in blue ovals. Arrows point from each term to its physical interpretation:

- $\frac{1}{g_i^2(M_Z)}$: gauge couplings at M_Z
- $-\frac{b_i}{8\pi^2} \ln\left(\frac{M_{\text{GUT}}}{M_Z}\right)$: one-loop running, with $b_i = \begin{cases} 33/5 \\ 1 \\ -3 \end{cases}$
- $\frac{1}{g_{i,\text{Thr}}^2}$: thresholds
- \dots : higher orders

- thresholds:

- heavy Standard Model fields (top, Higgs)
- supersymmetric particles
- GUT thresholds

- In this talk: gauge coupling unification **without GUT thresholds**

- Motivation:

↪ candidates for UV models with PGU

Hebecker, JHEP **0401** (2004), Hebecker et al., Nucl. Phys. **B713** (2005)

Błaszczak et al., Phys. Lett. **B683** (2010), Raby et al., Nucl. Phys. **B868** (2013)

↪ proton decay

↪ bottom up

- MSSM thresholds:

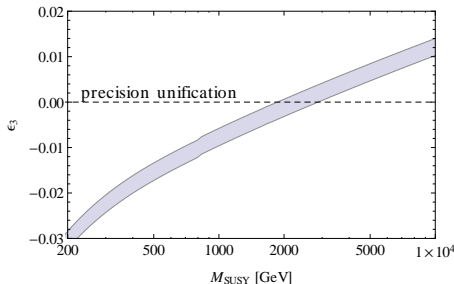
$$\frac{1}{g_{i,\text{Thr}}^2} = \sum_{\eta} \frac{b_i^{\eta}}{8\pi^2} \ln \left(\frac{M_{\eta}}{M_Z} \right) ,$$

b_i^{η} : contribution of field η to beta function coefficient b_i

A simple example

- Simple example: MSSM superpartners + heavy Higgs at M_{SUSY}

$$\frac{1}{g_{i,\text{Thr}}^2} = \frac{b_i - b_i^{\text{SM}}}{8\pi^2} \ln \left(\frac{M_{\text{SUSY}}}{M_Z} \right)$$

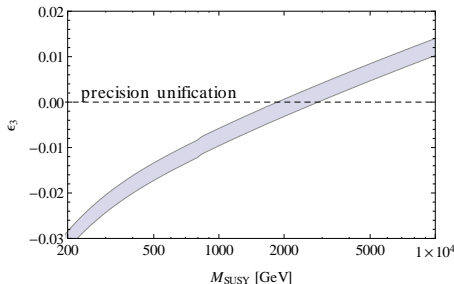


$$\epsilon_3 = \left. \frac{g_3^2 - g_{1,2}^2}{g_{1,2}^2} \right|_{M_{\text{GUT}}}$$

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precision unification for $M_{\text{SUSY}} \sim 2 \text{ TeV}$

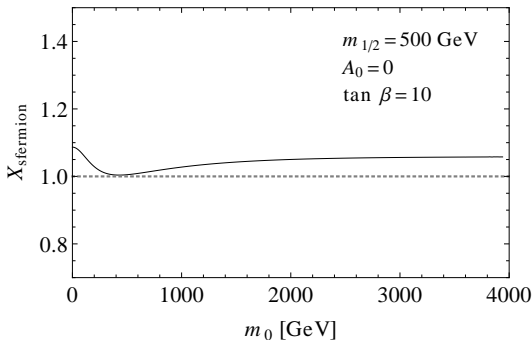
- superpartners not mass-degenerate: define effective M_{SUSY}
Carena et al., Nucl. Phys. **B406** (1993)
- interpretation: same effect on gauge couplings as if all superpartners had a common mass M_{SUSY} (up to changes of g , M_{GUT})

$$M_{\text{SUSY}} = \frac{m_{\tilde{W}}^{32/19} m_{\tilde{h}}^{12/19} m_H^{3/19}}{m_{\tilde{g}}^{28/19}} \underbrace{\prod_{i=1\dots 3} \left(\frac{m_{\tilde{L}_i}^{3/19}}{m_{\tilde{D}_i}^{3/19}} \right) \left(\frac{m_{\tilde{Q}_{Li}}^{7/19}}{m_{\tilde{E}_i}^{2/19} m_{\tilde{U}_i}^{5/19}} \right)}_{X_{\text{sfermion}}}$$

- $X_{\text{sfermion}} = 1$ if sfermions mass-degenerate among SU(5) multiplets
 \hookrightarrow Split SUSY does not destroy gauge unification

Arkani-Hamed et al., JHEP **0506** (2005)

- RGE running splits SU(5) multiplets
- but: effects on gauge coupling unification very small



- precision unification must be achieved in the gaugino / higgsino sector

- RGE running decreases $m_{\widetilde{W}}/m_{\widetilde{g}}$
 \hookrightarrow unfavorable for precision gauge unification
- for models with universal $m_{1/2}$

$$M_{\text{SUSY}} \simeq 0.3 m_{\widetilde{h}} \left(\frac{m_{1/2}^4 m_H^3}{m_{\widetilde{h}}^7} \right)^{1/19}$$

- $M_{\text{SUSY}} \sim 2 \text{ TeV}$ requires super-heavy higgsinos $m_{\widetilde{h}} \sim 10 \text{ TeV}$
- precision unification very unnatural in models with universal gaugino mass

A string-inspired mediation scheme

- Heterotic MiniLandscape

Lebedev et al., Phys. Lett. **B645** (2007), Phys. Lett. **B668** (2008), Phys. Rev. **D77** (2008)

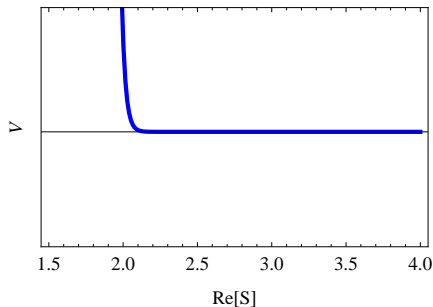
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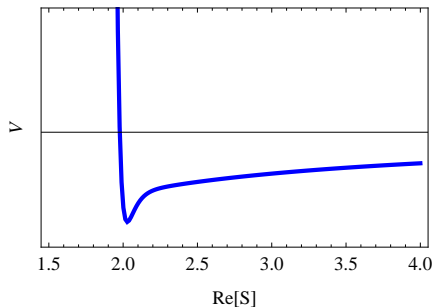
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$$+ \underbrace{\text{const}}$$

e.g. from approx R-symmetries

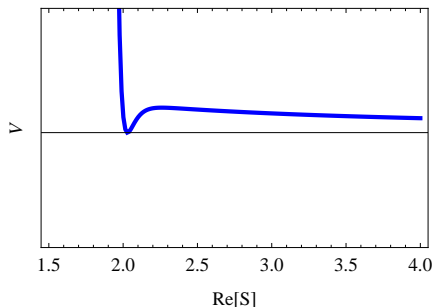
Brümmer et al., JHEP **1004** (2010)

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Brümmer et al., JHEP **1004** (2010)

$$+ \underbrace{W_{\text{up}}}_{\text{uplifting by matter F-term } F_X}$$

Lebedev et al., Phys. Lett. **B636** (2006)

- SUSY breaking pattern fixed by modulus stabilization
- gravitino mass

$$m_{3/2} \sim e^{-bs}$$

- subdominant modulus contribution to SUSY breaking

$$F_S \simeq \frac{F_X}{b} \quad \text{with} \quad b \sim \log \frac{M_P}{m_{3/2}}$$

- in heterotic models gauge kinetic function $f_{ab} \sim s$
- gaugino masses suppressed compared to $m_{3/2}$

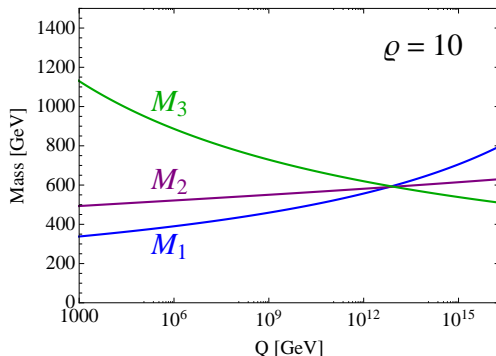
Compressed gaugino masses

- **mirage mediation:** mixed gravity / anomaly mediation

$$M_i = \frac{m_{3/2}}{16\pi^2} (\varrho + b_i g^2)$$

- occurs in various string constructions

KKLT: Choi et al., Nucl. Phys. **B718** (2005), Choi et al. JHEP **0509** (2005), Falkowski et al. JHEP **0511** (2005)
heterotic string: Lowen et al., Phys. Rev. **D77** (2005), Krippendorff et al., Phys. Lett. **B712** (2012)



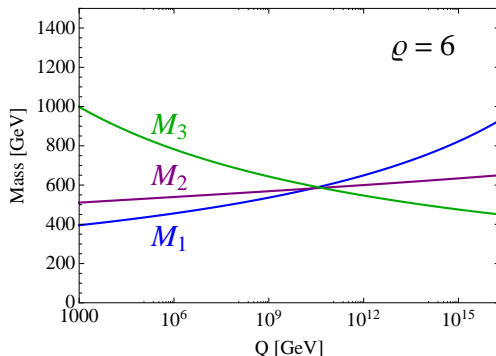
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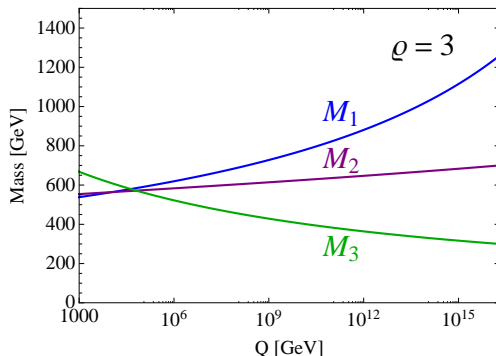
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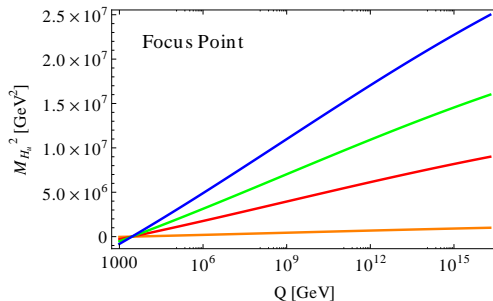
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Scalar masses in mirage mediation

- scalar masses model dependent, but hardly affect gauge unification
- natural mass scale $m_{3/2}$, unless sequestering
 \hookrightarrow hierarchy between gauginos and sfermions
- attractive possibility: universal scalar masses

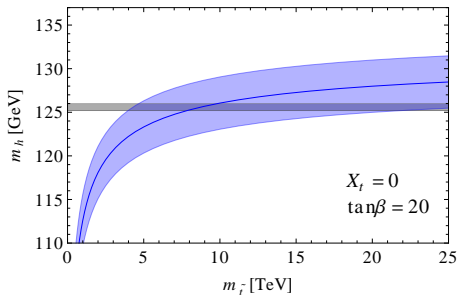


- Multi-TeV scalars not unnatural due to focus point

Feng et al., Phys. Rev. Lett. 84 (1999)

Fixing the scalar mass

- m_0 can be fixed by requiring $m_h = 126$ GeV



- large theoretical uncertainties in m_h

Feng et al., [arXiv:1306.2318](#) (2013)

$$m_0 = 3 \dots 30 \text{ TeV}$$

A comment on electroweak symmetry breaking

- $m_0 \gg m_{1/2}$ potential problem with electroweak symmetry breaking
- heavy right-handed neutrinos give a small threshold Δm^2

Asano et al., Phys. Lett. **B708** (2012), Moroi et al. [arXiv:1305.7357](#) (2013)

$$m_{H_u}^2 = m_0^2 - \Delta m^2 \quad (\text{at } M_{\text{GUT}})$$

allowing for successful EWSB with large universal m_0

- equivalent boundary conditions

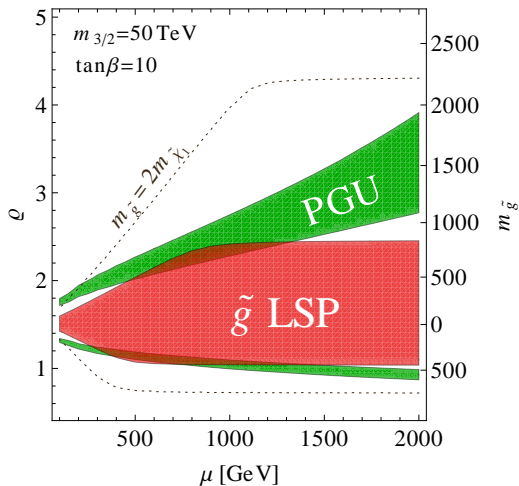
$$m_{\text{sfermion}} = m_H = m_0, \quad \mu \ll m_0$$

- after using the Higgs mass to fix m_0 , the free parameters are

$$m_{3/2}, \quad \varrho, \quad \mu, \quad \tan \beta$$

Precision gauge unification in mirage mediation

- searching regions with precision unification

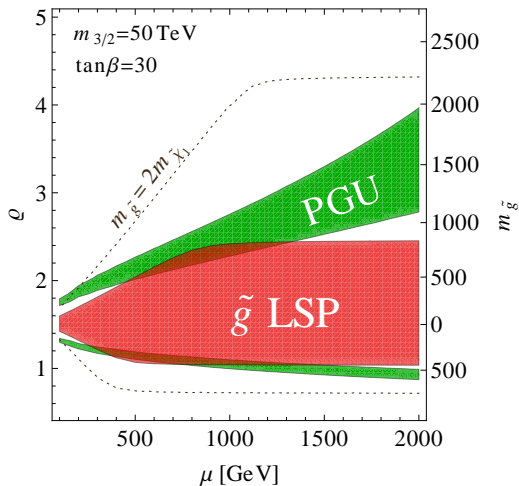


- precision unification with small μ
- **compressed spectrum**

$$m_{\tilde{g}} < 2 m_{\tilde{\chi}_1}$$

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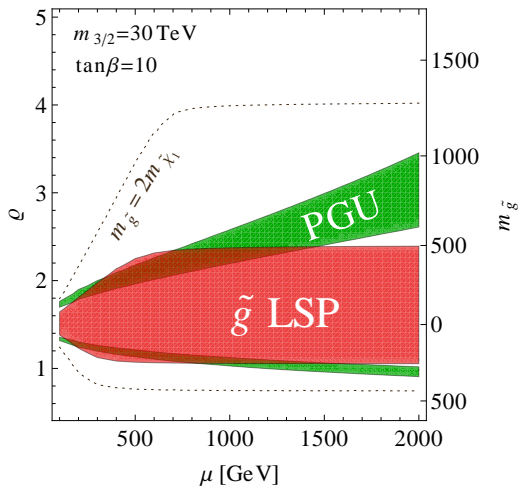


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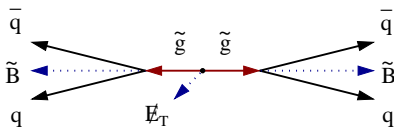
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Compressed spectrum at the LHC

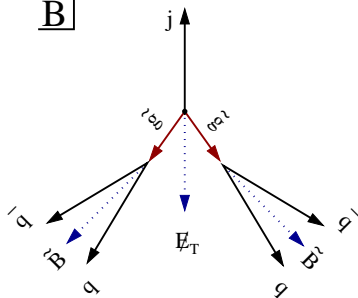
- Small mass splitting $\Delta m = m_{\tilde{g}} - m_{\tilde{\chi}_1}$
- difficult to detect: soft jets, low \cancel{E}_T

LeCompte et al., Phys. Rev. **D84** (2011), Dreiner et al., Europhys. Lett. **99** (2012)

A



B



Alwall et al., Phys. Lett. **B666** (2008)

- initial state radiation required

- simplified models to estimate constraints
- ATLAS, CMS performed searches for

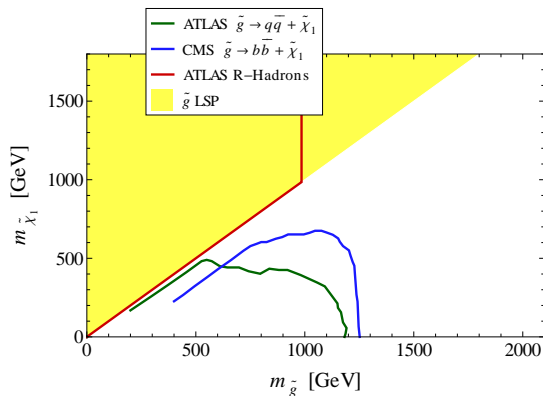
$$\begin{aligned} pp &\rightarrow \tilde{g}\tilde{g} \rightarrow 2 q\bar{q} + 2 \chi_1 && \text{or} \\ pp &\rightarrow \tilde{g}\tilde{g} \rightarrow 2 b\bar{b} + 2 \chi_1 \end{aligned}$$

ATLAS-CONF-2012-109, CMS-PAS-SUS-13-007

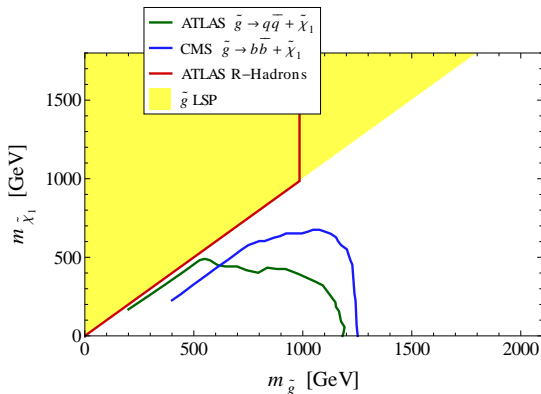
- for gluino LSP: searches for stable R-hadrons

ATLAS collaboration, Phys. Lett. **B720** (2013)

Precision Unification at the LHC



Precision Unification at the LHC

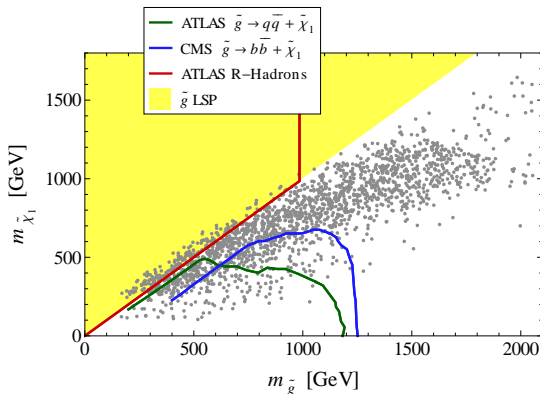


- we generated a large benchmark sample with random

$$\varrho = 0.5-30 \quad m_{3/2} = \frac{40 - 200 \text{ TeV}}{\varrho} \quad \mu = 0.1-2\text{TeV} \quad \tan \beta = 10-50$$

- required precision gauge coupling unification

Precision Unification at the LHC



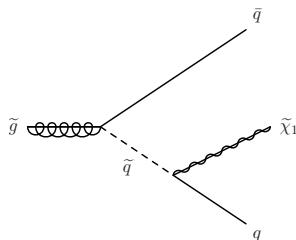
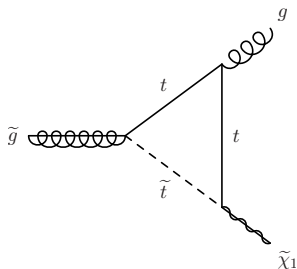
many viable points
with $m_{\tilde{g}} < 1 \text{ TeV}$

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Gluino decays

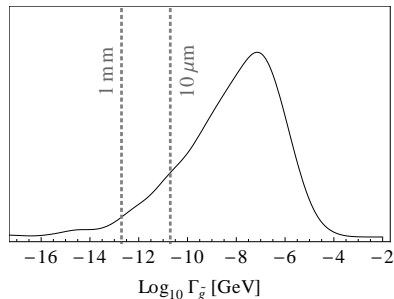
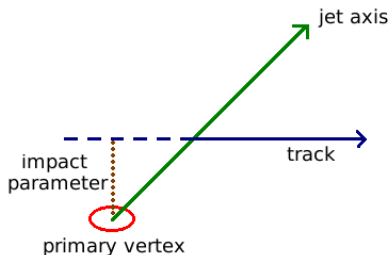


$$\Gamma \propto \text{higgsino fraction} \times \frac{\Delta m^3 m_t^2}{m_{\tilde{q}}^4}$$

$$\Gamma \propto \text{gaugino fraction} \times \frac{\Delta m^5}{m_{\tilde{q}}^4}$$

- gluino decay pattern encodes information about superpartner spectrum
- strong suppression of $\Gamma_{\tilde{g}}$ especially for gaugino-like LSP
- displaced vertices?

Displaced vertices



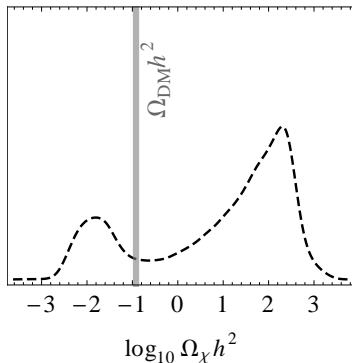
- distribution of $\Gamma_{\tilde{g}}$ among benchmark sample
- $\sim 10\%$ of benchmark points have $c/\Gamma_{\tilde{g}} = 10 \mu\text{m} - 10 \text{ mm}$
- $\Gamma_{\tilde{g}}$ very sensitive to squark sector, could even be larger
- possibly detectable (e.g. transverse impact parameter)
- may affect b -tagging

Dark Matter

- lightest neutralino very good dark matter candidate
- **BUT:** relic density typically **too large (bino)** or **too small (higgsino)**

Baer et al., JHEP **1010** (2010)

- compressed spectrum affects relic density



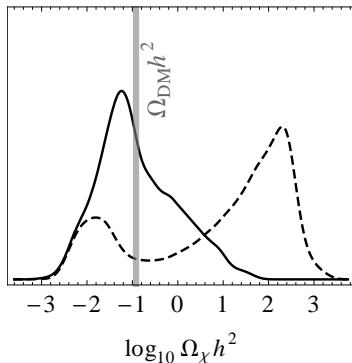
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↪ without PGU (dashed)

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- relic density with mirage boundaries
 - ↪ without PGU (dashed)
 - ↪ with PGU (solid)

- relic neutralino density

$$\Omega h^2 \simeq 0.1 \frac{\text{pb}}{\langle \sigma v \rangle}$$

- the problem:

$$\langle \sigma v \rangle_{\text{bino}} \sim 10^{-3\dots 4} \text{ pb} \quad \langle \sigma v \rangle_{\text{wino}} \sim 10^{+2\dots 3} \text{ pb}$$

- coannihilations for $m_{\tilde{W}} = m_{\tilde{B}} + \Delta m$ with small Δm

$$\begin{aligned} \langle \sigma v \rangle_{\text{eff}} &\simeq B \langle \sigma v \rangle_{\text{wino}} \\ B &\simeq \exp \left(-2 \frac{\Delta m}{T_F} \right) \end{aligned}$$

- PGU typically yields $\Delta m \sim 0.1 m_{\tilde{B}}$, which translates to $B \sim 10^{-2}$

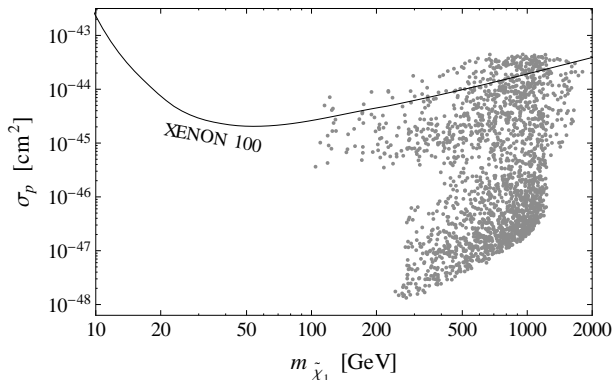
$$\langle \sigma v \rangle_{\text{eff}} \sim 1 \text{ pb}$$

Direct Detection

- neutralino-nucleon interactions dominated by Higgs exchange

$$\sigma_p \propto \text{higgsino-gaugino mixing}$$

- bino + wino coannihilation gives correct Ωh^2 but **tiny** σ_p
 \hookrightarrow hides from direct detection



Conclusion

- Main motivations for Supersymmetry are the **hierarchy problem**, **gauge coupling unification** and **dark matter**
- **But:** **non-observation of superpartners at the LHC**
gauge couplings typically miss by a few per cent
neutralino relic density too small or too large
- In mirage mediation, the reduced ratio $m_{\tilde{g}}/m_{\tilde{W}}$ improves gauge coupling unification
- Mirage mediation + **precision unification** predicts highly compressed gaugino spectrum, small $m_{\tilde{g}} - m_{\tilde{\chi}_1}$
- **LHC bounds relaxed**, $m_{\tilde{g}} \sim 500$ GeV ok
- Neutralino LSP has “automatically” the **correct relic density**