

# Constraining the Axion-Photon Coupling with Massive Stars

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*(Mostly) Based on:*

**Phys.Rev.Lett. 110 (2013) 061101**

with A. Friedland and M.Wise

**Desy, July 2013**

# Outline

- ❑ Stars and particle physics
- ❑ Axion and Axion Like Particles (ALPs)
- ❑ Cast and the Horizontal Branch bounds on the axion-photon coupling
- ❑ Bound on the axion-photon coupling from massive stars
- ❑ Conclusions

# Stars as Laboratories

For 50 years stars have been excellent laboratories for **light, weakly interacting particles**. In fact, stars are sensitive to very rare processes, e.g.,

$$\gamma \rightarrow \nu + \bar{\nu}$$

and

$$e^+ + e^- \rightarrow \nu + \bar{\nu}$$

which are extremely rare but play a fundamental role in stellar cooling.

Examples of models strongly constrained by stars include majorons, extra-dimensional photons, novel baryonic or leptonic forces, unparticles, etc.

J. Bernstein et al., Phys. Rev. 132, 1227 (1963)

H.M. Georgi, S.L. Glashow, and S. Nussinov, Nucl. Phys. B193, 297 (1981)

A. Friedland and M. Giannotti, Phys. Rev. Lett. 100, 031602 (2008)

Grifols and E. Masso, Phys. Lett. B 173, 237 (1986)

S. Hannestad, G. Raffelt, and Y.Y. Wong, Phys. Rev. D76, 121701 (2007)

A considerable improvement in astrophysical observations analysis is now leading to revision and improvement of stellar bounds

# Stars as Laboratories

For 50 years stars have been excellent laboratories for **light, weakly interacting particles**. In fact, stars are sensitive to very rare processes, e.g.,

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$$\gamma \rightarrow \nu + \bar{\nu}$$

and

$$e^+ + e^- \rightarrow \dots$$

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fundamental reactions

Recently, Raffelt and collaborators have been revising the old stellar bounds on neutrino magnetic moment and axion–electron coupling, providing more reliable bounds which include confidence levels.

[ *N. Viaux et. al. 2013, N. Viaux, Ph.D. thesis* ]

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Wong, Phys. Rev. D76, 121701 (2007)

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# Axion and Axion-Like Particles (ALPs)

A particularly interesting example of light, weakly interacting particle is the **axion**, hypothetical particle whose existence is a prediction of the Peccei-Quinn solution of the **Strong CP problem** and prominent dark matter candidate.

Peccei and Quinn (1977),  
Weinberg (1978),  
Wilczek (1978)

Preskill, Wise and Wilczek (1983)  
Abbott and Sikivie (1983)  
Dine and Fischler (1983)

In the *standard* (QCD) *axion models* the axion interaction with matter and photons and the axion mass are related through the Peccei-Quinn constant  $f_a$

$$L_{\text{int}} = -i \frac{C_i m_i}{f_a} a \bar{\psi} \gamma_5 \psi - \frac{1}{4} g_{a\gamma} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$m_a \approx \frac{6 \text{ eV}}{f_a / 10^6 \text{ GeV}}$$

Axion-photon  
coupling

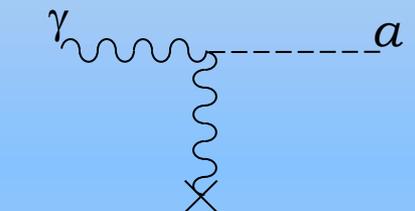
$$g_{a\gamma} = \xi \frac{\alpha_{em}}{2\pi f_a}$$

In more general models (**ALPs**), couplings and the mass are unrelated.

# Axions and Stellar Evolution

Light ALPs can be produced in stars through various mechanisms, e.g.

**Primakoff conversion**



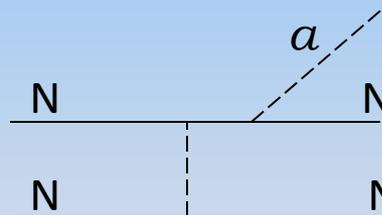
Relevant in He-burning stars

**Compton scattering**



Relevant in RG and WD

**Nucleon Bremsstrahlung**



Relevant in SN and neutron stars

The emission of axions could lead to an *overly efficient energy drain*, inconsistent with observations. This leads to bounds on the axion couplings with photons, electrons and nuclei.

# Axion detection: the role of the Axion-Photon Coupling

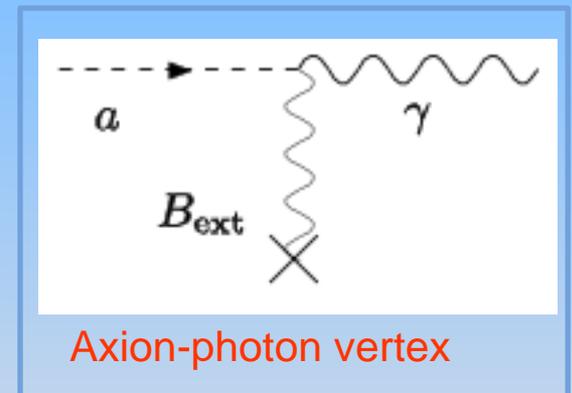
Most of the modern axion searches are based on the microwave cavity detection proposed by **P. Sikivie**, which relies on the axion-photon coupling. Axions can be converted into photons in an external magnetic field. These bounds depend on the axion mass.

A strong terrestrial bound on the axion (ALPs)-photon coupling comes from the Cern Axion Solar Telescope (**CAST**)

$$g_{a\gamma} \leq 0.88 \times 10^{-10} \text{GeV}^{-1}$$

The bound is however weakened at masses  $>0.02\text{eV}$  (in the QCD –axion region)

P. Sikivie, Phys.Rev.Lett.  
51, 1415 (1983)



# Experimental Axion (and ALPs) Search: CAST

## ■ Vacuum Phase

$$m_a \leq 0.02 \text{ eV}$$

- $g_{a\gamma}$  (95%)  $< 0.88 \times 10^{-10} \text{ GeV}^{-1}$
- Phys.Rev.Lett.94:121301, 2005
- JCAP 04 (2007) 010

## ■ $^4\text{He}$ Phase

$$0.02 \text{ eV} \leq m_a \leq 0.39 \text{ eV}$$

- $g_{a\gamma}$  (95%)  $< 2.2 \times 10^{-10} \text{ GeV}^{-1}$
- JCAP 02 (2009) 008

## ■ $^3\text{He}$ Phase: first results

$$0.39 \text{ eV} \leq m_a \leq 0.65 \text{ eV}$$

- $g_{a\gamma}$  (95%)  $< 2.3 \times 10^{-10} \text{ GeV}^{-1}$
- Phys.Rev.Lett. 107:261302, 2011

## ■ $^3\text{He}$ Phase: preliminary results

$$0.65 \text{ eV} \leq m_a \leq 1.18 \text{ eV}$$

- Publication in preparation

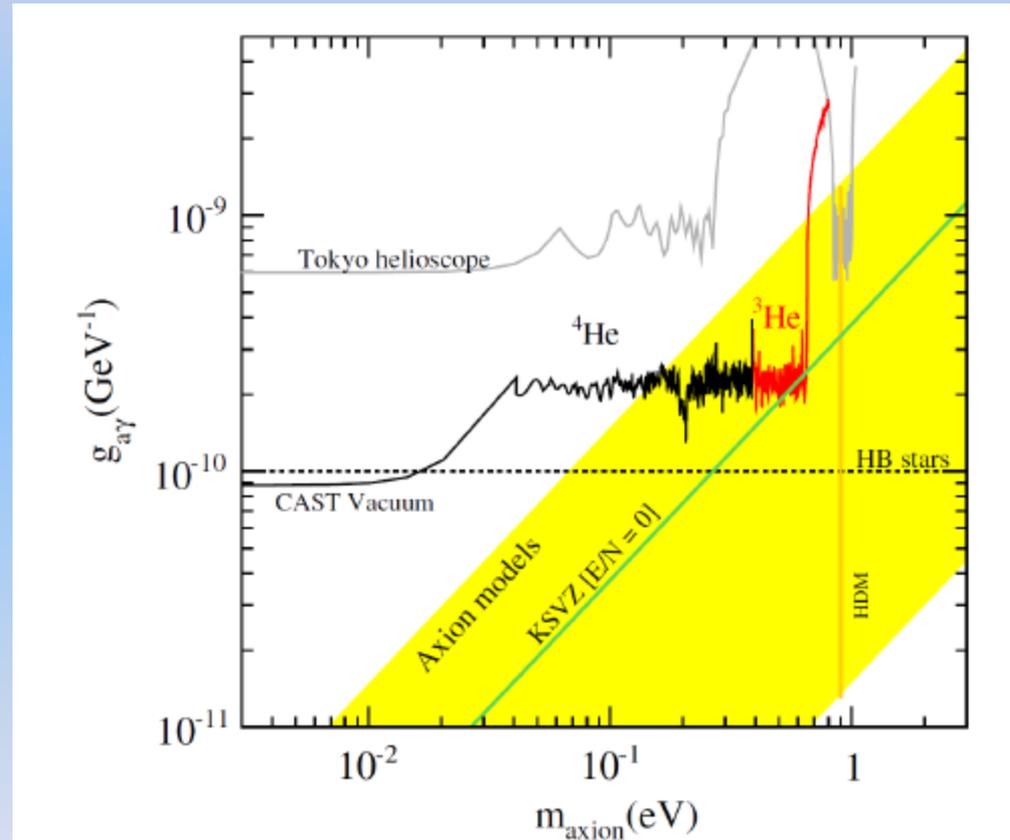
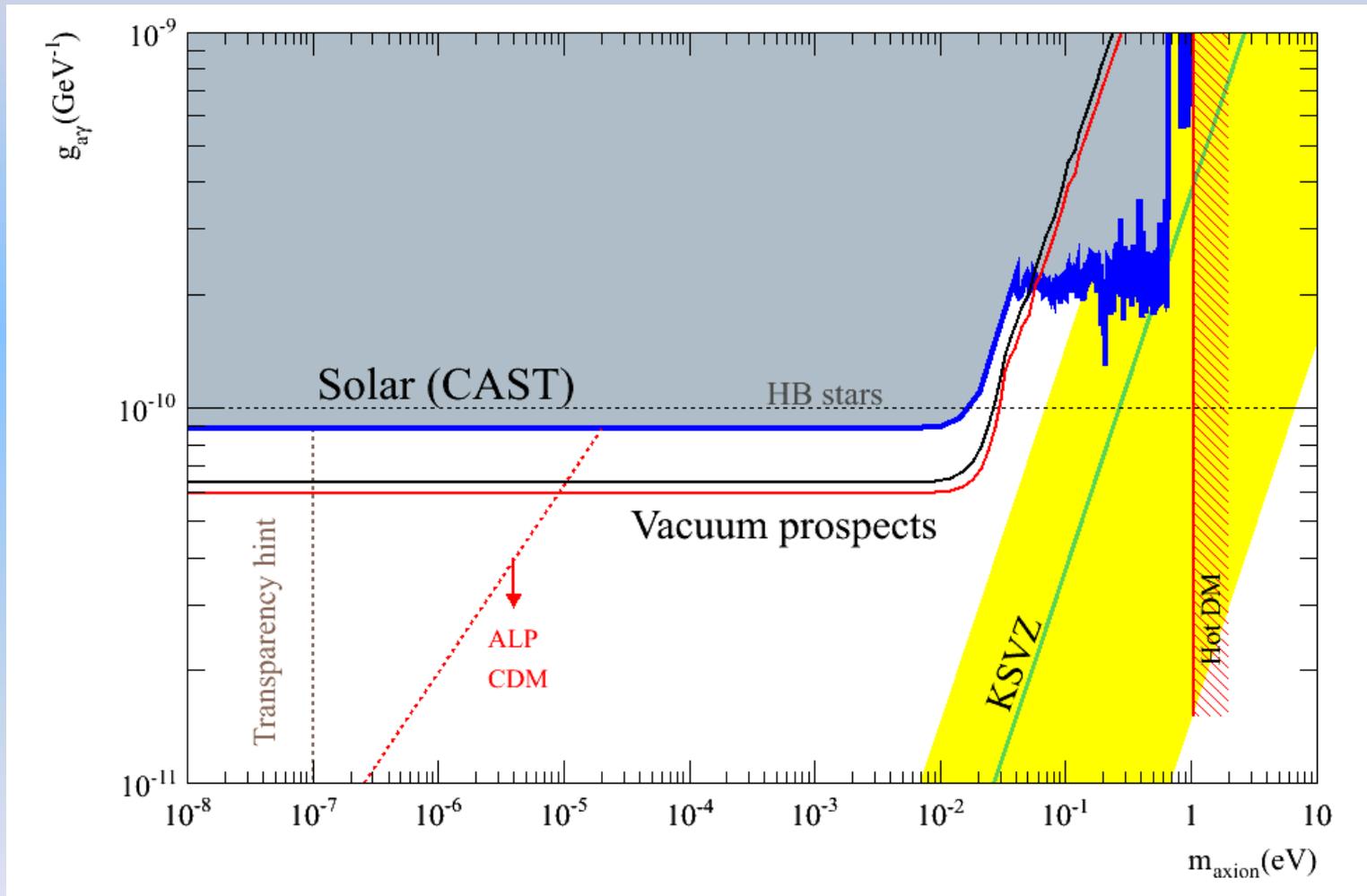


Image and data from Julia Vogel (LLNL), CAST

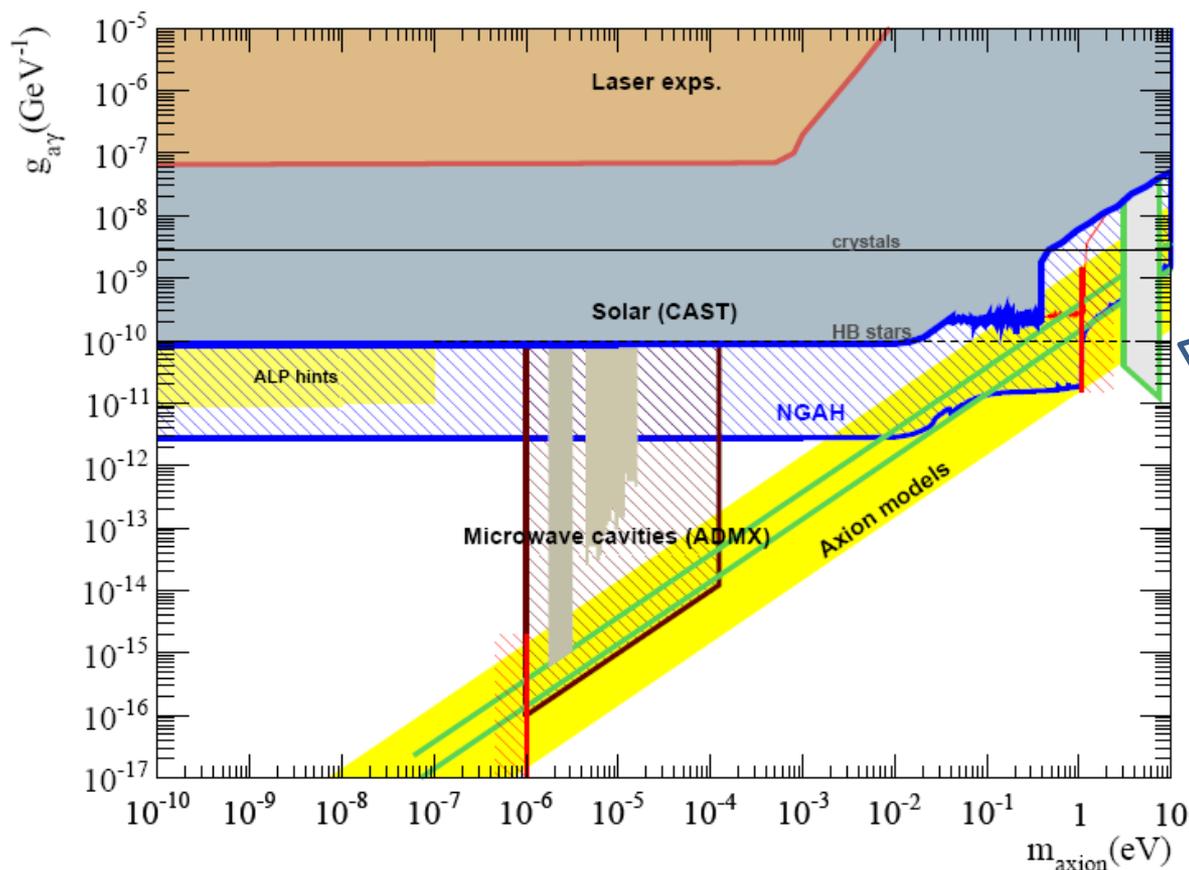
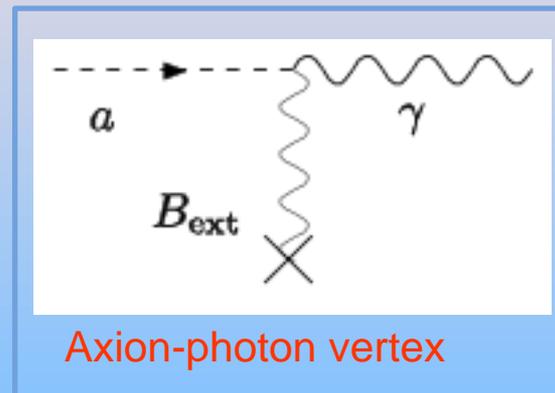
# Experimental Axion (and ALPs) Search: CAST



From Biljana Lakić, CAST, PATRAS 2013, Mainz

# Experimental Axion (and ALPs) Search

The **Next Generation Axion Helioscopes (NGAH)** are expected to improve the bounds by over an order of magnitude.



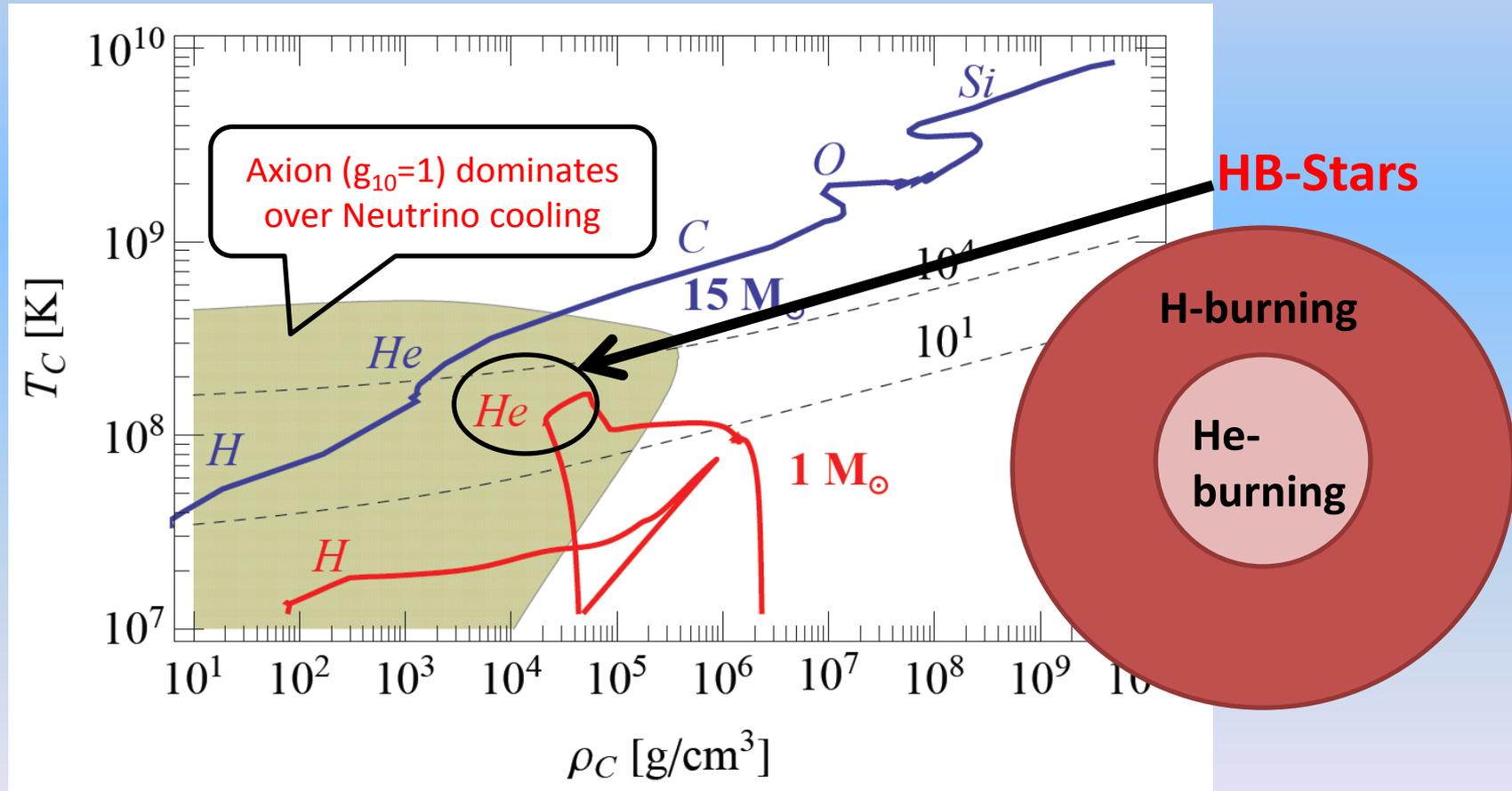
**HB-bound,  $g_{10}=1$ .**

Raffelt and Dearborn  
PRD 36 (1987)

From Biljana Lakić,  
CAST, PATRAS 2013,  
Mainz

# The HB bound on the axion-photon Coupling

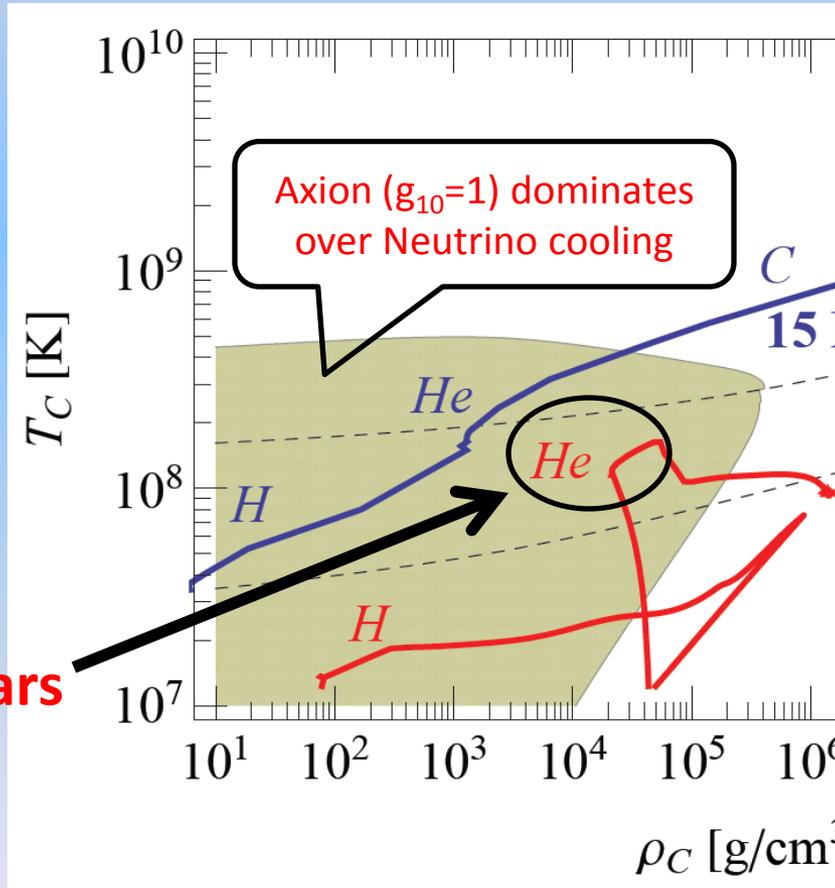
Axions can be produced in the core of a star from photons interacting with the electric field of the nuclei (**Primakoff process**).



Friedland, Giannotti, Wise, **Phys. Rev. Lett.** **110**, 061101 (2013)

# The HB bound on the axion-photon Coupling

Axions can be produced in the core of a star by interacting with the electric field of the



Friedland, Giannotti, Wise, *Phys. Rev.*

Axions coupled too strongly to photons would speed up the consumption of He in the HB star core and reduce the HB lifetime.

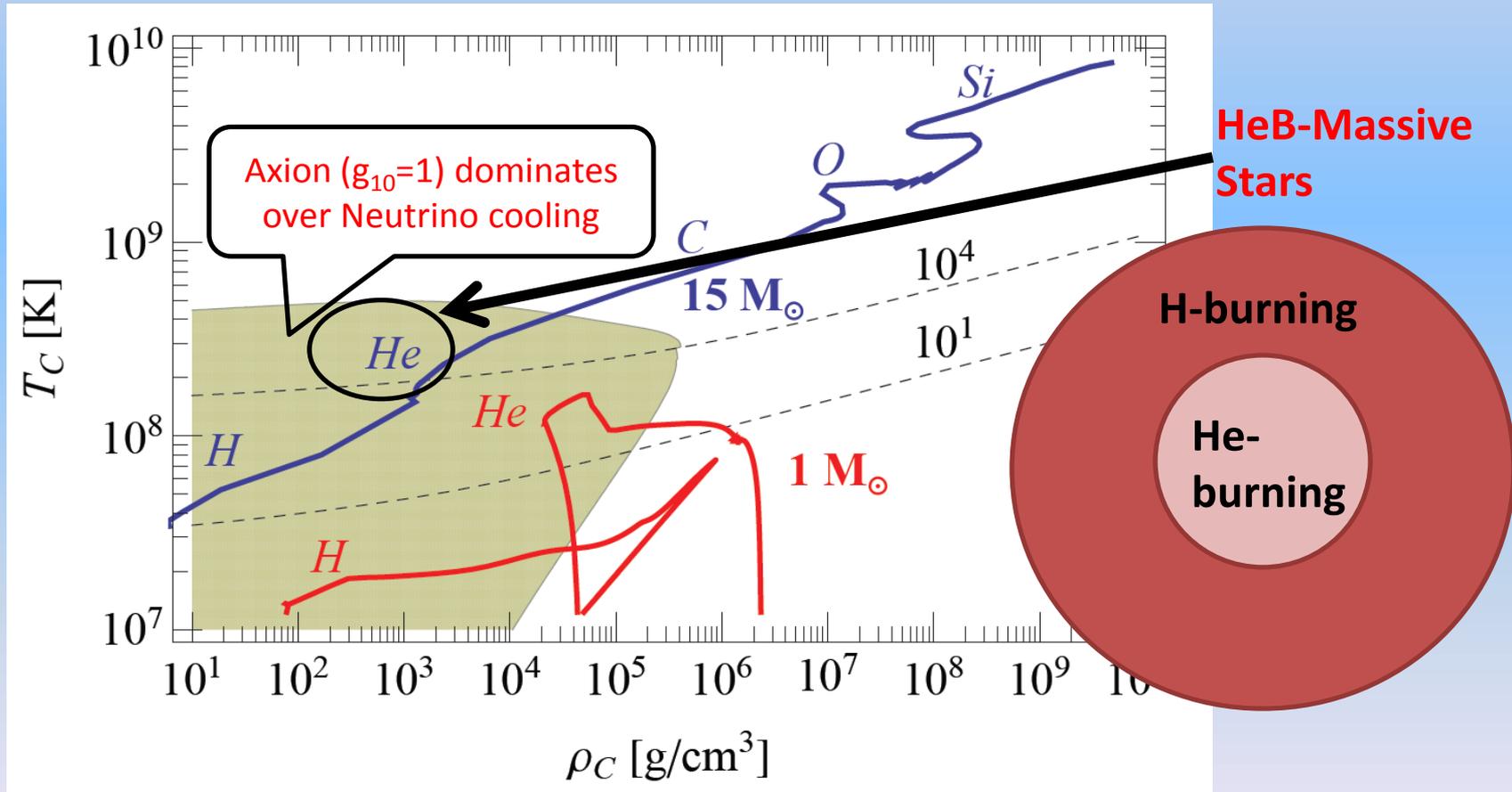
Observational constraints come from the comparison of the number of HB vs. RG stars.

The bound  $g_{ay} \leq 10^{-10} \text{GeV}^{-1}$  [G.G. Raffelt and D.S.P. Dearborn, *Phys. Rev. D36, 2211(1987)*] (which is the one currently reported in the PDG) corresponds to a reduction of the HB star lifetime by 30%.

The bound applies to axion masses  $< 30 \text{ keV}$  or so.

# Axion and He-burning massive stars

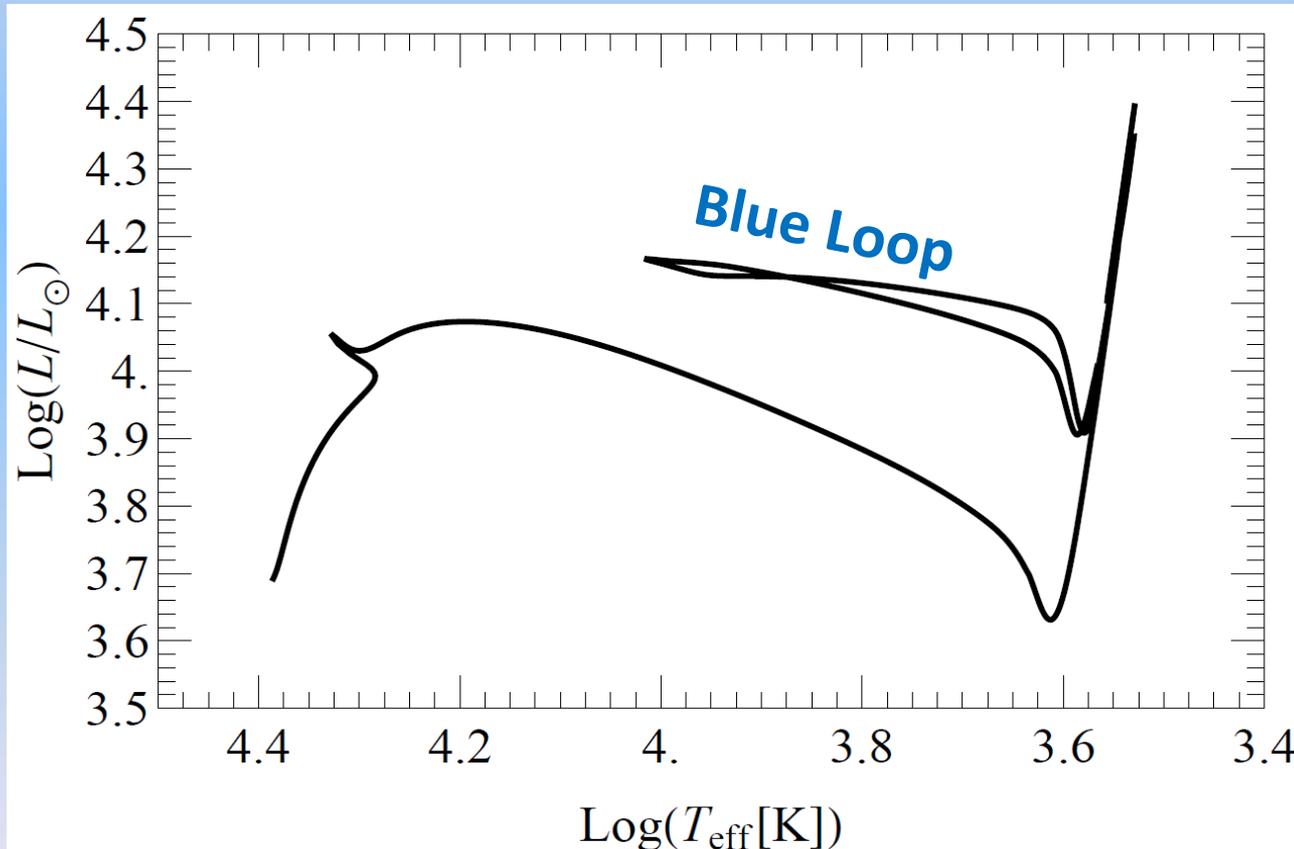
However, the currently strongest bound on the axion-photon coupling comes from the analysis of the evolution of He-burning massive stars



Friedland, Giannotti, Wise, **Phys. Rev. Lett.** **110**, 061101 (2013)

# The Blue Loop

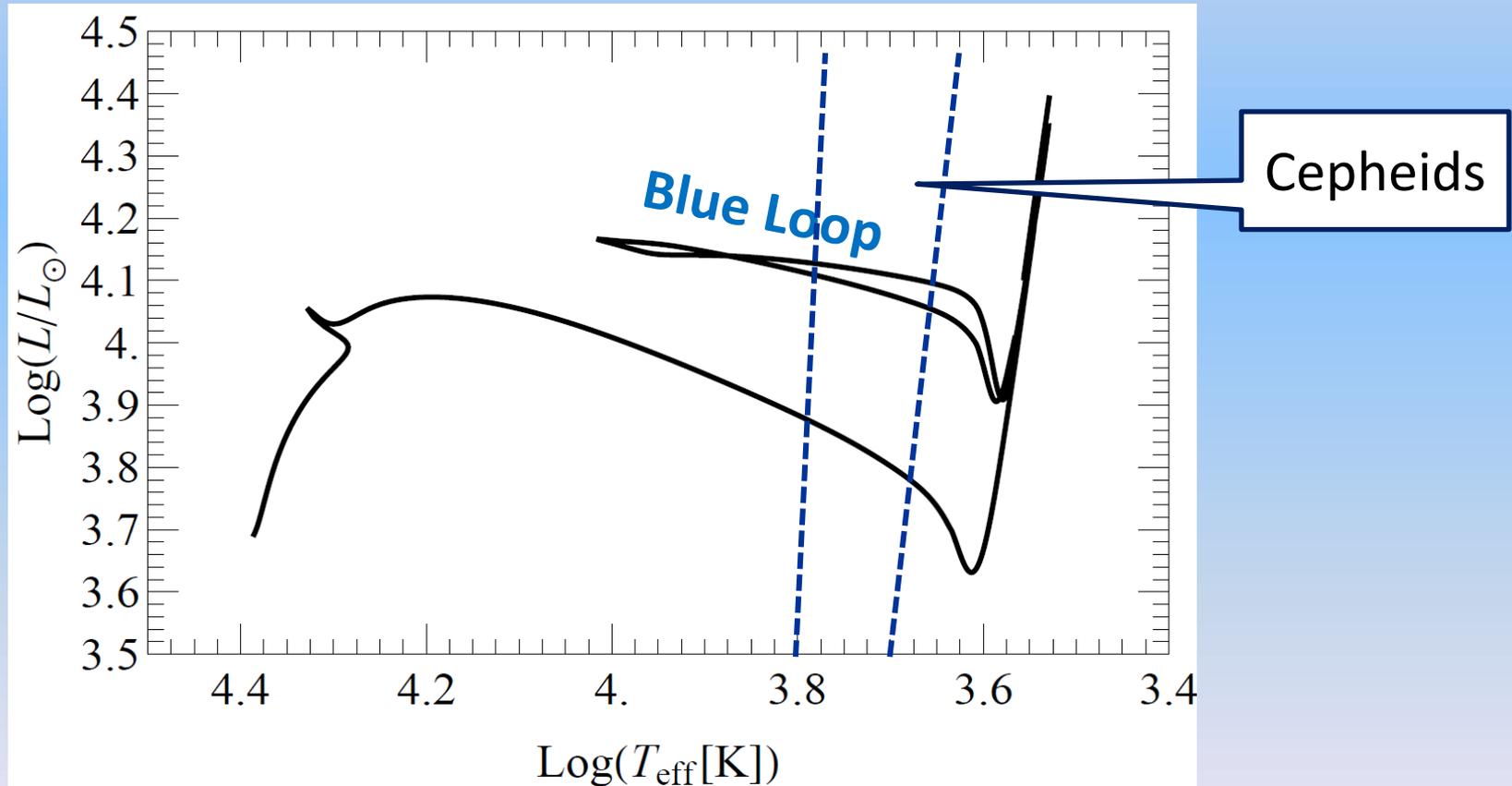
The HR diagram shows the luminosity v.s. surface temperature of a star. The **blue loop** is a prominent feature of the evolution of a massive star.



Simulations for a  $9.5M_{\odot}$ , solar metallicity, from main sequence to end of He-burning.  
**MESA** (Modules for Experiments in Stellar Astrophysics), Paxton et al. *ApJ Suppl.* **192** 3 (2011)  
[arXiv:1009.1622]

# The Blue Loop and the Cepheids

The blue loops are necessary to explain the existence of the Cepheids

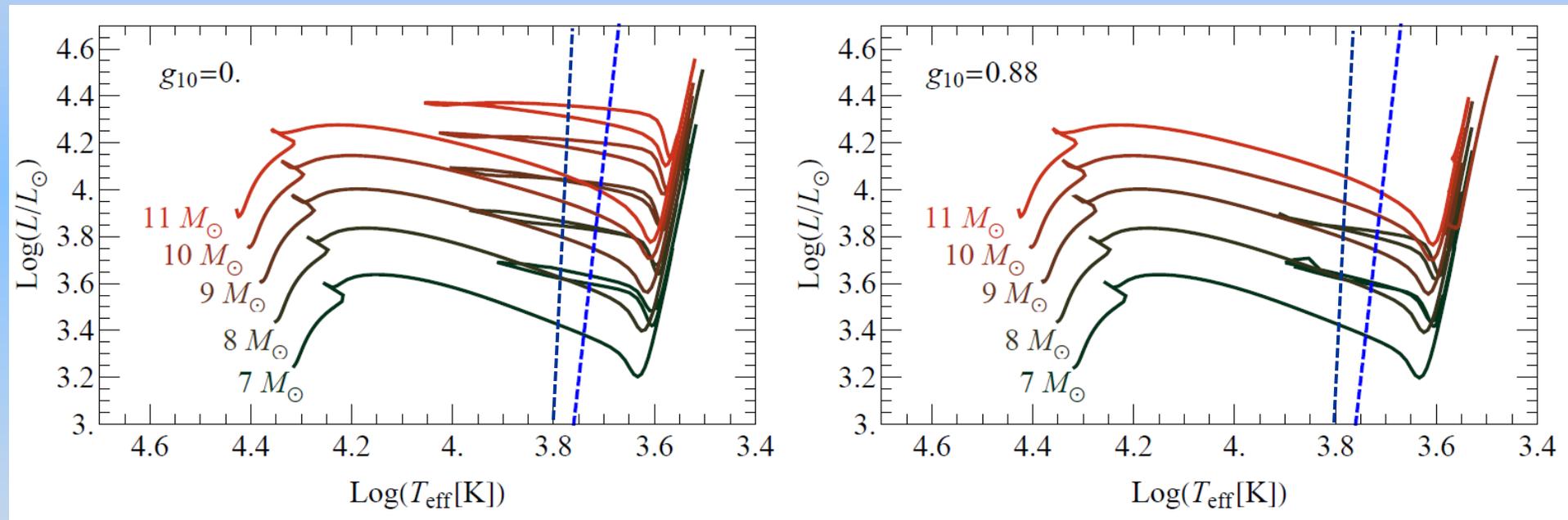


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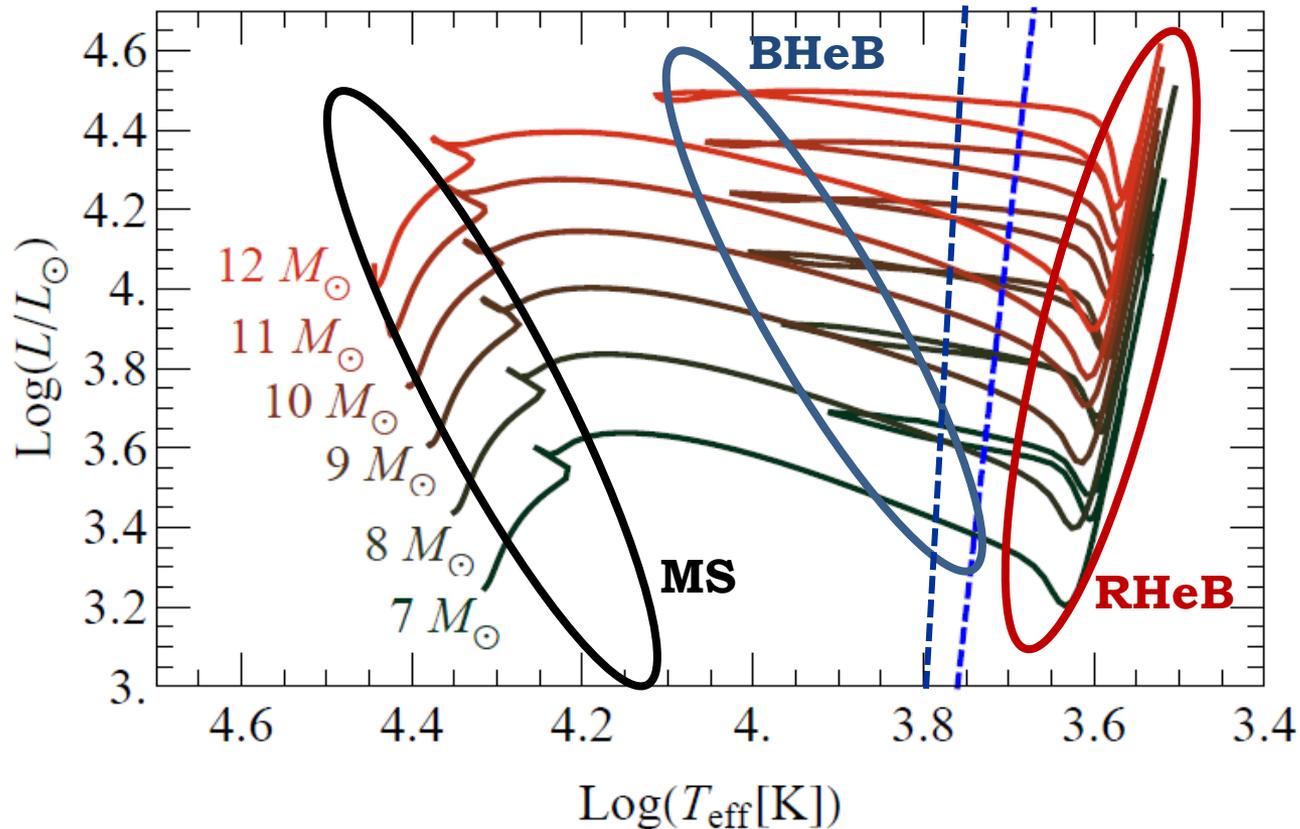
# Axions effects on the Blue Loop

The value  $g_{10}=0.88$  corresponds to the current CAST bound on the axion-photon coupling.



A value of  $g_{10}=0.8$  would provide qualitative changes in the stellar evolution. In particular, it would eliminate the blue loop stage of the evolution, leaving one without an explanation for the existence of Cepheid stars in a broad range of pulsation periods.

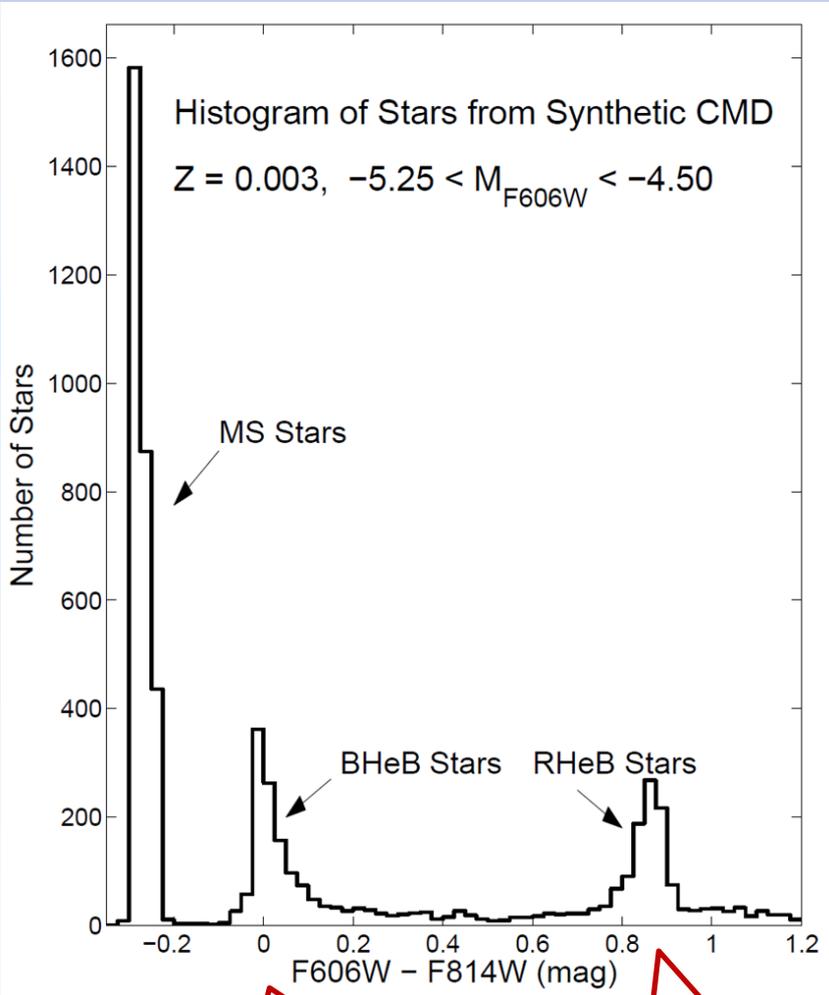
# Observations of massive stars: Main, Blue and Red Sequences



Most of the star life-time is spent in one of the three sequences: the **Main Sequence** (central H-burning), the **Red central He-Burning** sequence, and the **Blue central He-Burning** sequence

Simulations of evolution in H-R diagram of stars with solar metallicity, from main sequence to end of He-burning. [MESA]

# Observable evolutionary Phases: Central H- and He-burning



Blue stars have been observed for many decades and measurements are very accurate.

The contamination from MS stars transitioning to BHeB is conservatively estimated to be less than 10% (Dohm-Palmer & Skillman 2002).

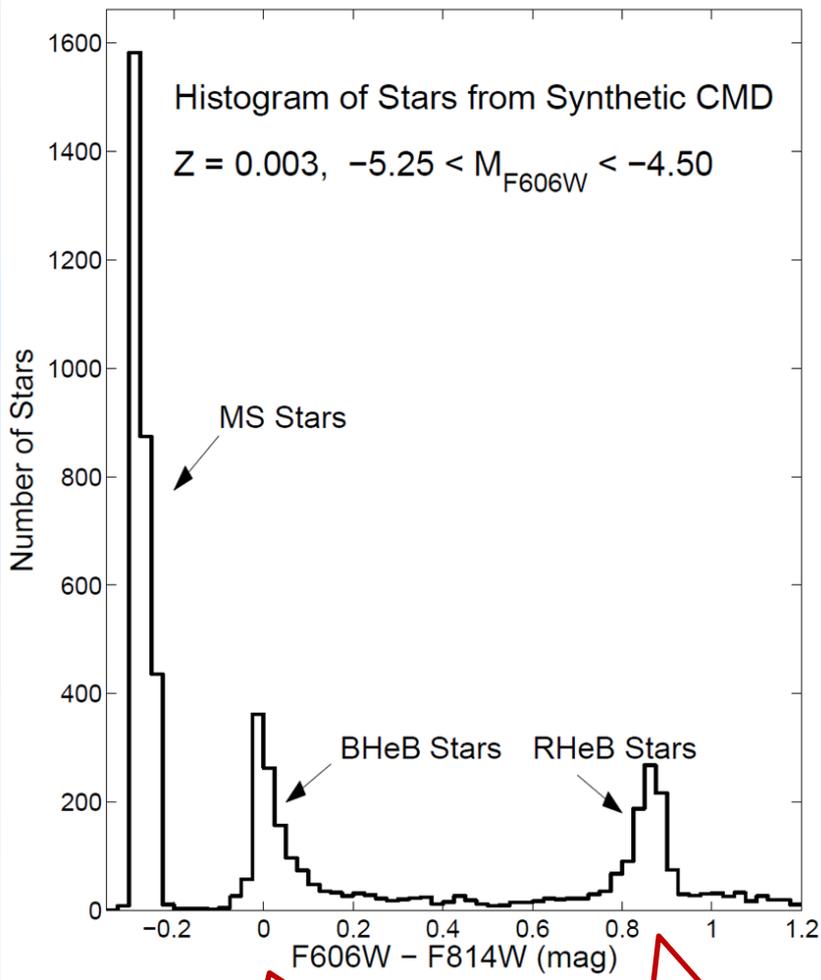
The complete disappearance of all blue stars in certain luminosity regions is physically unacceptable.

From Kristen B. W. McQuinn et. al.,  
*Astrophys.J.* 740 (2011)

Log T[k]=4

Log T[k]=3.7

# Observable evolutionary Phases: Central H- and He-burning



## Result:

A value of  $g_{10}$  above 0.8 would be incompatible with the current observations of HeB sequences.

This analysis provides the *strongest bound to date* on the axion-photon coupling.

A. Friedland, M.G., and M. Wise,  
**Phys. Rev. Lett. 110, 061101 (2013)**

See also

G. Raffelt, <http://physics.aps.org/articles/v6/14>

**Astrophys.J. 740 (2011)**

Log T[k]=4

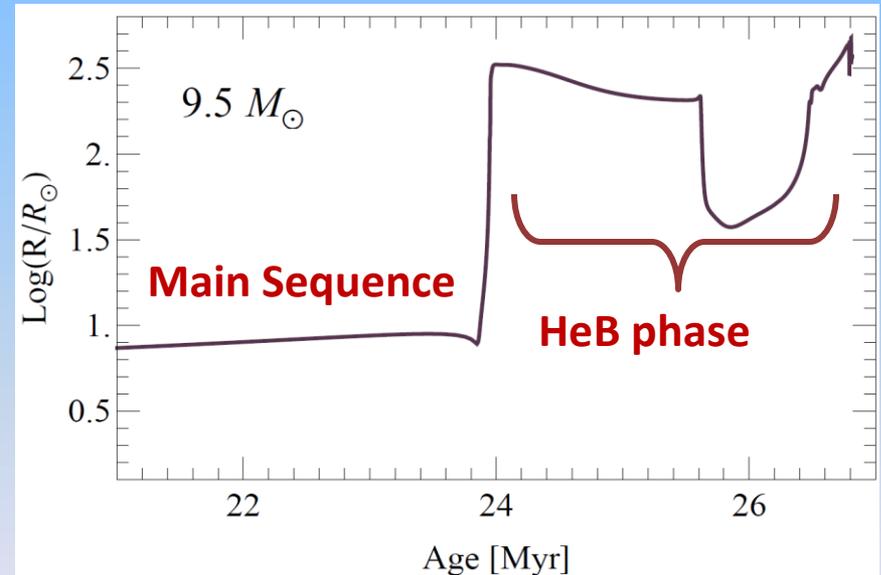
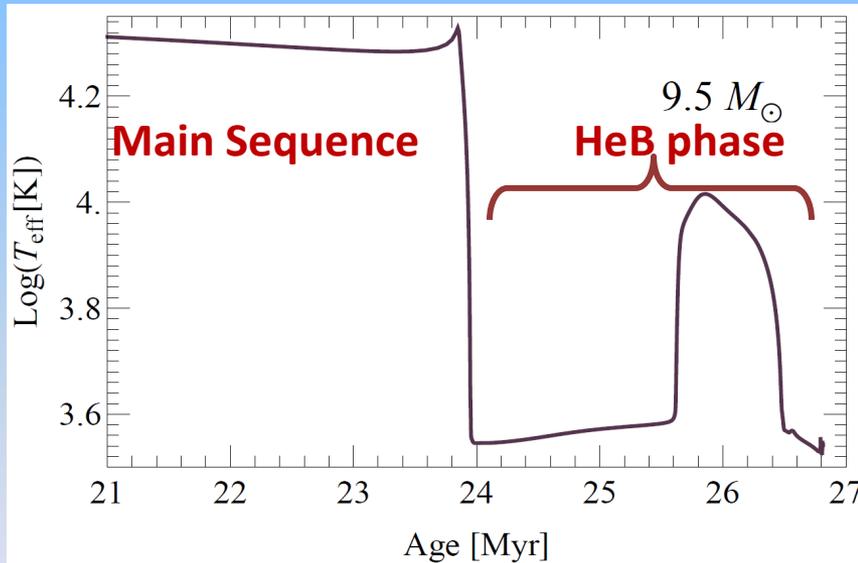
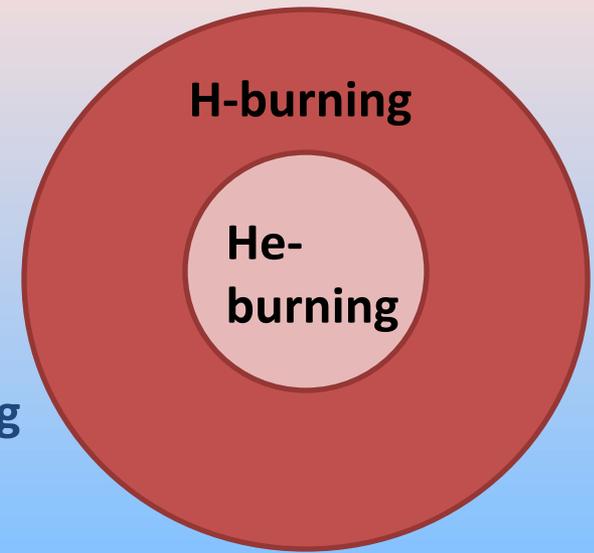
Log T[k]=3.7

# Massive stars and the Blue Loop



H-burning phase  
(Main Sequence)

He-burning phase

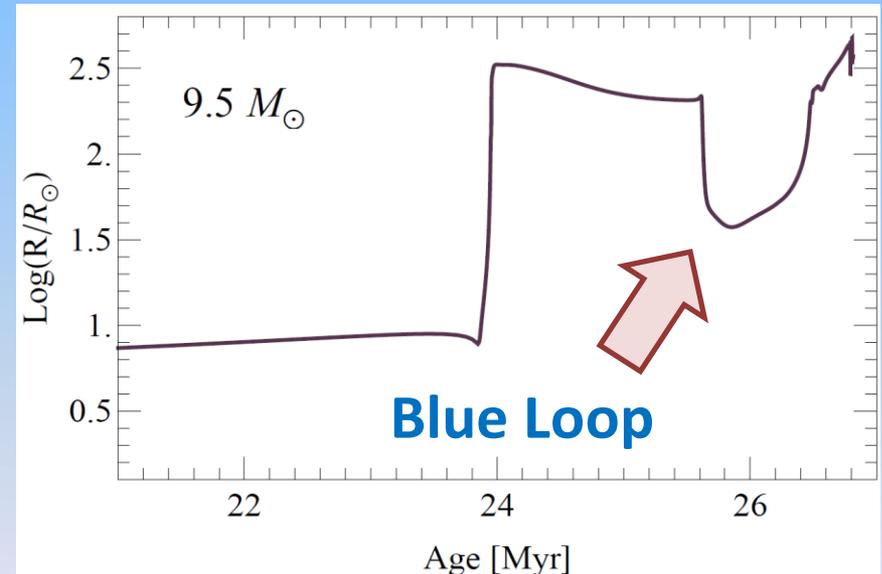
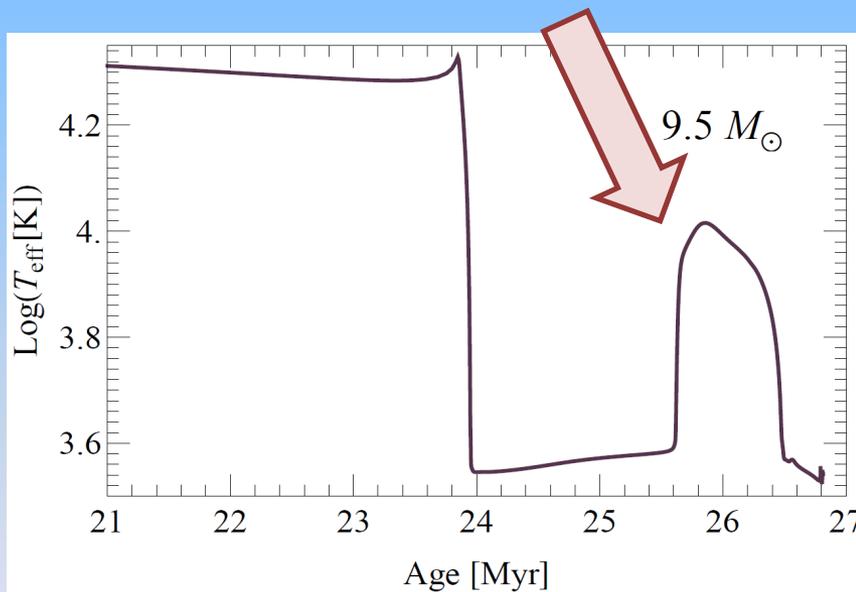


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[arXiv:1009.1622]

# The Blue Loop

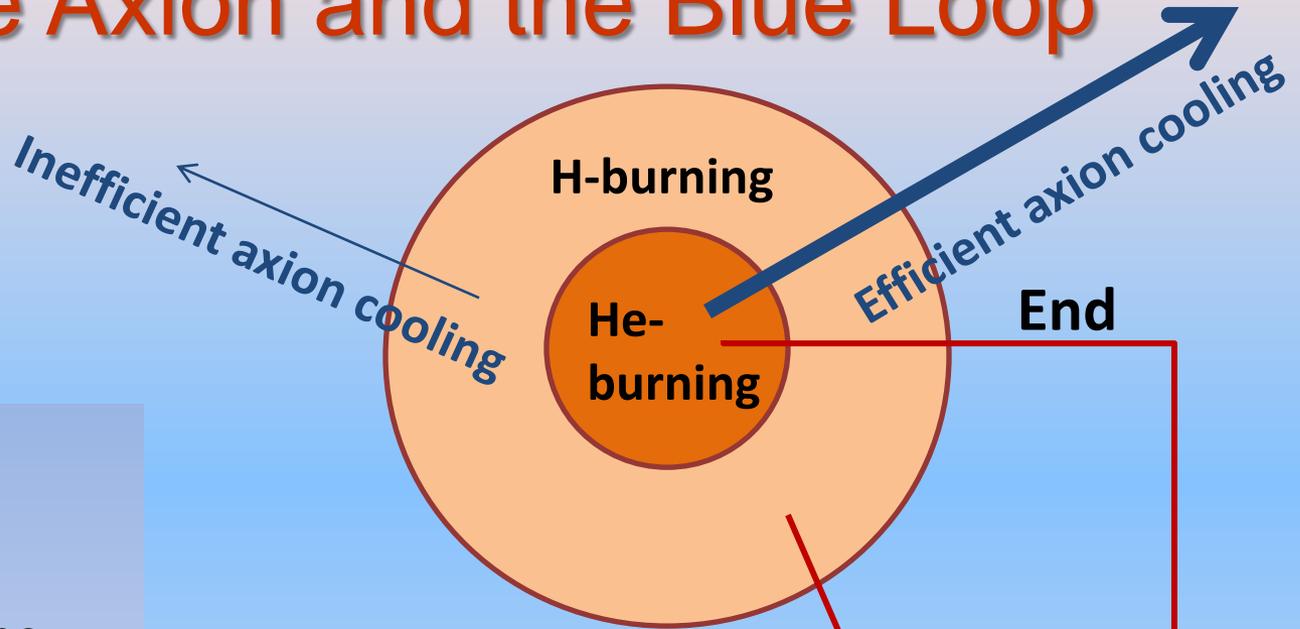
The journey of the star toward the hotter regions and back is called the **blue loop**. It happens for stars of a few solar masses during the He-burning stage.

## Blue Loop



Simulations for a  $9.5 M_{\odot}$ , solar metallicity, from main sequence to end of He-burning.  
**MESA** (Modules for Experiments in Stellar Astrophysics), Paxton et al. *ApJ Suppl.* **192** 3 (2011)  
[arXiv:1009.1622]

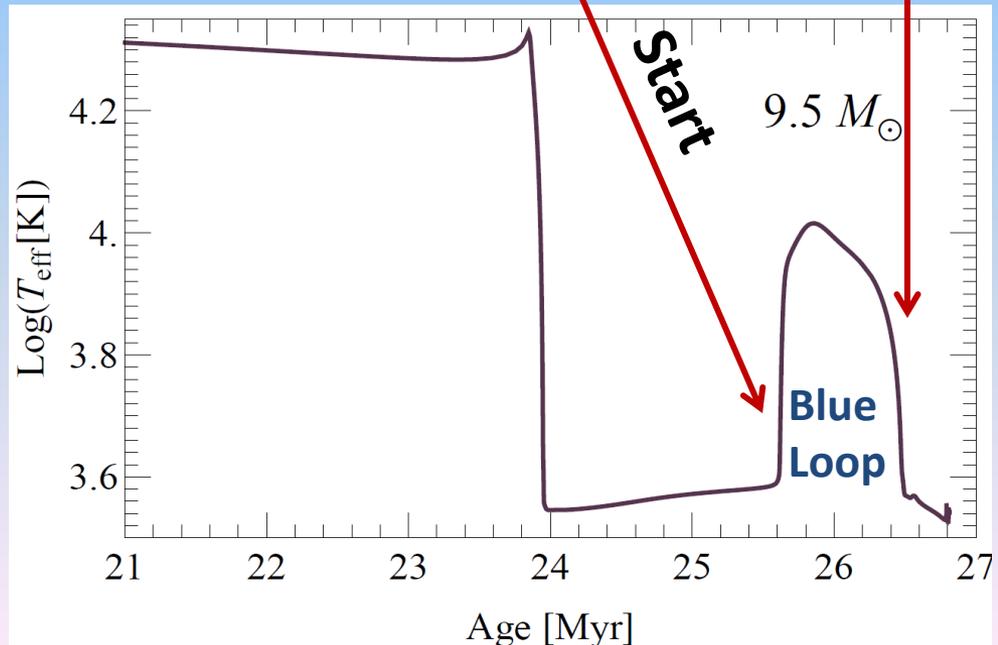
# The Axion and the Blue Loop



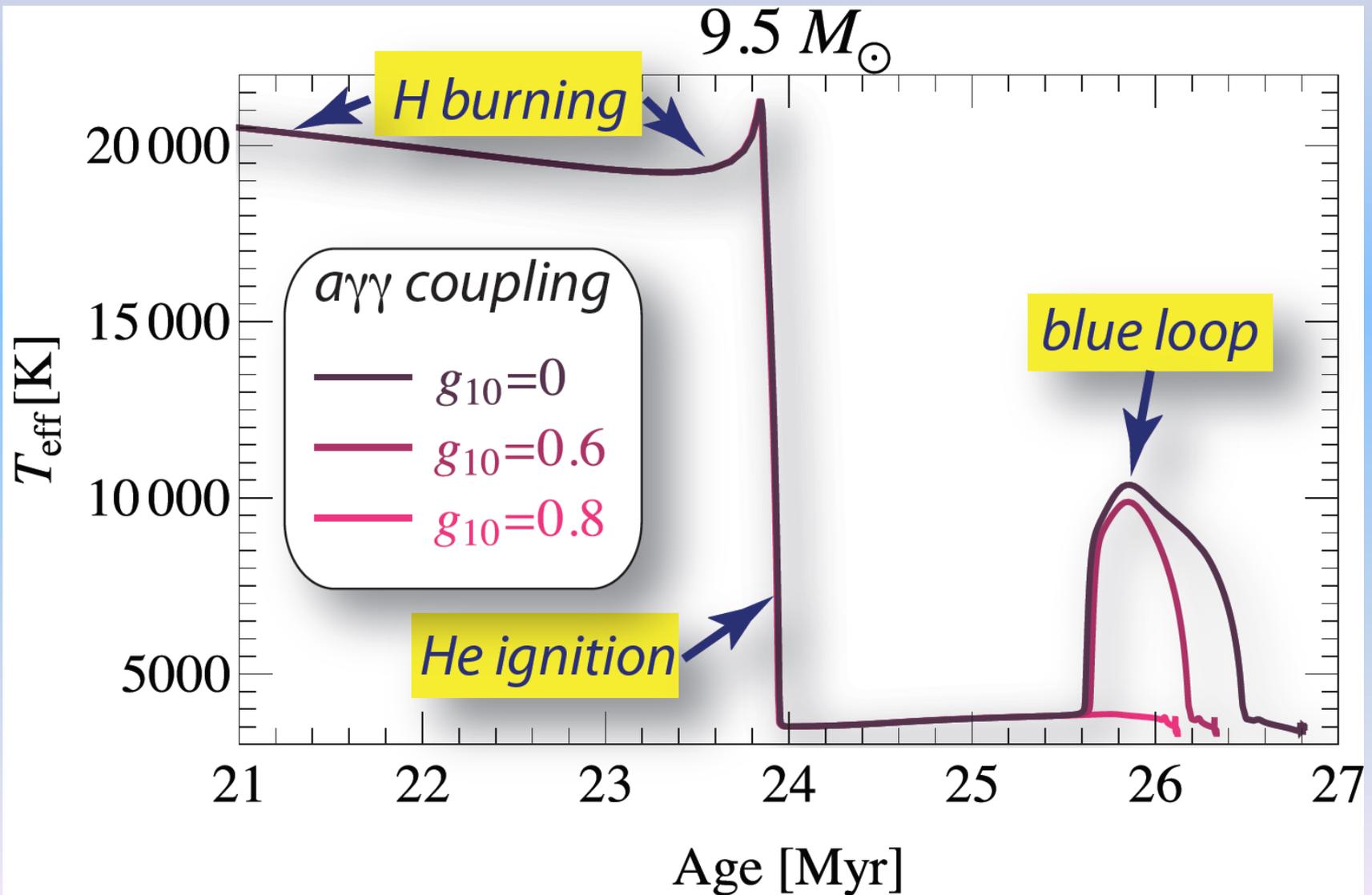
## Blue Loop:

the beginning of the blue loop is set by the H-burning shell time scale whereas the end is set by the He-burning core time scale

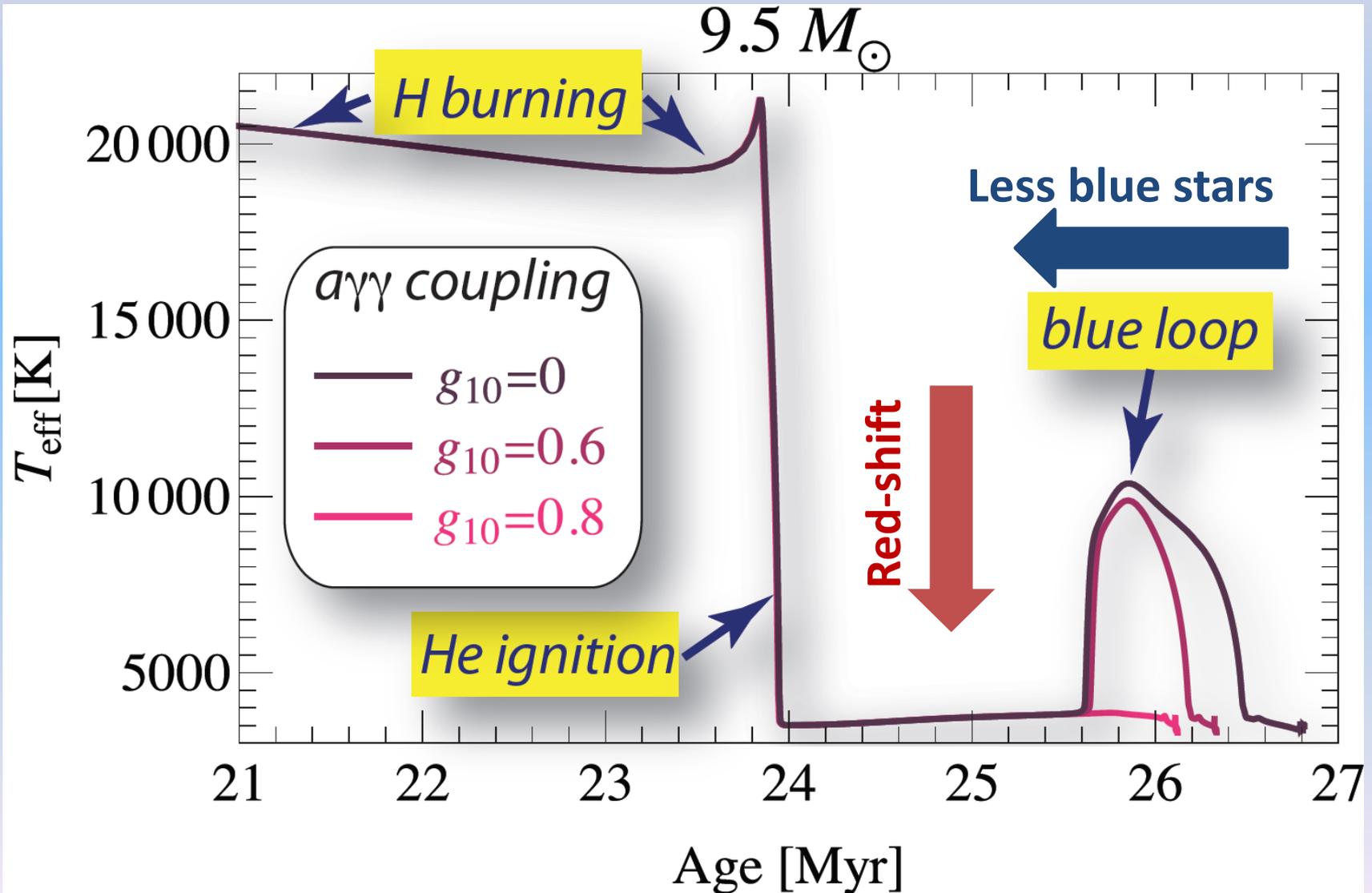
[Lauterborn et al., A&A 10, (1971),  
Kippenhahn and Weigert (1994)]



# The Axion and the Blue Loop



# The Axion and the Blue Loop

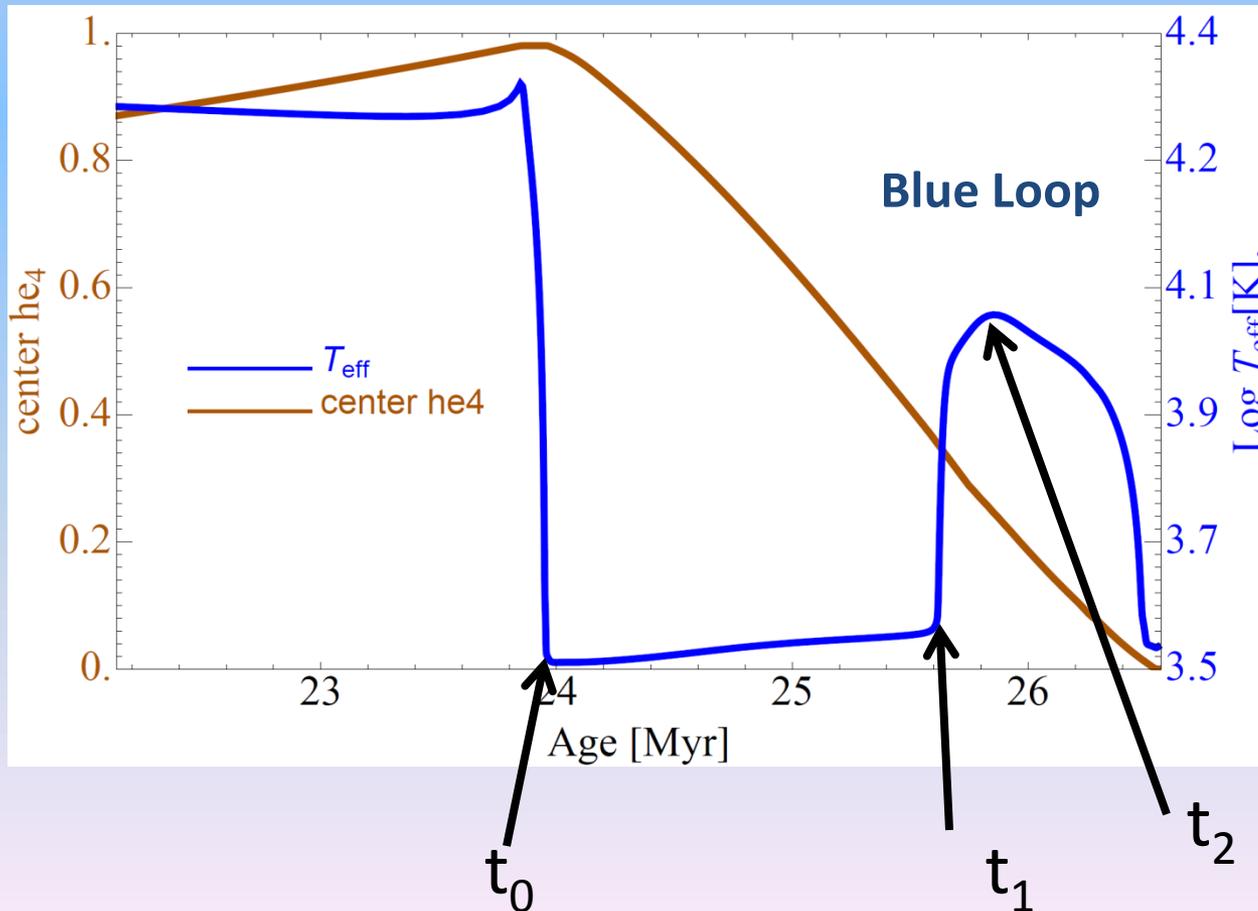


# Blue loop as a probe for Fundamental Physics

When the He content in the core reaches a certain lower value ( $t_2$ ), the surface temperature stops increasing and goes back to the red region of the HR diagram.

**Speeding up the core evolution would eliminate the blue loop phase**

[Lauterborn et al., A&A 10, (1971), E. Hofmeister, Z.F.A. 67, (1967)]



Characteristic times:

$t_0$ : beginning of the core He-burning phase

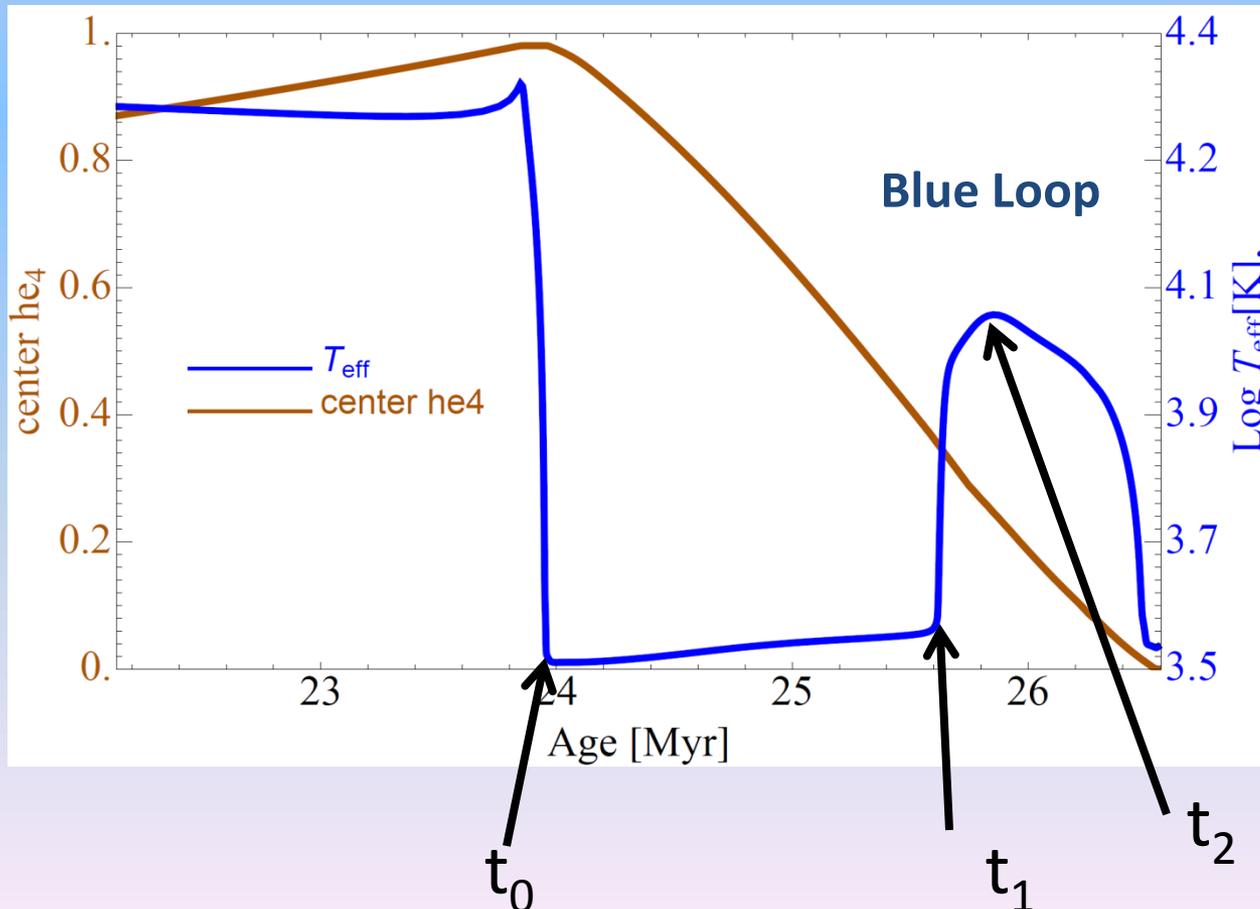
$t_1$ : beginning of the blue loop phase

$t_2$ : turning point

# Blue loop as a probe for Fundamental Physics

**A possible criteria:** *A novel cooling mechanism which changes the he-burning time from  $t_{\text{He}}$  to  $t'_{\text{He}}$  would eliminate the blue loop phase if:*

$$\frac{t'_{\text{He}}}{t_{\text{He}}} \leq \frac{t_1 - t_0}{t_2 - t_0}$$



This criteria is independent from convection prescriptions and other possible uncertainties.

A conservative requirement for stars around 9-11  $M_{\odot}$  is that

$$t'_{\text{He}} / t_{\text{He}} \leq 0.8$$

is not allowed by observations.

# Astro Labs: Massive vs. HB stars

HB and massive stars offer two different criteria for probing exotic processes which are efficient during the He-burning stage:

## HB stars:

Lower core temperature,  
higher core density;

Relatively low energy  
production during he-  
burning:

$$\varepsilon_x \approx 80 \text{ ergs g}^{-1} \text{ s}^{-1}$$

Observations require:

$$t'_{\text{He}} / t_{\text{He}} \geq 0.7$$

[cfr. PDG and G.Raffelt in Axions,  
Springer (2010), chapter 3]

## Massive stars:

Higher core temperature,  
lower core density;

High energy production  
during he-burning:

$$\varepsilon_x \approx \text{a few } 10^3 \text{ ergs g}^{-1} \text{ s}^{-1}$$

Observations require:

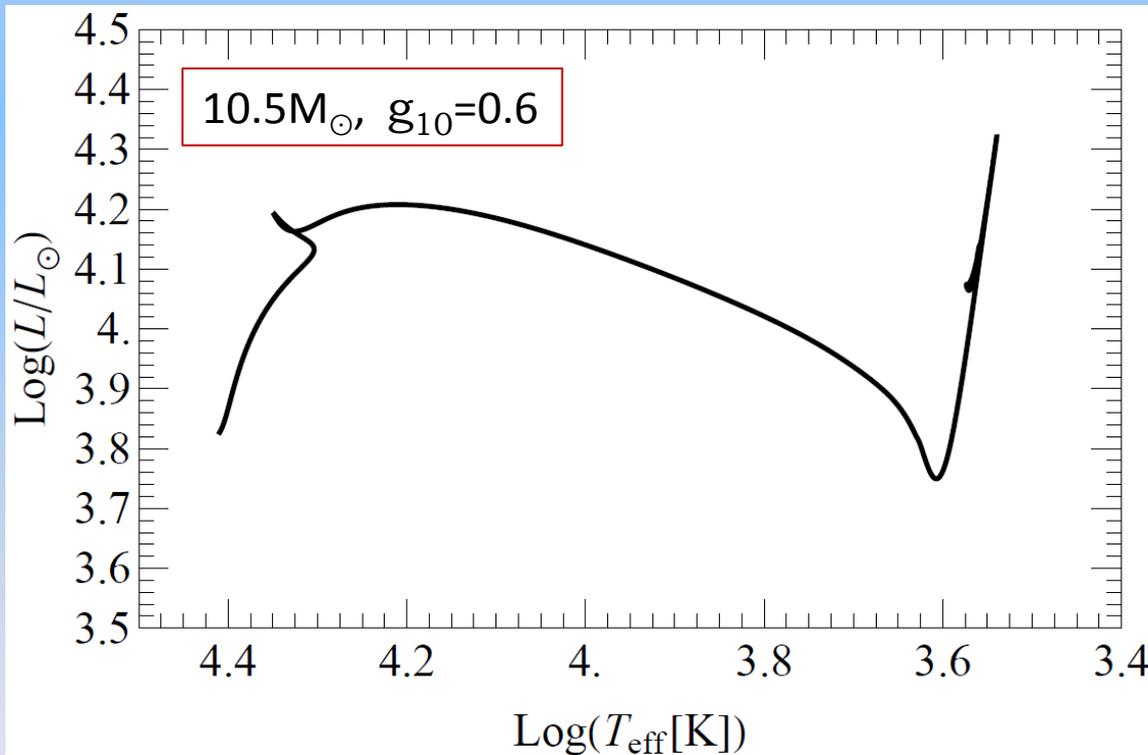
$$t'_{\text{He}} / t_{\text{He}} \geq 0.8$$

Massive stars are hotter and less dense. The maximal exotic cooling allowed by observation is larger than HB stars.

Exotic processes which are very sensitive to temperature are likely to be constrained more efficiently from massive stars. Otherwise, HB stars offer a better lab

# Even Stronger Bounds?

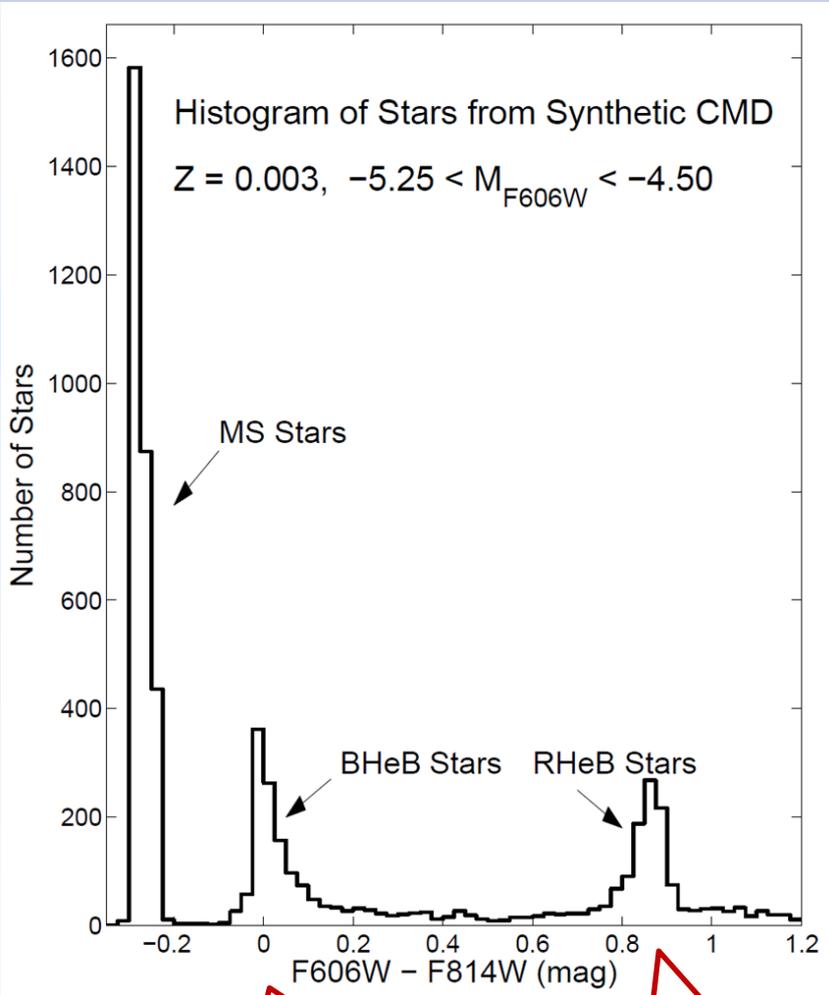
Other stars showed an even stronger response to axion cooling. (Friedland, M.G., Wise, work in progress).



Stars of mass between  $10.5M_{\odot}$  and  $11.5M_{\odot}$  don't show a loop already for  $g_{10}=0.6$ .

Stars of even larger mass are somewhat harder to model.

# Even Stronger Bounds?



Log T[k]=4

Log T[k]=3.7

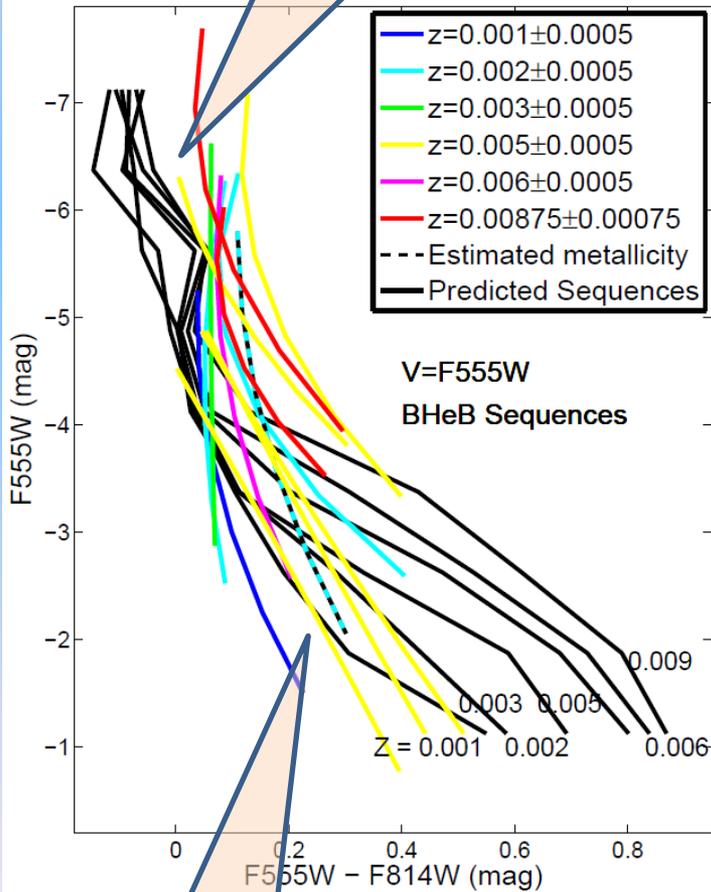
In addition, the requirement that we used to derive the  $g_{10}=0.8$  bound is conservative.

Given accurate counts, it may be possible to check whether the number of stars in the blue loop phase is reduced.

For example,  $g_{10} = 0.6$  would reduce the time a  $9.5 M_{\odot}$  star spends on the blue loop by a factor of two or so. To get the same sensitivity for  $g_{10}$  from solar-mass stars requires knowing the numbers of HB stars to a 10% precision.

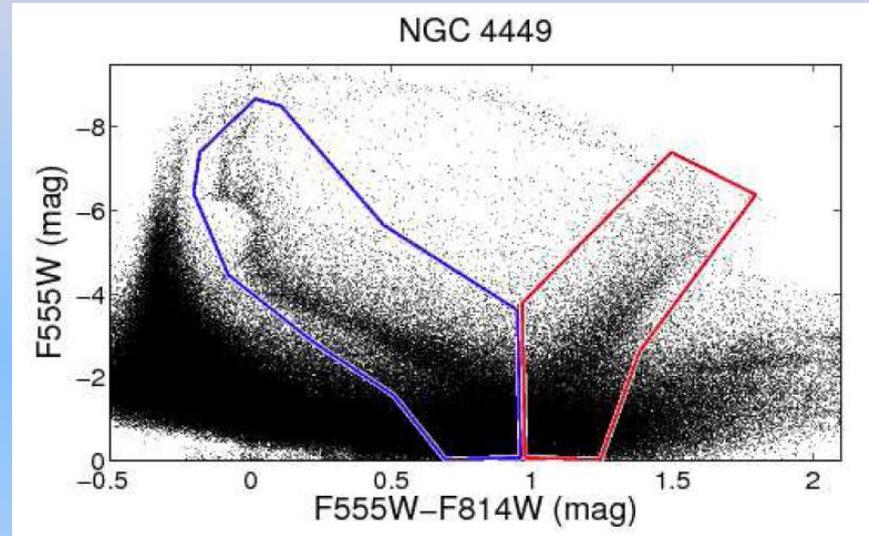
# Experimental Evidence for Blue Sequences

High mass  
(young cluster)

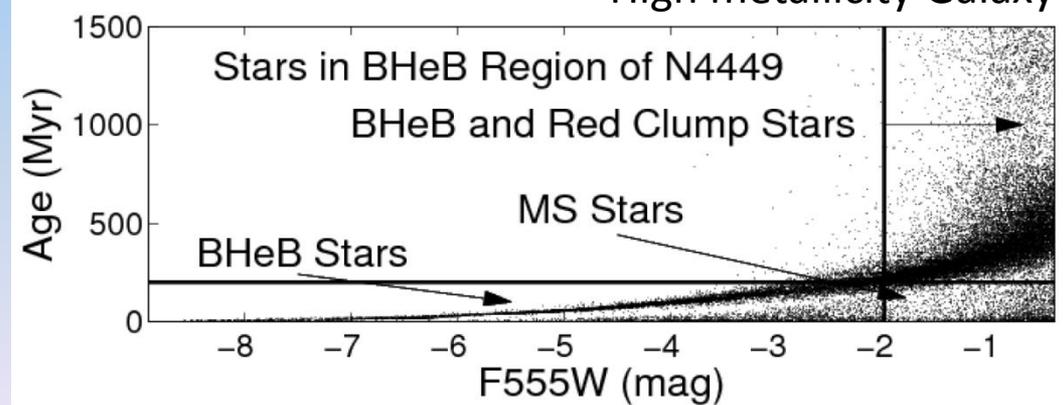


Low mass  
(old cluster)

V-I Color

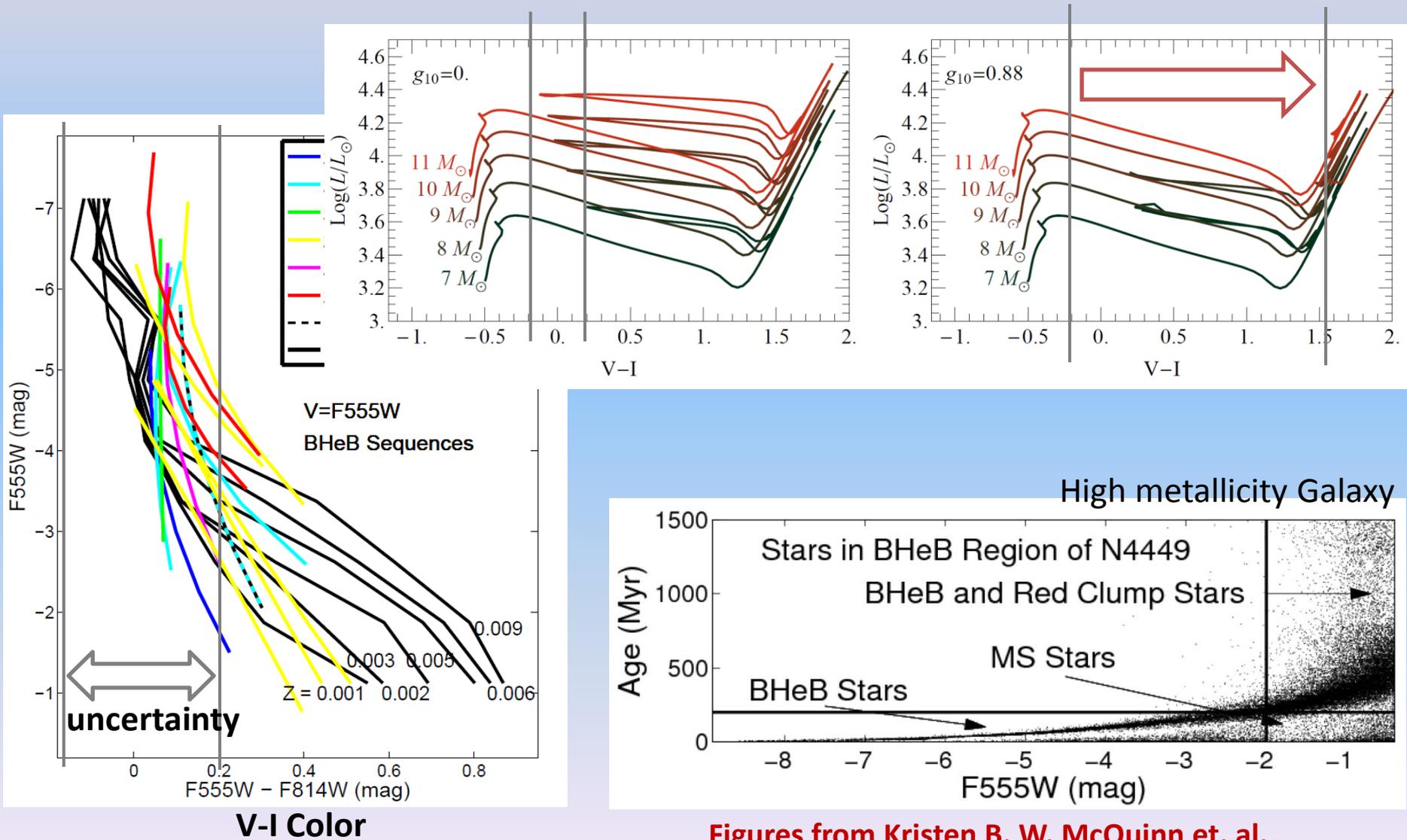


High metallicity Galaxy



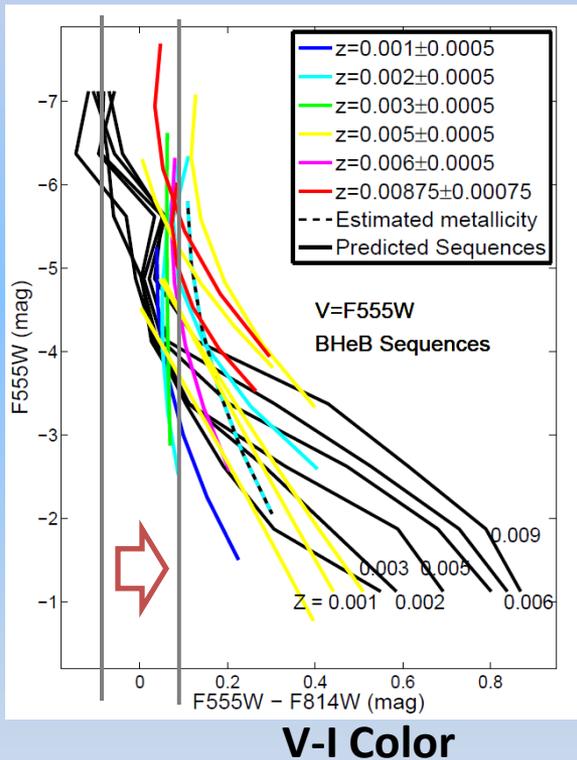
Figures from Kristen B. W. McQuinn et. al.,  
Astrophys.J. 740 (2011)

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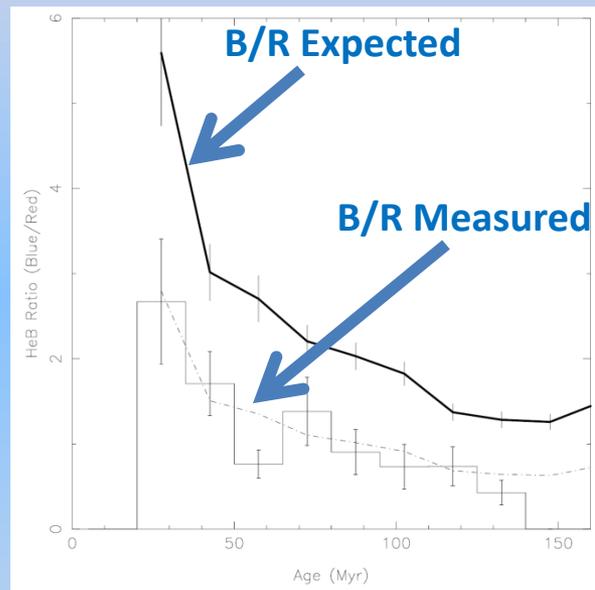


Figures from Kristen B. W. McQuinn et. al.,  
Astrophys.J. 740 (2011)

# Two interesting astrophysical problems: is there a hint to new physics?



Kristen B. W. McQuinn et. al.,  
Astrophys.J. 740 (2011)



R. C. Dohm-Palmer and E. D. Skillman,  
The Astronomical Journal, 123 (2002)

From Palmer, Skillman (2002):  
*“Note how well the functional form of the observations matches that of the model. However, the model values are twice as large as the observations.”*

Current observations show:

- 1) a small red-shift of the bluest point of the blue loop in the high luminosity region of the CMD and
- 2) too many blue stars (B/R problem).

An axion-photon coupling a little below the current bound would reduce (eliminate?) both problems.

# Conclusions and future directions

- Low and intermediate mass stars offer good laboratories to study models of weakly interacting particles.
- Axions have a strong impact on the evolution of low and intermediate mass stars during the core He-burning stage.
- Massive stars offer an efficient novel probe to test new physics. For axions they provide a bound of  $g_{10}=0.8$  (up to  $m_a < 40$  keV or so). A strong bound could possibly be derived from a more accurate analysis of observations.
- Observations of red and blue massive stars show some anomalies and it looks like an axion with  $g_{10}$  between  $\sim 0.1$  and  $\sim 0.8$  would help reducing these anomalies. Though it is still premature to draw precise conclusions, maybe new physics will be necessary to solve these problems.

Thank

You