

Studying the QGP in ultrarelativistic heavy-ion collisions

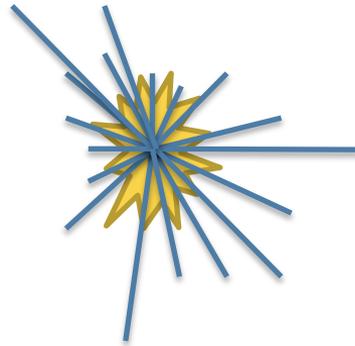
Carsten Greiner

PANIC, Hamburg, August 25th, 2014

- **Observables** at RHIC and **LHC** (elliptic flow, jets, heavy quarks)
- Dynamical description: “macro” or **fluid dynamics**
- Dynamical description: “micro” or **transport dynamics**
- **Transport properties** of the QGP

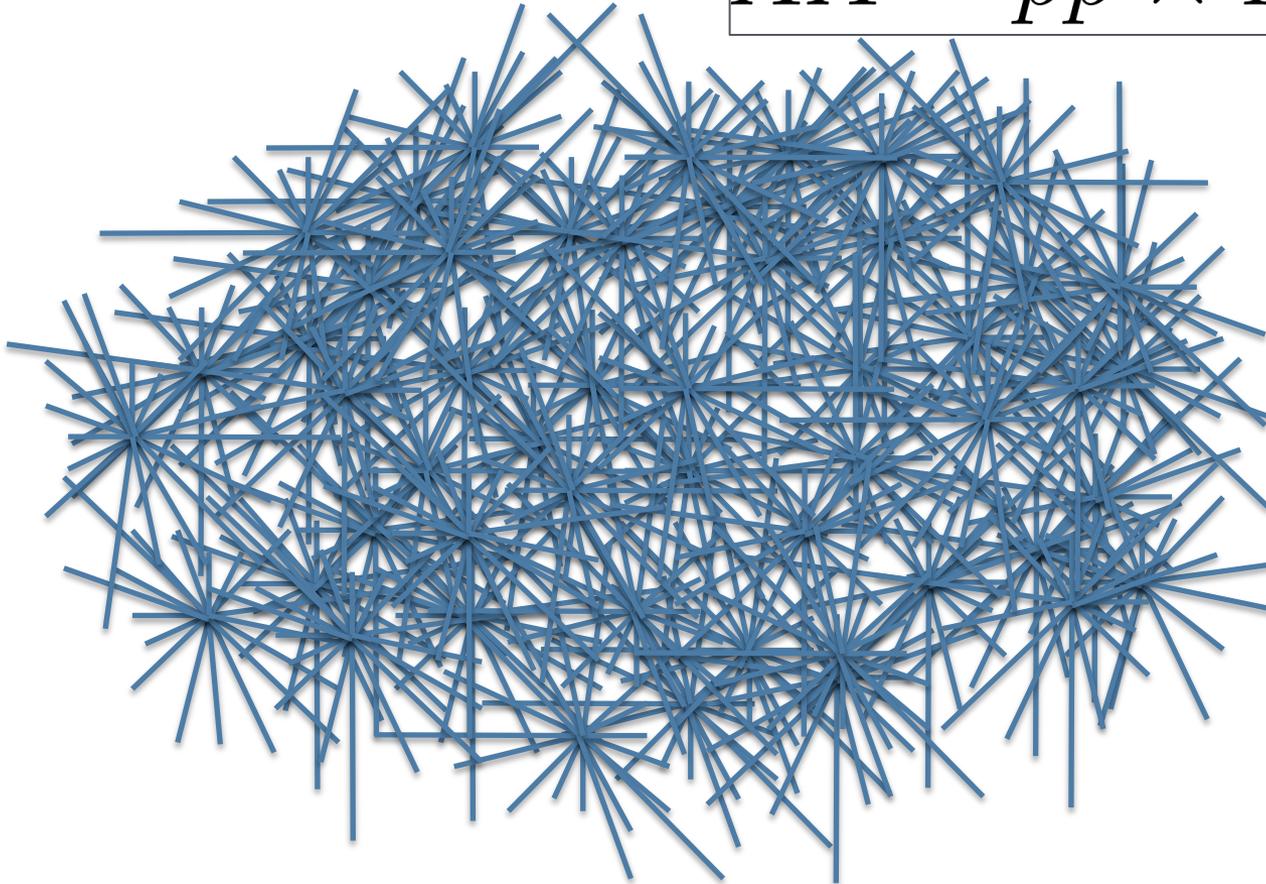
Sketch of heavy ion collision

Initial pp scattering

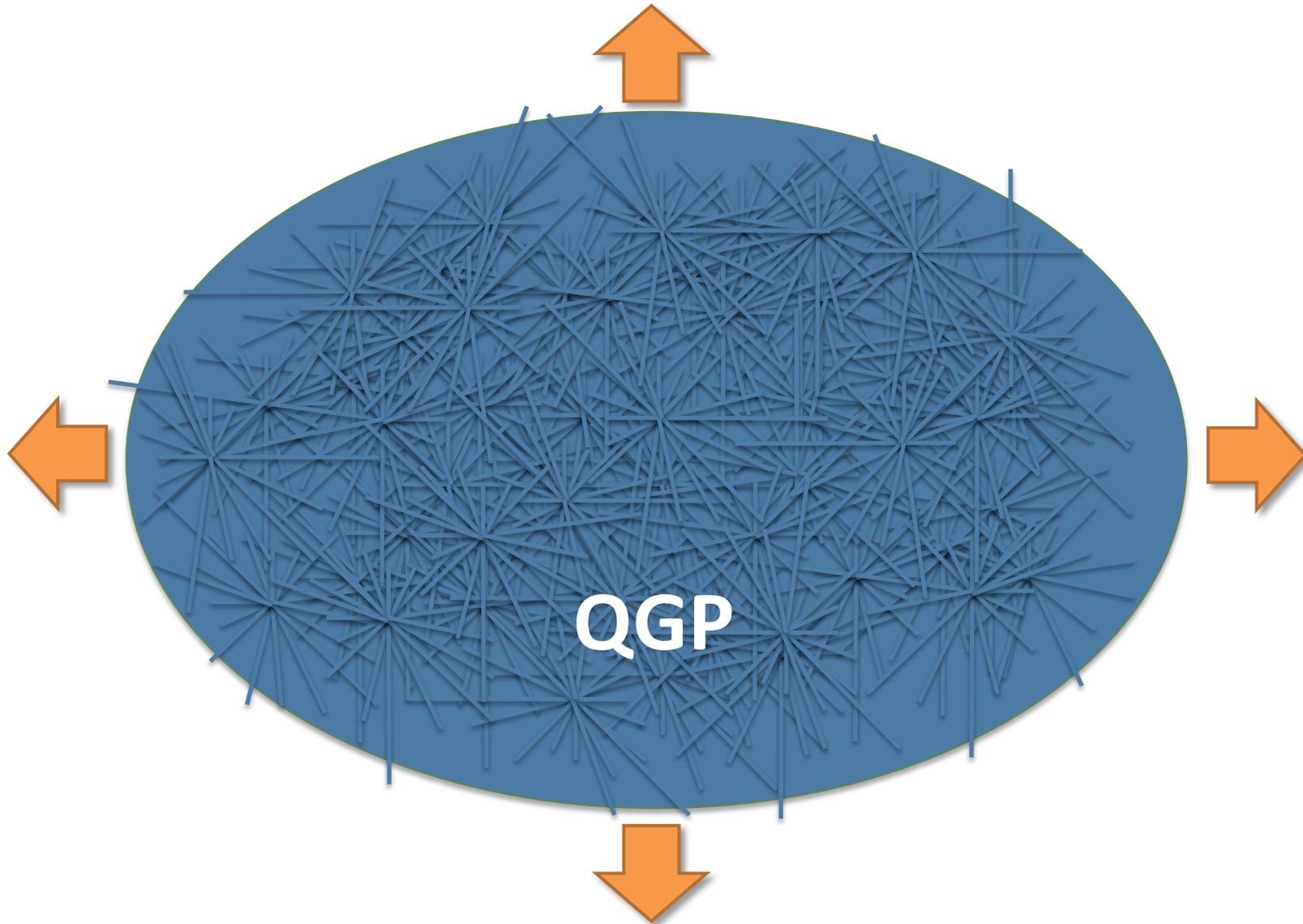


Sketch of heavy ion collision

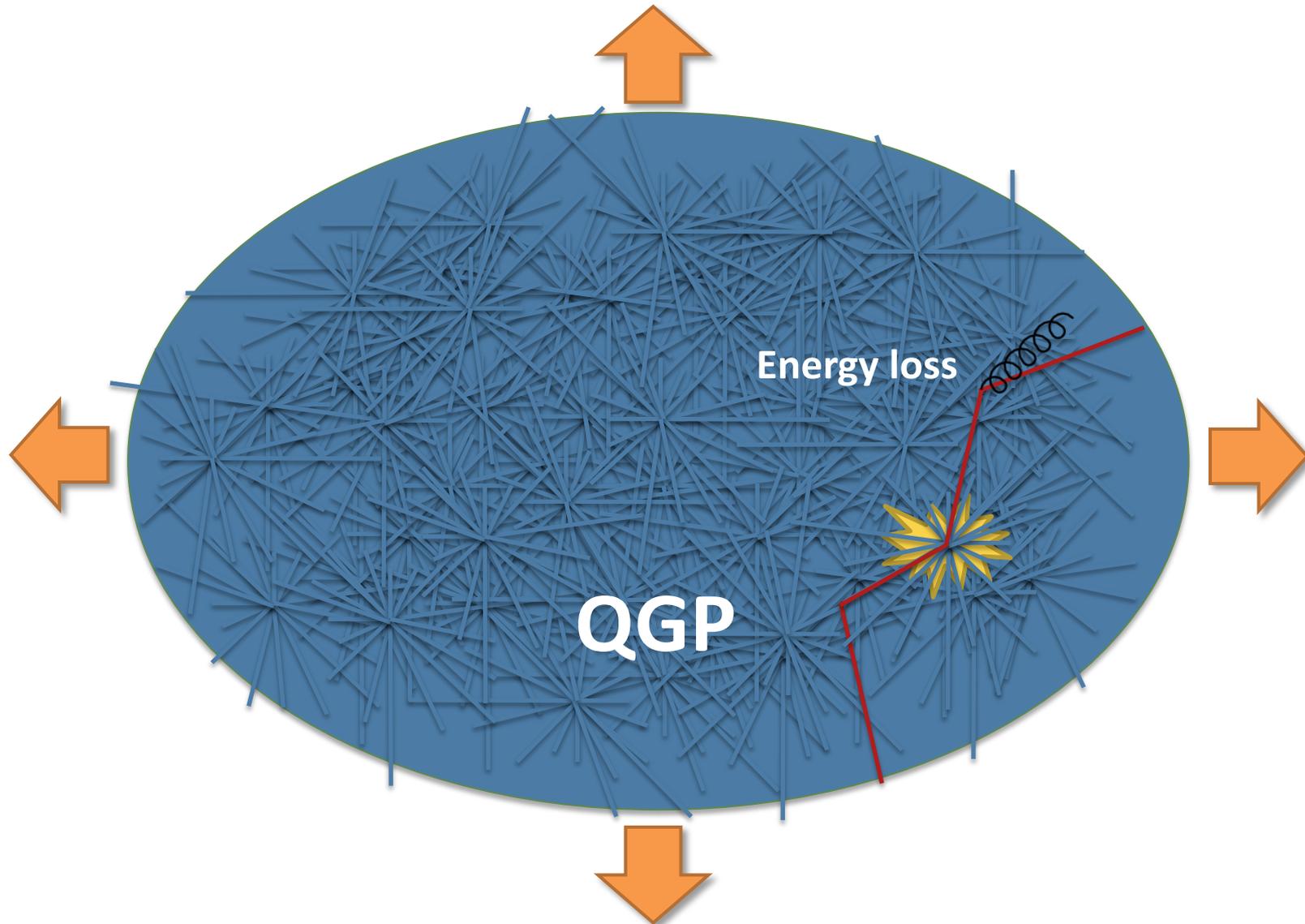
$$AA = pp \times N_{\text{binary}}$$



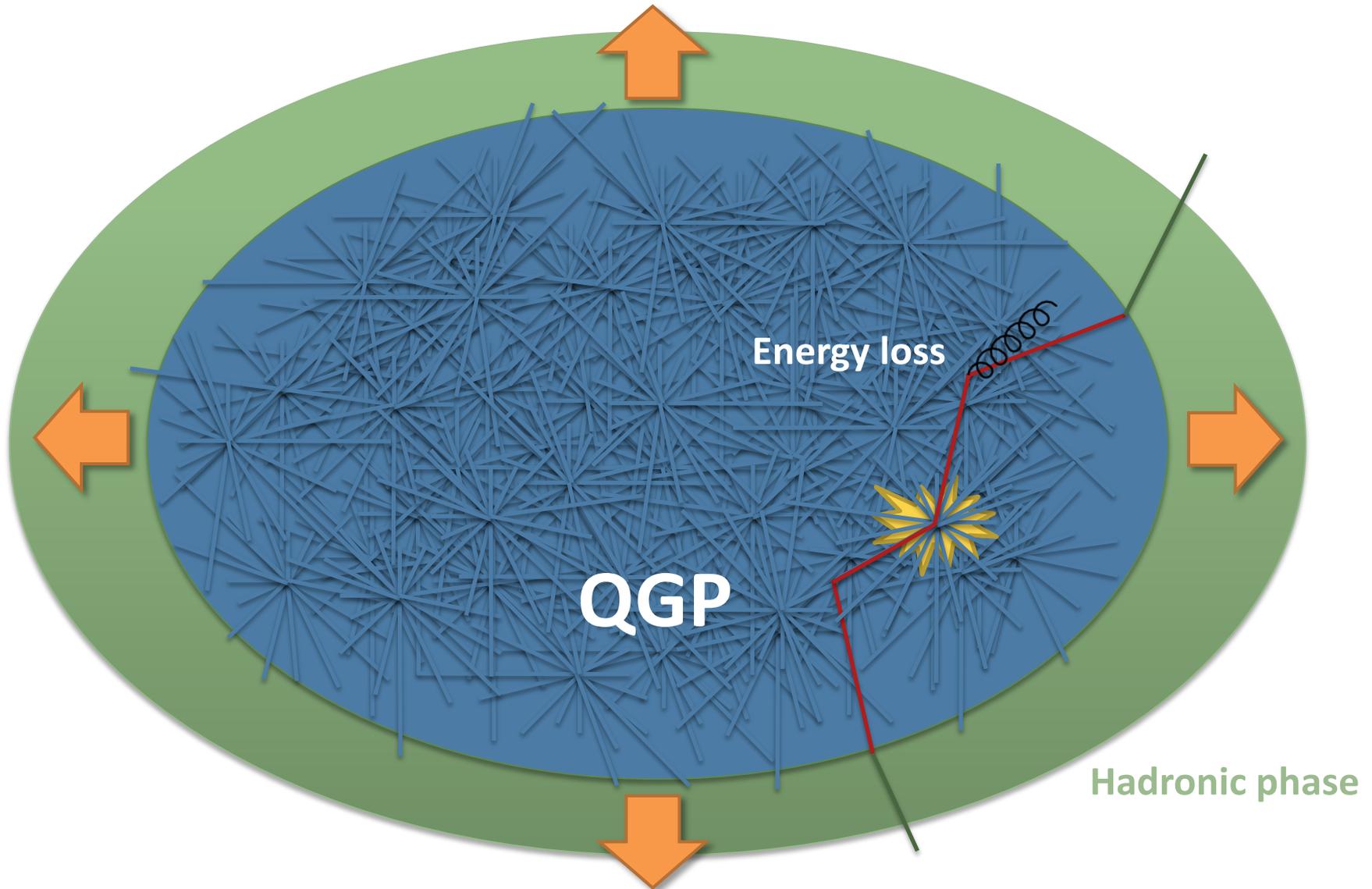
Sketch of heavy ion collision



Sketch of heavy ion collision

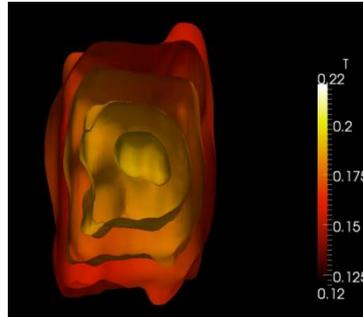
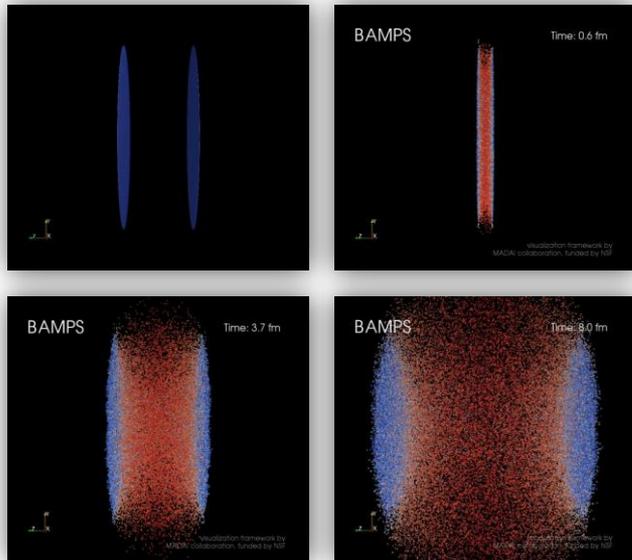


Sketch of heavy ion collision

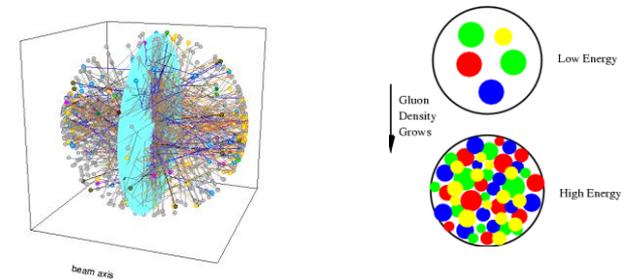


Heavy-ion collisions are complex !

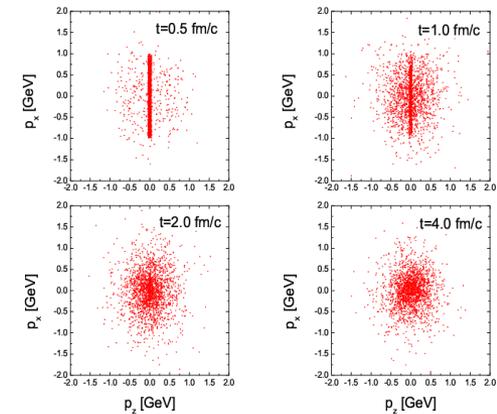
Dynamical bulk description



Glauber Gluon saturation



Early thermalisation

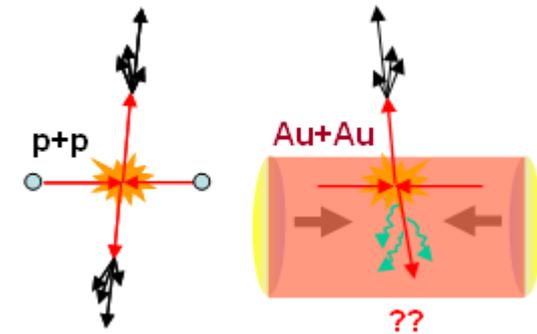
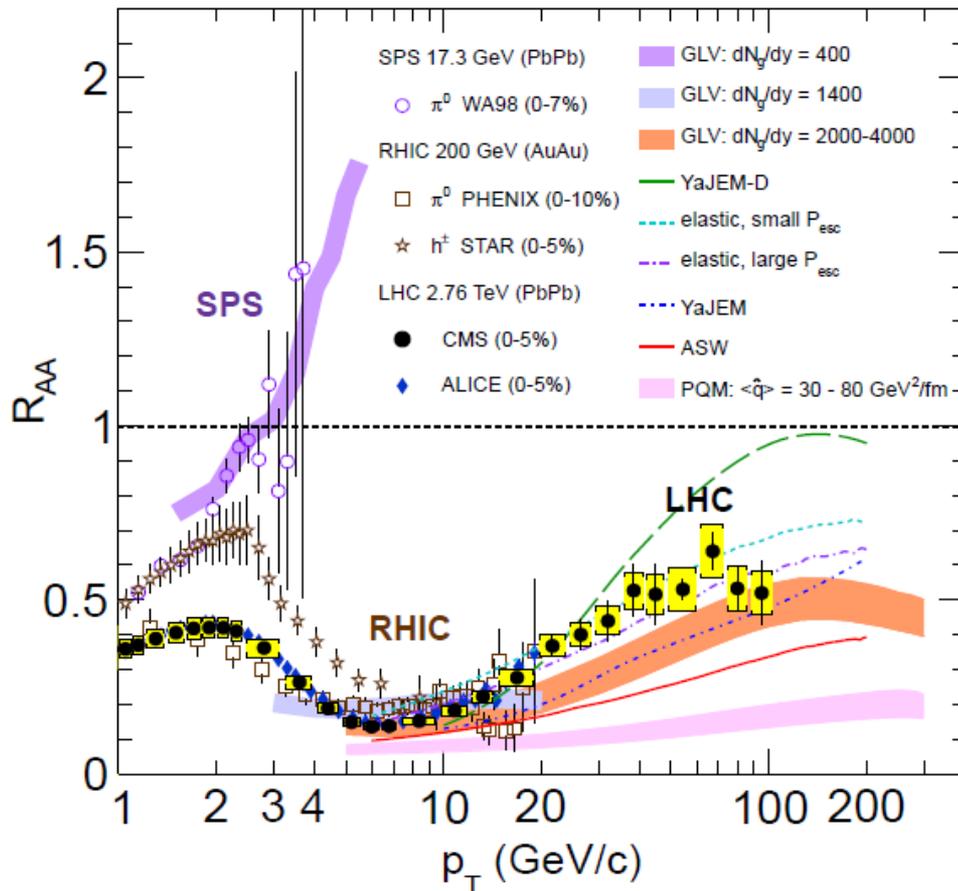


No model can describe all aspects of the QGP evolution

Nuclear modification factor R_{AA}

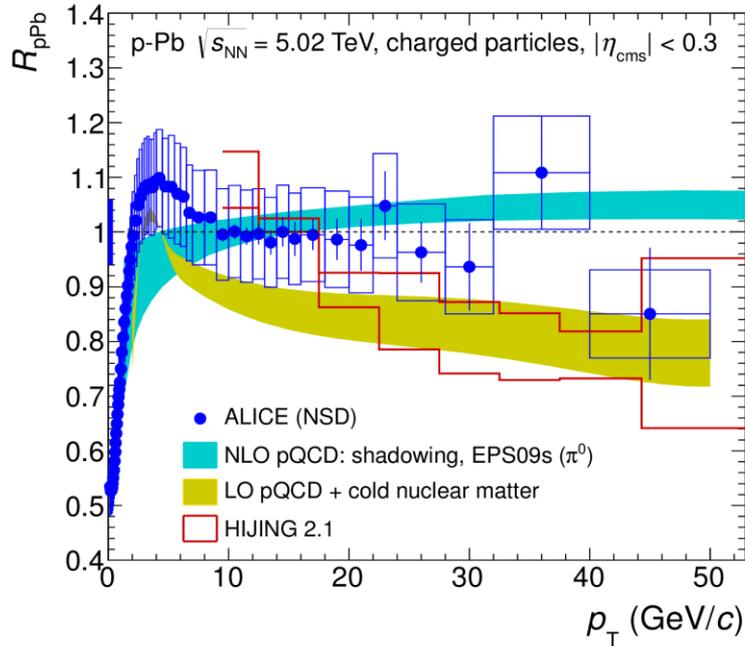
Nuclear modification factor

$$R_{AA} = \frac{dN/dp_T dy|_{A+A}}{N_{\text{bin}} dN/dp_T dy|_{p+p}}$$

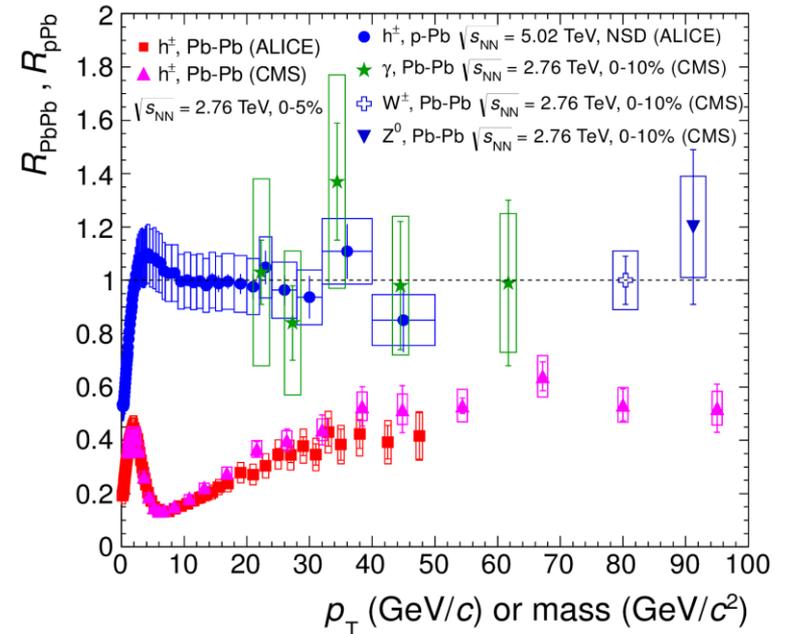


Successfully described
by perturbative QCD

CMS, Eur. Phys. J. C72 (2012)



ALICE, PRL 110 (2013) 082302

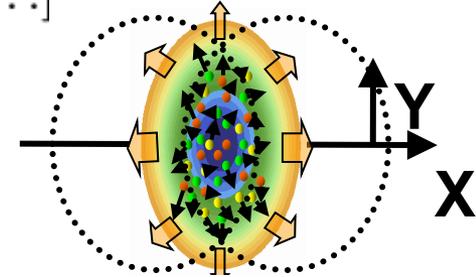
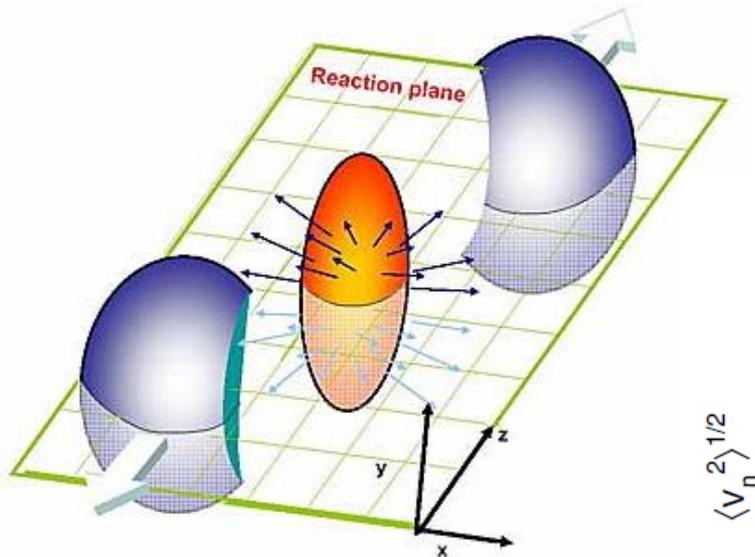


No indication of any departure from unity in p+Pb

... provide experimental demonstration that suppression =
parton energy loss

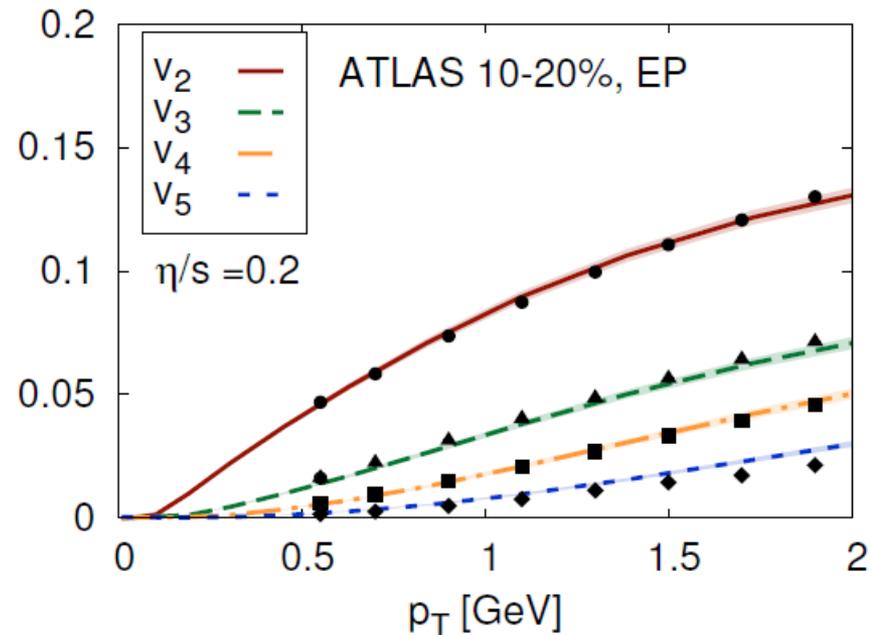
Elliptic flow v_2

$$\frac{d^3 N}{p_T dp_T dy d\phi}(p_T, y, \phi) = \frac{1}{2\pi} \frac{d^2 N}{p_T dp_T dy} [1 + 2v_2(p_T, y) \cos(2\phi) + \dots]$$



Successfully described by
(viscous) hydrodynamics

(... but no microscopic information)



Gale, Jeon, Schenke, Tribedy, Venugopalan,
Phys. Rev. Lett. 110 (2013)

Relativistic viscous fluid dynamics

- Basic equations: energy and momentum conservation

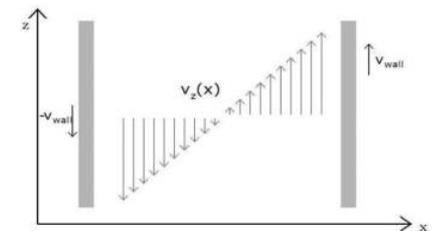
$$\partial_\mu T^{\mu\nu} = 0 \quad \text{with} \quad T^{\mu\nu} = (\overset{\text{energy density}}{\varepsilon} + \overset{\text{pressure}}{P}) \underset{\text{flow velocity}}{u^\mu} \underset{\text{viscous correction}}{\overset{\text{viscous correction}}{\Pi^{\mu\nu}}} - P g^{\mu\nu} + \Pi^{\mu\nu}$$

- Constituent equations for $\Pi^{\mu\nu}$

$$\Delta_\alpha^\mu \Delta_\beta^\nu (u \cdot \partial) \Pi^{\alpha\beta} = -\frac{1}{\tau_\pi} (\Pi^{\mu\nu} - S^{\mu\nu}) - \frac{4}{3} (\partial \cdot u) \Pi^{\mu\nu} \quad \frac{F_z}{A} = -\eta \frac{\partial v_z}{\partial x}$$

$$\text{with } S^{\mu\nu} = \eta \left(\nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} (\partial \cdot u) \right)$$

$$\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu \quad \text{and} \quad \nabla^\nu = \Delta^{\mu\nu} \partial_\nu$$

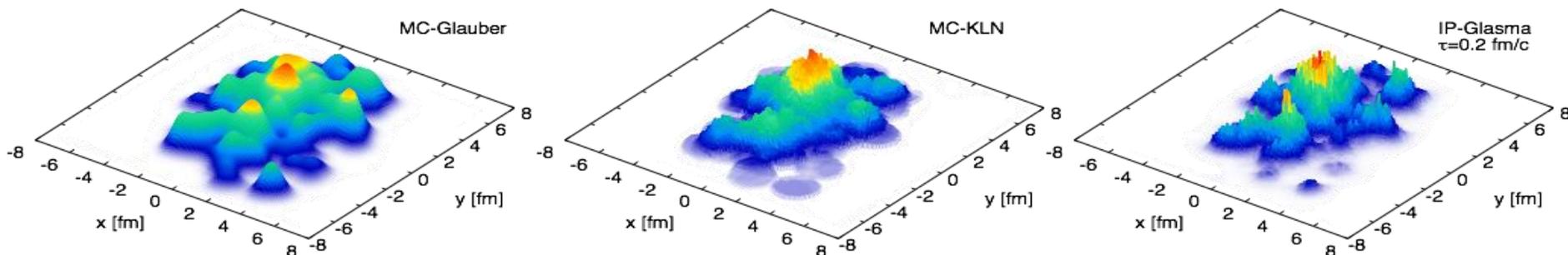


- Equation of state $P(\varepsilon)$ relates thermodynamic pressure to the energy density

η : shear viscosity

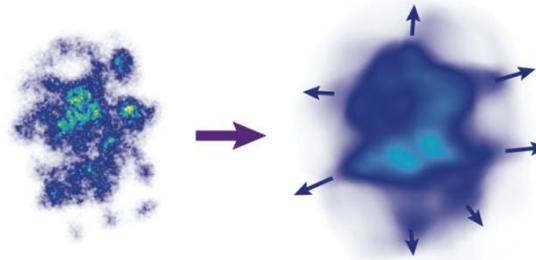
τ_π : relaxation time

- Models need to provide input for fluid dynamic simulations: initial energy density, flow velocities, shear stress tensor
- Initial conditions fluctuate from event to event
- Main source of fluctuations: nucleon positions
- Different models give different *energy density distributions*



C. Gale, S. Jeon, B. Schenke, Int.J.Mod.Phys. A28 (2013) 1340011

Anisotropic flow

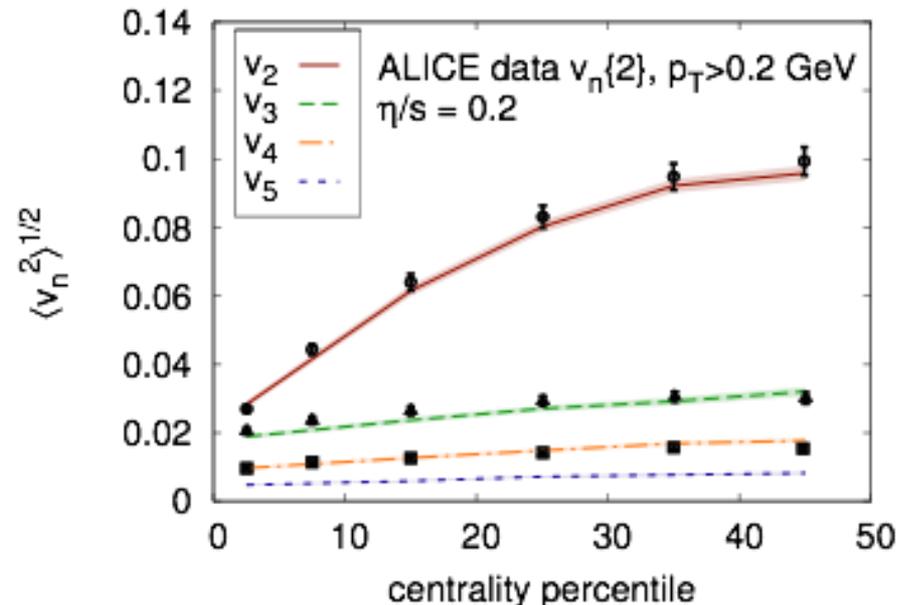


- System expands
- Initial shape is converted into particle momentum distribution
- Quantify its azimuthal anisotropy via Fourier expansion:
- Compare IP-Glasma+MUSIC simulation with ALICE data
- Extracted shear viscosity to entropy density ratio:

$$\eta/s = 0.2$$

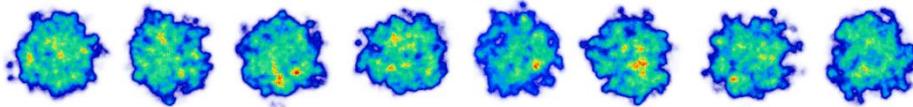
- Almost perfect fluidity!

$$\frac{dN}{d\phi} = \frac{N}{2\pi} \left(1 + \sum_n (2v_n \cos(n\phi)) \right)$$



Event-by-event fluctuations

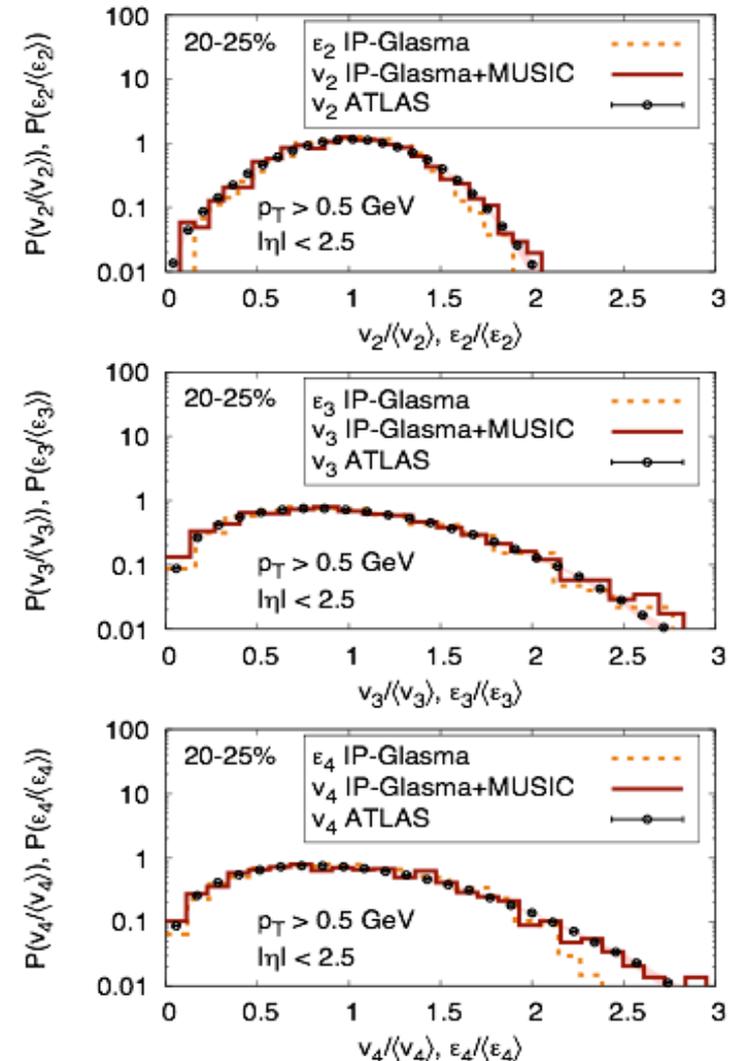
- Each event produces different v_n



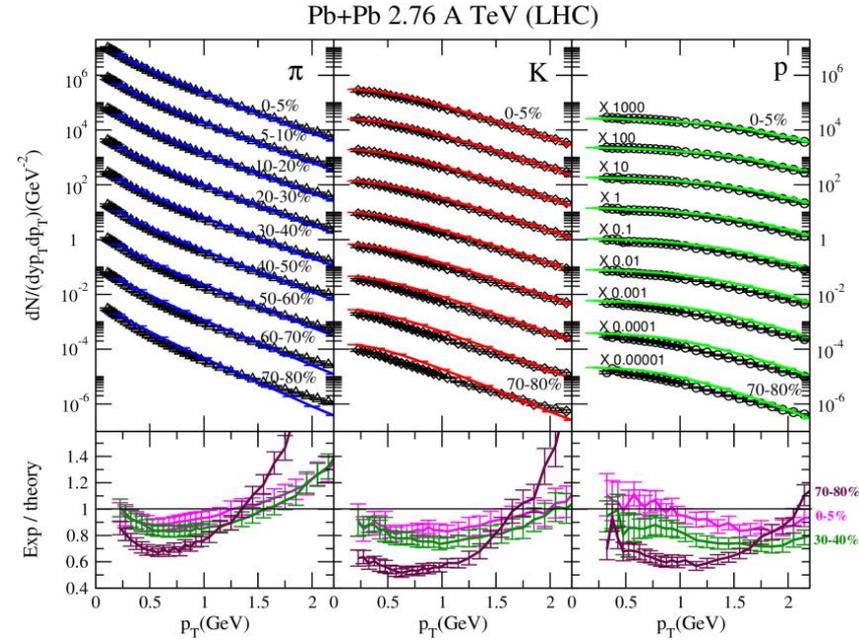
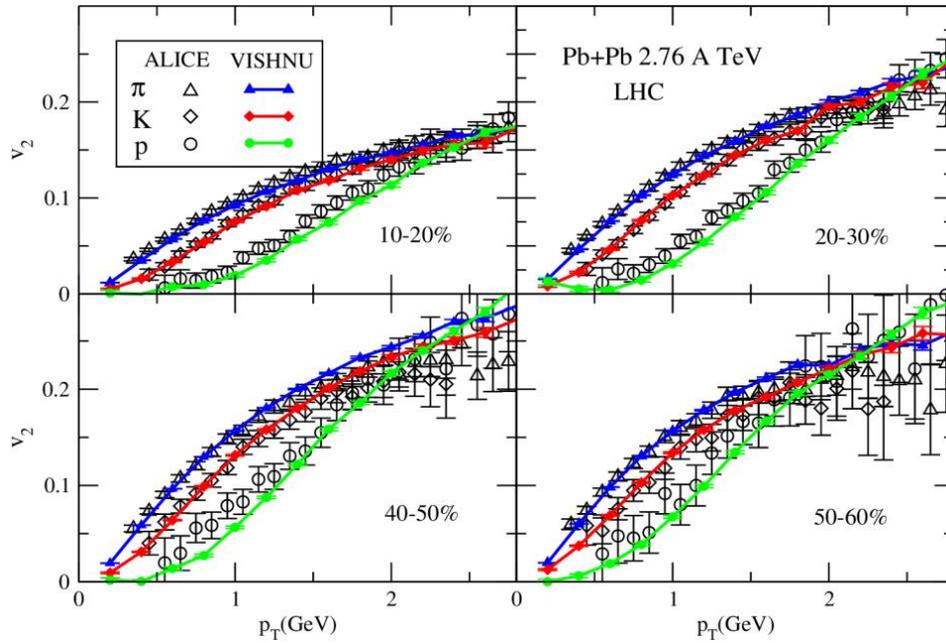
- Event-by-event distributions of v_n are also well reproduced by the simulations

- Fluid dynamics correctly describes the bulk dynamics of heavy-ion collisions

$$\eta/s = 0.2$$



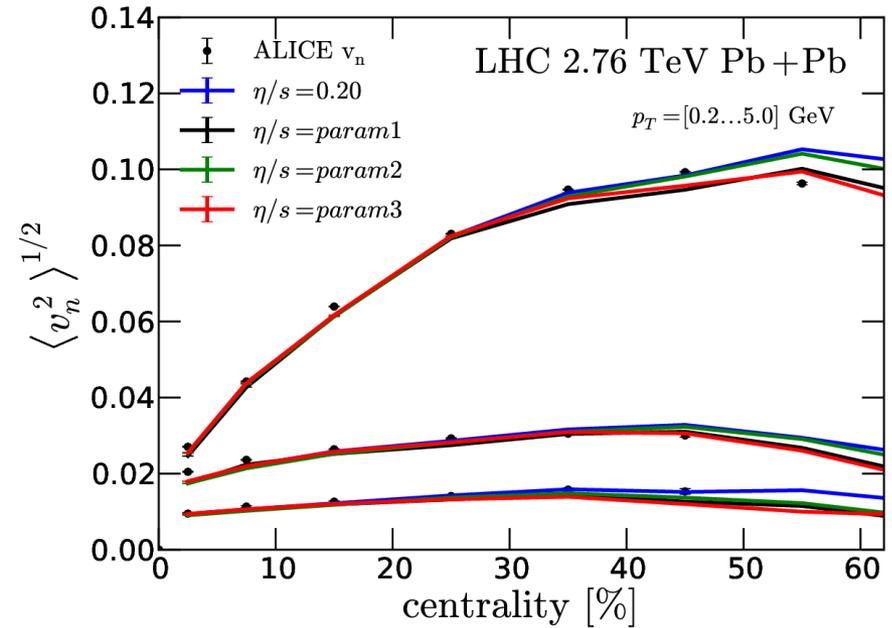
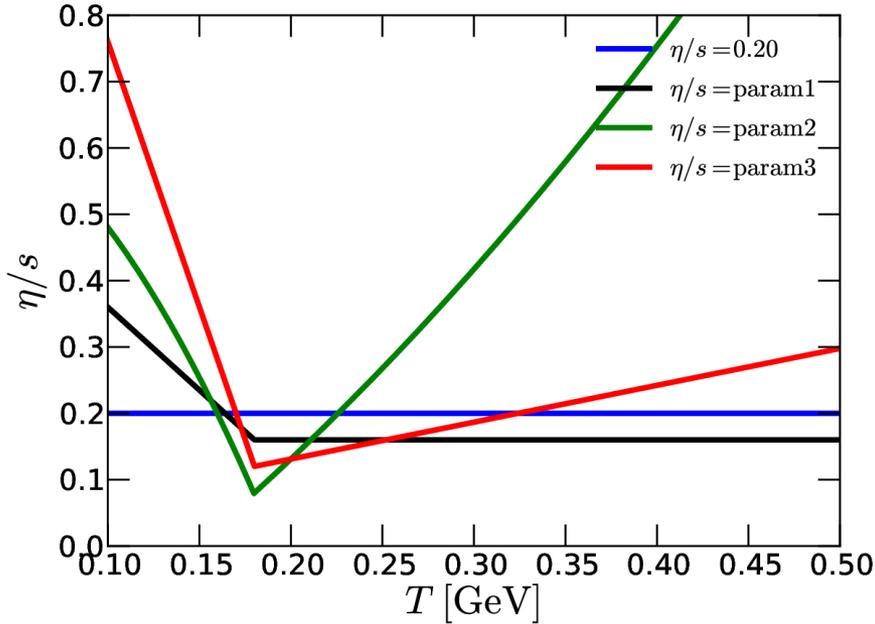
Identified hadrons π , K, p: Elliptic flow / p_t – spectra



H. Song, S. Bass and U. Heinz, Phys. Rev. C 89, 0349119 (2014)

(2+1)-D viscous hydro + UrQMD (VISHNU), $(\frac{\eta}{s})_{QGP} = 0.16$

$\eta/s(T)$ from LHC v_n s

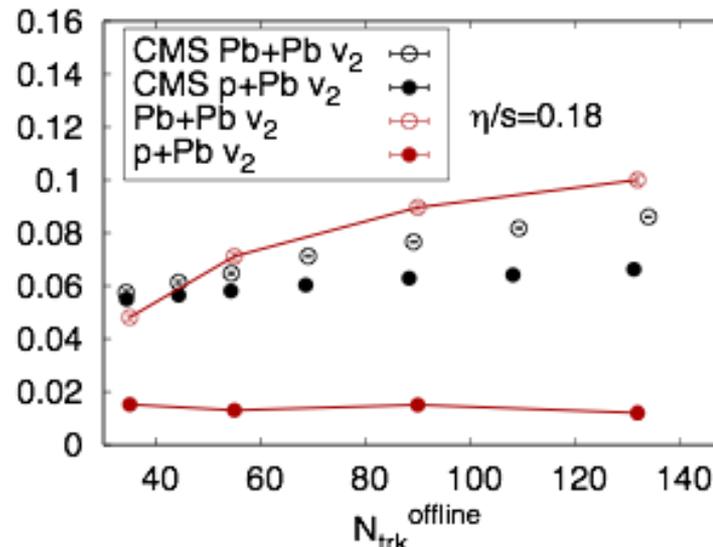
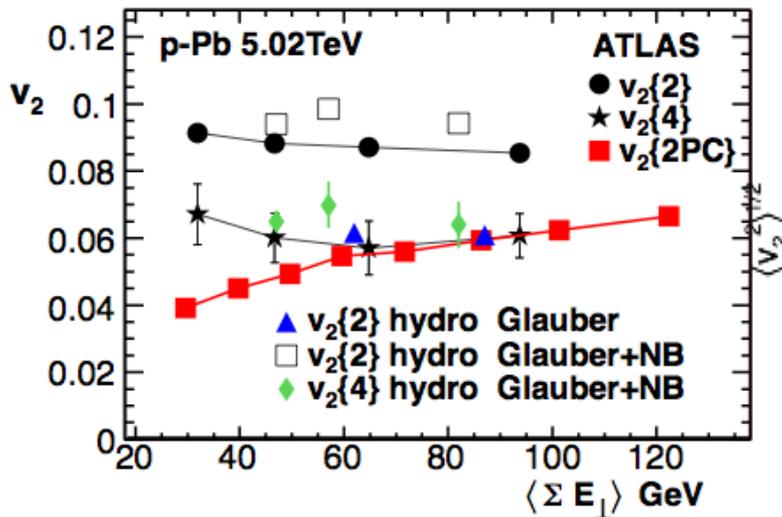
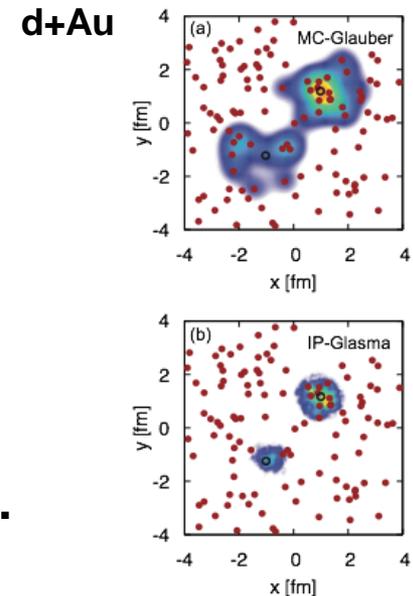


Paatelainen, Eskola, Niemi and Tuominen, Phys. Lett. B 731, 126 (2014)

... temperature dependence of the shear viscosity / entropy density ?

Evidence for fluid behavior in p+A collisions?

- Experiments find similar v_n in p+A as in A+A collisions CMS Collaboration, Phys.Lett. B724, 213 (2013)
- Can fluid dynamics work in such small systems? Viscous corrections become very large
- Initial state strongly depends on model
- Some models work, some do not. **Not yet settled...**

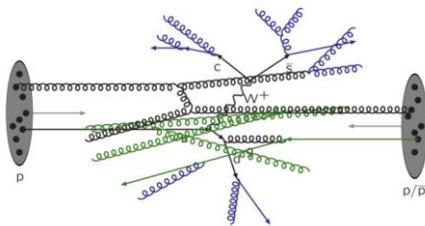


BAMPS: Boltzmann Approach to Multi-Parton Scatterings

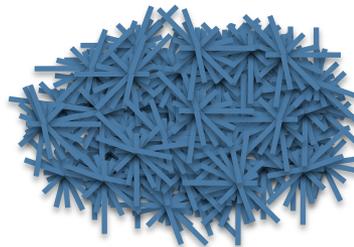
- 3+1 dimensional, fully dynamic parton transport model
- Boltzmann equations for on-shell partons with pQCD interactions

$$\left(\frac{\partial}{\partial t} + \frac{\mathbf{p}_i}{E_i} \frac{\partial}{\partial \mathbf{r}} \right) f_i(\mathbf{r}, \mathbf{p}_i, t) = C_i^{2 \rightarrow 2} + C_i^{2 \leftrightarrow 3} + \dots$$

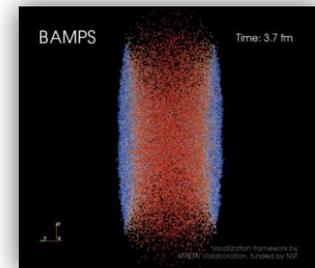
Z. Xu & CG,
Phys. Rev. C71 (2005)
Phys. Rev. C76 (2007)



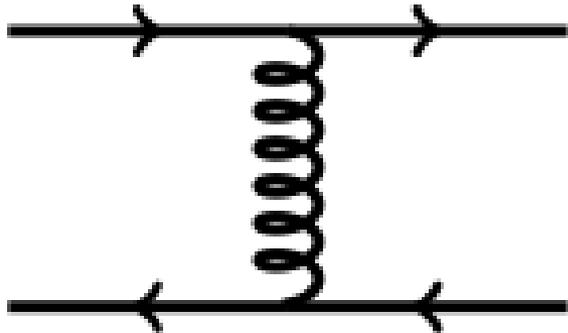
PYTHIA-Glauber



$$AA = pp \times N_{\text{binary}}$$



noneq. QGP



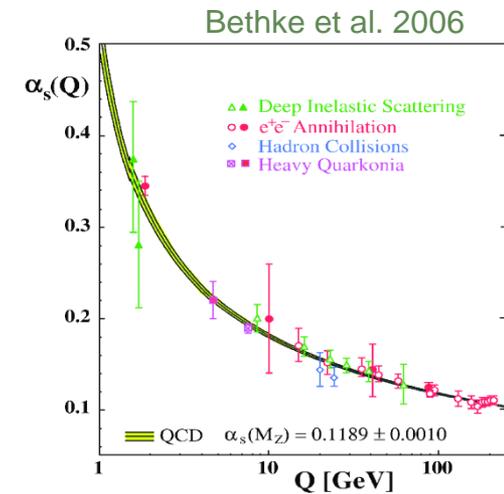
- Leading-order pQCD cross sections

$$|\overline{\mathcal{M}}_{qq' \rightarrow qq'}|^2 = \frac{64\pi^2}{9} \alpha_s^2(t) \frac{u^2 + s^2}{[t - m_D^2(\alpha_s(t))]^2}$$

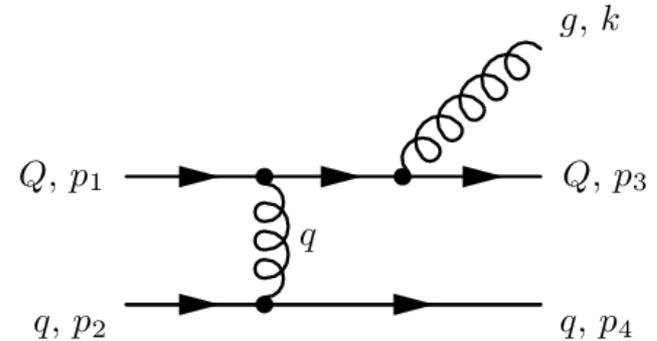
- Divergences screened by Debye mass

- Running coupling

$g g \rightarrow g g$		
$g g \rightarrow q \bar{q}$		
$q \bar{q} \rightarrow g g$	and	$q \bar{q} \rightarrow q' \bar{q}'$
$q g \rightarrow q g$	and	$\bar{q} g \rightarrow \bar{q} g$
$q \bar{q} \rightarrow q \bar{q}$		
$q q \rightarrow q q$	and	$\bar{q} \bar{q} \rightarrow \bar{q} \bar{q}$
$q q' \rightarrow q q'$	and	$q \bar{q}' \rightarrow q \bar{q}'$



$$\begin{aligned}
 &gg \leftrightarrow ggg \\
 &qg \leftrightarrow qgg \quad \text{and} \quad \bar{q}g \leftrightarrow \bar{q}gg \\
 &q\bar{q} \leftrightarrow qq\bar{q} \\
 &qq \leftrightarrow qqg \quad \text{and} \quad \bar{q}\bar{q} \leftrightarrow \bar{q}\bar{q}g \\
 &qq' \leftrightarrow qq'g \quad \text{and} \quad q\bar{q}' \leftrightarrow q\bar{q}'g
 \end{aligned}$$



Improved Gunion-Bertsch matrix element:

$$|\overline{\mathcal{M}}_{X \rightarrow Y+g}|^2 = |\overline{\mathcal{M}}_{X \rightarrow Y}|^2 P_g$$

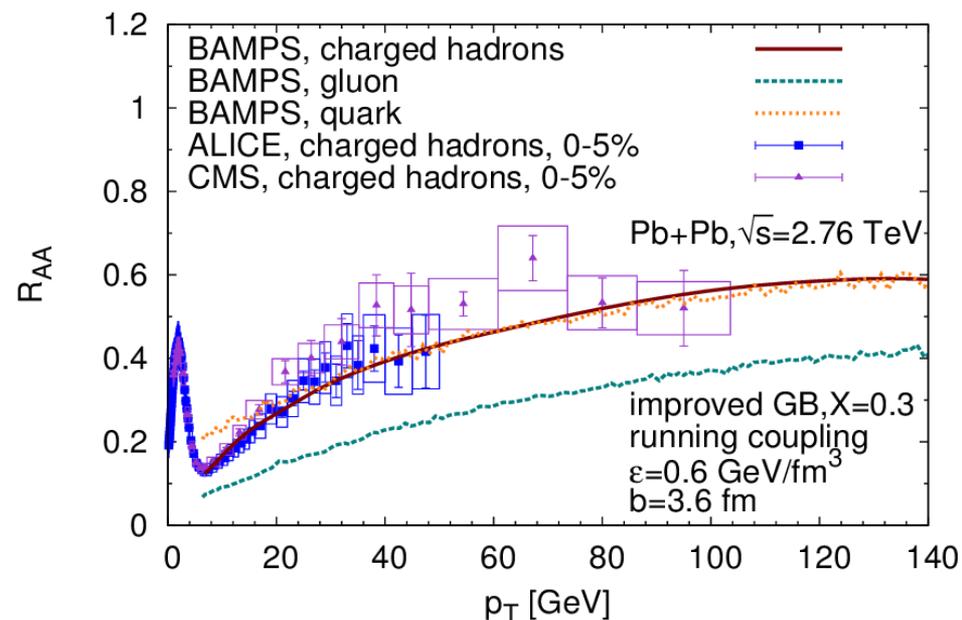
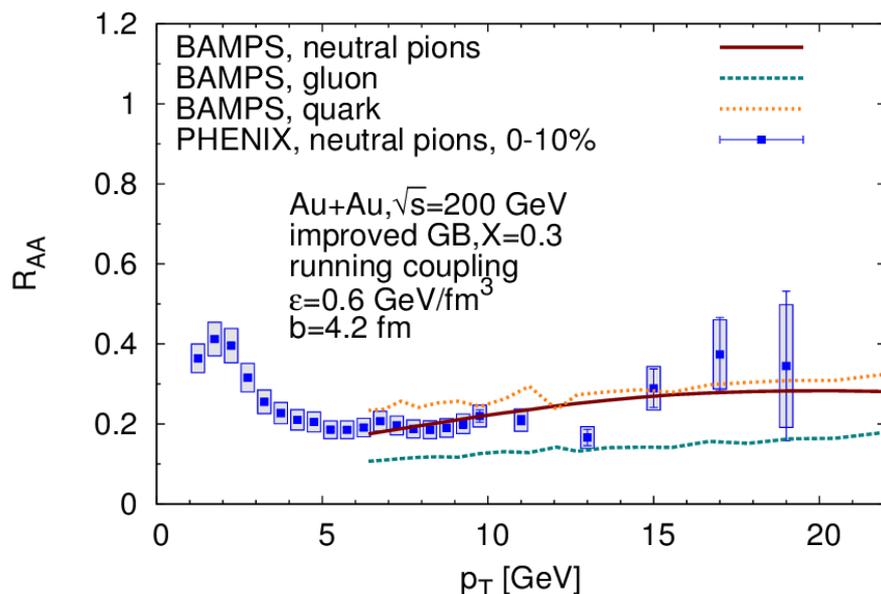
$$\begin{aligned}
 P_g &= 48\pi\alpha_s(k_\perp^2) (1 - \bar{x})^2 \\
 &\times \left[\frac{\mathbf{k}_\perp}{k_\perp^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + m_D^2 (\alpha_s(k_\perp^2))} \right]^2
 \end{aligned}$$

Effective QCD LPM effect:

$$\text{Mean free path} \quad \lambda > X \tau \quad \text{Gluon formation time}$$

Fochler, Uphoff, Xu, CG
Phys. Rev. D88 (2013)

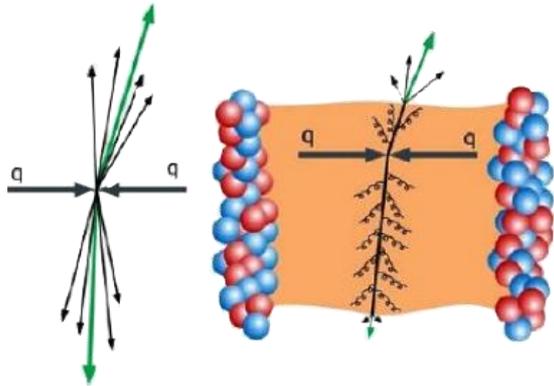
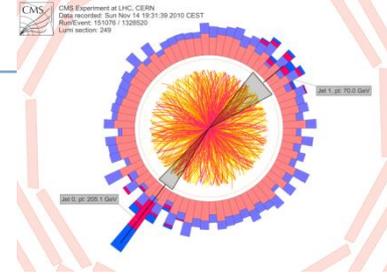
Nuclear modification factor R_{AA}



Uphoff et al, arXiv:1401.1364

- Hadronization of high p_t partons with AKK fragmentation functions
- LPM parameter fixed by comparison to RHIC data
- Realistic suppression both for RHIC and LHC

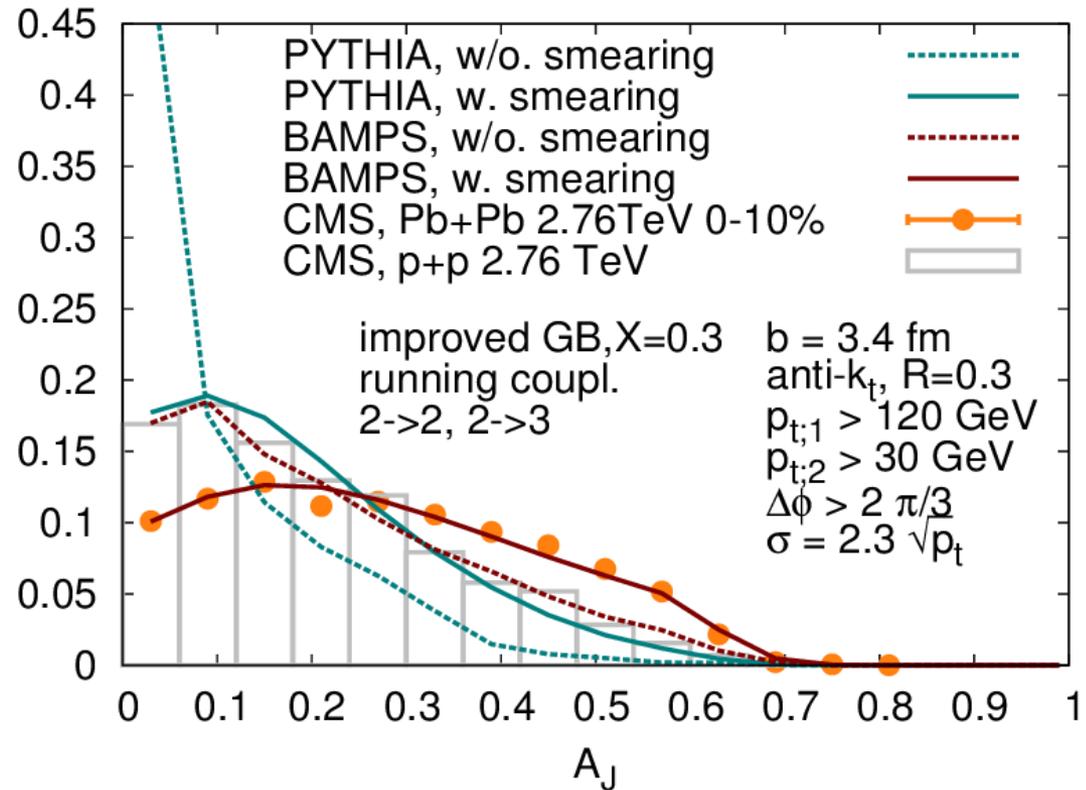
Reconstructed jets



Momentum imbalance

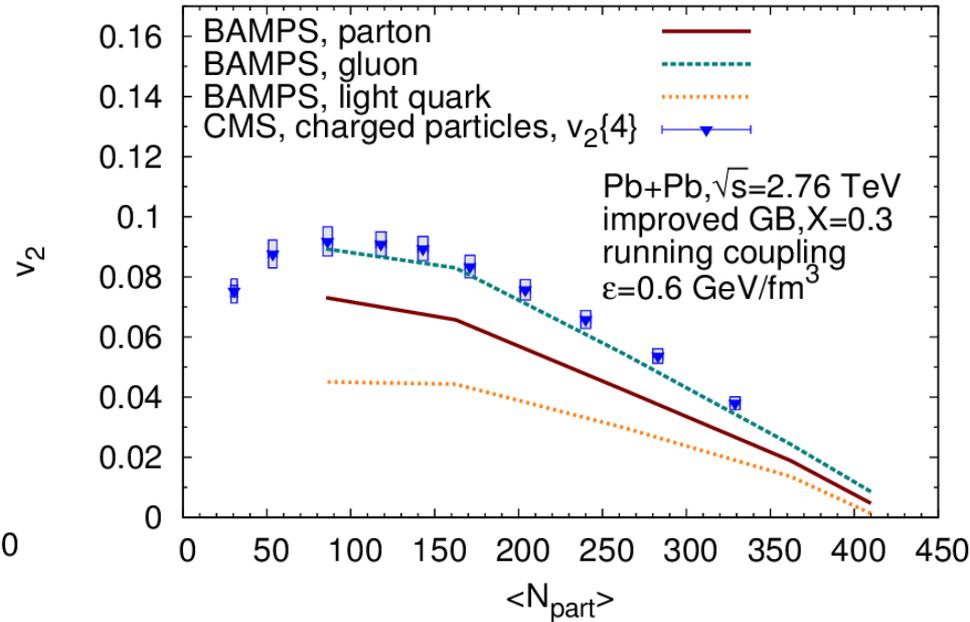
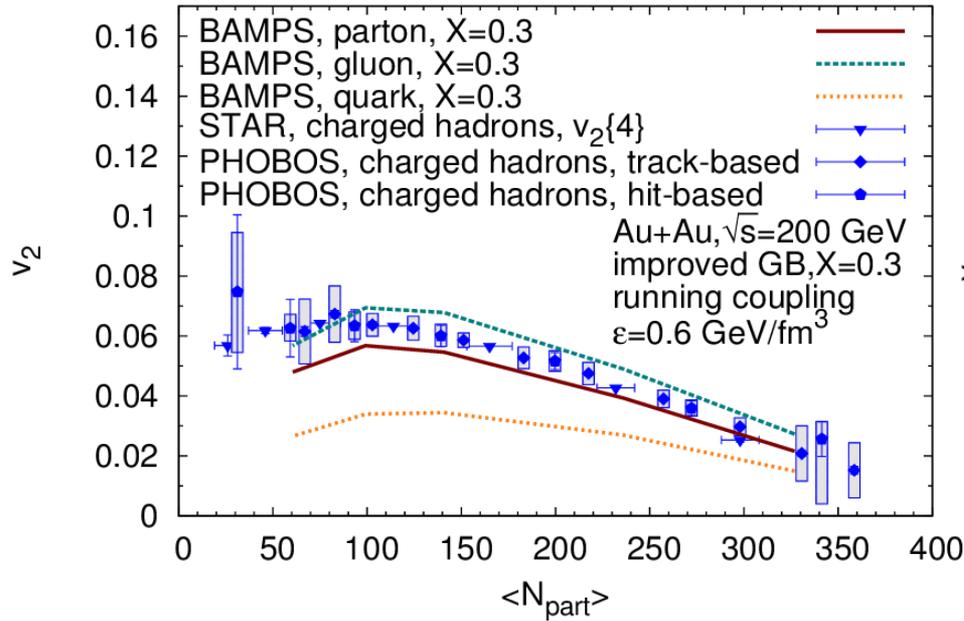
$$A_J = \frac{p_{t;leading} - p_{t;subleading}}{p_{t;leading} + p_{t;subleading}}$$

$P(A_J)$



Senzel et al, arXiv:1309.1657

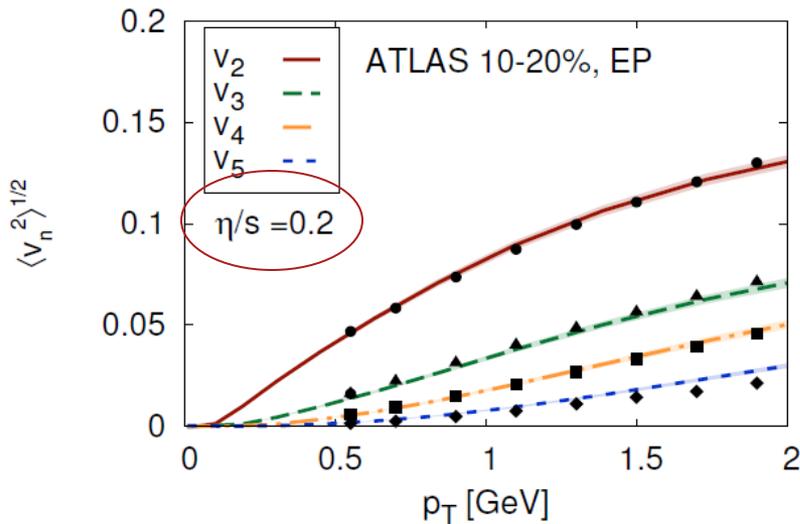
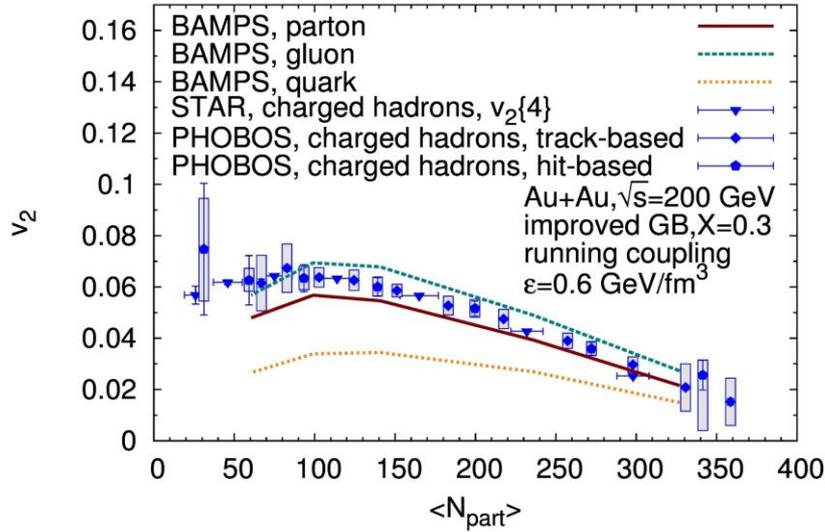
Elliptic flow v_2



Uphoff et al, arXiv:1401.1364

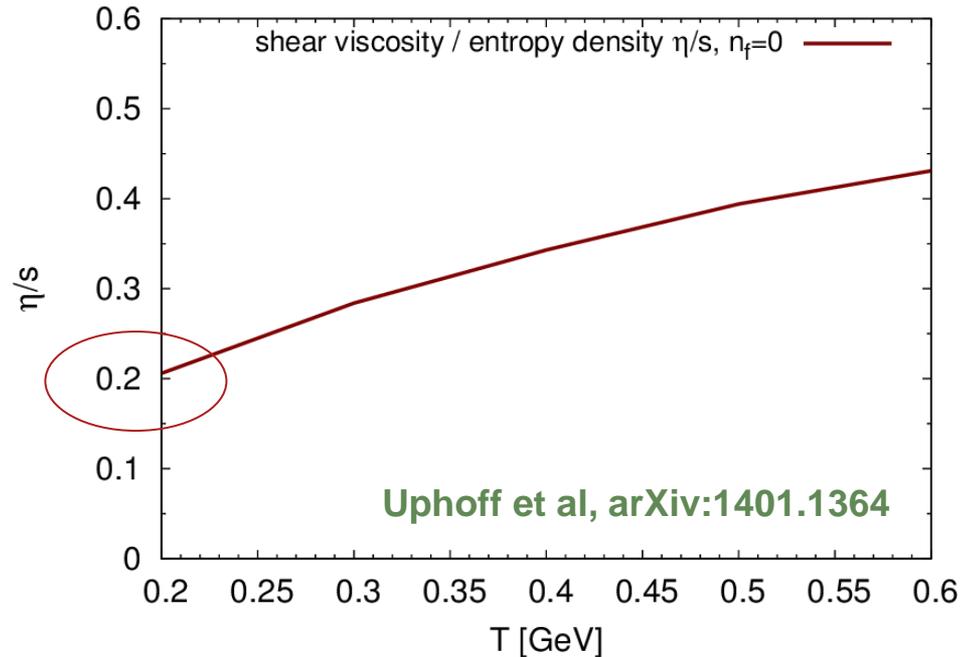
- Same pQCD interactions lead to a sizeable elliptic flow for bulk medium
- No hadronization for bulk medium \rightarrow no hadronic after-burner

Shear viscosity as QGP transport parameter



Reason for large elliptic flow:
Small shear viscosity to
entropy density ratio

From parameters to calculations:

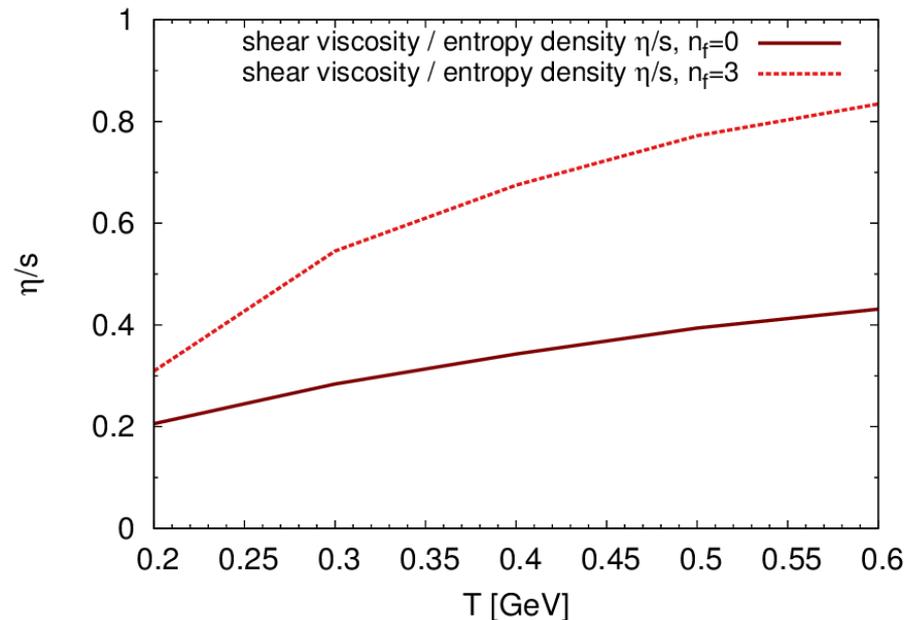
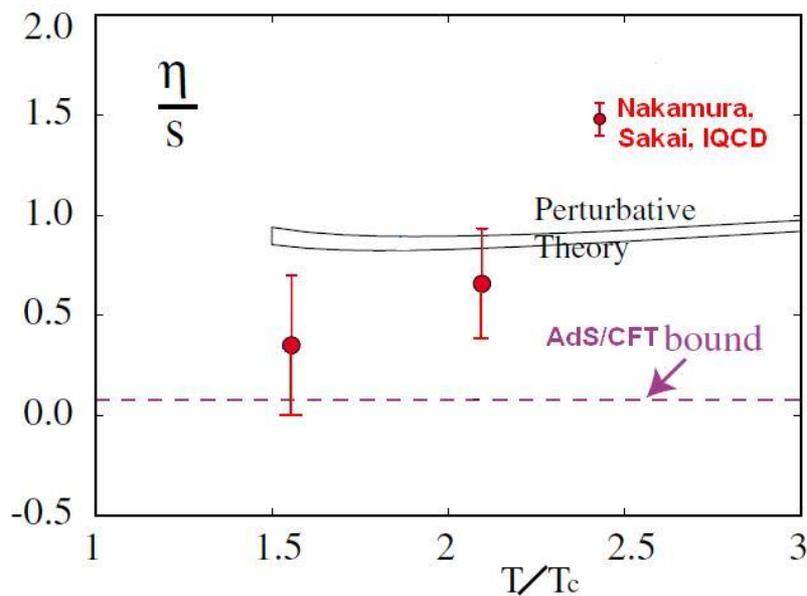


- Shear viscosity η
- bulk viscosity ζ
- heat conductivity κ
- electric conductivity σ

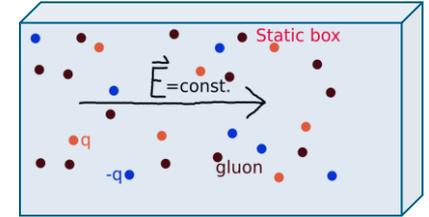
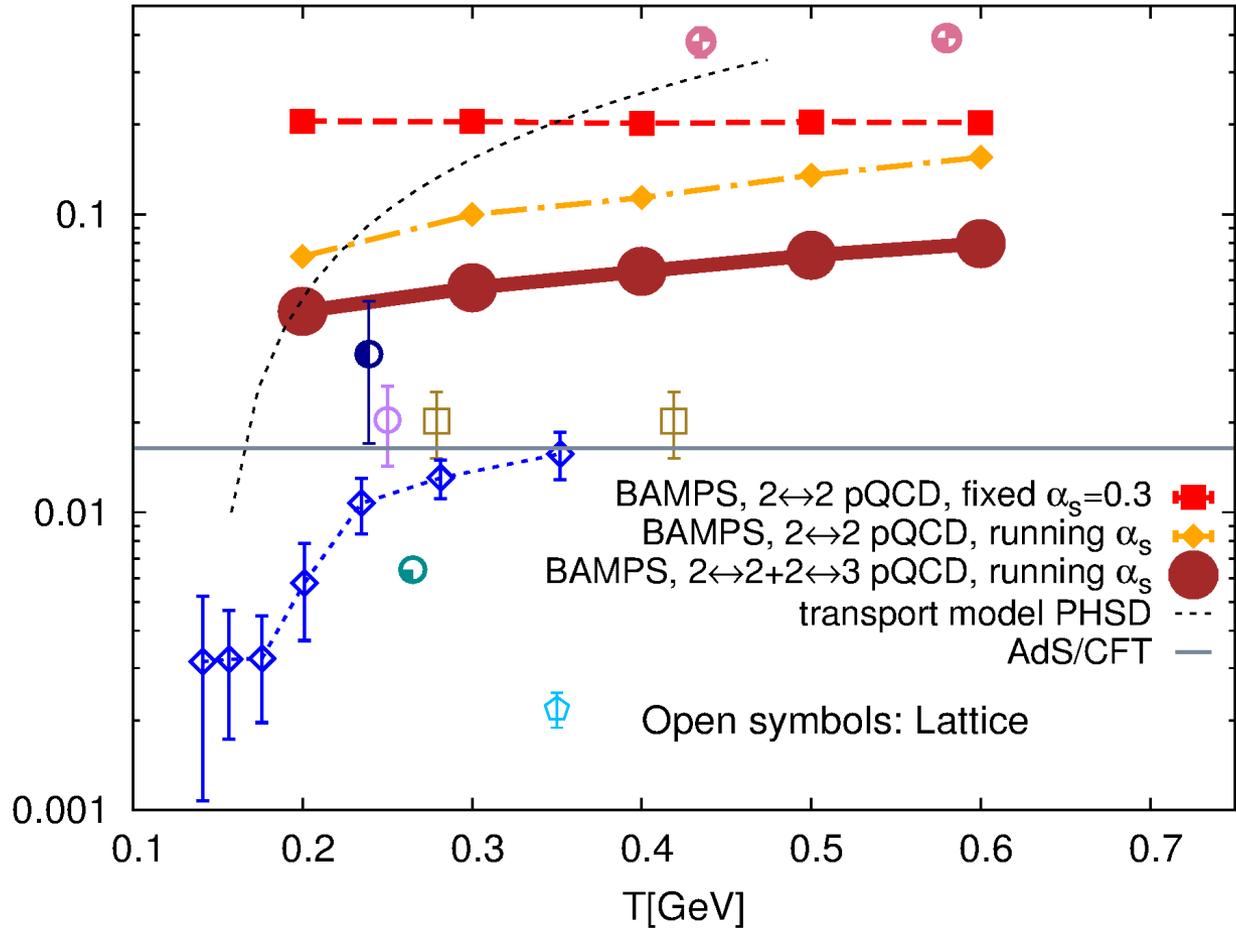
Other coefficients of interest:
Heavy quark diffusion constants, susceptibilities,...

Green-Kubo relation:

$$\eta = \frac{1}{10T} \int_0^\infty dt \int d^3r \langle \pi^{ij}(r,t) \pi^{ij}(0,0) \rangle$$



... electric conductivity



$$j = \sigma_{el} E$$

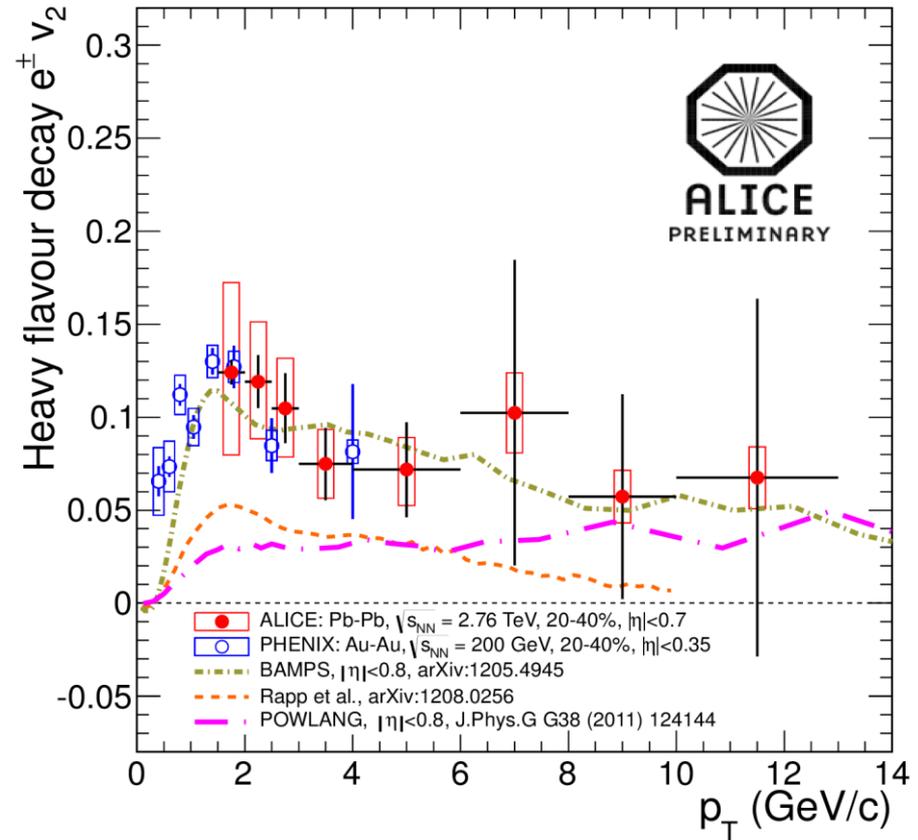
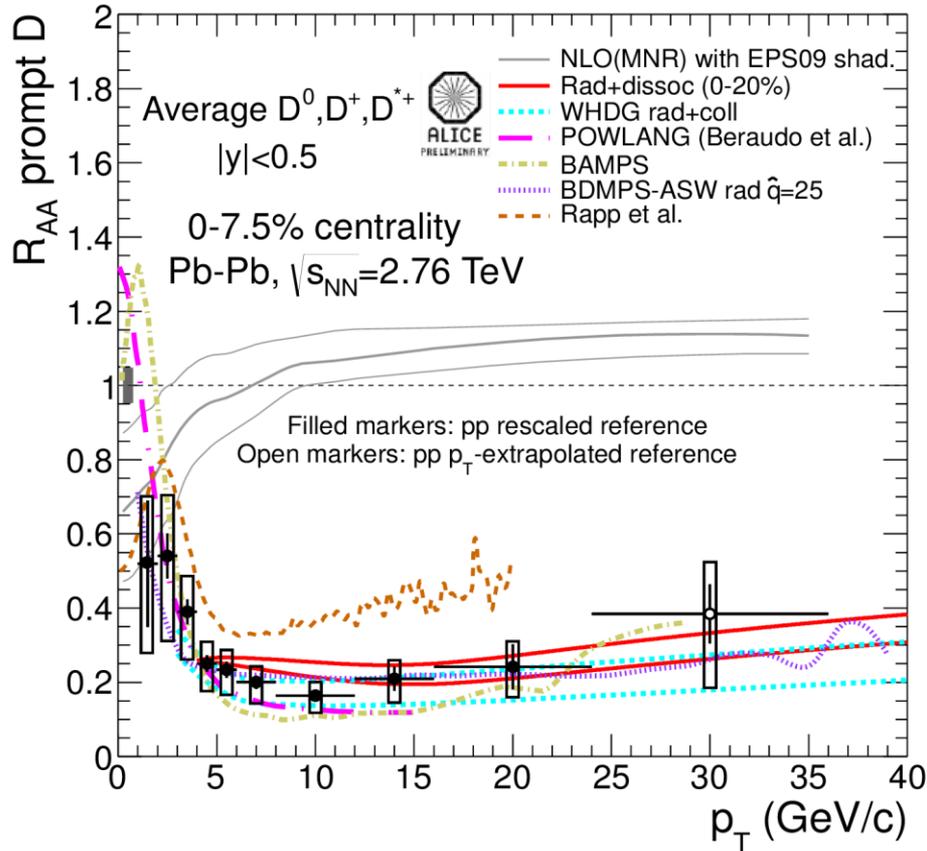
Green-Kubo relation:

$$\sigma_{el} = \frac{V}{T} \int_0^{\infty} \langle j_x(0) j_x(t) \rangle dt$$

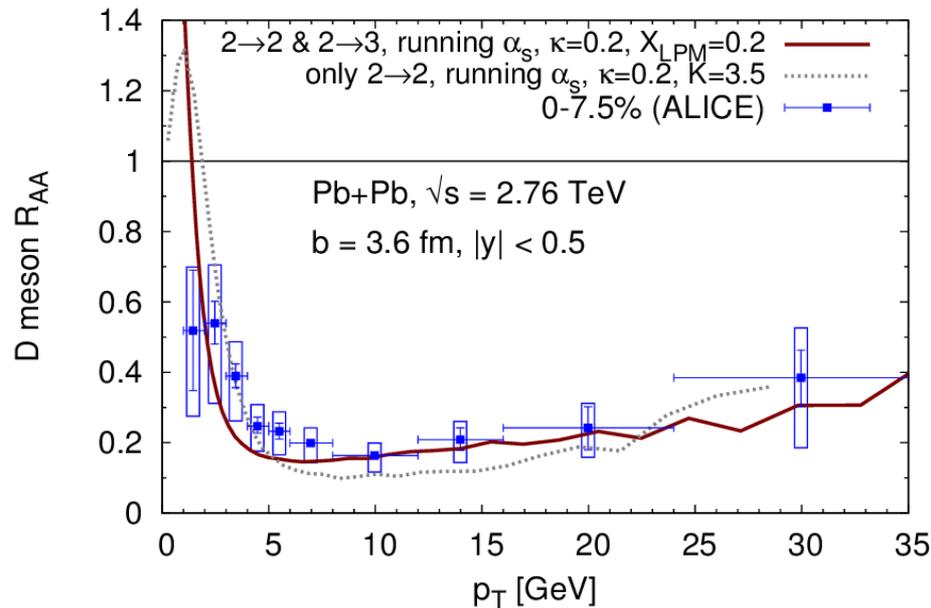
Lattice theory versus transport theory

... and learn about the (strongly) interacting QGP

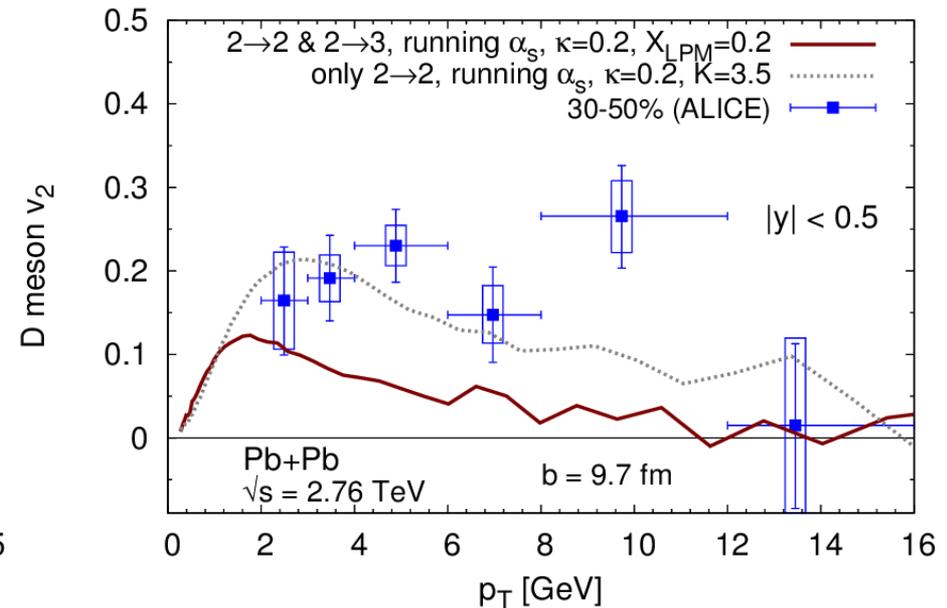
Heavy flavor R_{AA} and v_2 at LHC



D meson R_{AA} and v_2 at LHC



ALICE data, QM12



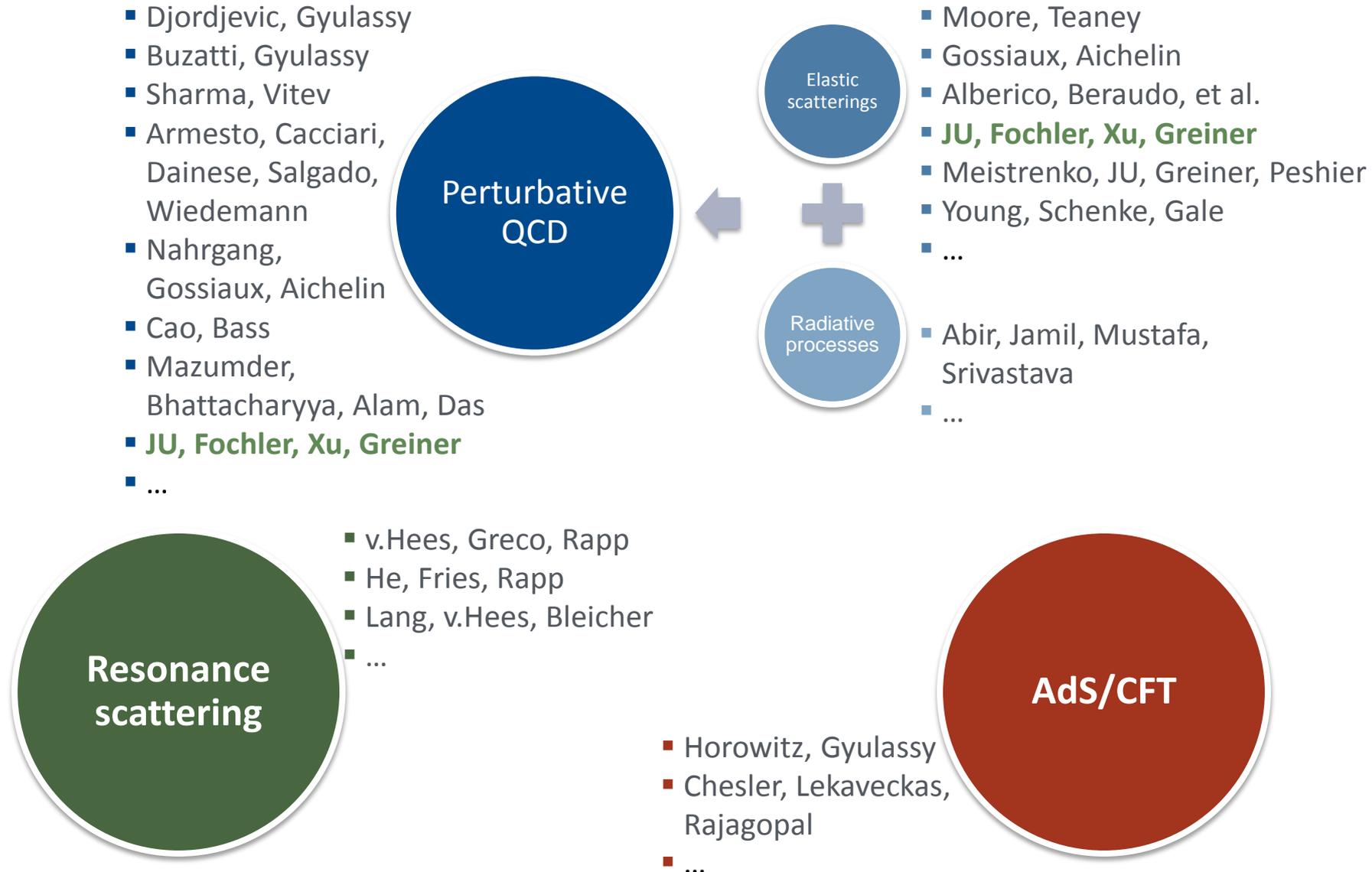
Uphoff, Fochler, Xu, CG, Phys. Lett. B 717 (2012),
 Uphoff, Fochler, Xu, CG, arXiv: 1408.2964

Heavy flavor R_{AA} and v_2 simultaneously seems difficult

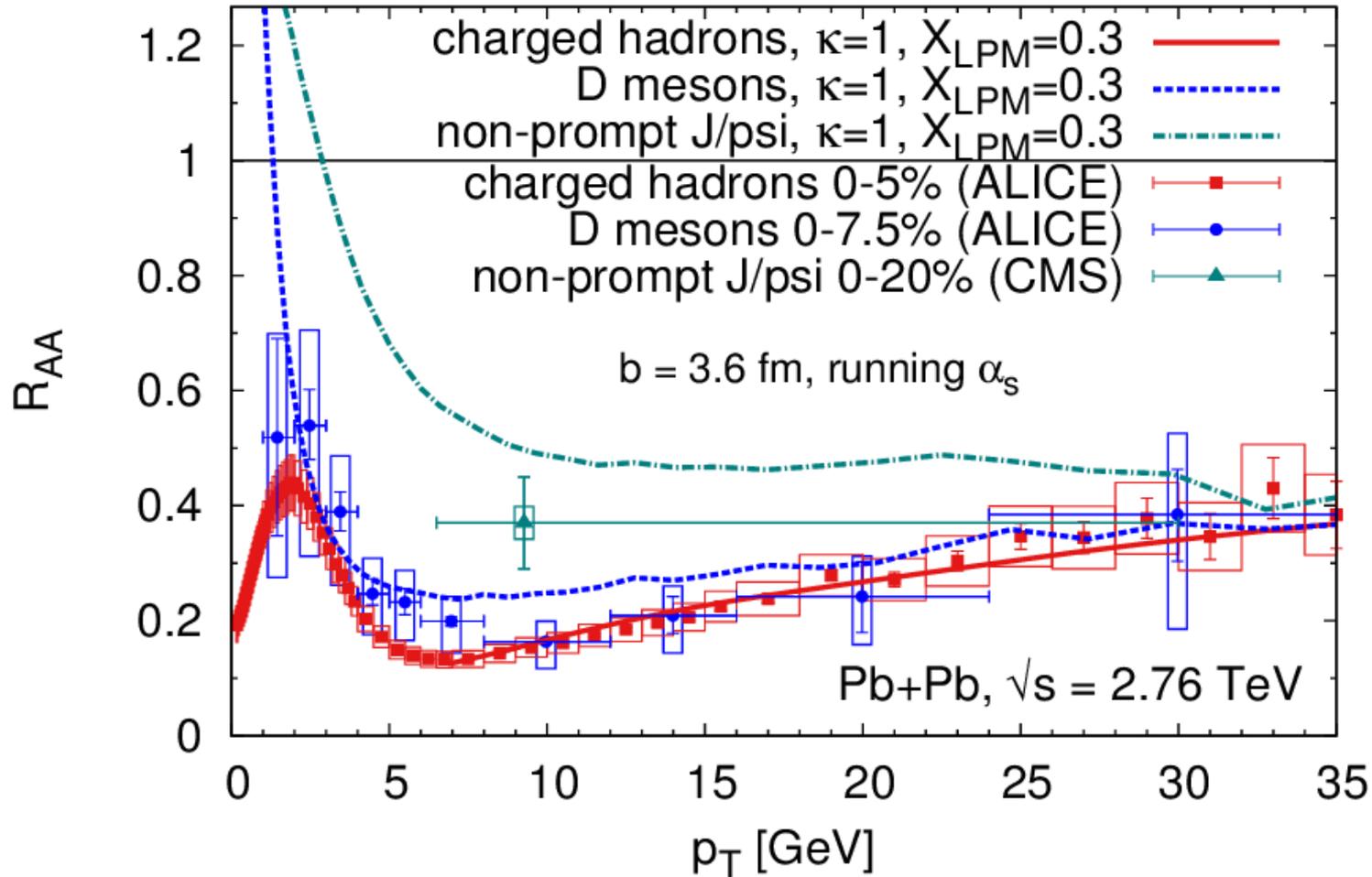
- Heavy-ion collisions provide an unique opportunity for investigating QCD matter under extreme conditions
- Strong collective behavior of the QGP is successfully described by fluid or transport dynamics
- Both at RHIC and the LHC, hard probes (high p_t and heavy flavor) are quenched while traversing the QGP
- Transport coefficients allow a connection between dynamical models and lattice QCD
- What can we learn for heavy-ion collisions from smaller systems (p+p, p+A, ...)?

Thank you for your attention.

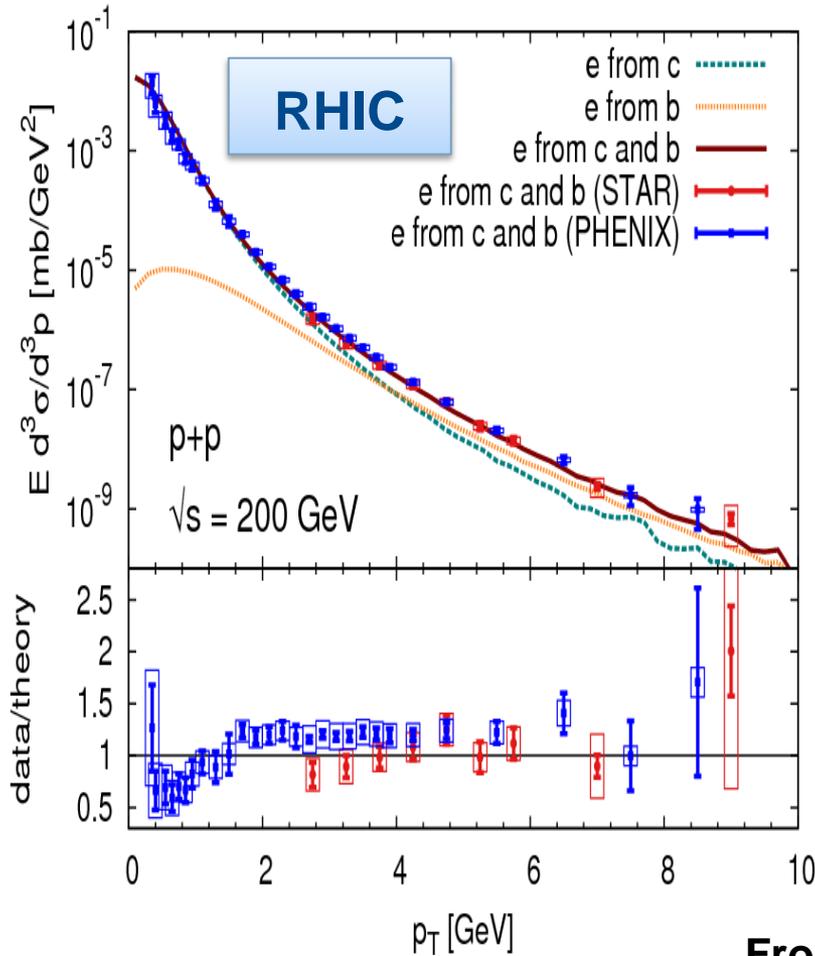
Heavy quark energy loss mechanism



Heavy flavor and charged hadron R_{AA} at LHC

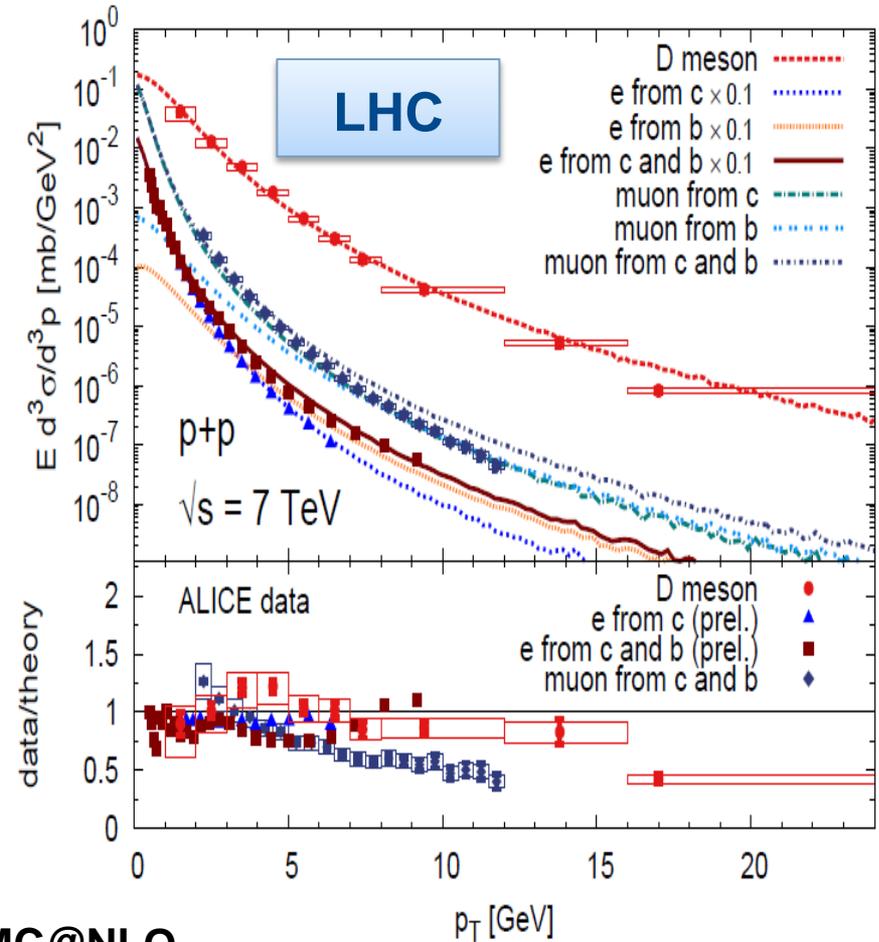


Initial heavy flavor spectrum



From MC@NLO

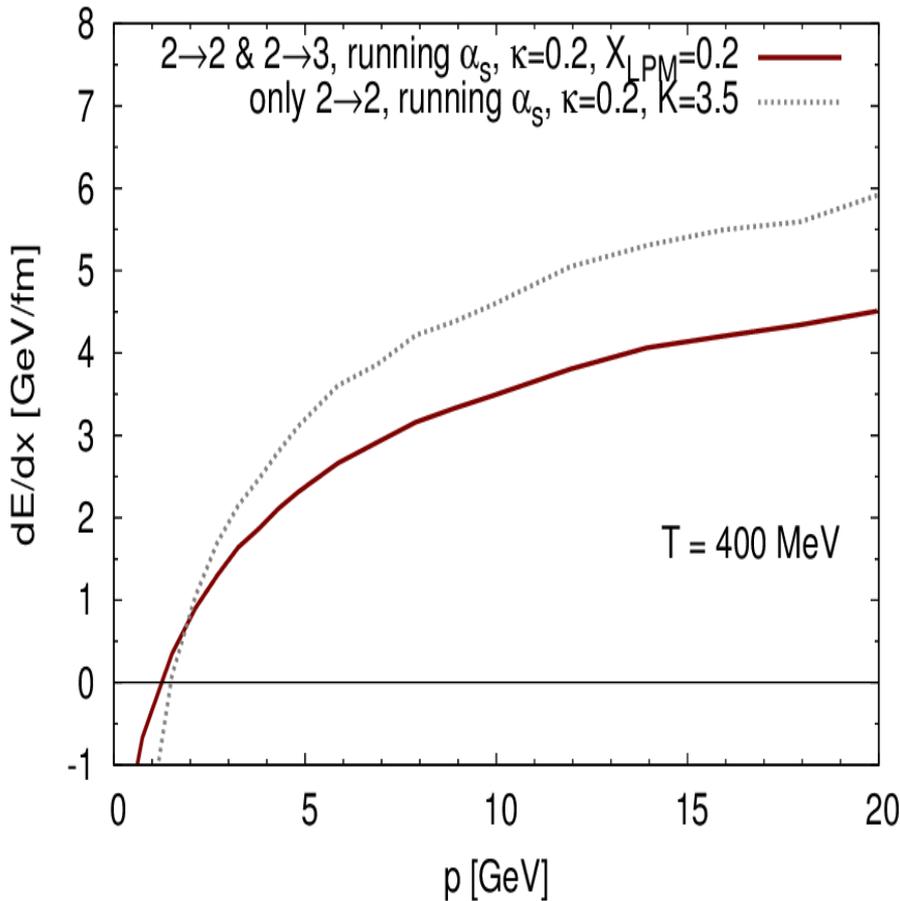
JU, Fochler, Xu, Greiner
Phys. Rev. C84 (2011)



JU, Fochler, Xu, Greiner
Phys. Lett. B 717 (2012)

Energy loss and transport cross section

Energy loss in static medium



Transport cross section in static medium

