

Global Fits Beyond the SM

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in collaboration with

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based on [PRL 109 (2012) 241802], [JHEP07 (2013) 118]

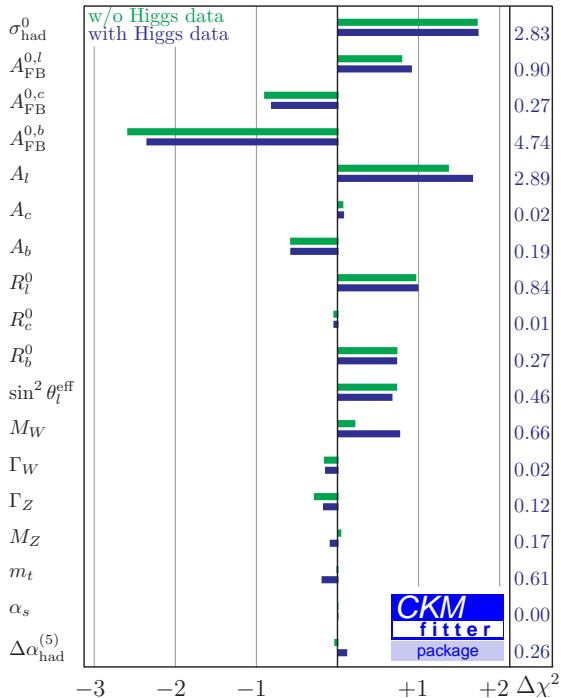


Helmholtz-Alliance Meeting, Karlsruhe, December 2013

SM Electroweak Fit

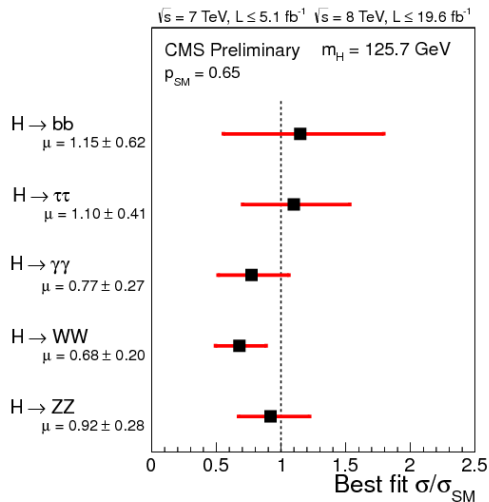
- Fits were used in the LEP era to determine the free parameters of the SM.
- With the Higgs discovery the last free parameter of the SM has been measured directly.
- The SM electroweak fit determines these Parameters from the electroweak precision observables (EWPOs). →
- The influence of this measurement on the SM electroweak fit is relatively small.

[PRL 109 (2012) 241802]



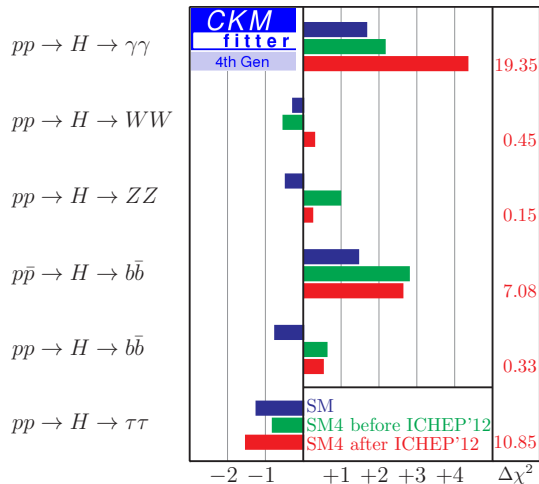
Higgs Signal Strengths

- The EWPOs also **severely constrain** many **models beyond the SM**.
- Constraints from the observed **Higgs signal strengths** can already be **equally powerful**.
- **Combined fits** of both types of data can (and should) be used to study models beyond the SM, but...
- the **statistical interpretation** of the fits can be (conceptually and technically) **less trivial** for models beyond the SM.



The Case of the 4th Generation

- Direct searches push us in a region where the 4G fermions have **large Yukawa couplings**.
- **Virtual corrections** to the Higgs couplings **modify** the Higgs **signal strengths**.
[Denner, Dittmaier, Mück, Passarino, Spira, Sturm, Uccirati, Weber; arXiv:1111.6395]
- **EWPOs** constrain the **mass splittings** inside the $SU(2)$ doublets.
- Combined fit gives $\Delta\chi^2 = 38$.



Chi-squares and p -values

- The usual analytic relation between $\Delta\chi^2$ and the p -value (Wilks' theorem) requires **nested models**.

- The additional fermions of the SM4 **do not decouple**.

⇒ You cannot obtain the SM3 as a limiting case of the SM4.

⇒ The computation of the p -value requires a **very expensive** numerical **simulation**, which is **unfeasible without special simulation methods**.

- These methods were implemented in the public code **myFitter** (<http://myfitter.hepforge.org>) and documented in [CPC 184 (2013) 2438].

- The SM4 is excluded at **5.3 standard deviations**. (Wilks' theorem gives 3.5 standard deviations.)

2HDM Fits

The 2HDM of type II

- Two scalar SU(2) doublets.
- A **softly broken \mathbb{Z}_2 symmetry** which forbids FCNCs.
- No Higgs-sector CP violation.
- Scalar particle content: h, H, A, H^\pm .
- Independent sets of **real parameters** are
 - $v_2/v_1 \equiv \tan\beta, m_{12}^2, \lambda_1, \lambda_2, \lambda_3, \lambda_4, \lambda_5$
 - $\tan\beta, \beta - \alpha, m_{12}^2, m_h, m_H, m_A, m_{H^\pm}$
- This time we have a **decoupling limit**:
 $\beta - \alpha = \pi/2, m_H, m_A, m_{H^\pm} \gg m_h$.
- The map between the two parametrisations **near the decoupling limit** is **not very smooth**.

Theoretical Constraints

- The Higgs potential must be **bounded from below**

$$\lambda_1 > 0 \quad , \quad \lambda_2 > 0 \quad , \quad \lambda_3 > -\sqrt{\lambda_1 \lambda_2} \quad , \\ |\lambda_5| < \lambda_3 + \lambda_4 + \sqrt{\lambda_1 \lambda_2} \quad .$$

- ‘Our’ minimum of the Higgs potential must be the **global minimum**

$$m_{12}^2(m_{11}^2 - m_{22}^2\sqrt{\lambda_1 \lambda_2})(\tan\beta - (\lambda_1/\lambda_2)^{1/4}) > 0 \quad .$$

- The Higgs self-couplings must be **perturbative**:

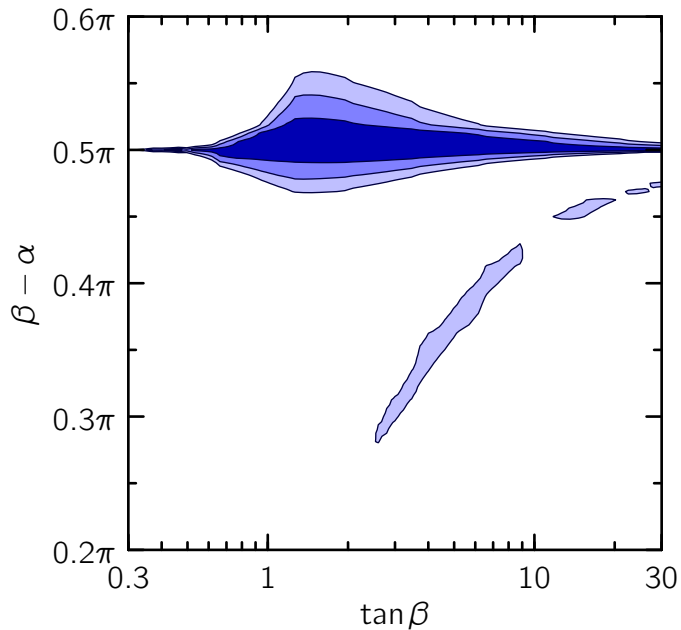
$$|\text{Eigenvalues of } \phi\phi \rightarrow \phi\phi \text{ scattering matrix}| < 2\pi \quad .$$

⇒ Fit requires **optimisation under non-linear constraints**.

Experimental Constraints

- Full set of **electroweak precision observables** (no S , T , U).
- **Signal strengths** of the light Higgs boson (including correlations between different production mechanisms).
- Limits on heavy $H \rightarrow WW, ZZ$ and $H \rightarrow \tau\tau$ resonances.
- **Flavour observables** relevant for the low $\tan\beta$ region:
 Δm_{B_s} and $\text{Br}(\bar{B} \rightarrow X_s \gamma)$.

$\tan\beta$ vs. $\beta - \alpha$



SM Fits

The Fourth
Generation

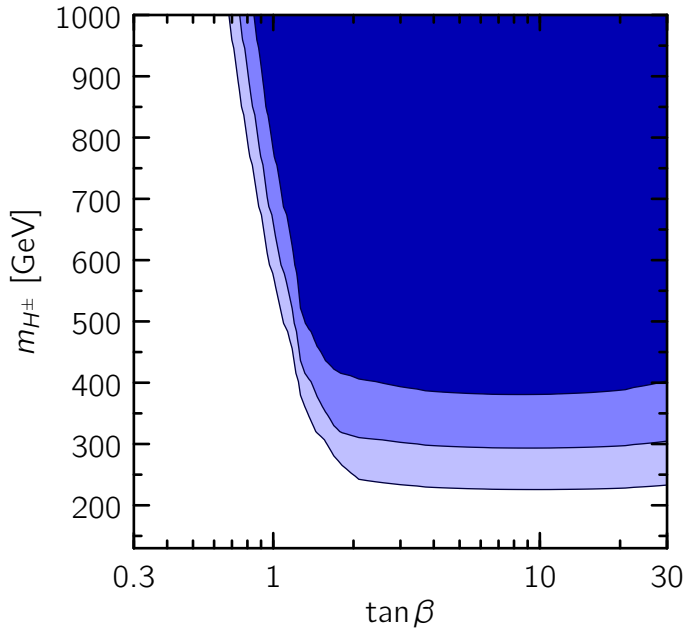
2HDM Fits

Constraints

Fit results

Conclusions

$\tan\beta$ vs. m_{H^\pm}



SM Fits

The Fourth
Generation

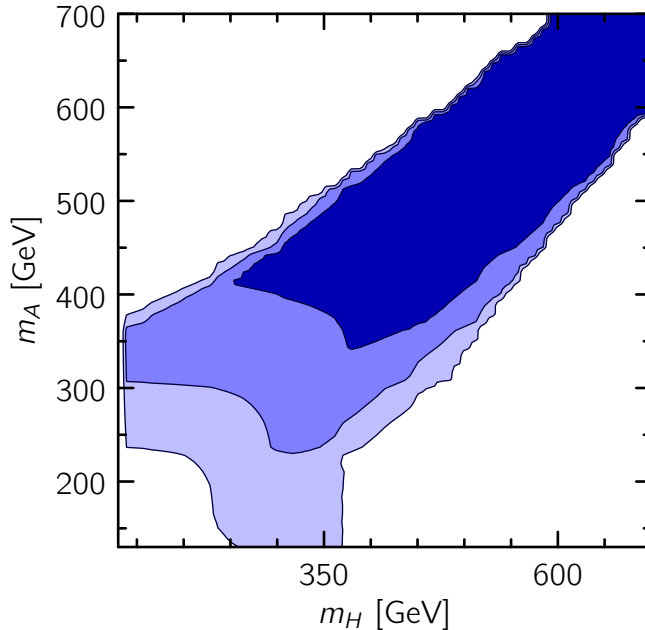
2HDM Fits

Constraints

[Fit results](#)

Conclusions

m_H vs. m_A , m_{H^\pm} free



SM Fits

The Fourth
Generation

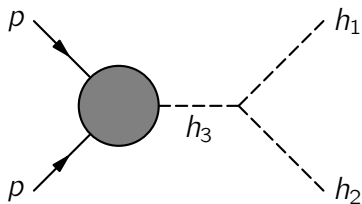
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Triple Higgs Couplings

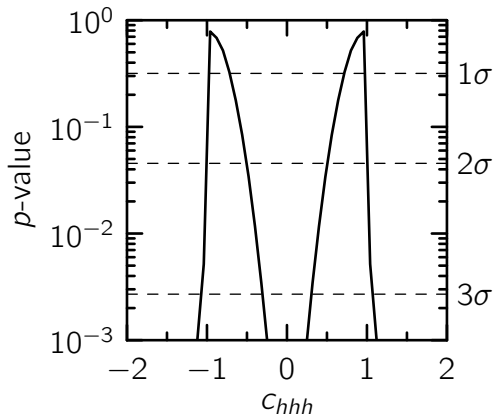


$$C_{h_1 h_2 h_3} = \frac{g_{h_1 h_2 h_3}^{2\text{HDM}}}{g_{HHH}^{\text{SM}}|_{m_H=126 \text{ GeV}}}$$

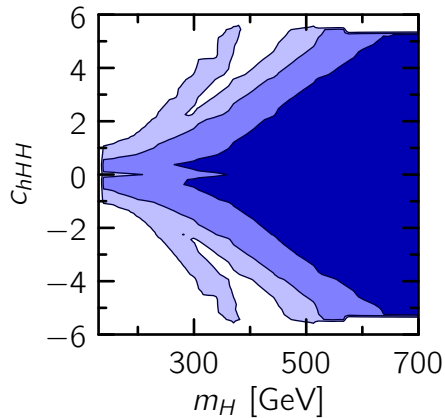
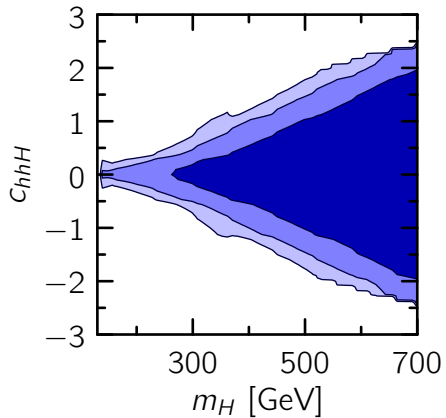
For $\beta - \alpha = \pi/2$:

$$C_{hhh} = 1 \quad ,$$

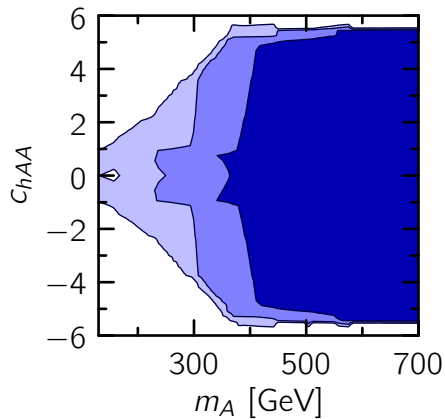
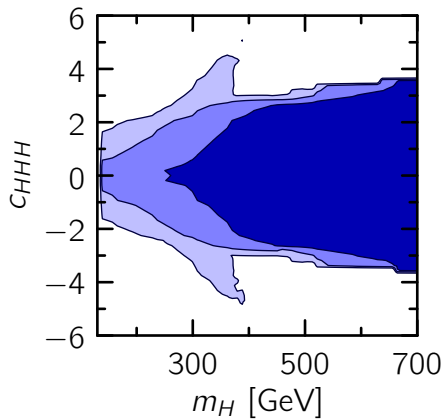
$$C_{hhH} = 0 \quad .$$



Triple Higgs Couplings



Triple Higgs Couplings



Conclusions (SM4)

- SM with a sequential **fourth generation** is **ruled out** by a combination of **Higgs and electroweak precision data**.
- Computation of **p -values** in **non-decoupling models** is **non-trivial** and requires numerical simulations which become unfeasible for small p -values.
- **Importance sampling** techniques as implemented in ***myFitter*** can **speed things up considerably**.

Conclusions (2HDM)

- Best-fit scenario of the type-II 2HDM is the decoupling limit.
- Scenarios with non-SM-like h couplings are allowed by Higgs data but disfavoured (at 2σ) by flavour observables (for tight perturbativity bound).
- Scenarios with m_H and m_A below 300 GeV are ruled out at 2σ .
- hhh coupling can only be reduced.
- $c_{hhH} < 1$ for $m_H \lesssim 500$ GeV.
- $c_{hHH}, c_{HHH} < 2$ for $m_H \lesssim 400$ GeV.
- $c_{hAA} < 5$ for $m_A \lesssim 400$ GeV.
- Be careful with purely scan-based analyses. They don't necessarily give you the full picture.