

# Chiral limit of $N=4$ SYM and ABJM and integrable Feynman graphs

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GATIS Closing Workshop - DESY

based on work to appear soon with O. Gurdogan and V. Kazakov

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All of these preserve **Integrability!**



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**Idea:** play with these new  
parameters, to kill interacting  
terms in the Lagrangian



Simpler field theories,  
with easier perturbation theory.  
Easier to prove integrability!



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$$[A, B] \rightarrow [A, B]_\gamma = e^{i\gamma_k Q_i^A Q_j^B \epsilon^{ijk}} AB - e^{-i\gamma_k Q_i^A Q_j^B \epsilon^{ijk}} BA$$

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Cartan charges of the  $SU(4)$  R-symmetry group

R-symmetry is broken:  $SU(4) \rightarrow U(1)^3$

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Deformation based on conserved charges  $\rightarrow$  all-loop integrability is equally preserved



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Lagrangian is now a function of  $g$  and  $\gamma_k$

$$\begin{aligned} \mathcal{L}_{int} = N_c \text{tr} & \left[ \frac{g}{4} \{ \phi_i^\dagger, \phi^i \} \{ \phi_j^\dagger, \phi^j \} - g e^{-i\epsilon^{ijk} \gamma_k} \phi_i^\dagger \phi_j^\dagger \phi^i \phi^j \right. \\ & - e^{-\frac{i}{2} \gamma_j^-} \bar{\psi}_j \phi^j \bar{\psi}_4 + e^{+\frac{i}{2} \gamma_j^-} \bar{\psi}_4 \phi^j \bar{\psi}_j + i\epsilon_{ijk} e^{\frac{i}{2} \epsilon_{jkm} \gamma_m^+} \psi^k \phi^i \psi^j \\ & \left. - e^{+\frac{i}{2} \gamma_j^-} \psi_4 \phi_j^\dagger \psi_j + e^{-\frac{i}{2} \gamma_j^-} \psi_j \phi_j^\dagger \psi_4 + i\epsilon^{ijk} e^{\frac{i}{2} \epsilon_{jkm} \gamma_m^+} \bar{\psi}_k \phi_i^\dagger \bar{\psi} \right] \end{aligned}$$

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Double scaling limit [\[Gurdogan, Kazakov'15\]](#)

$$q_i \equiv e^{-i \frac{\gamma_i}{2}} \rightarrow \infty$$

$$g \rightarrow 0$$

with  $\xi_i \equiv g q_i$  finite

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Chiral limit: parameters are complex hence resulting theories will be non-unitary

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$$\begin{aligned}\mathcal{L}_{\text{int}} = & N_c \text{tr} [\xi_1^2 \phi_2^\dagger \phi_3^\dagger \phi_2 \phi_3 + \xi_2^2 \phi_3^\dagger \phi_1^\dagger \phi_3 \phi_1 + \xi_3^2 \phi_1^\dagger \phi_2^\dagger \phi_1 \phi_2 \\ & + i\sqrt{\xi_2 \xi_3} (\psi^3 \phi^1 \psi^2 + \bar{\psi}_3 \phi_1^\dagger \bar{\psi}_2) \\ & + i\sqrt{\xi_1 \xi_3} (\psi^1 \phi^2 \psi^3 + \bar{\psi}_1 \phi_2^\dagger \bar{\psi}_3) \\ & + i\sqrt{\xi_1 \xi_2} (\psi^2 \phi^3 \psi^1 + \bar{\psi}_2 \phi_3^\dagger \bar{\psi}_1) ]\end{aligned}$$

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$\xi_i = \xi$  strong  $\beta$ -deformation with  $\mathcal{N} = 1$  SUSY

$$\xi_1 = \xi_2 = 0, \quad \xi_3 = \xi$$

$$\mathcal{L} = \frac{N_c}{2} \text{tr} (\partial^\mu \phi_1^\dagger \partial_\mu \phi_1 + \partial^\mu \phi_2^\dagger \partial_\mu \phi_2 + 2\xi^2 \phi_1^\dagger \phi_2^\dagger \phi_1 \phi_2)$$

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In ABJM, the twisted theory depend on the coupling  $\lambda = \frac{N}{k}$  and 3 twist parameters.

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One gets a  $\phi^6$  type scalar Lagrangian

$$\mathcal{L} = N_c \text{Tr} \left[ -\partial_\mu Y_1^\dagger \partial^\mu Y^1 - \partial_\mu Y_2^\dagger \partial^\mu Y^2 - \partial_\mu Y_4^\dagger \partial^\mu Y^4 \right. \\ \left. + (4\pi)^2 \frac{\xi_1 \xi_2}{\xi_3} Y^1 Y_4^\dagger Y^2 Y_1^\dagger Y^4 Y_2^\dagger \right]$$



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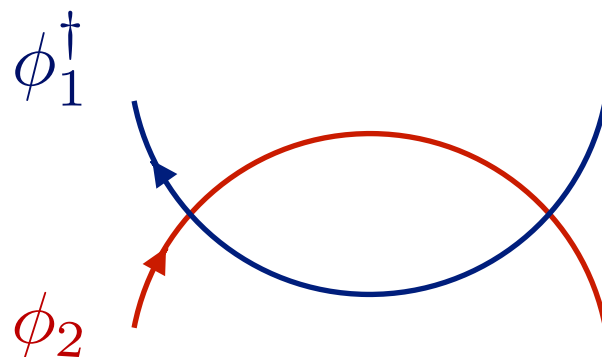
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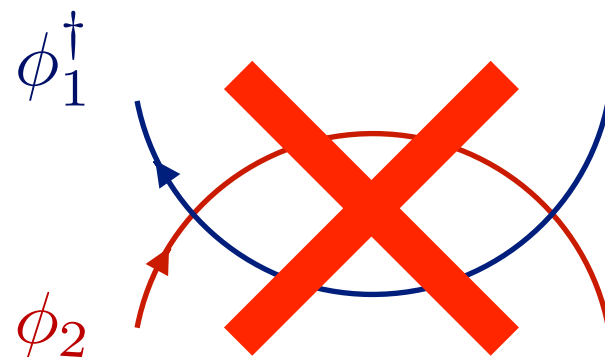




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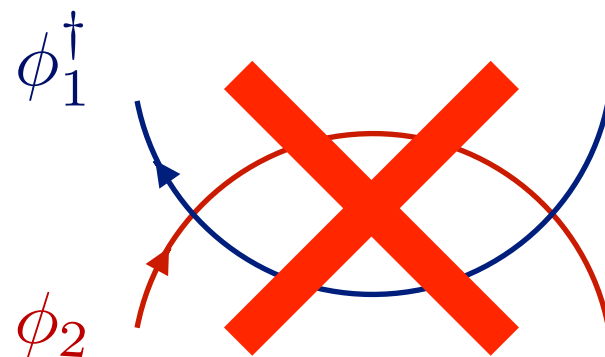


at large  $N$ ,  
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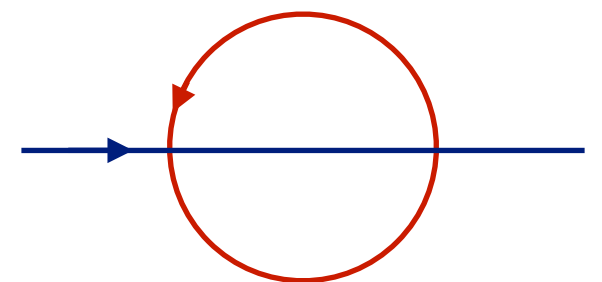
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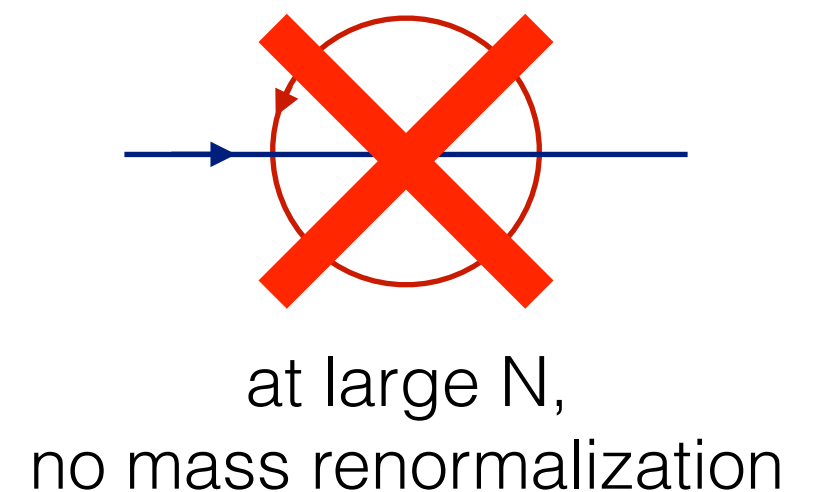
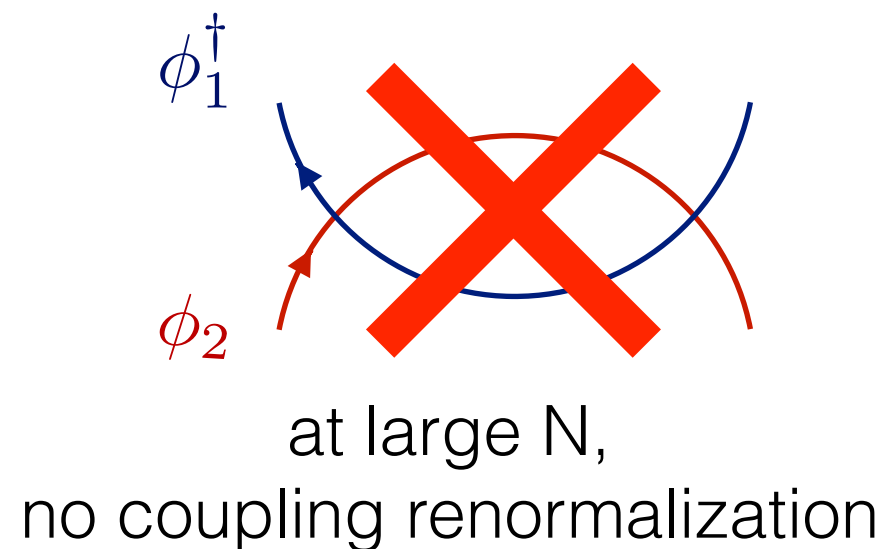
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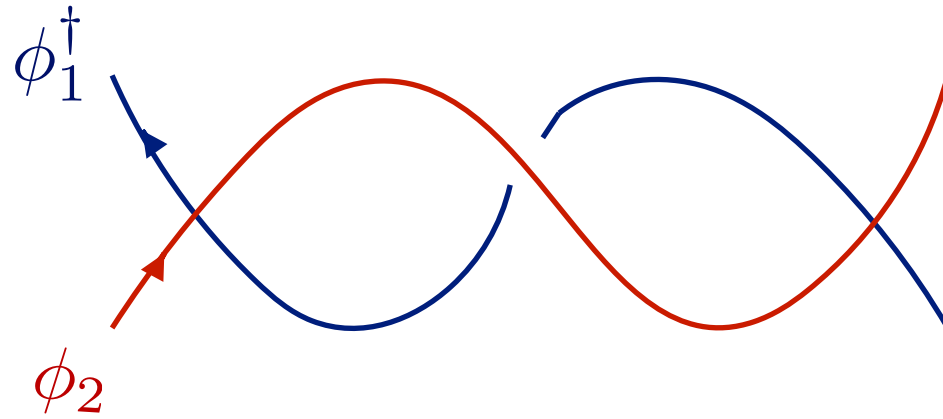
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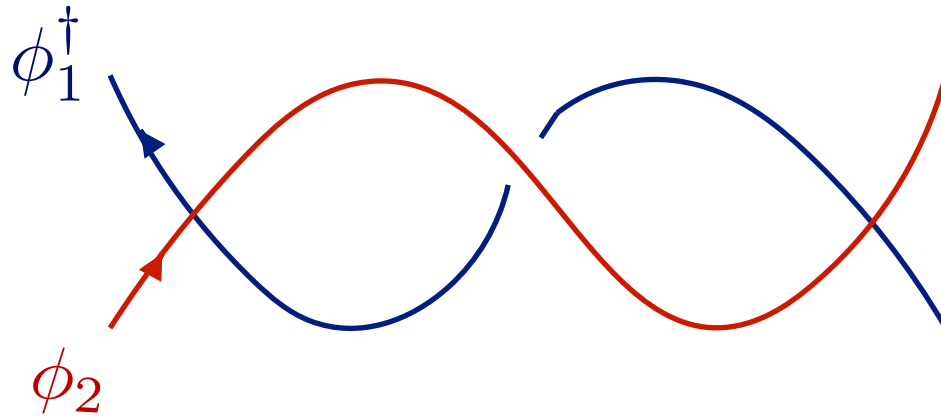
**$A_t \mid N \dots$**

# $A \propto 1/N \dots$



allowed, of order  $1/N$

# At $1/N \dots$

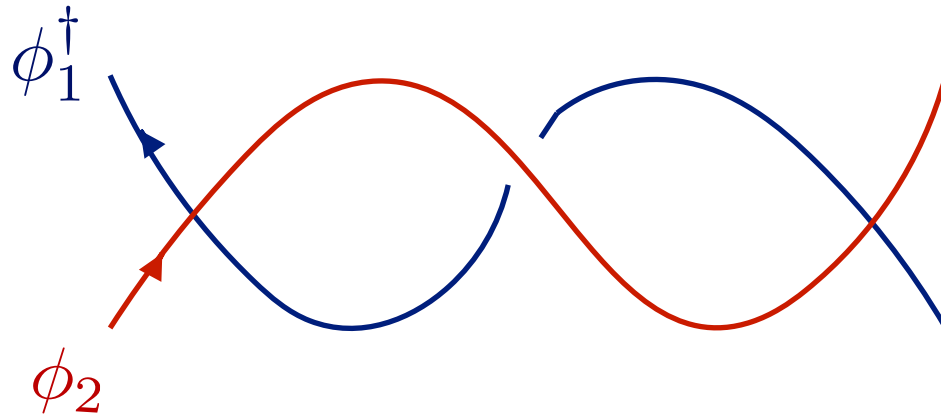


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—> New divergences that cannot be absorbed in the renormalization of operators. **Need counterterms of the form trace-trace:** [Fokken, Sieg, Wilhelm'13,16]

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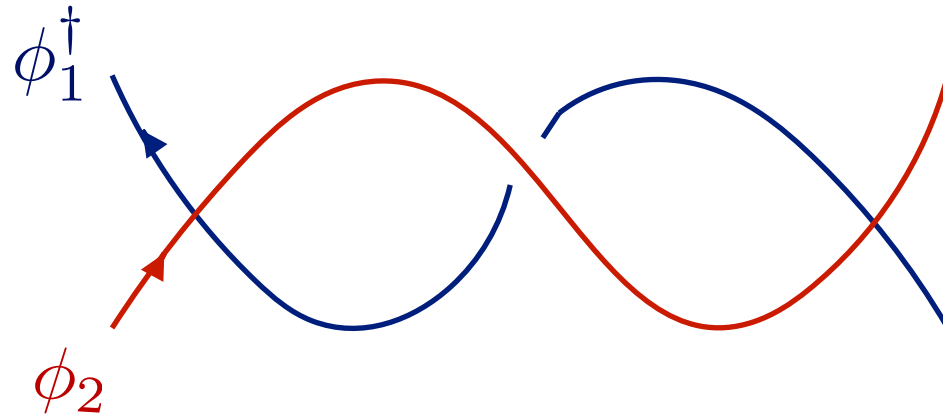
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Conformal symmetry broken even in the planar limit!



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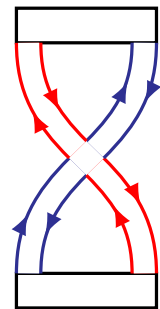
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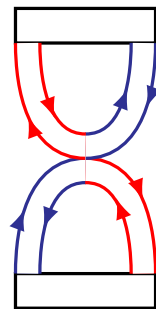
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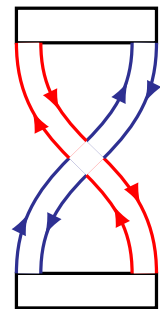
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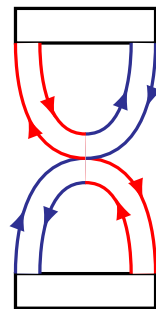
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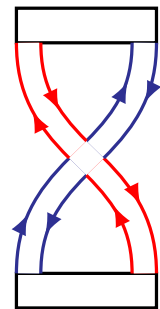
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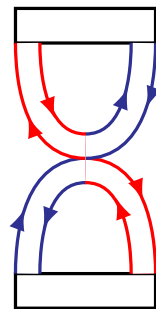
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These  $L=2$  states play the role of **tachyons** in  $\gamma$ -twisted dual string theory. At strong coupling, these tachyons have been identified. [\[Rastelli, Pomoni'08\]](#)

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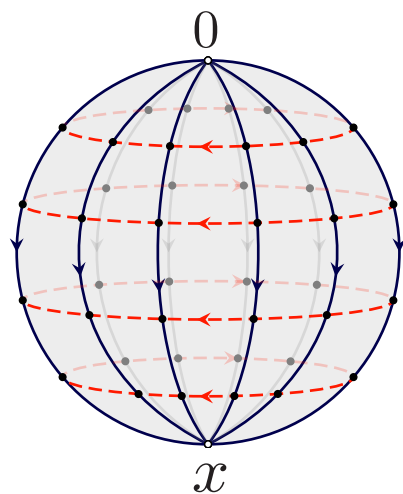
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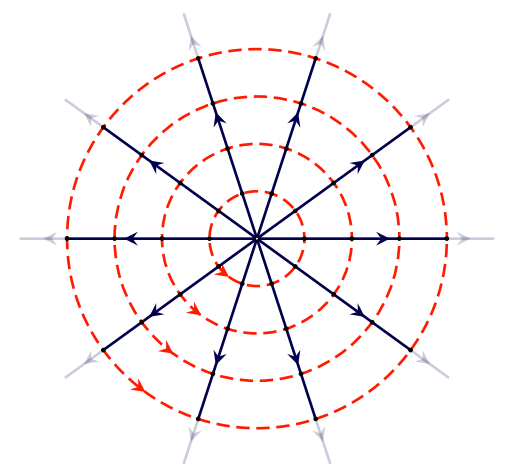
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amputation of one external operator



wheel graphs



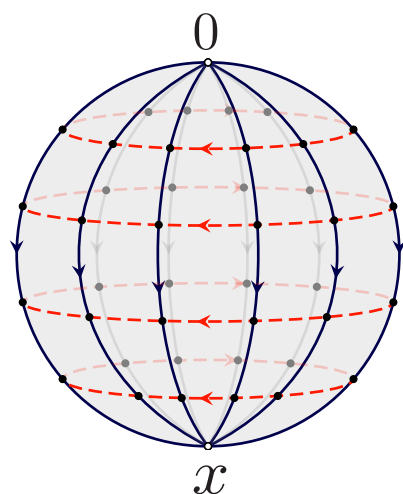
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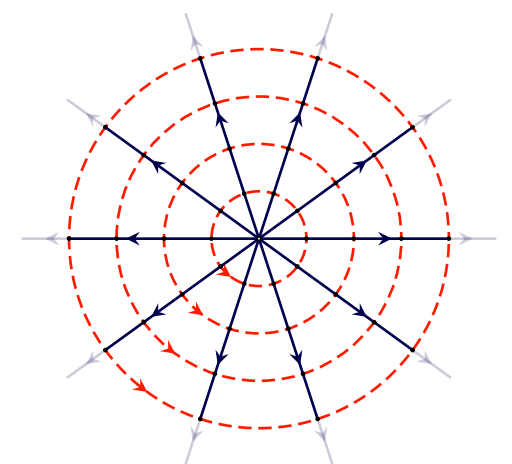
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- 1 wheel [Broadhurst, '80]

- 2 wheels from the integrability based solution of  $\gamma$  twisted

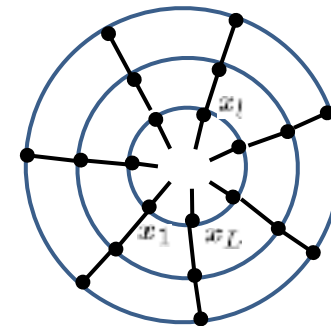
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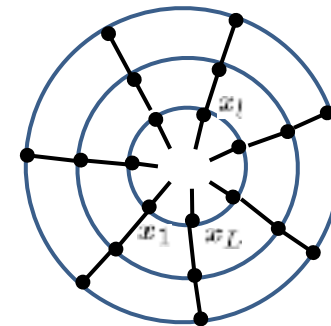
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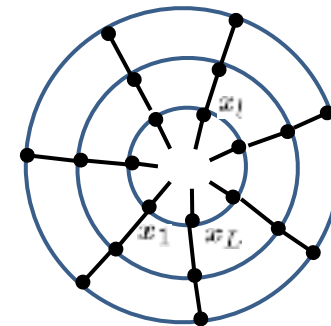


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# A glimpse on the origin of Integrability

- ‘Hamiltonian’ generating the wheel graphs (fishnet lattice)

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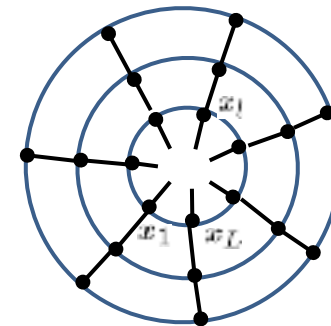
physical space in principal series irreps of  $\text{SL}(4)$

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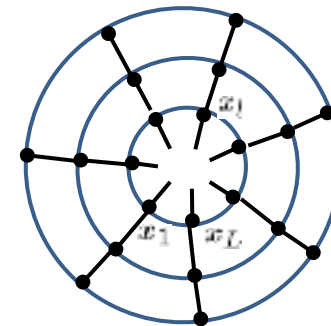
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- Other operators: modification of boundary conditions.
- Still fishnet in the bulk!



# Excited states

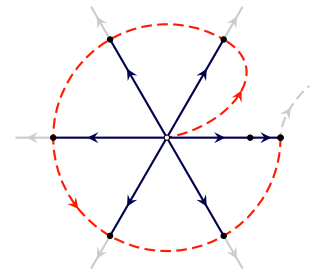
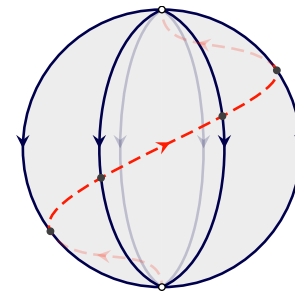
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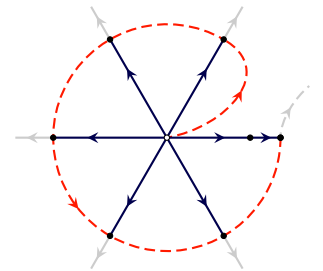
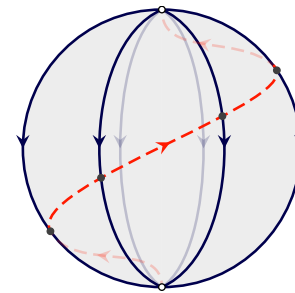


ladder diagrams

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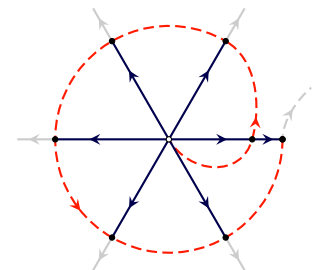
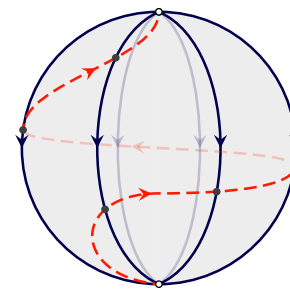
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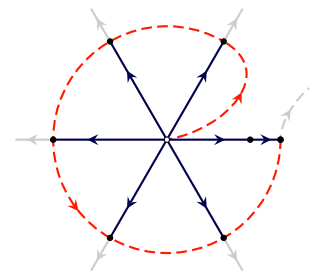
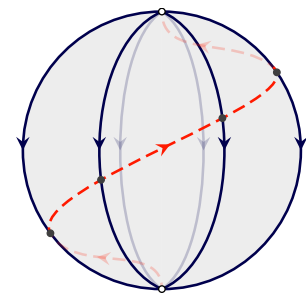


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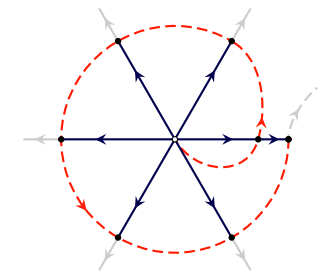
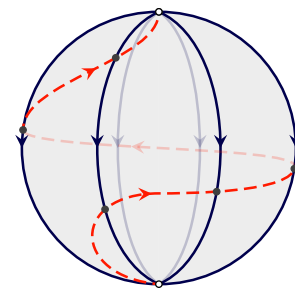
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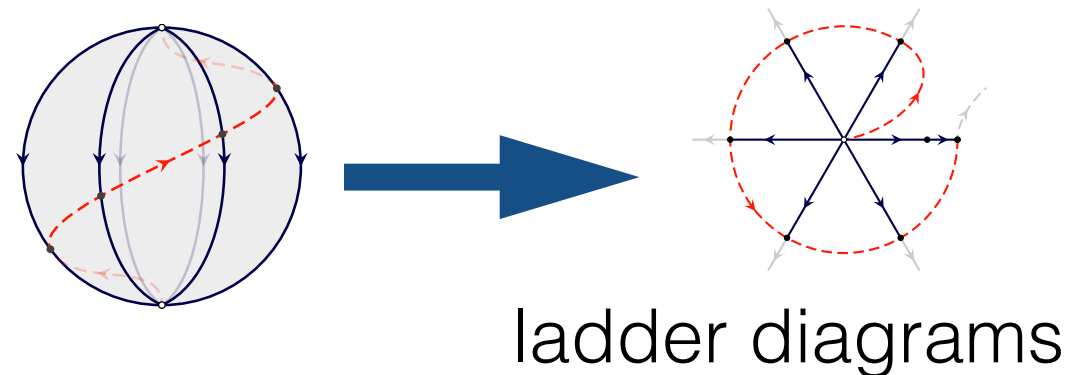
Asymptotic Bethe Ansatz:  $e^{ipL} = q_3^{-2L}$

$$\gamma = -\frac{1}{2} + \frac{1}{2} \sqrt{1 + 4g^2 \sin(p/2)^2} \rightarrow -\frac{1}{2} + \frac{1}{2} \sqrt{1 - 4\xi_3^2}$$

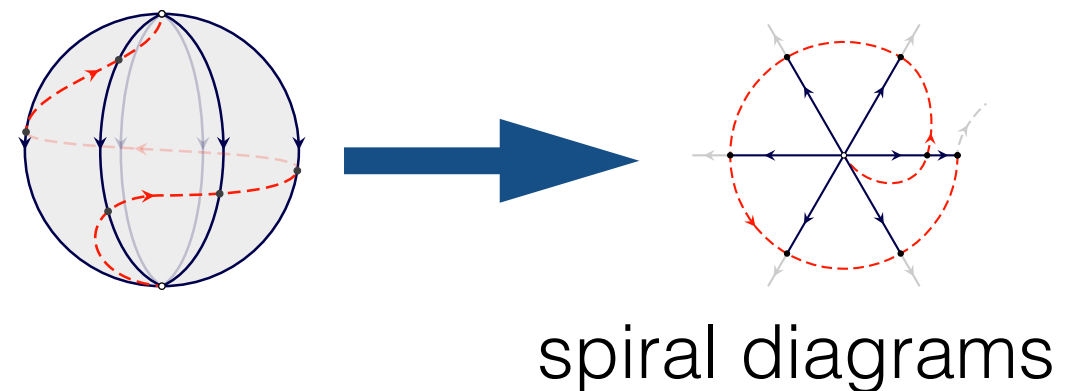
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Resummation of ladder diagrams [Broadhurst, '93, Gross, Mikhailov, Roiban' 02]

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- The anomalous dimension is:  $\gamma = \sum_{k=1}^N i \left(u_k + \frac{i}{2}\right)$

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Say at four loops for L=5:

$$\mathcal{G}_{\alpha\beta} \big|_{\xi^8} = \left[ \begin{array}{c|c} \text{diagram 1} & \text{diagram 2} \\ \hline \text{diagram 3} & \text{diagram 4} \end{array} \right]$$

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Spectrum + input of lower loop integrals + 

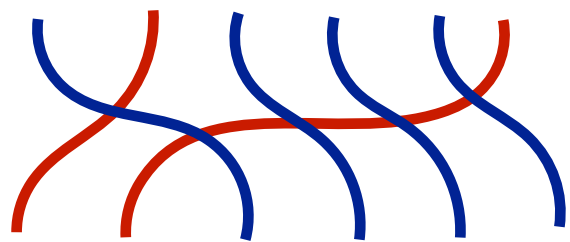


fixes completely  and  up to  $1/\epsilon$  term

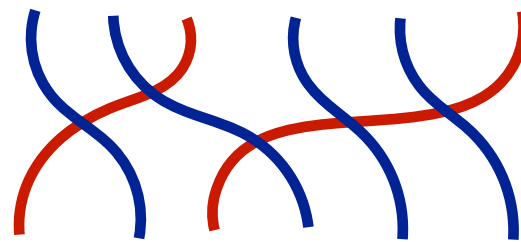
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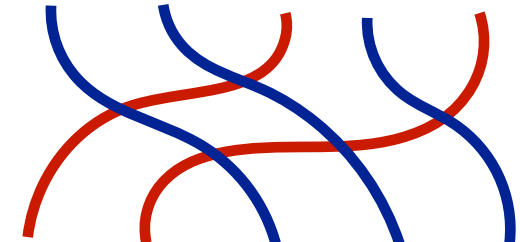
Spectrum provides relations between these 5-loop integrals  
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$I_a$



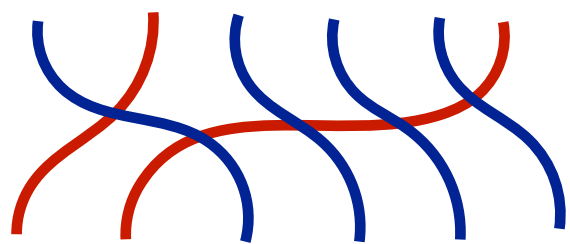
$I_b$



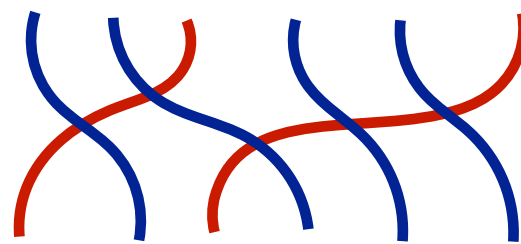
$I_c$

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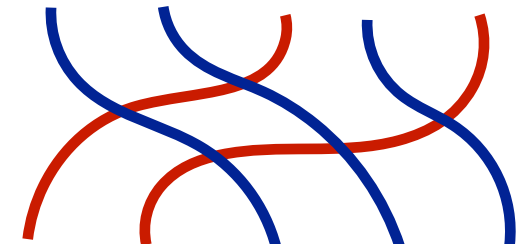
Spectrum provides relations between these 5-loop integrals and lower loop ones



$I_a$



$I_b$



$I_c$

Spectrum + input lower loop integrals fixes everything up to a constant. The nontrivial part are  $1/\epsilon$  terms:

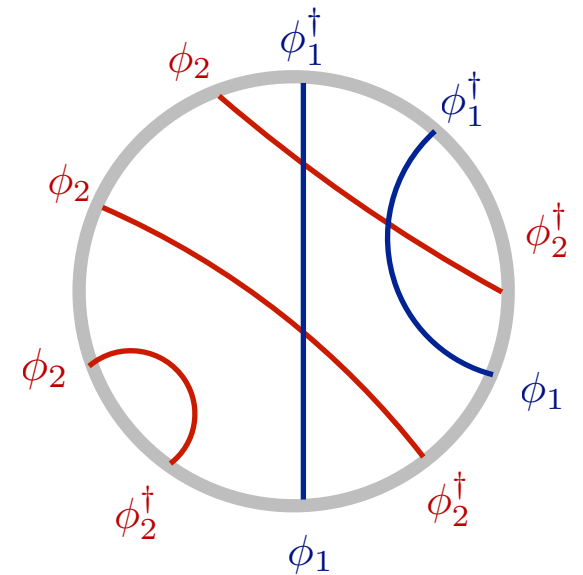
$$I_b|_{1/\epsilon} = -I_a|_{1/\epsilon} - \frac{160\zeta(3)}{9} + \frac{53\pi^4}{72} - \frac{187}{5} - \frac{25\pi^2}{12}$$

$$I_c|_{1/\epsilon} = I_a|_{1/\epsilon} + \frac{418\zeta(3)}{45} + \frac{121\pi^4}{360} + \frac{2\pi^2}{9} - \frac{112}{5}$$

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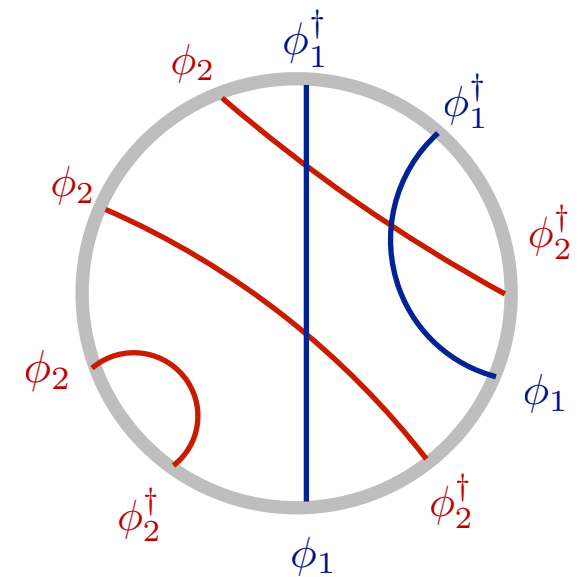
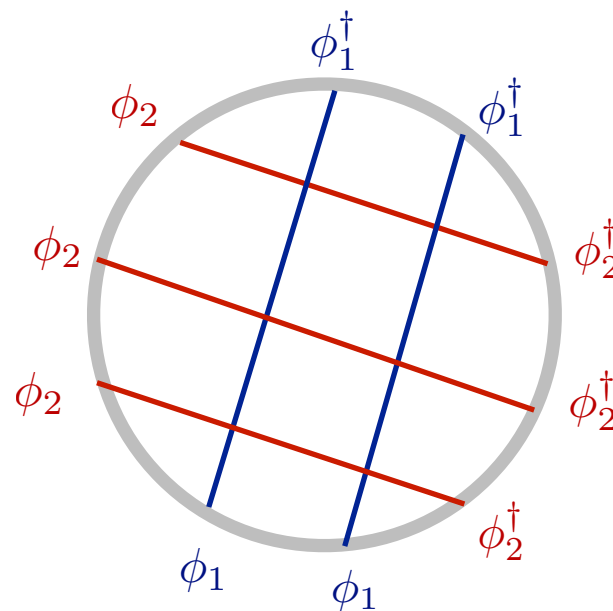
**Scattering amplitudes:** for a given ordering of external fields in the disk  
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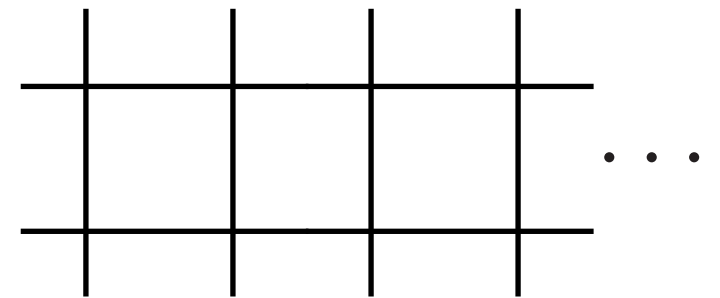
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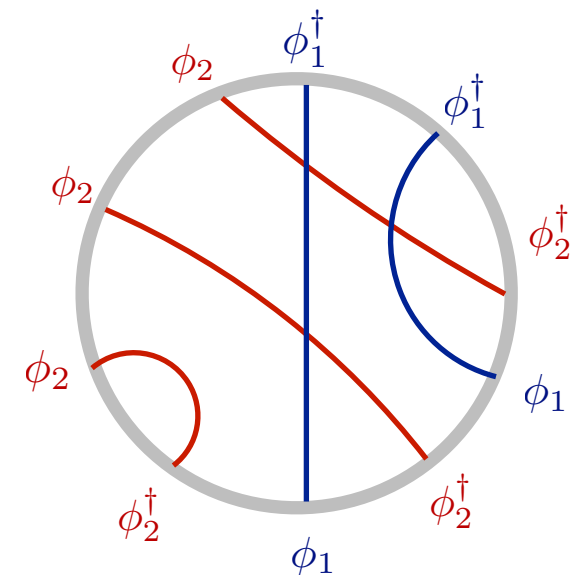
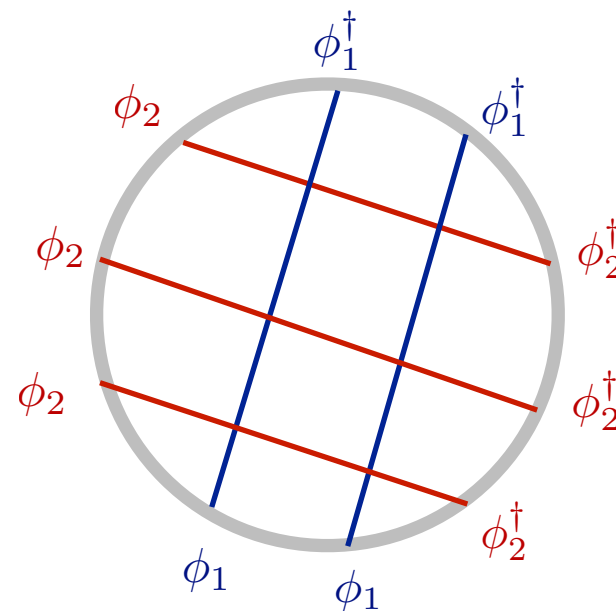
Only boxes are allowed



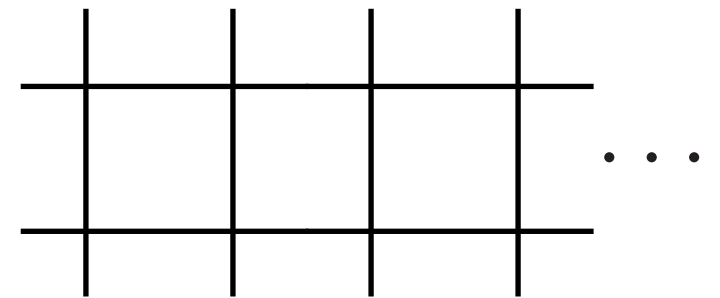
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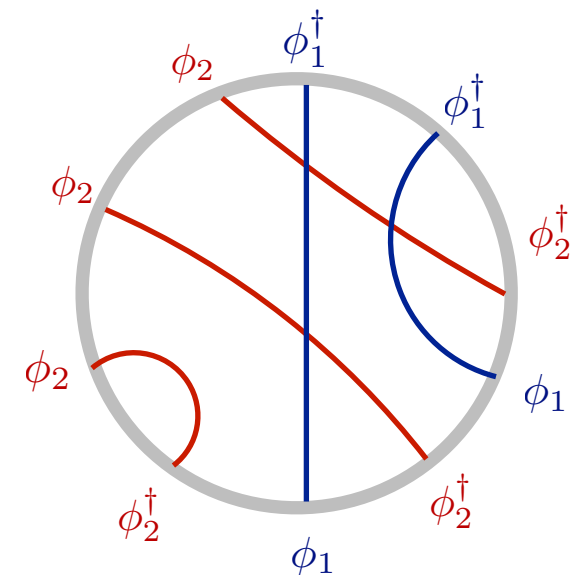
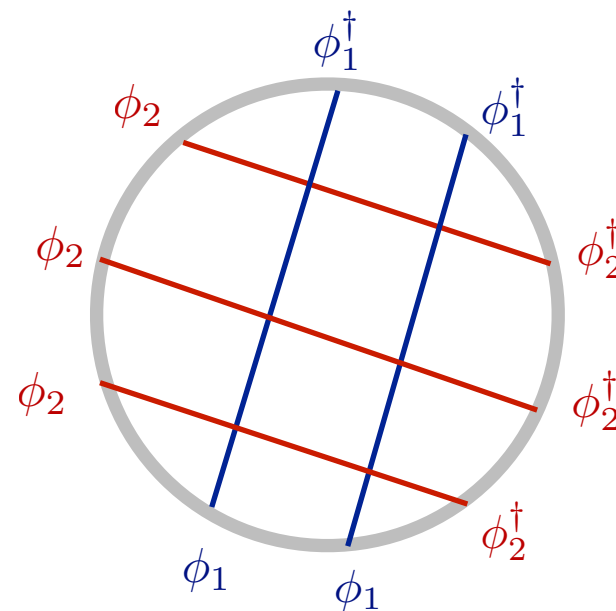
Amplitudes are **finite** and manifestly **dual conformal invariant!**



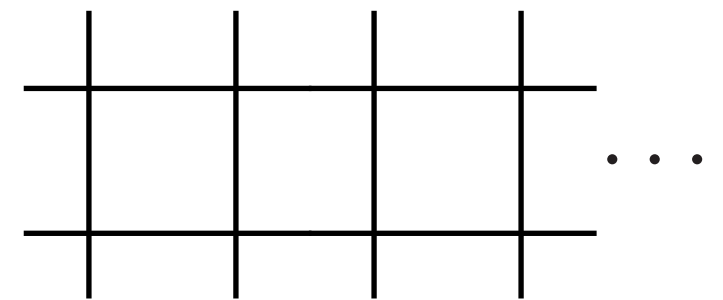
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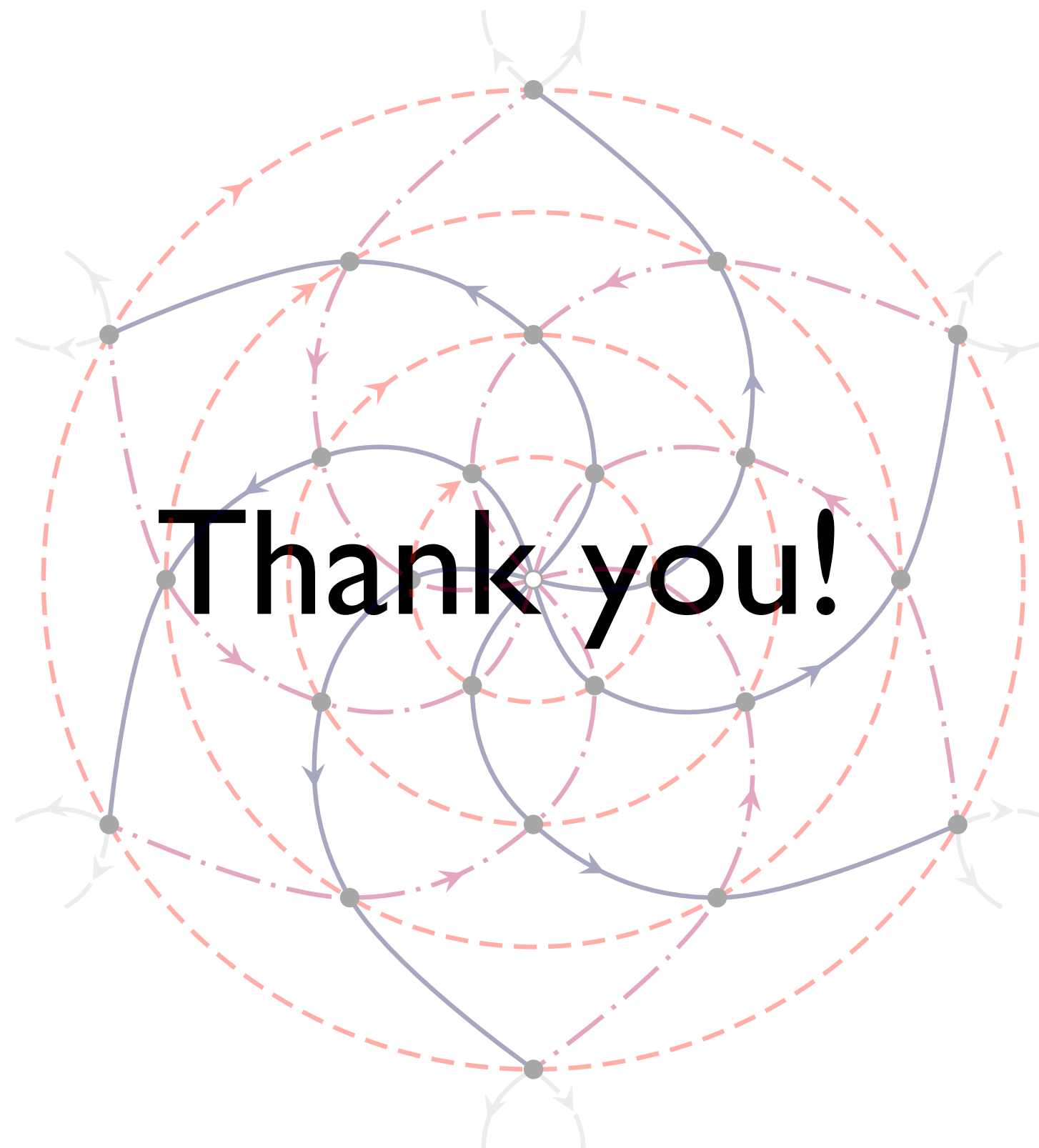


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**Three point functions:** simple perturbation theory. Hexagon program for bi-scalar model? [Basso, Komatsu, Vieira, '15]

# Conclusions

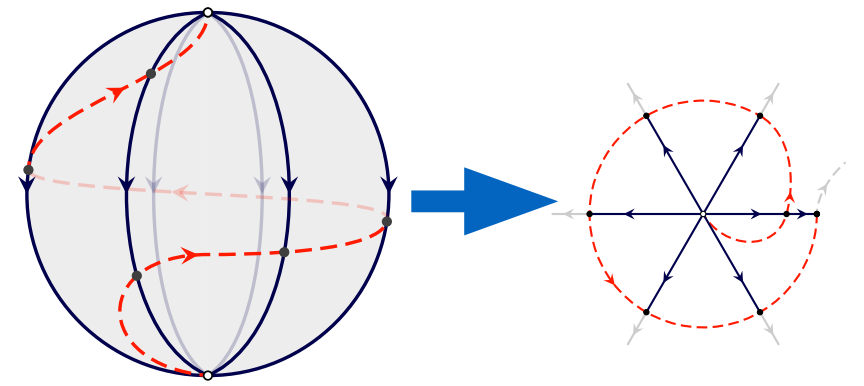
- New 4D and 3D integrable field theories.
- Shed light on the origins of Integrability. First explicit demonstration of all-loop integrability in a corner of  $N=4$  SYM - related to fishnet graphs.
- How to take the DS limit at the level of the Quantum Spectral Curve?
- Can we derive from first principles Thermodynamic Bethe Ansatz and/or Quantum Spectral Curve?
- Structure constants, scattering amplitudes, four-point functions?
- Can integrability further help in computing more integrals?
- Integrability to compute beta function?
- Dual string theory?



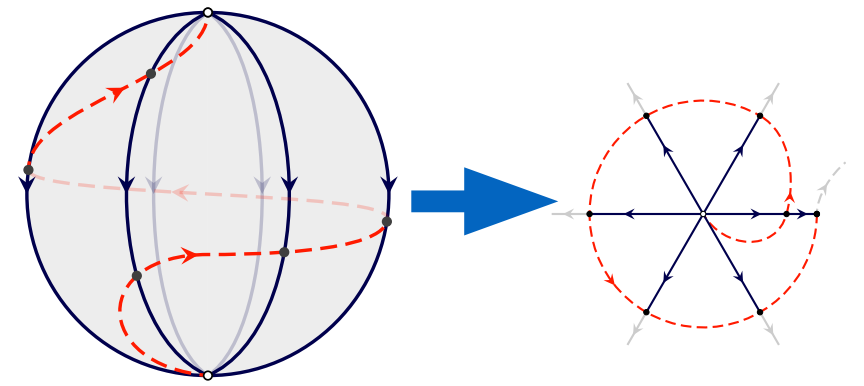
**Thank you!**

# Extra Slides

# Spiral graphs



# Spiral graphs

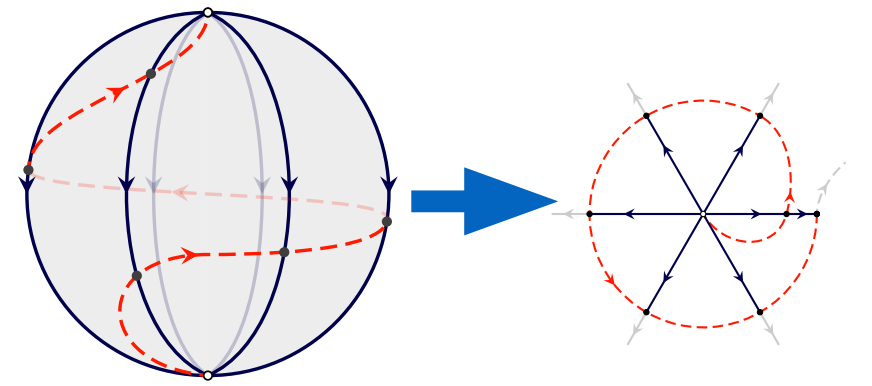


For these operators,  $\text{Tr}[\phi_1 \phi_1 \dots \phi_1 \phi_2 \phi_1 \dots \phi_1]$

strongly  $\beta$ -deformed = strongly  $\gamma$ -deformed

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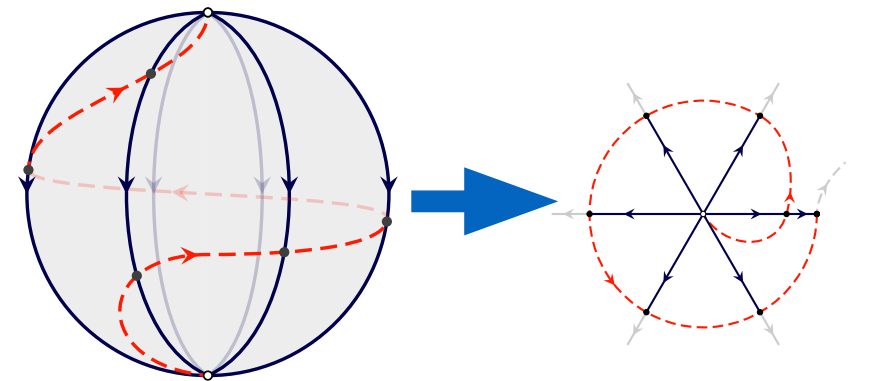
Anomalous dimension is exactly known

(including wrapping) in  $\beta$ -deformed theory by TBA:

[Gromov, Levkovich-Maslyuk, '10]

$$(\delta\gamma)_{\text{wrap}} = 4\xi^{2L} \sum_{k=3}^{L-1} \binom{2(-k + L[k/2] - 1)}{L - k} \zeta(2(-k + L + [k/2] - 1) + 1)$$

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We can easily predict this integral in dim-reg

wrapping integral = ladder integral up to  $1/\epsilon$  (by exponentiation)

+  $1/\epsilon$  term fixed by anomalous dimension from TBA



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- The anomalous dimension is then given by the sum of  $N$  dispersion relations evaluated on these solutions.

$$\gamma = ig \sum_j \left( \frac{1}{x^+(u_j)} + \frac{1}{x^-(u_j)} \right)$$

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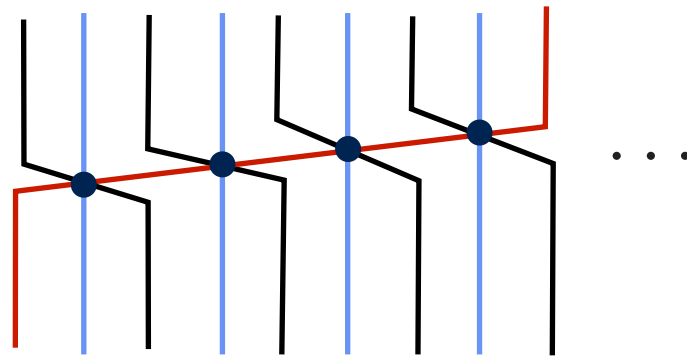
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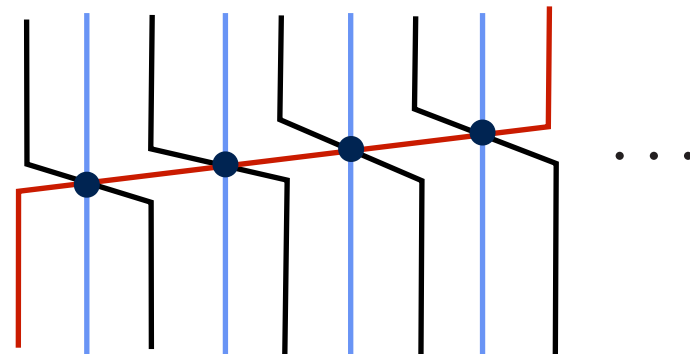
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Resums 3D ladders



Also Bethe Ansatz for the states  $\text{Tr}[(Y^1 Y_4^\dagger)^{L-N} (\textcolor{red}{Y}^2 Y_4^\dagger)^N]$

$$\xi^{-L} (u_j^2 + 1/4)^L = \prod_{j \neq k}^N \left[ \frac{u_k - u_j + i}{u_k - u_j - i} \sigma_0^{m,m}(u_k, u_j) \right]$$

Can be used for 3D Feynman integrals