New results for precision Higgs Physics

Zoltán Trócsányi

University of Debrecen and MTA-DE Particle Physics Research Group

in collaboration with

A. Kardos, M.V. Garzelli

Loops and Legs in Quantum Field Theory, Weimar
April 29, 2014

LHCphenOnet

Precision tools for Higgs Physics with `PowHel`

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Outline

- Motivation

- Method

- Predictions

- Conclusions and Plans

# Motivation

Higgs boson discovered at the LHC

- m_H [GeV]=$125.5 \pm 0.2_{stat} \pm 0.6_{syst}$ (ATLAS 2013)
  $125.7 \pm 0.3_{stat} \pm 0.3_{syst}$ (CMS 2013)

- All measured properties are consistent with SM expectations within experimental uncertainties
  - branching ratios as predicted
  - spin zero
  - parity +
  - couples to masses of W and Z (with c_v=1 within experimental uncertainty)

t-quark: potential tool for discovery

- The t-quark is heavy, Yukawa coupling ~1
  m_t [GeV]=173.34±0.64 (LHC+TeVatron, 2014)
  ($\Rightarrow y_t$=0.997±0.003)

  $\Rightarrow$ plays important role in Higgs physics
  (more tantalizing: $m_t m_Z = (125.7±0.3)^2$ GeV^2)

- y_t cannot be measured in $H \rightarrow t\bar{T}$ decay ($m_t > m_H$)

How to measure y_t?

- $H \to \gamma\gamma$ is sensitive to y_t through t-quark loop,
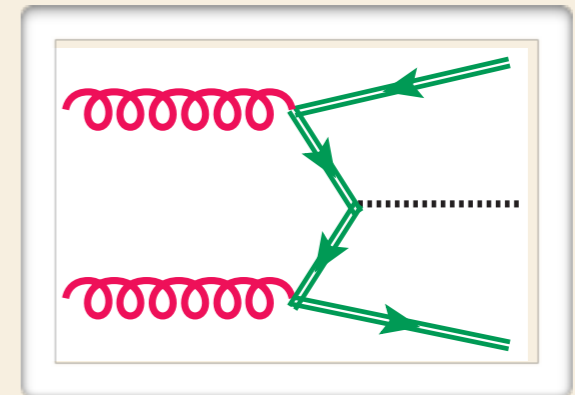
  but rates are small and W loop also contributes
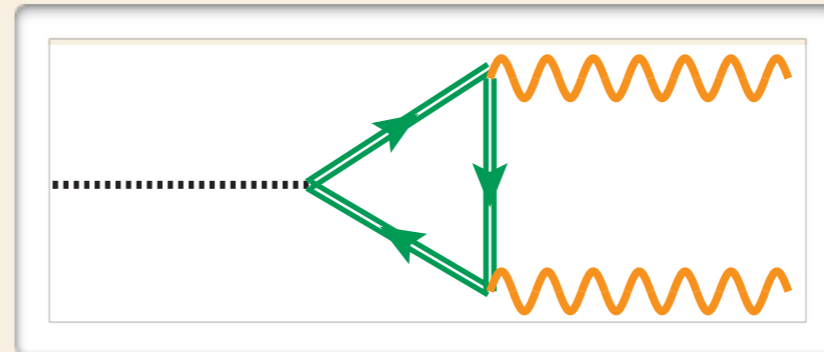
- $gg \to H$ is sensitive to y_t through t-quark loop

  if only SM model particles contribute (so far xsec is consistent with SM)
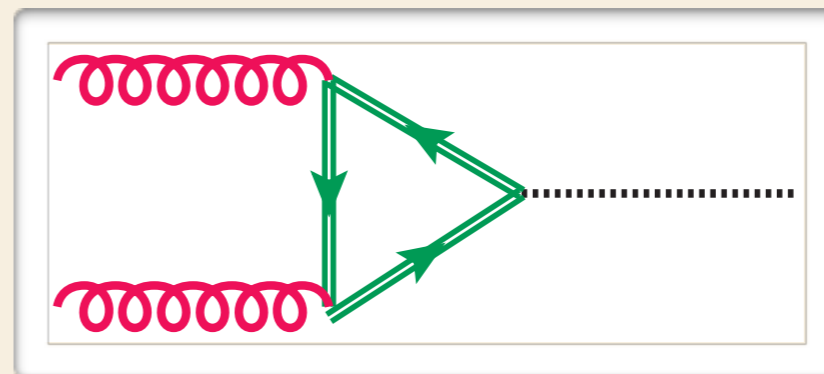
- $gg \to H$ is sensitive to BSM physics

  if y_t is measured separately

tTH hadroproduction

○ y_t can be measured in pp $\rightarrow$ tTH through many decay channels (all very difficult):

- ○ hadrons with single lepton: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}jj, H \rightarrow b\bar{b}$
- ○ hadrons with dilepton: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}\ell\nu, H \rightarrow b\bar{b}$
- ○ hadrons with hadronic tau: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}jj, H \rightarrow \tau_h^+ \tau_h^-$
- ○ diphoton with lepton: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}jj, H \rightarrow \gamma\gamma$
- ○ diphoton with hadrons: $t \rightarrow b\, jj, \bar{t} \rightarrow \bar{b}jj, H \rightarrow \gamma\gamma$
- ○ same sign dilepton: $t \rightarrow b\, jj, \bar{t} \rightarrow \bar{b}jj, H \rightarrow \ell\nu\ell[\nu]$
- ○ 3 leptons with di, trilepton: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}jj, H \rightarrow \ell[\nu]\ell[\nu]$
- ○ 4 lepton with di, trilepton: $t \rightarrow b\ell\nu, \bar{t} \rightarrow \bar{b}\ell[\nu], H \rightarrow \ell[\nu]\ell[\nu]$

The importance of being top

These require precise predictions of distributions at hadron level for
pp →tT+hard X, X = H,W,Z,γ,j,bB,2j...

...with decays: the t-quark is not detected because it decays before hadronization

$$|V_{tb}|^2 \gg |V_{ts}|^2,\ |V_{td}|^2$$
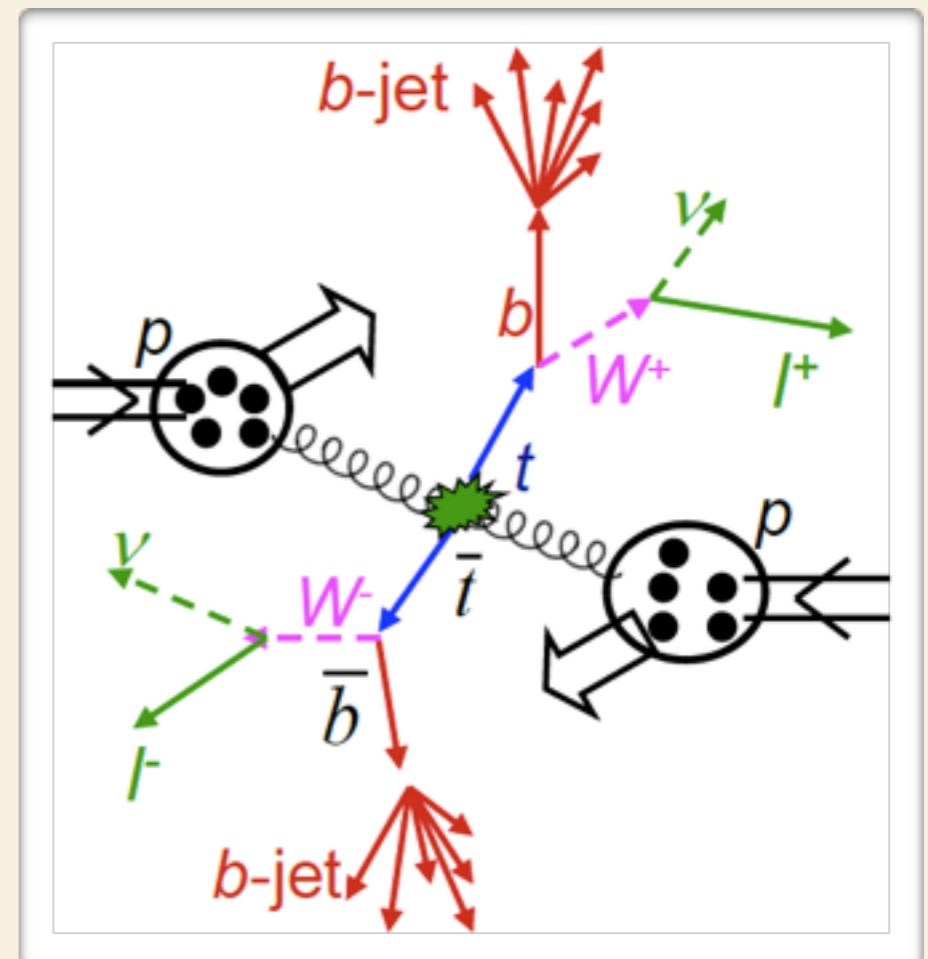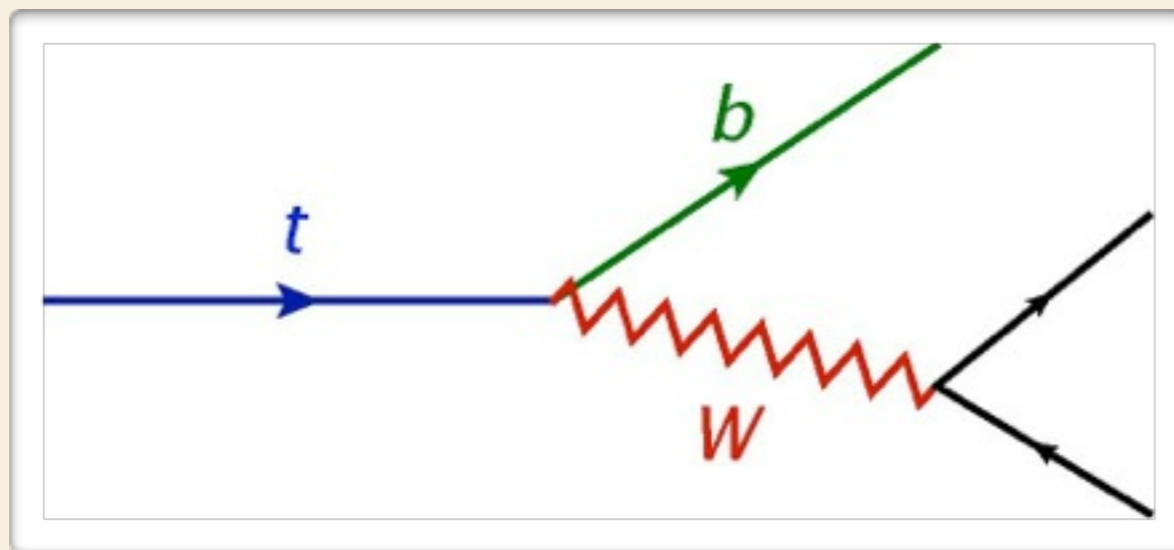
...to distributions, full of pitfalls & difficulties

Method

Cerro Torre Patagonia, courtesy of V Del Duca

There is a long way from loops and legs...

From standard SMC to POWHEG MC

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

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- Standard MC first emission:

$$\mathrm{d}\sigma_{\mathrm{SMC}} = B(\Phi_n)\mathrm{d}\Phi_n \left[\Delta_{\mathrm{SMC}}(t_0) + \Delta_{\mathrm{SMC}}(t) \underbrace{\frac{\alpha_{\mathrm{s}}(t)}{2\pi} \frac{1}{t} P(z)}_{} \Theta(t - t_0)\, \mathrm{d}\Phi_{\mathrm{rad}}^{\mathrm{SMC}} \right]$$

$$\int B(\Phi_n)\mathrm{d}\Phi_n = \sigma_{\mathrm{LO}}$$

$$= \lim_{k_\perp \to 0} R(\Phi_{n+1})/B(\Phi_n)$$

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- POWHEG MC first emission:

$$\mathrm{d}\sigma = \bar{B}(\Phi_n)\mathrm{d}\Phi_n \left[\Delta(\Phi_n, p_\perp^{\mathrm{min}}) + \Delta(\Phi_n, k_\perp)\frac{R(\Phi_{n+1})}{B(\Phi_n)}\Theta(k_\perp - p_\perp^{\mathrm{min}})\mathrm{d}\Phi_{\mathrm{rad}} \right]$$

$$\bar{B}(\Phi_n) = B(\Phi_n) + V(\Phi_n) + \int \left[R(\Phi_{n+1}) - A(\Phi_{n+1}) \right]\mathrm{d}\Phi_{\mathrm{rad}}$$

[Frixione, Nason, Oleari
arXiv: 0709.2092]

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$$\int B(\Phi_n)\mathrm{d}\Phi_n = \sigma_{\mathrm{LO}}$$

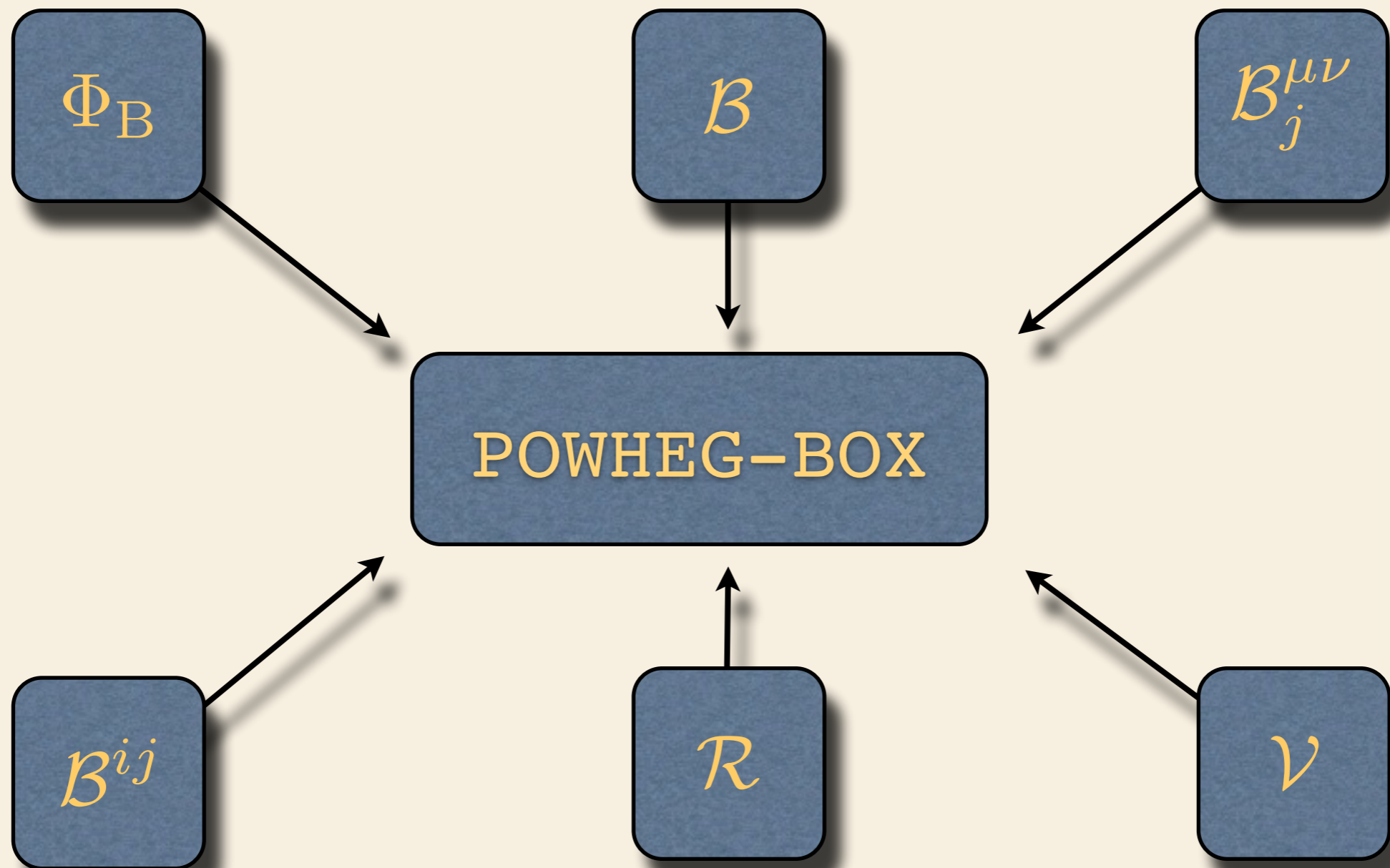- POWHEG MC first emission:

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$$\int \bar{B}(\Phi_n)\mathrm{d}\Phi_n = \sigma_{\mathrm{NLO}}$$

[Frixione, Nason, Oleari arXiv: 0709.2092]

POWHEG-BOX framework

PowHel framework

HELAC-NLO

[Bevilaqua et al, arXiv: 1110.1499]

POWHEG-BOX

[Alioli, Nason, Oleari, Re, arXiv: 1002.2581]

PowHel

RESULT of PowHel:

Les Houches file of Born and Born+1st radiation events (LHE) ready for processing with SMC followed by almost arbitrary experimental analysis

Why should we care about NLO + PS?

- Hadrons in final state
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  event generation is, faster than an NLO computation

  (once the code is ready!)

Why should we care about NLO + PS?

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  ...but we deliver the events on request

tTbB production

Choice of scales

▸ QCD corrections are

  ▸ large with scales $\mu_0 = m_t$ or $m_t + m_{b\bar{b}}/2$ (about 80%)

  ▸ moderate with dynamical scale $\mu_0 = (m_t^2\, p_{T,b} p_{T,\bar{b}})^{1/4}$ (about 25%) (proposed by Bredenstein et al in arXiv:1001.4006), implying better convergence by emulating higher order effects through CKKW-type scale choice

Choice of scales

▸ QCD corrections are

    ▸ large with scales $\mu_{fix} = m_t$ or $m_t + m_{b\bar{b}}/2$ (about 70%)

    ▸ moderate with dynamical scale $\mu_{dyn} = (m_t^2 \, p_{T,b} p_{T,\bar{b}})^{1/4}$ (about 25%) (proposed by Bredenstein et al in arXiv:1001.4006), implying better convergence by emulating higher order effects through CKKW-type scale choice,

<div align="center">but</div>

    ▸ we simulate higher order effects through the PS: μ_{dyn} is too small near threshold where cross section is largest, even for a b with $p_T = 100$ GeV and another b with $p_T = 20$ GeV $\mu_{dyn} = 90$ GeV $\ll m_t$ resulting in an artificially large xsection at LO
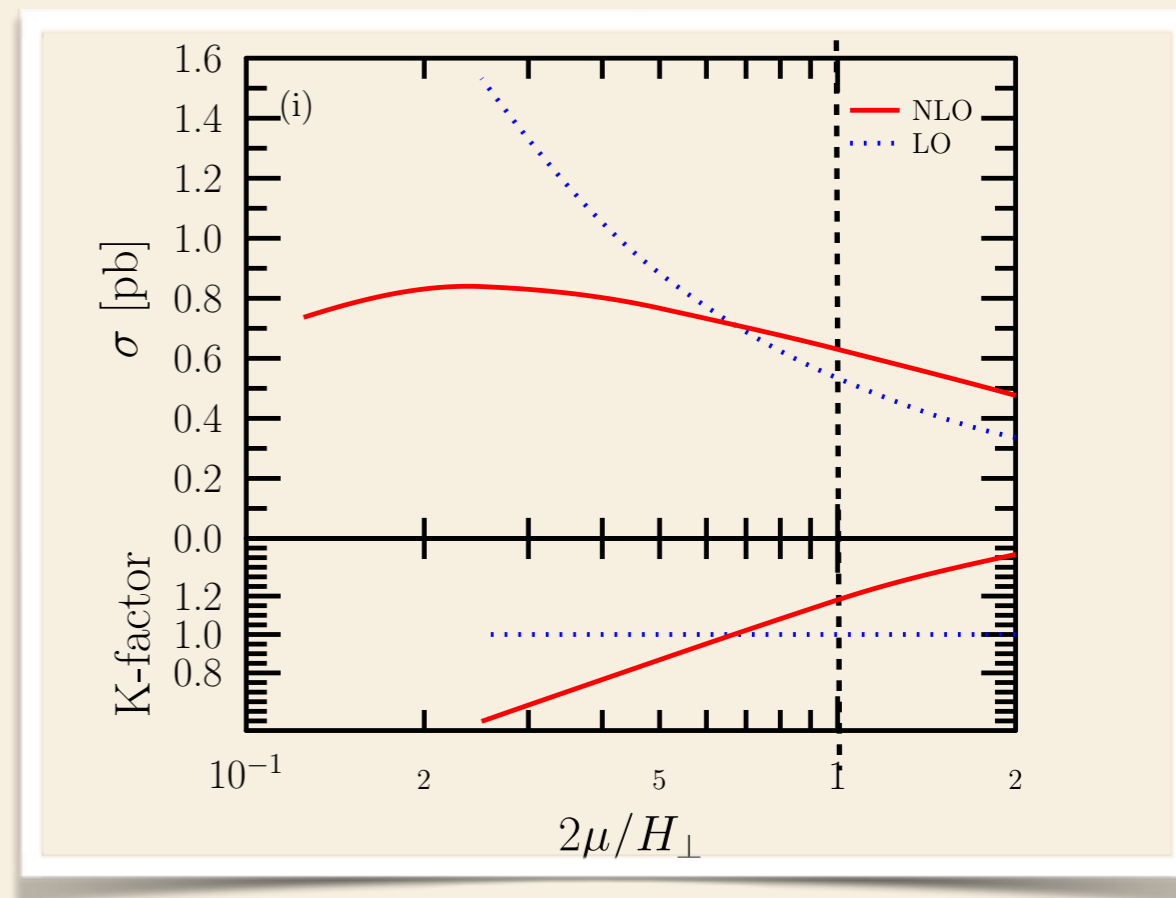
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With this scale

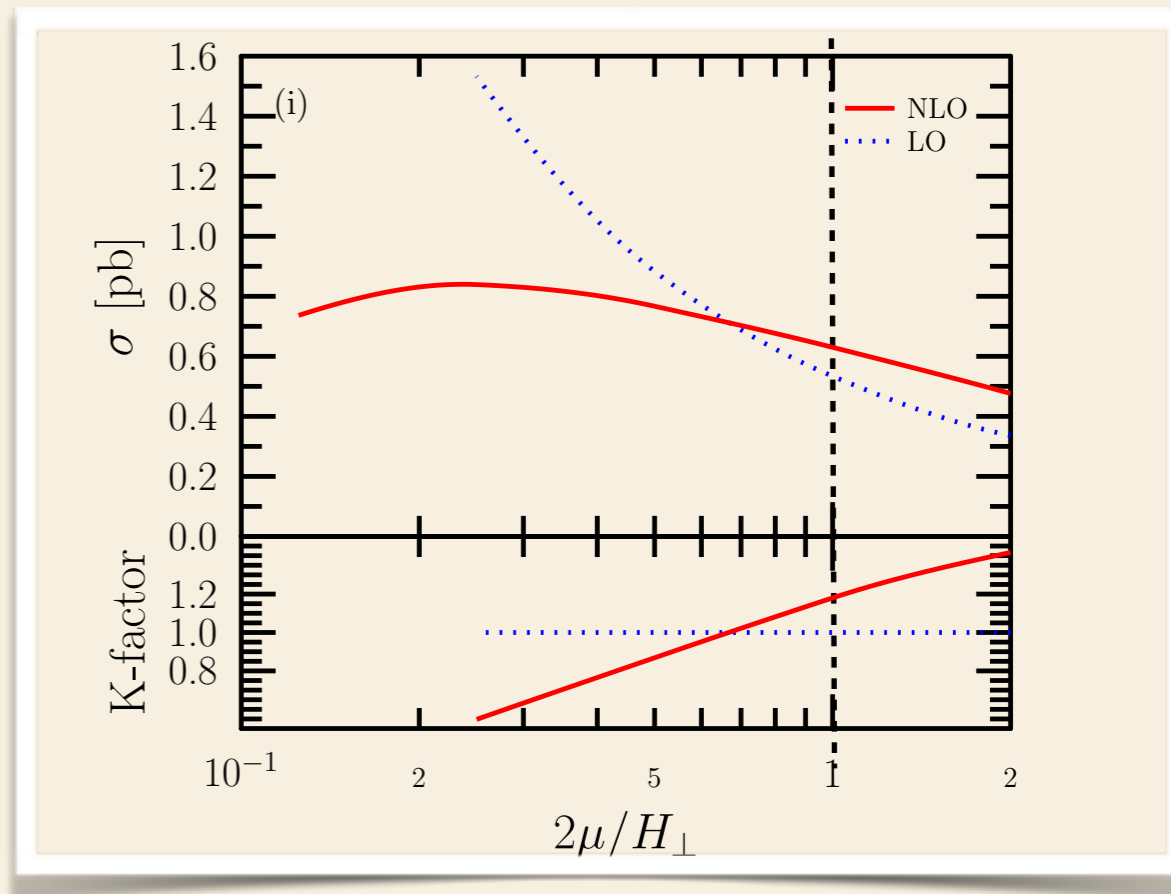✓ the K factor is even smaller, implying good convergence

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With this scale

✓ the K factor is even smaller, implying good convergence

✓ the cross sections are smaller (with BDDP cuts):

σ_{LO} = 534 fb, σ_{NLO} = 630 fb, K = 1.18

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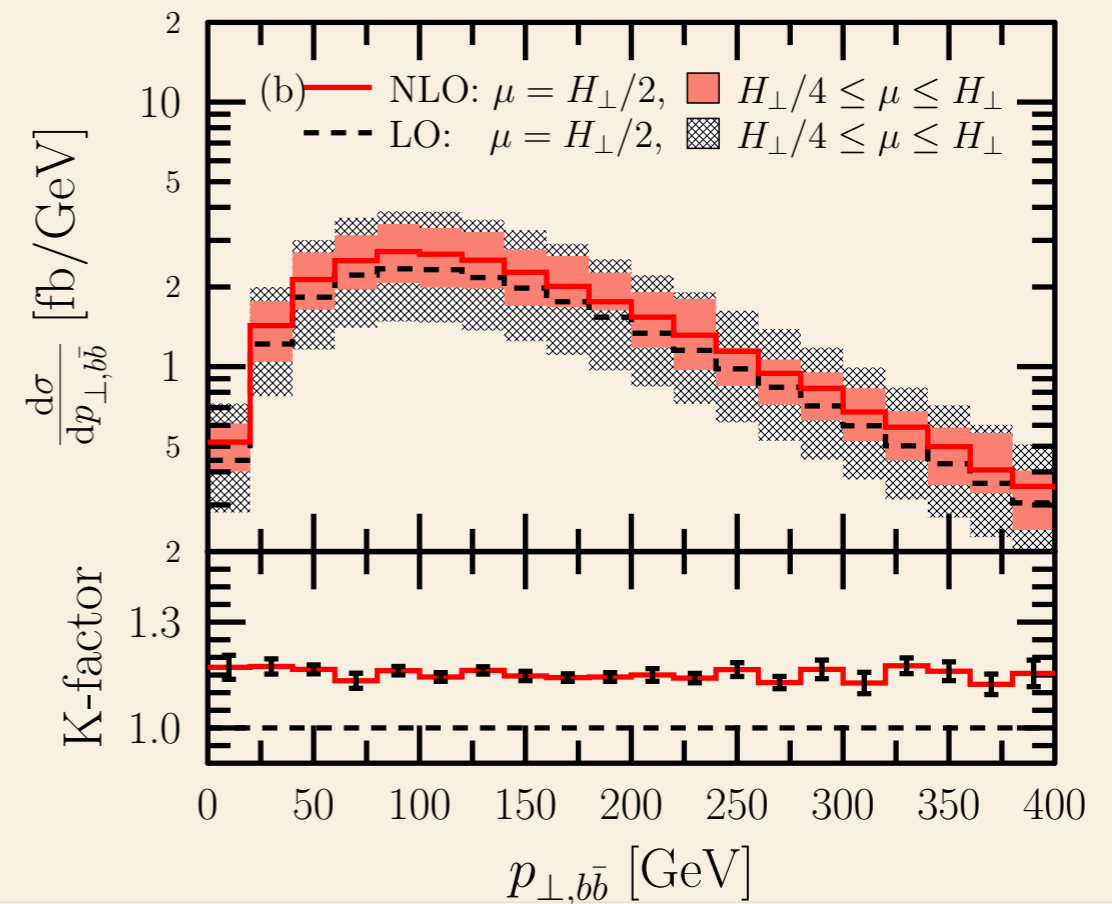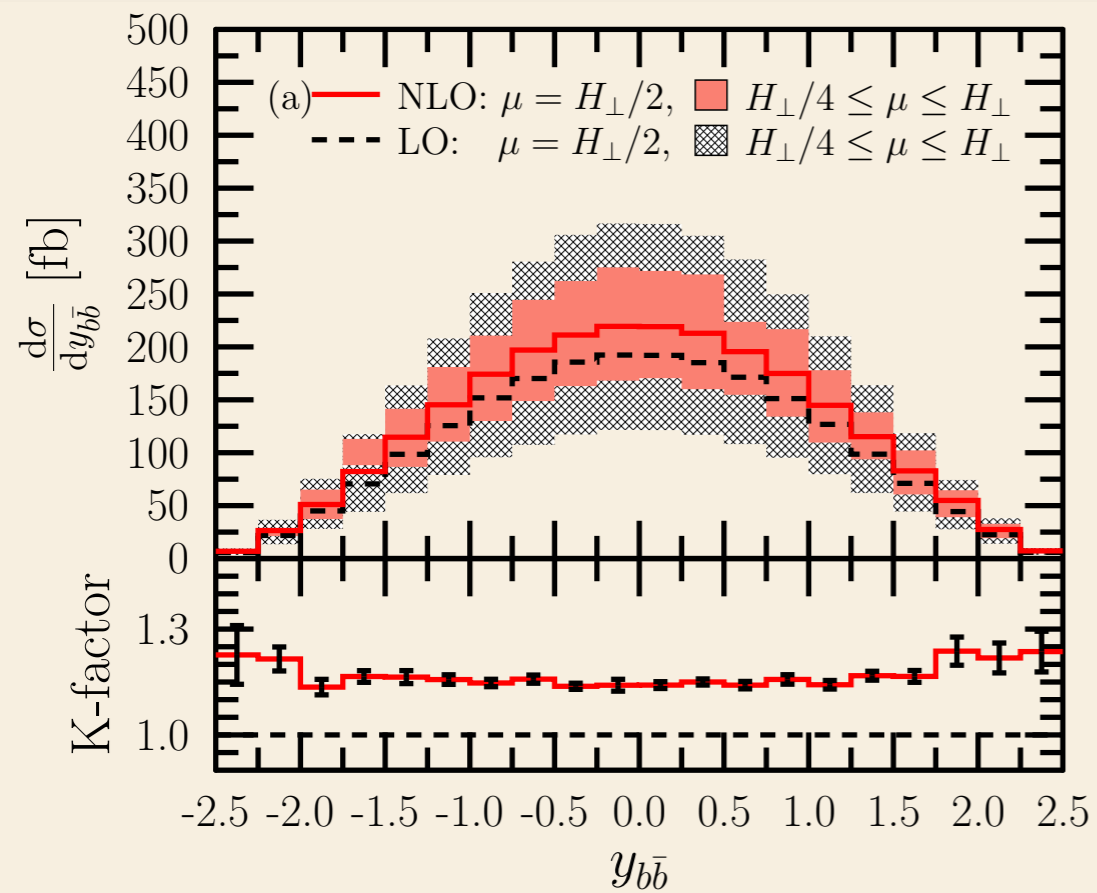✓ the K factor is even smaller, implying good convergence

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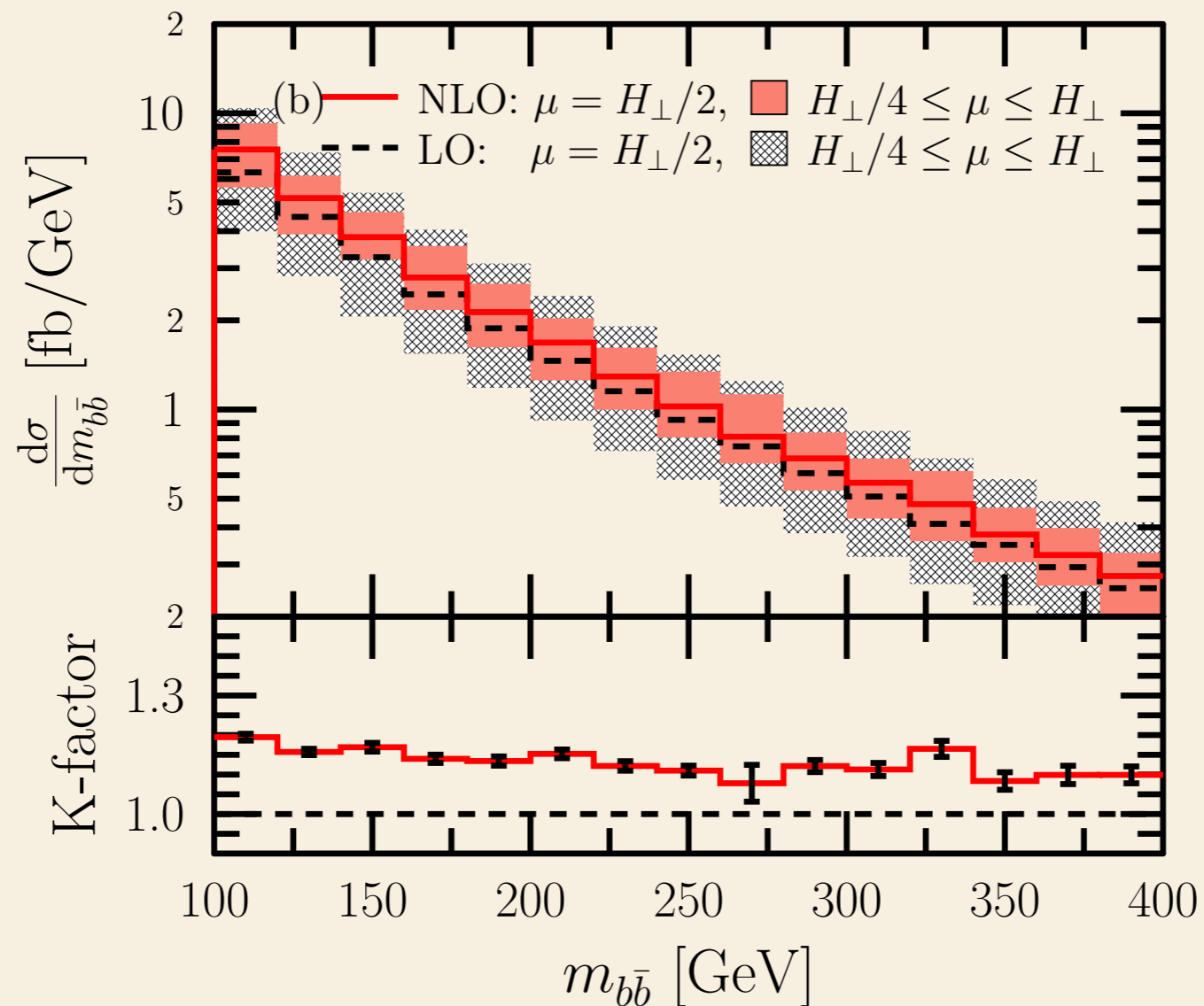
scale dependence: $^{+32\%}_{-22\%}$, largest if $\mu_R = \mu_F = \mu_{dyn}$
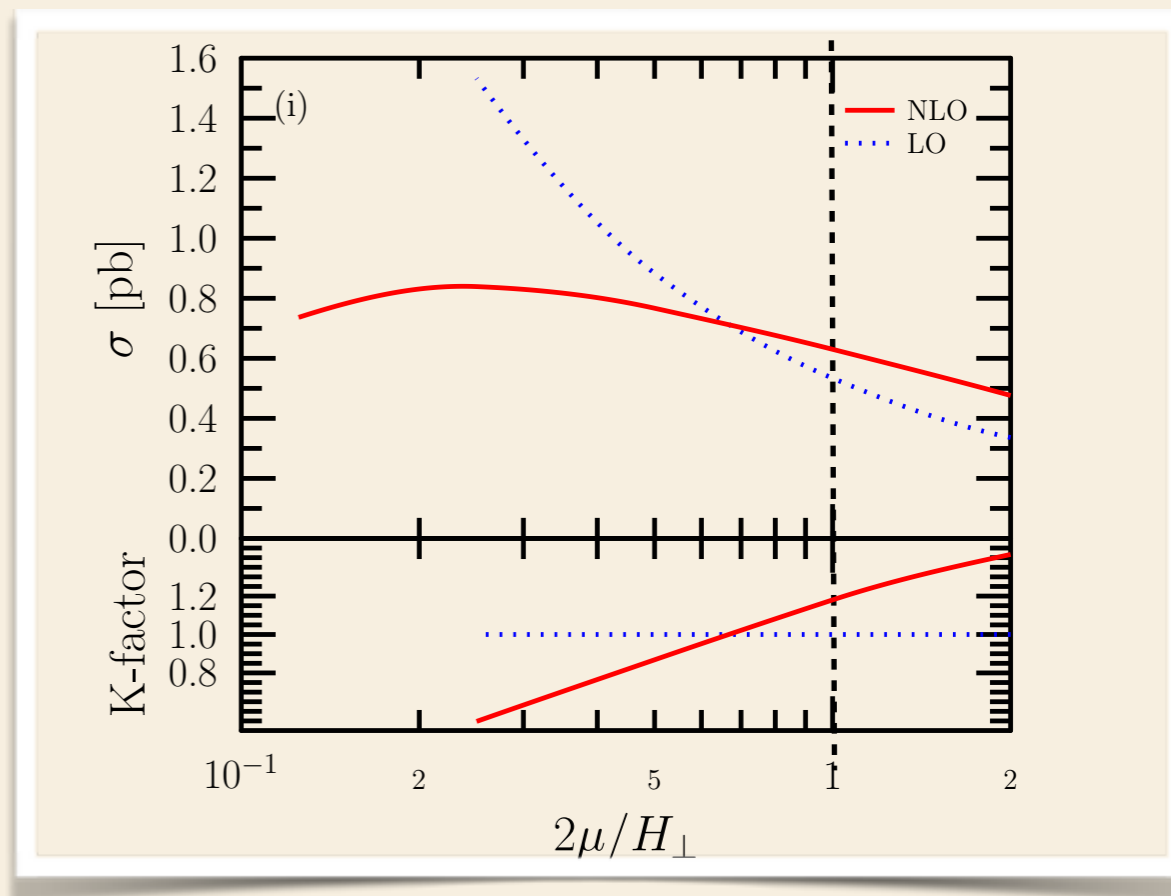
Small changes in shapes of distributions

Small changes in shapes of distributions

Formal accuracy of the POWHEG MC

$$\langle O \rangle = \int \mathrm{d}\Phi_{\mathrm{B}} \widetilde{B} \left[\Delta(p_{\perp,\mathrm{min}}) O(\Phi_{\mathrm{B}}) + \int \mathrm{d}\Phi_{\mathrm{rad}} \Delta(p_{\perp}) \frac{R}{B} O(\Phi_{\mathrm{R}}) \right] =$$

...

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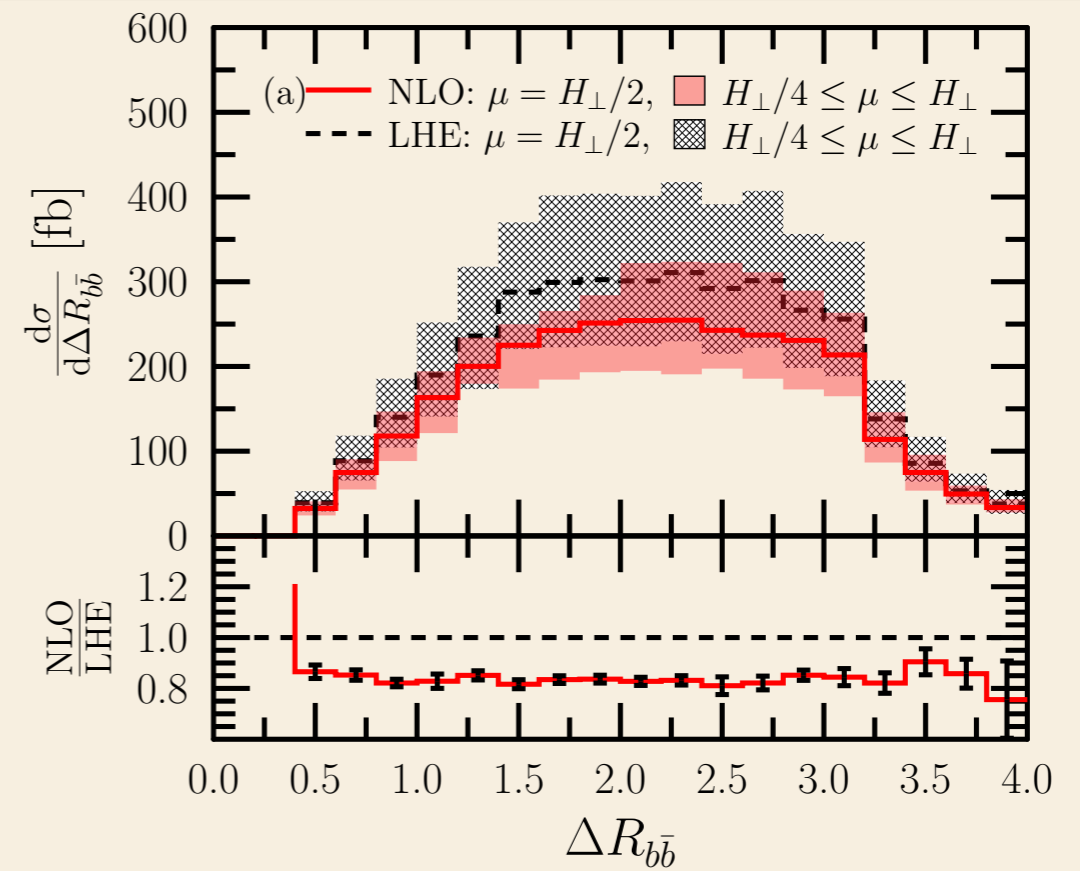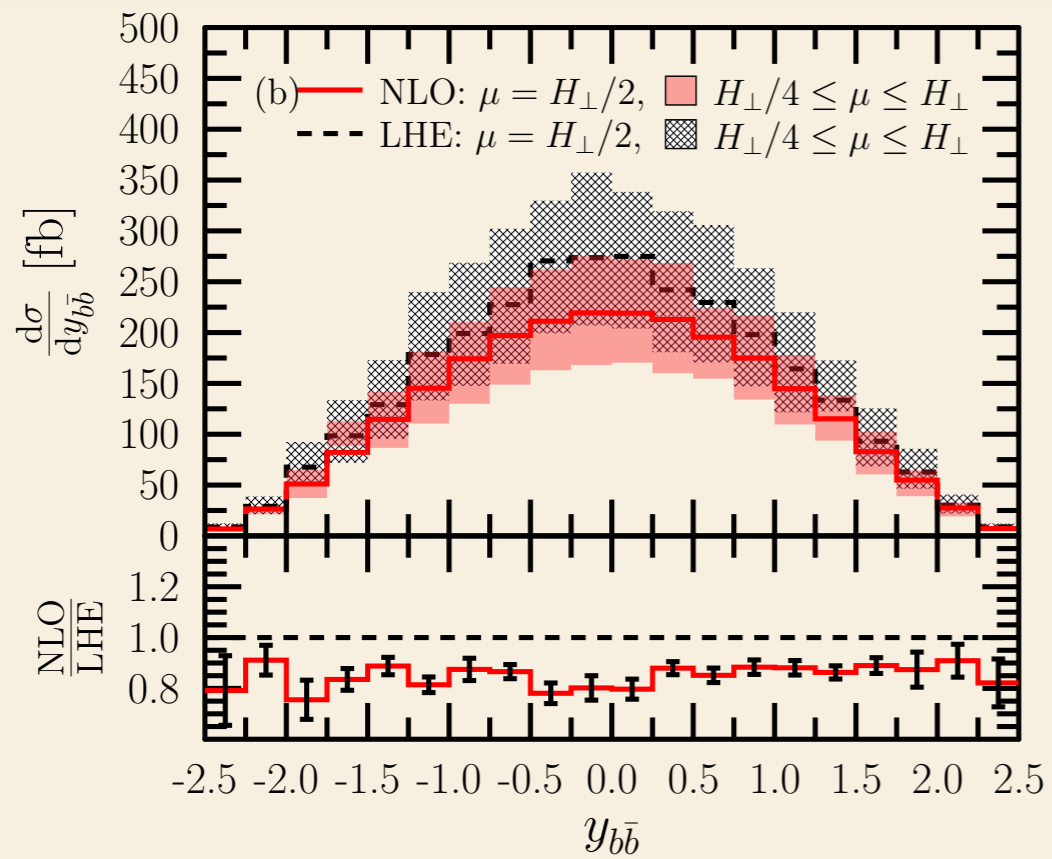
$$\cdots$$

$$= \left\{ \int d\Phi_B \left[B + V \right] O(\Phi_B) + \int d\Phi_R R O(\Phi_R) \right\} (1 + \mathcal{O}(\alpha_S))$$
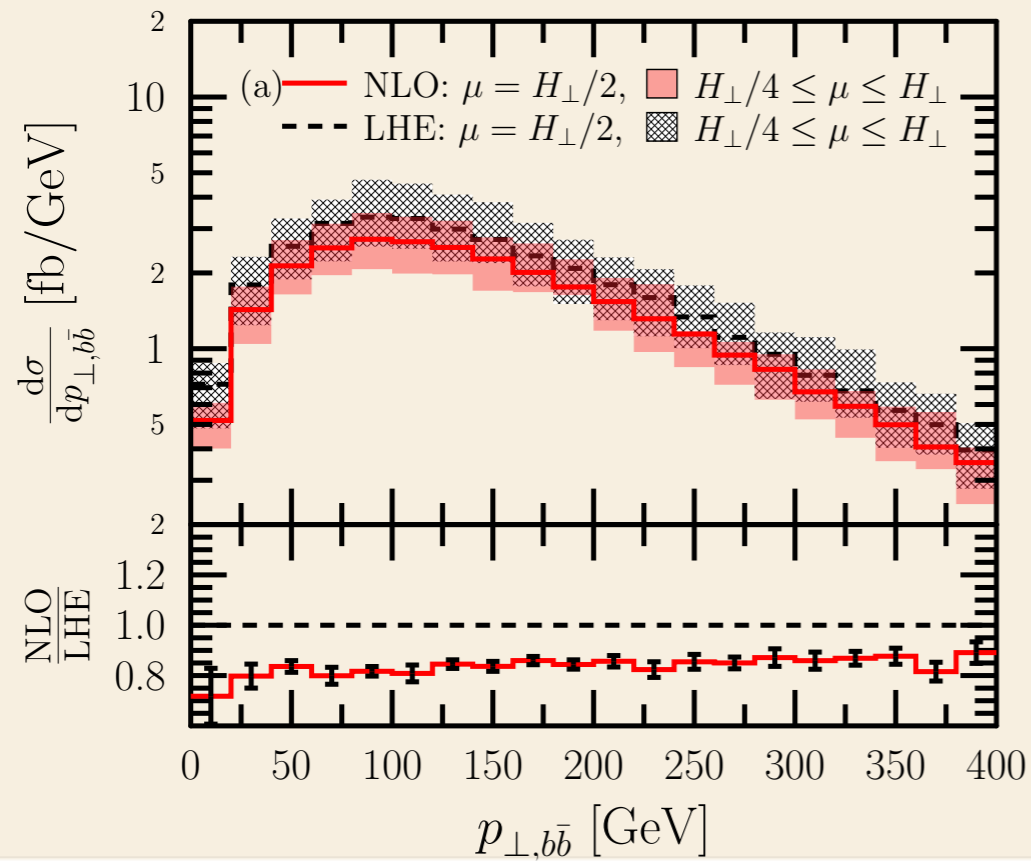
$\langle O \rangle_{\text{NLO}}$

Useful for checking

LHE vs. NLO
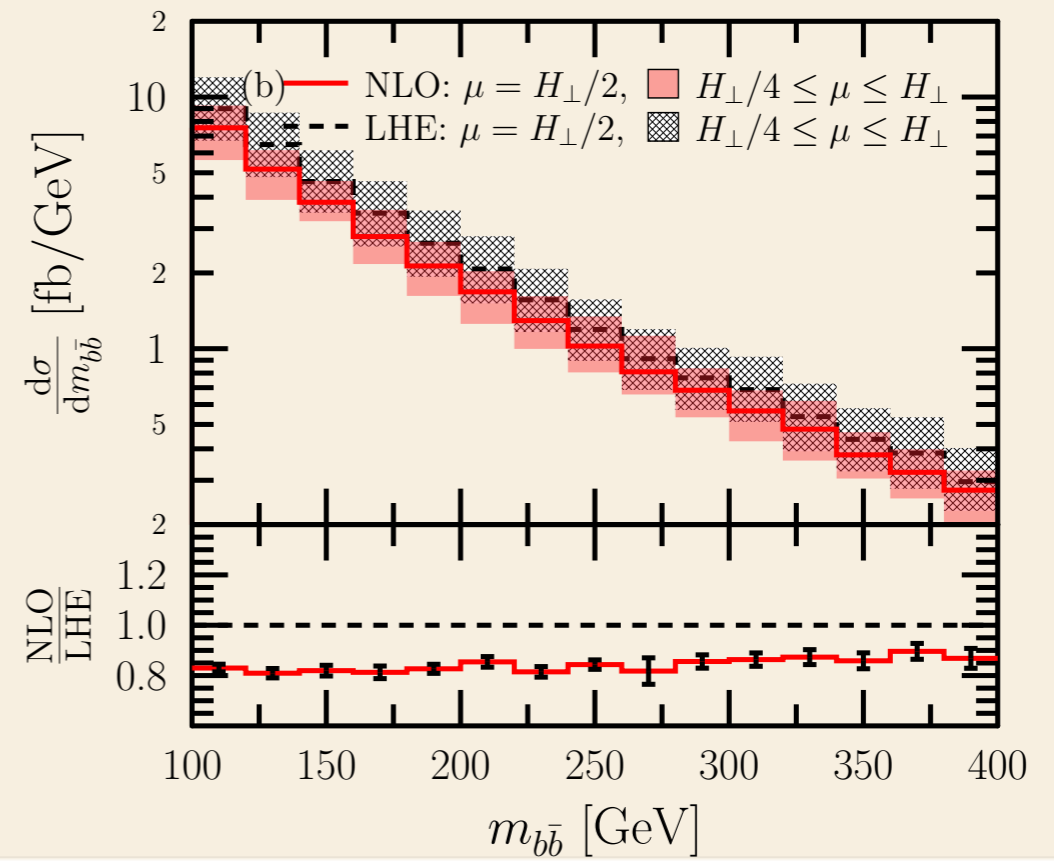
Message:
we can trust the LHE's, so can make

Predictions

Four possible forms of predictions

LHE: distributions from events at BORN+1st radiation

Decay: on-shell decays of heavy particles (t-quarks), shower and hadronization effects turned off

PS: parton showering (PYTHIA or HERWIG) included (t-quarks kept stable)

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Full SMC: decays, parton showering and hadronization are included by using PYTHIA or HERWIG

Four possible forms of predictions

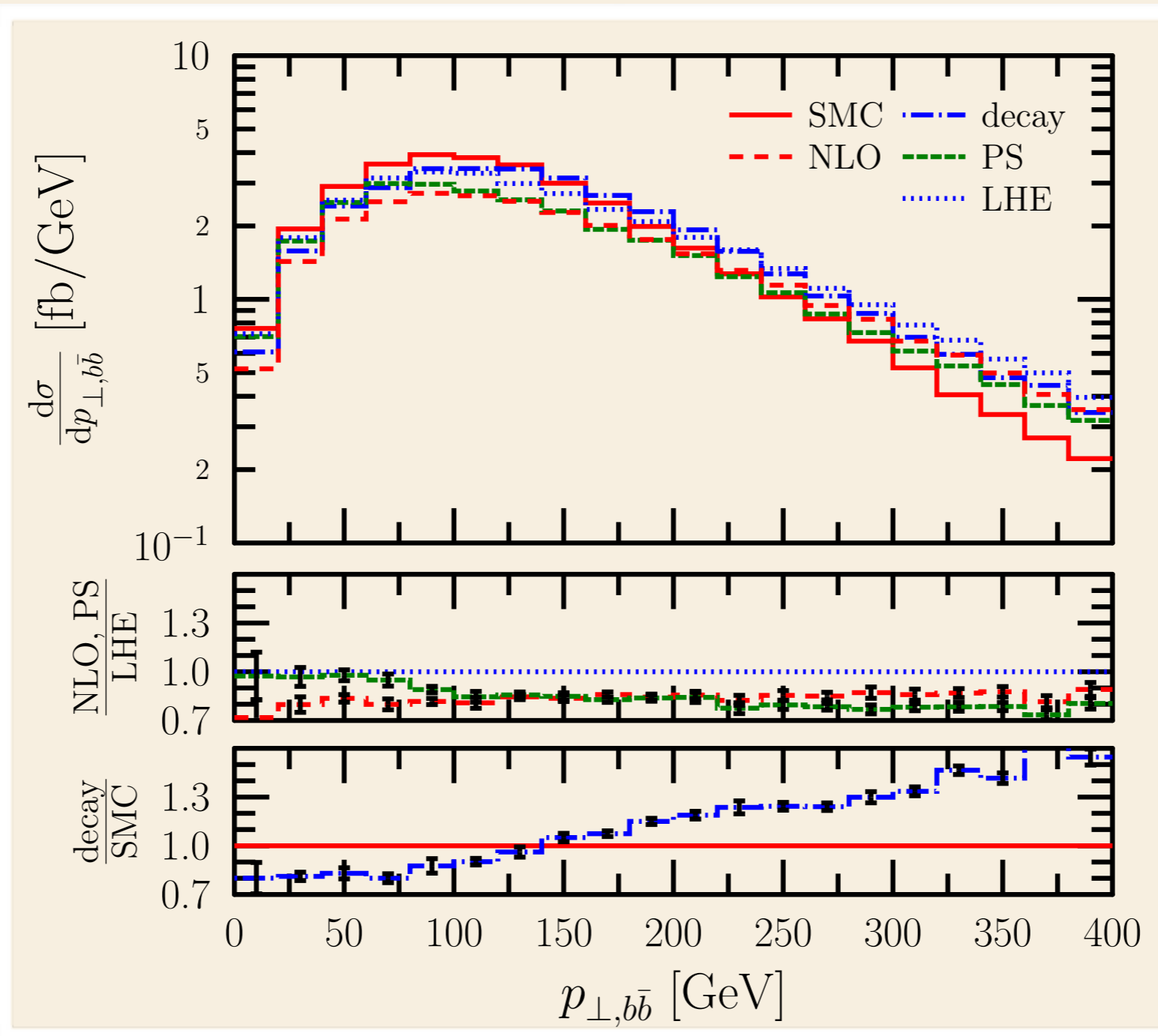**LHE:** distributions from events at BORN+1st radiation

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Number and type of particles are very different => to study the effect of SMC we employ selection cuts to keep the cross section fixed
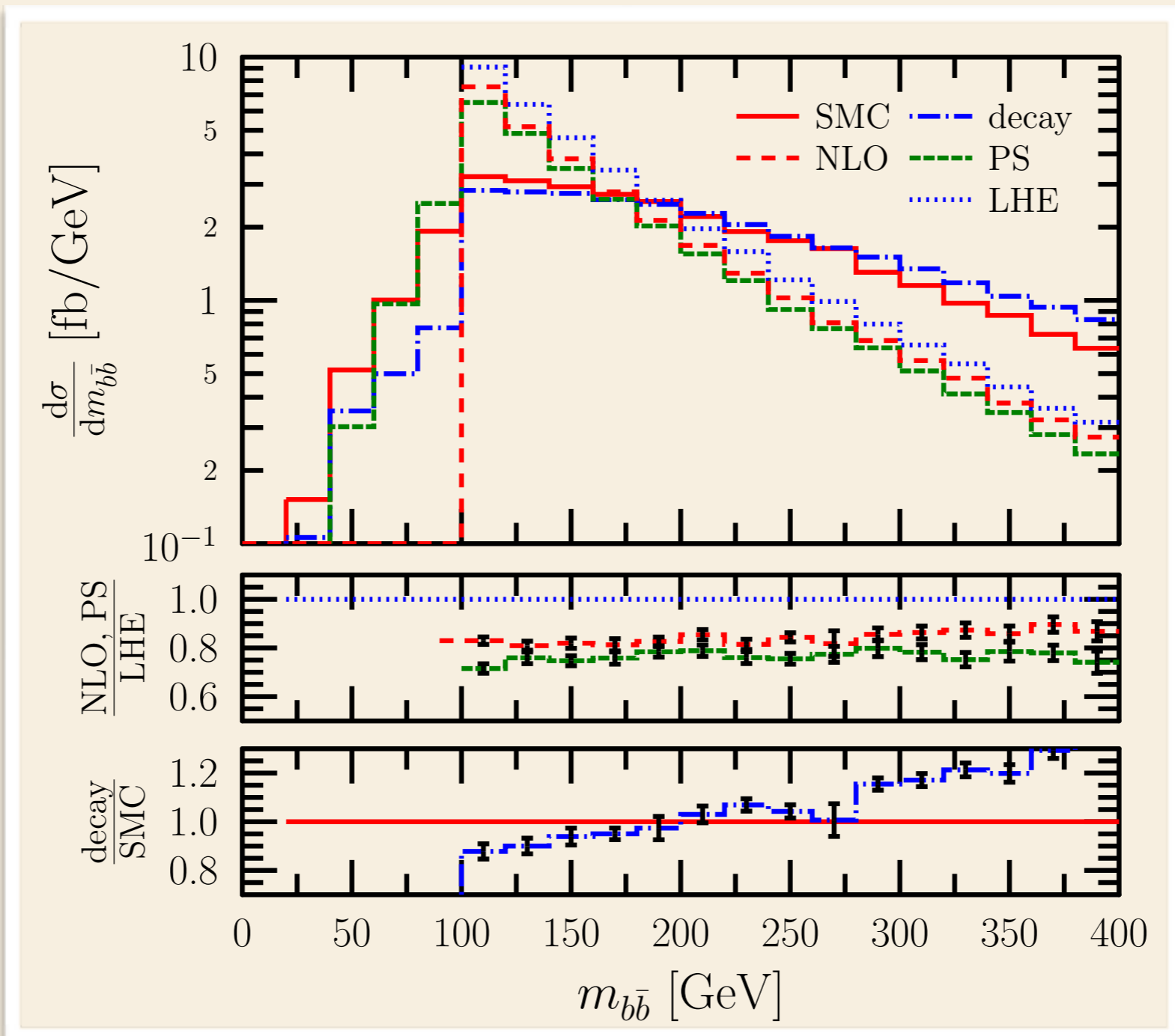
NLO vs. PS and decay vs. full SMC at 14TeV, μ = H_T/2

NLO vs. PS and decay vs. full SMC at 14TeV, μ = $H_T/2$
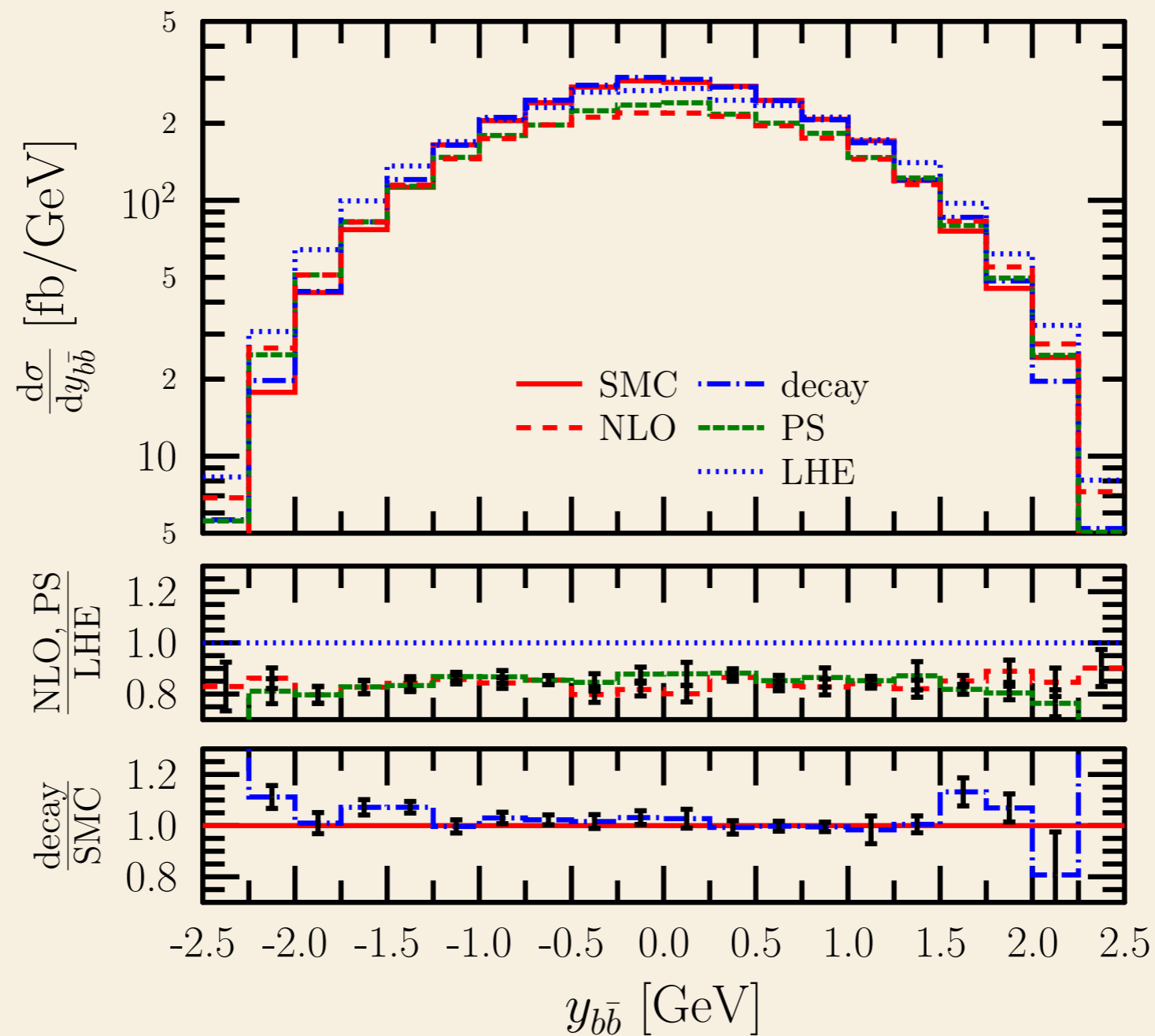
Measurement of the Cross Section Ratio $\sigma(t\bar{t}b\bar{b})/\sigma(t\bar{t}jj)$ in pp Collisions at $\sqrt{s} = 8$ TeV

in dilepton decay mode

The CMS Collaboration

| Final state	ee	$\mu\mu$	$e\mu$	All
$t\bar{t} + b\bar{b}$	18.1 ± 0.8	26.8 ± 1.0	60.9 ± 1.5	105
$t\bar{t} + b$	34.3 ± 1.1	51.4 ± 1.4	111	196
$t\bar{t} + c\bar{c}$	13.4 ± 0.9	20.5 ± 1.0	47.0 ± 1.6	80.9 ± 2.4
$t\bar{t} + LF$	244	359	822	1,425
$t\bar{t}$ others	20.5 ± 1.1	25.6 ± 1.1	63.7 ± 1.9	109
multijet	< 0.1	1.4 ± 1.2	1.4 ± 1.2	2.9 ± 2.2
$W + jets$	< 0.1	< 0.1	< 0.1	< 0.1
VV	< 0.1	0.3 ± 0.1	< 0.1	0.4 ± 0.7
Single top-tW	7.9 ± 2.0	11.6 ± 2.5	25.1 ± 3.7	44.7 ± 4.1
$Z/\gamma* \to ll$	5.6 ± 4.3	5.7 ± 3.9	2.9 ± 3.2	14.4 ± 5.7
Total expected	351	512	1,159	2,023
Data	367	506	1,145	2,018

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t̄tH signal on t̄tbb̄ background
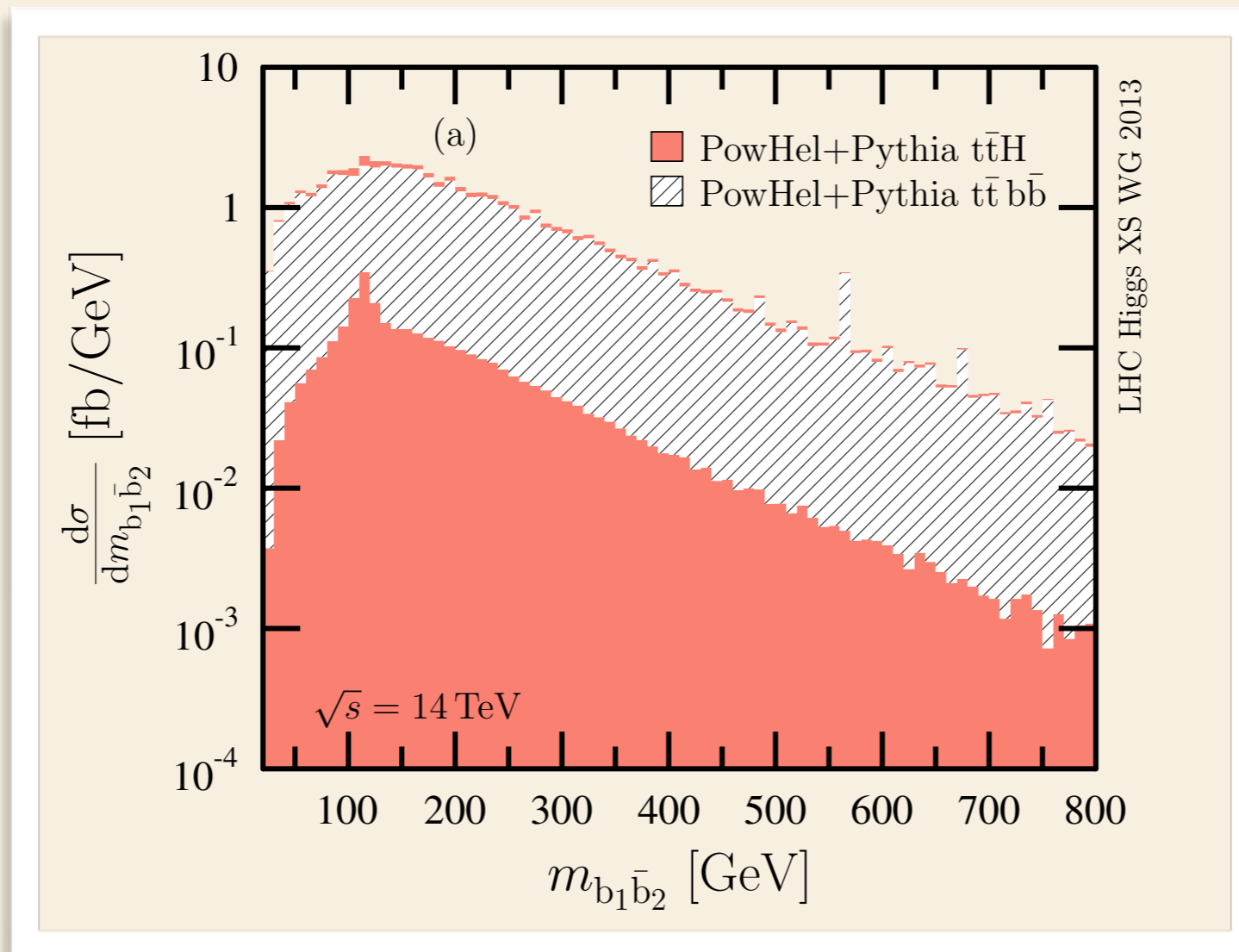
Distribution of the invariant mass of the hardest bb̄ jet pair in pp → t̄t H and t̄t bb̄ at LHC (14 TeV)

Conclusions and outlook

Conclusions

- First computation of pp → $t\bar{t}b\bar{b}$ at NLO + SMC accuracy [A. Kardos and Z.T. arXiv:1303.6291, Cascioli et al arXiv:1309.5912, Meierhofer this morning]

- NLO cross sections agree with published predictions and with CMS results

- Effects of decay of t-quarks could be important

- LHE event files for pp → $t\bar{t}$, $t\bar{t}H$, $t\bar{t}W$, $t\bar{t}Z$, $t\bar{t}jet$, $t\bar{t}b\bar{b}$ processes available, to put into SMC and perform experimental analyses on events with hadrons

Processes available in PowHel

✓ tT   [Kardos et al, arXiv:
✓ tT + Z   1111.0610,1111.1444,
✓ tT + W   1208.2665,
✓ tT + H/A   1108.0387,
✓ tT + j   1101.2672,
✓ WWbB   PoS LL2012 057
✓ tT + bB   1103.6291]

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✓WWbB   PoS LL2012 057

✓tT + bB   1103.6291]

✓tT +... (2 more processes coming soon)

The end
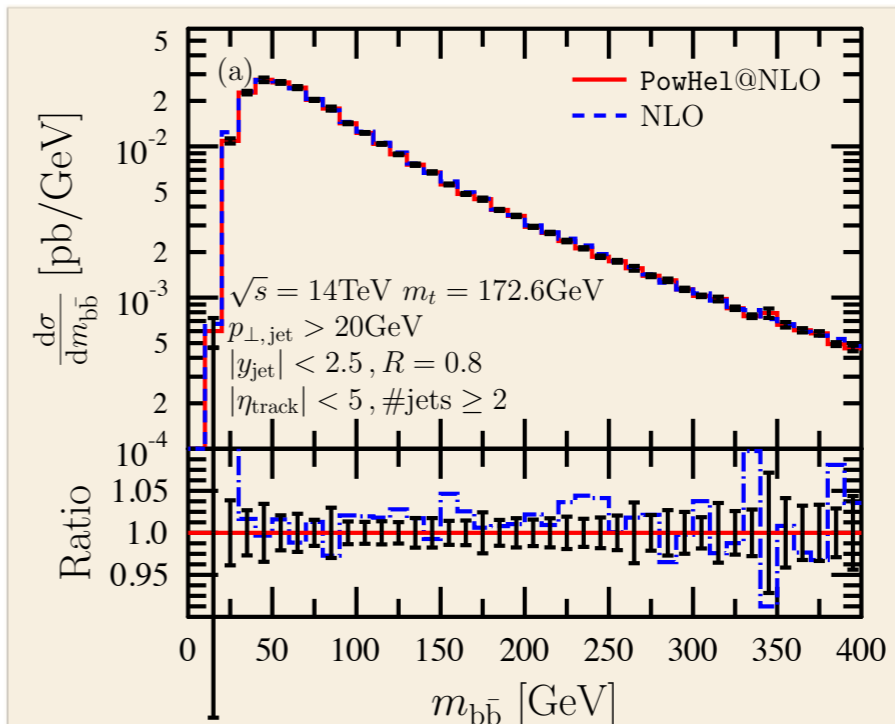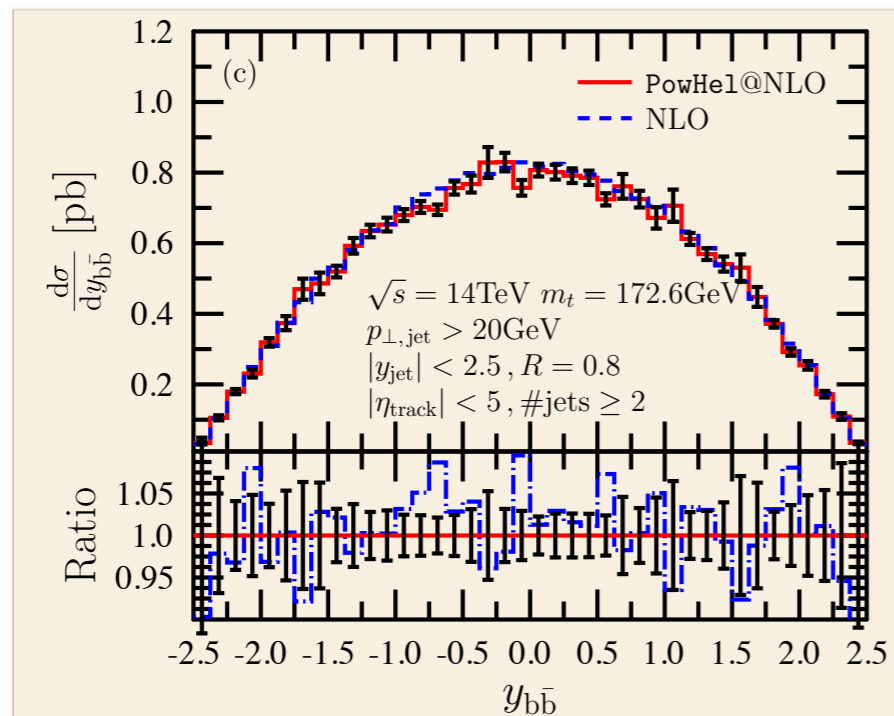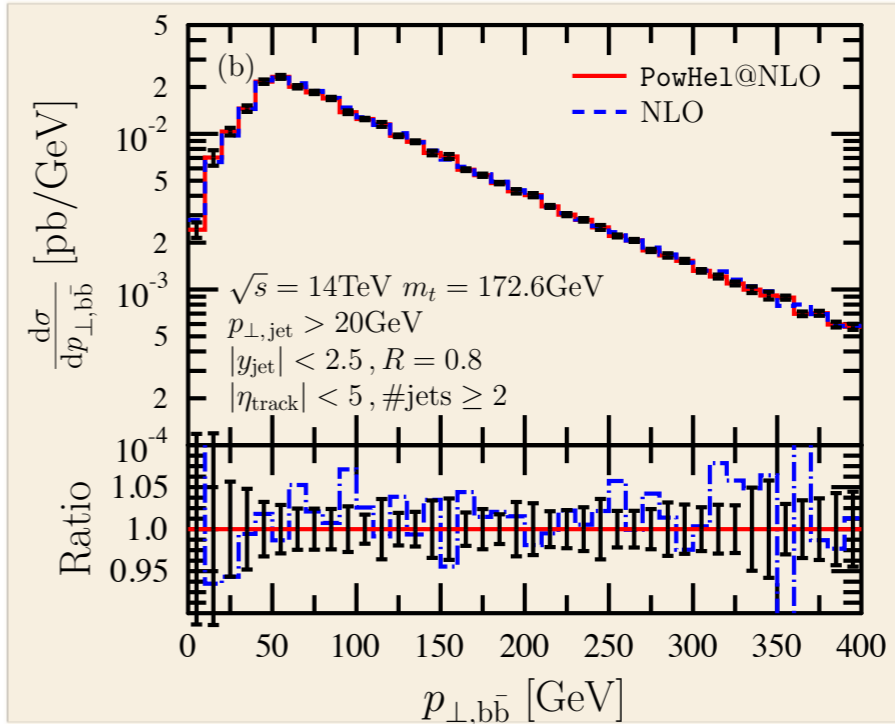
Appendix

Selection cuts in NLO studies

Cuts employed by Bevilacqua et al in arXiv:0907.4723

▸ A track was considered as a possible jet constituent if $|\eta^{track}|<5$, t-quarks were excluded from the set of possible tracks, jets were reconstructed with the k_T-algorithm using R=0.4

▸ Events with invariant mass of the $b\bar{b}$-jet pair below $m^{min}_{b\bar{b}}$ = 20 GeV were discarded

▸ We require $p_{Tmin,j}$ = 20 GeV and

▸ at least two, one b- and one $\bar{b}$-jet, with $|y_{b(\bar{b})}|<2.5$
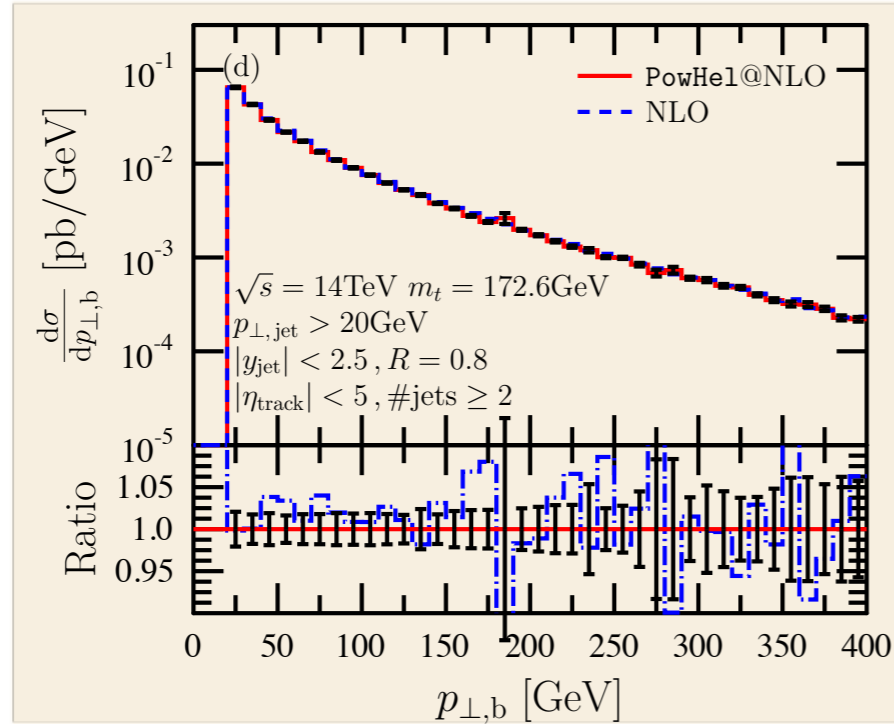
Selection cuts for decay vs. SMC

▸ **Applied on the LHE's:**

  ▸ A track was considered as a possible jet constituent if $|\eta^{track}|<5$, t-quarks were excluded from the set of possible tracks. Jets were reconstructed with the anti-k_T algorithm using R=0.4.

  ▸ Events with invariant mass of the $b\bar{b}$-jet pair below $m^{min}_{b\bar{b}}$ = 100 GeV were discarded.

▸ **Applied on LHE's and checked also on the existing particles at different stages of evolution:**

  ▸ we require $p_{Tmin,j}$ = 25 GeV and

  ▸ at least two, one b- & one $\bar{b}$-jet with $|\eta_{b(\bar{b})}|<2.5$.

Comparison to CMS PAS TOP-13-010

▸ at least one pair of isolated (with R=0.3, I_{rel} = 0.15) opposite sign leptons with $p_{Tmin,\ell}$ = 20 GeV/c, $|\eta_\ell|$<2.4, 12 GeV < $m_{\ell\ell}c^2$ ($\notin$[77, 107] GeV if ee or μμ)

▸ p_T^{miss} = 30 GeV/c if ee or μμ

▸ jets reconstructed with the anti-k_T algorithm using R=0.4, with $p_{Tmin,j}$ = 20 (40) GeV and $|\eta_j|$<2.5

▸ at least four well separated jets with ΔR > 0.5 both from leptons and jets

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$t\bar{t} + c\bar{c}$	3.1 ± 0.4	4.6 ± 0.5	9.5 ± 0.7	17.3 ± 1.6
$t\bar{t} + LF$	62.2 ± 2.0	94.1 ± 2.3	211 ± 3.6	368
$t\bar{t}$ others	9.5 ± 0.8	11.8 ± 0.8	28.3 ± 1.3	49.7 ± 2.2
multijet	< 0.1	0.3 ± 0.6	0.3 ± 0.6	0.7 ± 1.6
$W + jets$	< 0.1	< 0.1	< 0.1	< 0.1
VV	< 0.1	< 0.1	< 0.1	< 0.1
Single top-tW	2.8 ± 1.24	2.7 ± 1.2	4.4 ± 1.7	9.9 ± 2.7
$Z/\gamma* \to ll$	2.2 ± 3.02	< 0.1	2.9 ± 3.2	5.2 ± 3.9
Total expected	100	139	315	555
Data	90	148	311	549

Measurement of the Cross Section Ratio $\sigma(t\bar{t}b\bar{b})/\sigma(t\bar{t}jj)$ in pp Collisions at $\sqrt{s} = 8$ TeV

in dilepton decay mode

The CMS Collaboration

Final state	ee	$\mu\mu$	$e\mu$	All
$t\bar{t} + b\bar{b}$	4.0 ± 0.4	5.9 ± 0.5	13.3 ± 0.7	23.3 ± 1.5
$t\bar{t} + b$				
$t\bar{t} + c\bar{c}$	4.1 ± 0.2	4.3 ± 0.3	12.3 ± 0.3	20.7 ± 0.5
$t\bar{t} + LF$	62.2 ± 2.0	94.1 ± 2.3	211 ± 3.6	368
$t\bar{t}$ others	9.5 ± 0.8	11.8 ± 0.8	28.3 ± 1.3	49.7 ± 2.2
multijet	< 0.1	0.3 ± 0.6	0.3 ± 0.6	0.7 ± 1.6
$W + jets$	< 0.1	< 0.1	< 0.1	< 0.1
VV	< 0.1	< 0.1	< 0.1	< 0.1
Single top-tW	2.8 ± 1.24	2.7 ± 1.2	4.4 ± 1.7	9.9 ± 2.7
$Z/\gamma* \to ll$	2.2 ± 3.02	< 0.1	2.9 ± 3.2	5.2 ± 3.9
Total expected	100	139	315	555
Data	90	148	311	549

Cuts for background study for tTH

Applied after full SMC

▶ a track was considered as a possible jet constituent if $|\eta^{track}|<5$, jets were reconstructed with the anti-k_T algorithm using R=0.4

we require

▶ at least six jets with $p_{Tmin,j}$ = 20 GeV and $|\eta_j|<5$

▶ at least two b-jets & two $\bar{b}$-jets with $|\eta_{b(\bar{b})}|<2.7$, with MCTRUTH tagging

▶ at least one isolated (with R=0.4) lepton with $p_{Tmin,\ell}$ = 20 GeV and $|\eta_\ell|<2.5$

▶ p_T^{miss} = 15 GeV

to disentangle background in the semileptonic $t\bar{t}$ decay

t̄tH signal on t̄tb̄b background