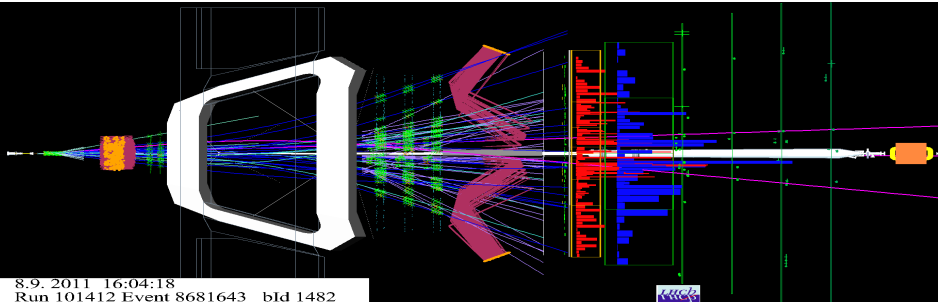


Rare B meson decays to NNLO

Matthias Steinhauser | TTP Karlsruhe

Loops and Legs, Weimar, April 30, 2014



■ $B_s \rightarrow \mu^+ \mu^-$

- NNLO QCD
- NLO EW

[Bobeth,Gorbahn,Hermann,Misiak,Stamou,Steinhauser'13]

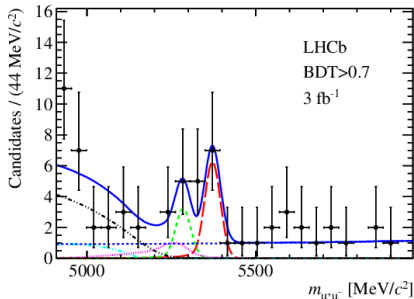
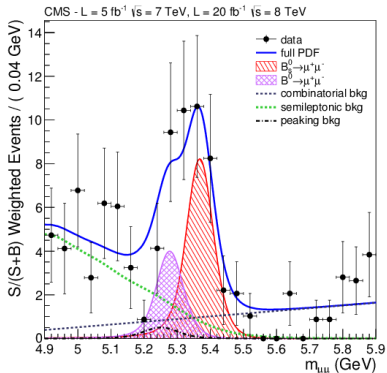
■ $\bar{B} \rightarrow X_s \gamma$

[Misiak et al.]

Experiment: invariant ($\mu^+ \mu^-$) mass

CMS

LHCb



$$\overline{B}_{S\mu} = (2.9 \pm 0.7) \times 10^{-9}$$

≥ 2015: uncertainty below 10% within the next 10 years

- NLO QCD corrections [Buchalla,Buras'93'99; Misiak,Urban'99]
- leading- m_t NLO electroweak corrections [Buchalla,Buras'98]
- uncertainty (from higher orders): $\approx 7\%$

$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)$: why interesting?

- SM prediction for $\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)$: $\mathcal{O}(10^{-9})$
nevertheless: experimentally accessible
- \Leftrightarrow “new physics” contribution can be sizeable
- hope: easy to see deviations
- but: $\overline{\mathcal{B}}^{\text{exp}} \approx \overline{\mathcal{B}}^{\text{SM}}$
- \Leftrightarrow constraints on “new physics”
 \Leftrightarrow precision test of SM

$$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)$$

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QCD} \times \text{QED}} + \frac{V_{tb}^* V_{ts} G_F^2 M_W^2}{\pi^2} \sum_n C_n Q_n + \text{h.c.}$$

$$Q_A = (\bar{b} \gamma_\alpha \gamma_5 s)(\bar{\mu} \gamma^\alpha \gamma_5 \mu) \quad Q_S = (\bar{b} \gamma_5 s)(\bar{\mu} \mu) \quad Q_P = (\bar{b} \gamma_5 s)(\bar{\mu} \gamma_5 \mu)$$

SM: only Q_A

BSM: Q_S and Q_P important

$$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-) = \frac{|N|^2 M_{B_s}^3 f_{B_s}^2}{8\pi \Gamma_H^S} \beta \left[|r C_A - u C_P|^2 F_P + |u \beta C_S|^2 F_S \right] + \mathcal{O}(\alpha_{em})$$

(time-integrated, CP-averaged, [De Bruyn et al.'12])

$$\text{SM: } F_P = 1, F_S = 1 - \Delta \Gamma^S / \Gamma_L^S \quad r = \frac{2m_\mu}{M_{B_s}}, \beta = \sqrt{1 - r^2}, u = \frac{M_{B_s}}{m_b + m_s}$$

Aim: compute C_A to 3 loops in QCD + NLO EW

$$\overline{B}(B_s \rightarrow \mu^+ \mu^-)$$

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QCD} \times \text{QED}} + \frac{V_{tb}^* V_{ts} G_F^2 M_W^2}{\pi^2} \sum_n C_n Q_n + \text{h.c.}$$

$$Q_A = (\bar{b} \gamma_\alpha \gamma_5 s)(\bar{\mu} \gamma^\alpha \gamma_5 \mu) \quad Q_S = (\bar{b} \gamma_5 s)(\bar{\mu} \mu) \quad Q_P = (\bar{b} \gamma_5 s)(\bar{\mu} \gamma_5 \mu)$$

SM: only Q_A

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$$\overline{B}(B_s \rightarrow \mu^+ \mu^-) = \frac{|N|^2 M_{B_s}^3 f_{B_s}^2}{8\pi \Gamma_H^S} \beta r^2 |C_A(\mu_b)|^2 + \mathcal{O}(\alpha_{em})$$

(time-integrated, CP-averaged, [De Bruyn et al.'12])

$$N = V_{tb}^* V_{tq} G_F^2 M_W^2 / \pi^2$$

$$r = \frac{2m_\mu}{M_{B_s}}, \beta = \sqrt{1 - r^2}, u = \frac{M_{B_s}}{m_b + m_s}$$

Aim: compute C_A to 3 loops in QCD + NLO EW

$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)$: missing $\mathcal{O}(\alpha_{em})$

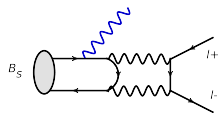
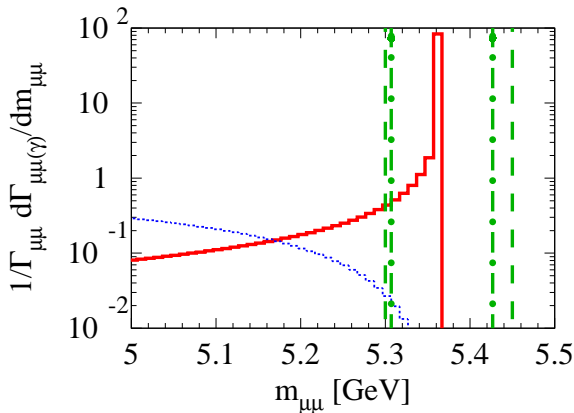
$$\overline{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-) = \frac{|N|^2 M_{B_s}^3 f_{B_s}^2}{8\pi \Gamma_H^S} \beta r^2 |C_A(\mu_b)|^2 + \mathcal{O}(\alpha_{em})$$

- no enhancement factor (like $\frac{1}{\sin^2 \theta_W}$, $\frac{m_t^2}{M_W^2}$ or $\ln^2 \frac{M_W^2}{\mu_b^2}$)
- contribution from other operators
 e.g., $(\bar{b}\gamma_\alpha \gamma_5 q)(\bar{\ell}\gamma^\alpha \ell)$ or $(\bar{b}\gamma_\alpha P_L c)(\bar{c}\gamma^\alpha P_L s)$
 with γ connecting quark and lepton lines
- non-perturbative effects beyond f_{B_s}
- compensates μ_b dependence of $C_A(\mu_b)$ ($\Leftrightarrow \approx 0.3\%$)
- soft Bremsstrahlung

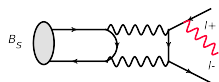
$\overline{B}(B_s \rightarrow \mu^+ \mu^-)$: missing $\mathcal{O}(\alpha_{em})$

$$\overline{B}(B_s \rightarrow \mu^+ \mu^-) = \frac{|N|^2 M_{B_s}^3 f_{B_s}^2}{8\pi \Gamma_H^S} \beta r^2 |C_A(\mu_b)|^2 + \mathcal{O}(\alpha_{em})$$

- soft Bremsstrahlung: $B_s \rightarrow \mu^+ \mu^- + (n\gamma) \ (n = 0, 1, 2, \dots)$

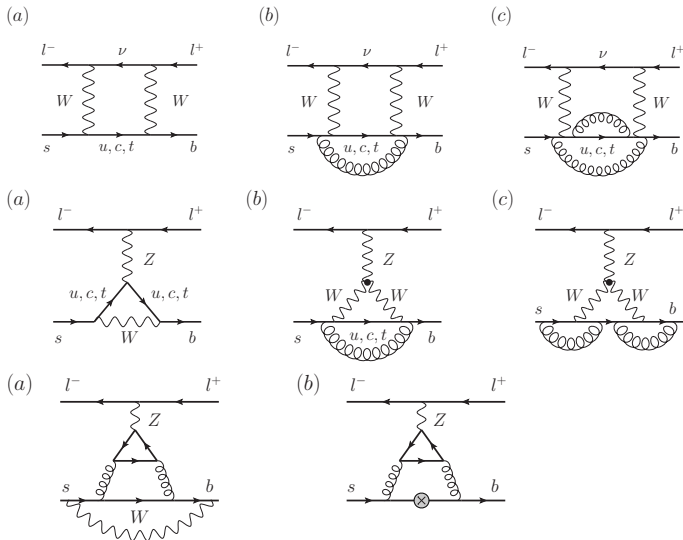


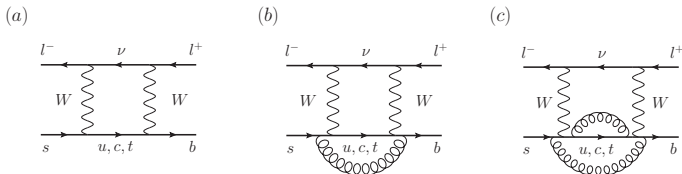
bkg for theory and exp.



taken into account by
LHCb and CMS

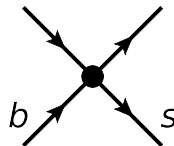
C_A to 3 loops: Feynman diagrams



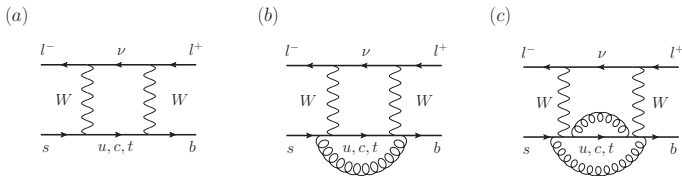


■ matching in $d = 4 - 2\epsilon$ dimensions

- effective theory: only tree-level contribution
- evanescent operators
- vacuum integrals, 2 masses: M_W and m_t
 - ↔ consider expansions: $M_W \approx m_t$; $M_W \ll m_t$



W box

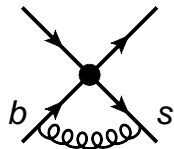
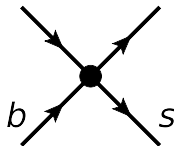


■ matching in $d = 4 - 2\epsilon$ dimensions

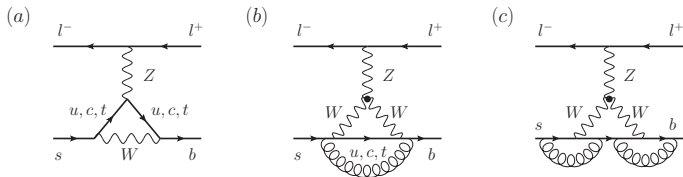
- effective theory: only tree-level contribution
- evanescent operators
- vacuum integrals, 2 masses: M_W and m_t
 - ⇒ consider expansions: $M_W \approx m_t$; $M_W \ll m_t$

■ matching in 4 dimensions

- NO evanescent operators
- loop diagrams in effective theory
- IR regulator needed: $m_s \ll m_b \ll M_W, m_t$
 - ⇒ nested asymptotic expansion



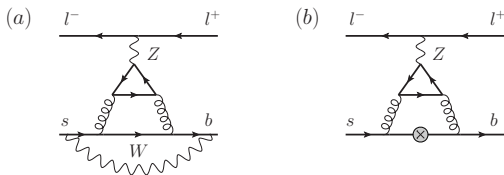
Z penguin



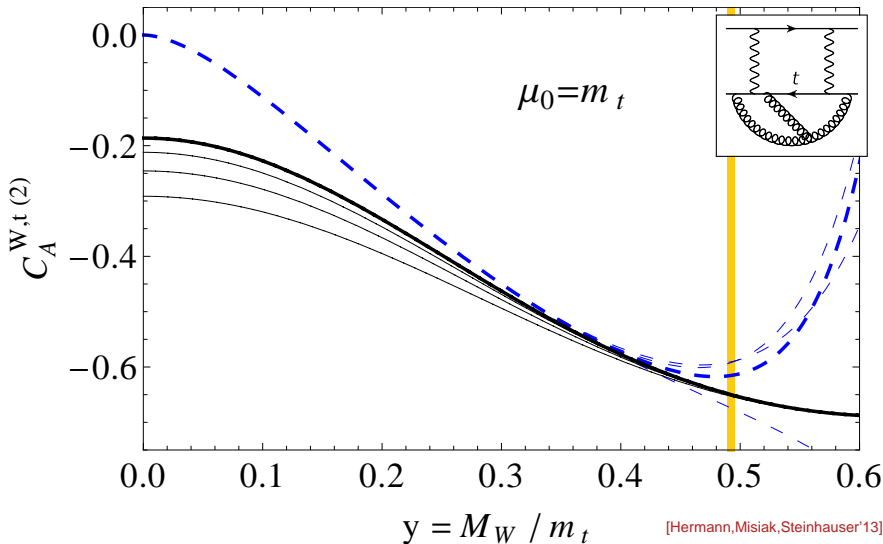
- matching in $d = 4 - 2\epsilon$ dimensions
- NO evanescent operators
- effective theory: only tree-level contribution
- triangle contributions

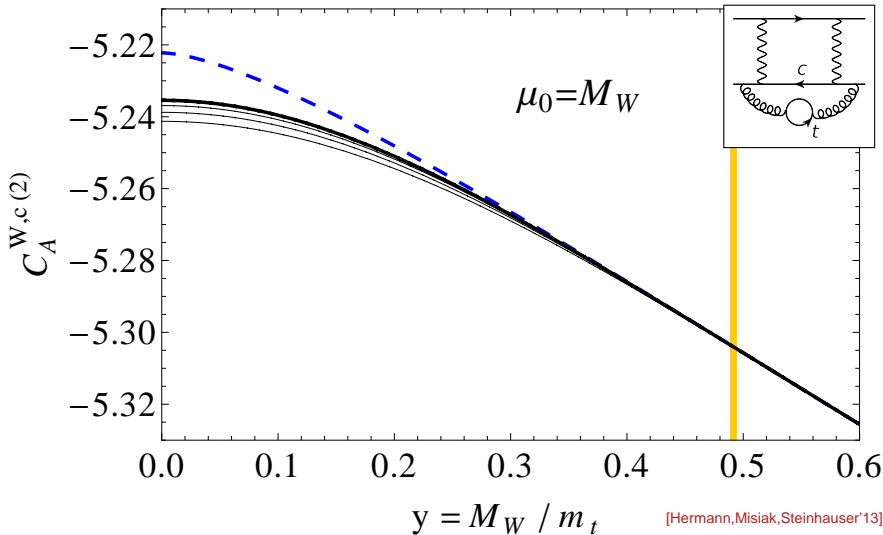
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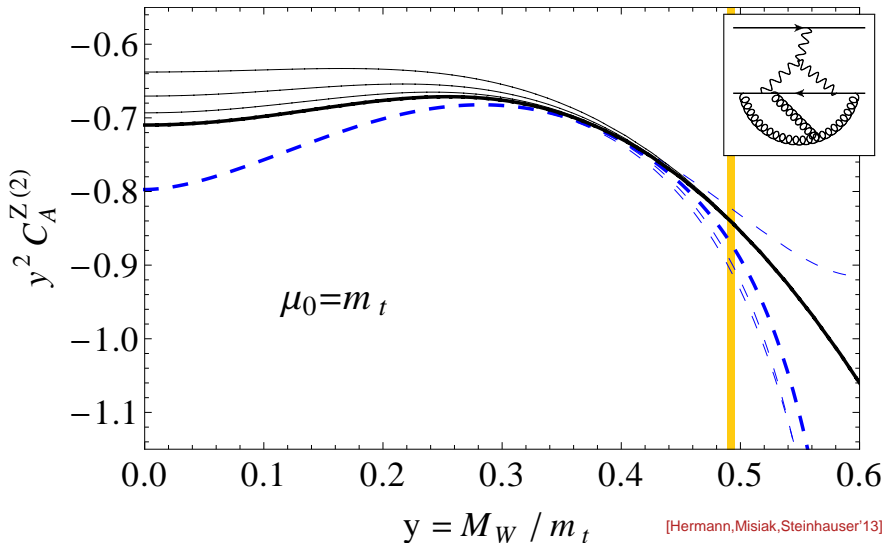
- matching in $d = 4 - 2\epsilon$ dimensions
- NO evanescent operators
- effective theory: only tree-level contribution
- triangle contributions



- $u - d$ and $c - s \leftrightarrow$ cancel
- non-zero contribution from $t - b$
- non-naive treatment of γ_5 (e.g. [Larin'93])

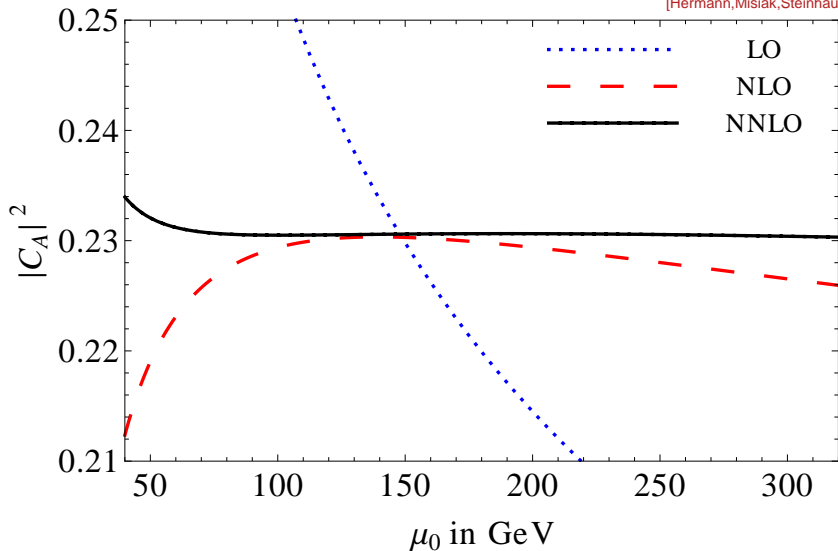


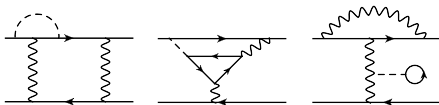
C_A W box: charm



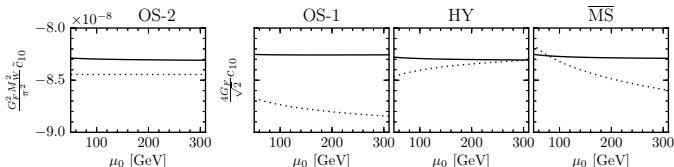
Results: μ dependence

[Hermann,Misiak,Steinhauser'13]





- different renormalization schemes
($G_F \alpha_{em}$ vs. $G_F^2 M_W^2$; OS, $\overline{\text{MS}}$, hybrid)



- RGE running (not in QCD!): $C_A(\mu_0 \sim M_W, m_t) \rightarrow C_A(\mu_b \sim m_b)$
- uncertainty: 7% $\implies < 1\%$

Results: parametrization of \mathcal{B}

[Bobeth, Gorbahn, Hermann, Misiak, Stamou, Steinhauser'13]

$$\overline{\mathcal{B}}_{S\mu} = (3.65 \pm 0.06) R_{t\alpha} R_S \times 10^{-9} = 3.65 \pm 0.23 \times 10^{-9}$$

$$R_{t\alpha} = R_t^{3.06} R_\alpha^{-0.18} = \tilde{R}_t^{3.02} R_\alpha^{0.032}$$

$$R_S = \left(\frac{f_{B_S} [\text{MeV}]}{227.7} \right)^2 \left(\frac{|V_{cb}|}{0.0424} \right)^2 \left(\frac{|V_{tb}^* V_{ts} / V_{cb}|}{0.980} \right)^2 \frac{\tau_H^S [\text{ps}]}{1.615}$$

$$R_\alpha = \frac{\alpha_s(M_Z)}{0.1184}$$

$$R_t = \frac{M_t^{\text{OS}}}{(173.1 \text{ GeV})}$$

$$\tilde{R}_t = \frac{m_t^{\overline{\text{MS}}}}{(163.5 \text{ GeV})}$$

Results: uncertainties

$$\overline{B}_{S\mu} = (3.65 \pm 0.06) R_{t\alpha} R_S \times 10^{-9} = 3.65 \pm 0.23 \times 10^{-9}$$

$$R_S = \left(\frac{f_{B_S} [\text{MeV}]}{227.7} \right)^2 \left(\frac{|V_{cb}|}{0.0424} \right)^2 \left(\frac{|V_{tb}^* V_{ts} / V_{cb}|}{0.980} \right)^2 \frac{\tau_H^S [\text{ps}]}{1.615}$$

f_{B_S} : [FLAG], V_{cb} : [Gambino,Schwanda'13], $|V_{tb}^* V_{ts} / V_{cb}|$: [CKMfitter,UTfit], τ_H^S : [HFAG],

	f_{B_S}	CKM	τ_H^S	M_t	α_S	other param.	non-param.	Σ
\overline{B}_{sl}	4.0%	4.3%	1.3%	1.6%	0.1%	< 0.1%	1.5%	6.4%

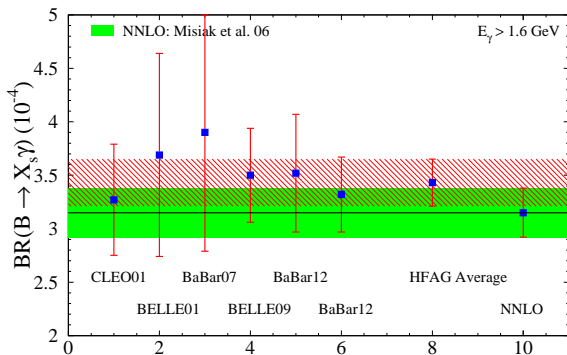
non-parametric uncertainties

- μ_b variation
- $\mathcal{O}(\alpha_S^3, \alpha_{em}^2, \alpha_{em}\alpha_S)$ to C_A
- M_B^2 / M_W^2
- $M_t^{\text{OS}} \rightarrow m_t^{\overline{\text{MS}}}$ transition

Status:

$$\mathcal{B}(\bar{B} \rightarrow X_s \gamma) |_{E_\gamma > 1.6 \text{ GeV}} = (3.15 \pm 0.23) \times 10^{-4}$$

[Misiak, Steinhauser'06; Misiak et al.'06]



5% non-perturbative

3% parametric

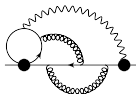
α_s, m_c, \dots

3% m_c interpolation

3% higher order

- update input parameters; fit to semi-leptonic B decay data in kinetic scheme [Gambino,Schwanda'13]

- m_c interpolation: $m_c = 0$ result for \longrightarrow



[Czakon,Fiedler,Huber,Misiak,Schutzmeier,Steinhauser]

- 3- and 4-body final state contributions to the NNLO interferences among Q_1 , Q_2 and Q_8 (BLM approximation) [Ligeti,Luke,Manohar,Wise'99;

Ferrogli,Haisch'10; Misiak,Poradziński'11]

- Four-loop $Q_1, \dots, Q_6 \rightarrow Q_8$ anomalous dimensions [Czakon,Haisch,Misiak'07]

- LO $b \rightarrow sq\bar{q}\gamma$ from $Q_{1,\dots,6}$ [Kamiński,Misiak,Poradziński'12]

- non-pert. $\mathcal{O}(\alpha_s\Lambda^2/m_b^2)$ to (Q_7, Q_7) interference [Ewerth,Gambino,Nandi'10]

- non-perturbative contributions [Benzke, Lee, Neubert, Paz'10]

- ...

\Leftrightarrow shifts of order $\leq \pm 2\%$

Updated result for $\mathcal{B}(\bar{B} \rightarrow X_s\gamma)$ soon

- $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$ to NNLO QCD and NLO EW
- theory uncertainty negligible
- main uncertainty source: f_{B_s} and CKM
- $\overline{\mathcal{B}}_{s\mu}^{\text{exp}} = (2.9 \pm 0.7) \times 10^{-9}$
 $\overline{\mathcal{B}}_{s\mu}^{\text{th}} = (3.65 \pm 0.23) \times 10^{-9}$
 $\overline{\mathcal{B}}_{d\mu}^{\text{exp}} = (3.6_{-1.4}^{+1.6}) \times 10^{-10}$
 $\overline{\mathcal{B}}_{d\mu}^{\text{th}} = (1.06 \pm 0.09) \times 10^{-10}$
- also predictions for $\overline{\mathcal{B}}_{se}, \overline{\mathcal{B}}_{s\tau}, \overline{\mathcal{B}}_{de}, \overline{\mathcal{B}}_{d\tau}$