Forward-backward multiplicity correlations in pp collisions at high energy in Monte Carlo model with string fusion

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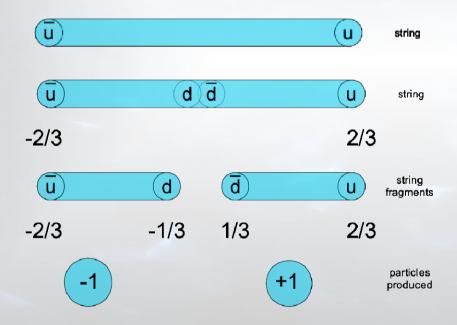


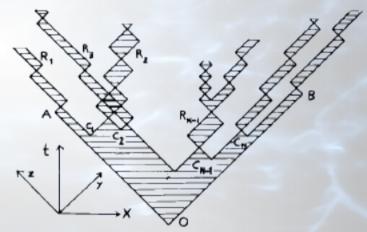




Overview

- The soft QCD processes is not described by usual perturbation theory
- The model of quark-gluon strings, stretched between projectile and target partons
- semiphenomenological approach to the multiparticle production

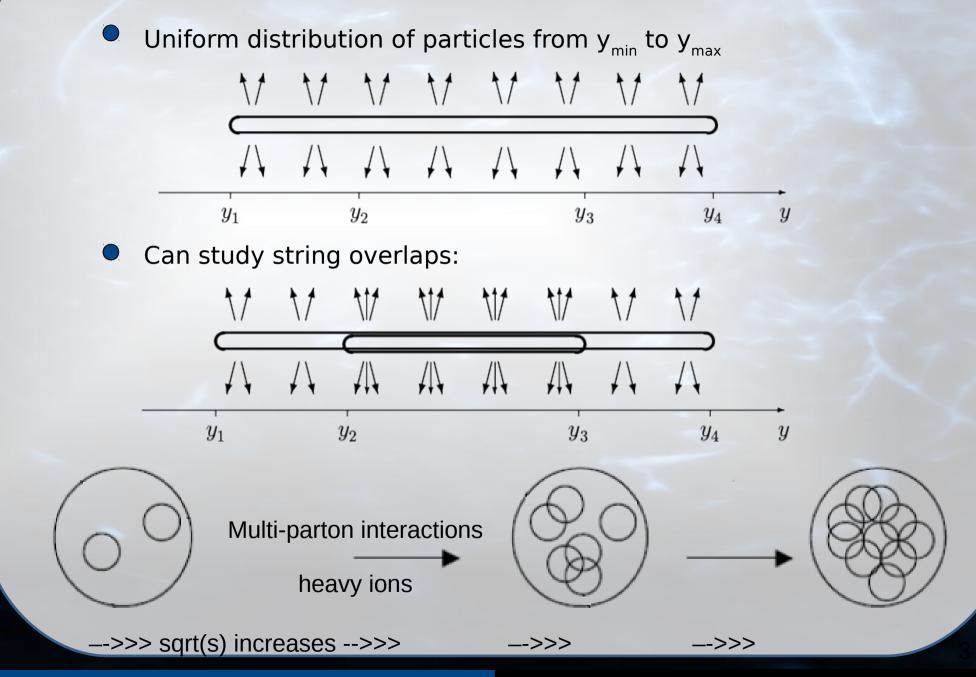




X. Artru and G. Mennessier, Nucl Phys B 70 (1974) 93 "String Model and Multiproduction",

- Correlations play crucial role:
- causality requires appearance of long-range correlations if they exist at the very early stages between particles detected in separated rapidity intervals

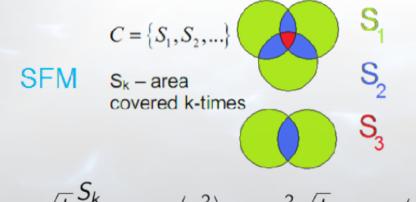
String in rapidity space



String fusion

$$Q^{2}(n) = \left(\sum_{i=1}^{n} \overrightarrow{Q_{i}}(1)\right)^{2} = \sum_{i=1}^{n} Q_{i}^{2}(1) + \sum_{i \neq j} \overrightarrow{Q_{i}}(1) \cdot \overrightarrow{Q_{i}}(1)$$
$$\left\langle Q^{2}(n) \right\rangle = nQ^{2}(1)$$

overlaps



String fusion mechanism predicts:

- decrease of multiplicity
- increase of p_T
- growth of p_T with multiplicity
- in pp, pA and AA collisions

 growth of strange particle yields

results are in a good agreement with the experiment

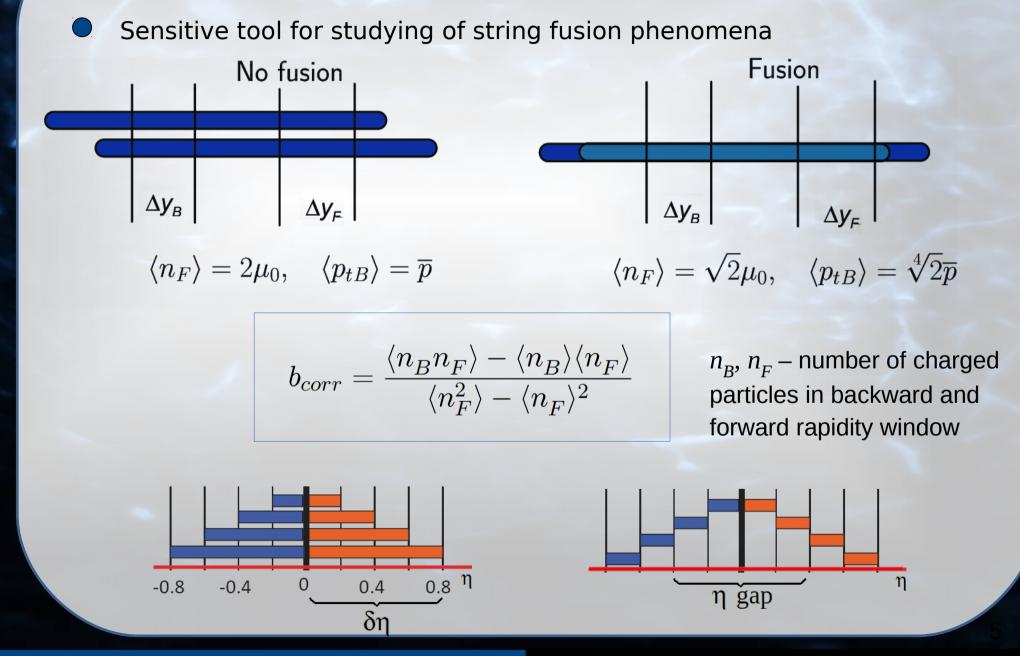
$$\langle \mu \rangle_k = \mu_0 \sqrt{k} \frac{S_k}{\sigma_0} \qquad \langle p_t^2 \rangle_k = p_0^2 \sqrt{k} \qquad \langle p_t \rangle_k = p_0 \sqrt[4]{k}$$

 S_k – area, where k strings are overlapping, σ_0 single string transverse area, μ_0 and p_0 – mean multiplicity and transverse momentum from one string

M. A. Braun, C. Pajares, Nucl. Phys. B 390 (1993) 542.

M. A. Braun, R. S. Kolevatov, C. Pajares, V. V. Vechernin, Eur. Phys. J. C 32 (2004) 535.
N.S. Amelin, N. Armesto, C. Pajares, D. Sousa, Eur.Phys.J.C22:149-163 (2001), arXiv:hep-ph/0103060
G. Ferreiro and C Pajares J. Phys. G: Nucl. Part. Phys. 23 1961 (1997)

Long-range (forward-backward) correlations



Monte Carlo model

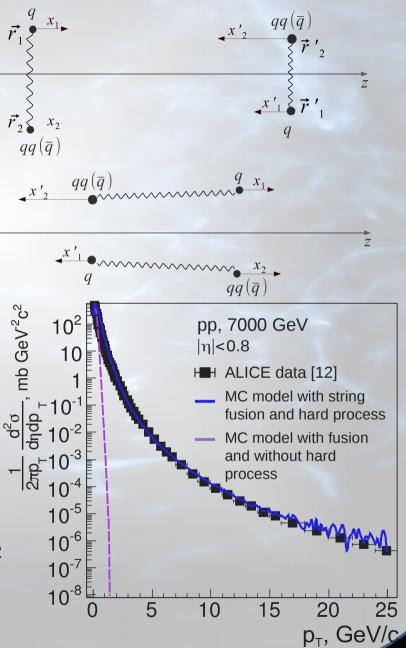
- Partonic picture based on dipole interaction
- Energy and angular momentum conservation in the initial state
- The probability amplitude depends on transverse coordinates:

$$f = \frac{\alpha_s^2}{2} \ln^2 \frac{|\vec{r}_1 - \vec{r}_1'| |\vec{r}_2 - \vec{r}_2'|}{|\vec{r}_1 - \vec{r}_2'| |\vec{r}_2 - \vec{r}_1'|}$$

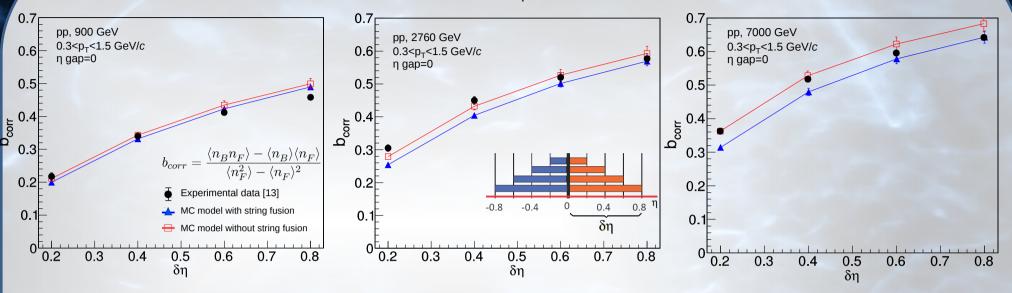
• The **hardness** of the elementary collisions is defined by transverse size of dipoles:

 $d_{1i} = |\vec{r_1} - \vec{r_2}|, d_i' = |\vec{r_1}' - \vec{r_2}'|$ • Transverse momentum of a cluster of strings:

$$p_1^4 = \sum_{i}^{k} p_{T stri}^4$$
, $p_{T stri}^2 = \frac{1}{d_i^2} + \frac{1}{d_i'^2} + p_0^2$

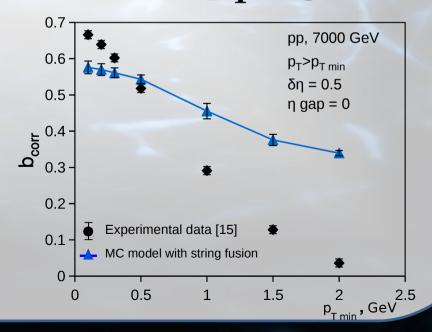


Dependence of b_{corr} **on window size**

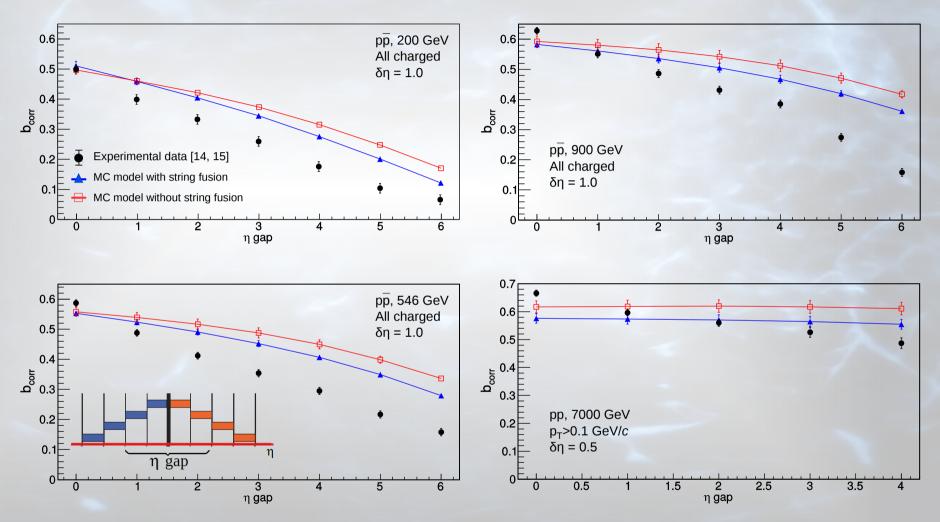


- The general trends (growth of b_{corr} with energy and pseudorapidity windows width) are well described by the model
- The data does not allow to distinguish between fusion/no fusion
- Qualitative agreement of the dependence on the lower bound of the p_T range

and on p_T region



Dependence of b_{corr} **on pseudorapidity gap**



• The model describes the behavior of b_{corr} in wide energy range

- At small η gap the data is contaminated by short-range effects
- The model with string fusion is better supported by the data

Summary and conclusions

- The forward-backward multiplicity correlations in pp collisions are studied in the Monte Carlo model with string formation and fusion
- The Monte Carlo model reasonably describes the main features of the correlation coefficient in a wide energy range:
 - The general growth of the correlation coefficient with collision energy
 - Growth of b_{corr} with increase of the pseudorapidity window size and decrease of b_{corr} with increase of the gap between windows
 - Decrease of b_{corr} with increase of lower p_T bound.
- The model with string fusion is better supported by the data, while the case without string fusion is disfavored

End Of Presentation

Thank you!



Backup



Forward-backward multiplicity correlations in pp collisions at high energy in Monte Carlo model with string fusion



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pp. 200 GeV

Abstract and motivation

The magnitude of the correlations between multiplicities in two separated rapidity windows, proposed as a tool for study of the string fusion and percolation phenomena [1], is studied in the framework of the Monte Carlo string-parton model [2-3]. The model is based on the picture of strings formation in elementary collisions of color dipoles. The hardness of the elementary collisions is defined by a transverse size of the interacting dipoles.

In the framework of the model the charged particles spectra with the account of string interaction in the transverse plane is calculated. The interaction of string is is realized in the accordance with the string fusion model [4-6] prescriptions by the introduction of the lattice in the impact parameter plane and taking into account the finite rapidity length of strings. The parameters of the model were fixed with the experimental data on total inelastic cross section and charged multiplicity. The dependencies of the forward-backward correlation strength on the width and position of the pseudorapidity windows and on the transverse momentum range of observed particles were studied. The detailed modeling of the charged particles spectra allowed to make a direct comparison to the results of experimental measurements at different energies.



• Partonic picture of nucleons inte	eraction [2, 3].	
 Energy and angular momentum 	conservation i	in the initial state
of a nucleon.	2 ° ×1.	$x', qq(\bar{q})$
 The probability of dipoles 	15	P'2
interaction depends on their transverse coordinates [7-8] with effective coupling :	\vec{F}_{2} \vec{x}_{2} $qq(\bar{q})$	<u></u>
$f = \frac{\alpha_s^2}{2} \ln^2 \frac{ \vec{r}_1 - \vec{r}_1' \vec{r}_2 - \vec{r}_2' }{ \vec{r}_1 - \vec{r}_2' \vec{r}_2 - \vec{r}_1' }$	x', qq(q)	
$j = \frac{1}{2} \prod_{i=1}^{n} \vec{r}_1 - \vec{r}_2' \vec{r}_2 - \vec{r}_1' $	**:•~~~~~ q	$qq(\bar{q})$

Multiplicity and transverse momentum are obtained in the approach of <u>colour strings</u>, stretched between projectile and target partons.

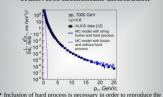
• The interaction of strings is realized in the accordance with the string fusion model prescriptions [4-6]. Mean multiplicity and the mean trans verse momentum of the particles produced from a cluster of strings are:

 $\begin{array}{l} \langle \mu \rangle_k = \mu_j \sqrt{k} \frac{S_k}{\sigma_j} \qquad \langle p_t^2 \rangle_k = p_j^2 \sqrt{k} \qquad \langle p_t \rangle_k = p_j \sqrt{k} \\ S_k - \text{ area, where } k \text{ strings overlapping, } \sigma_i = \pi r_{str}^2 - \text{ single string transverse area, } \mu_i \text{ and } p_i - \text{ mean multiplicity and } \end{array}$

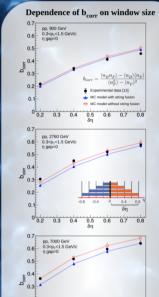
transverse momentum from one string. Multiplicity from one string is distributed according to Poisson distribution, with Gaussian transverse momentum spectra. The **hardness** of the elementary collisions is either constant (p_0) – in case of no hard subprocess; or it is defined by a transverse size of the interacting dipoles, similarly to DIPSY [9, 10]: $a_{ij} = [\vec{r}_i - \vec{r}_i], a_j' = [\vec{r}_i' - \vec{r}_i']$

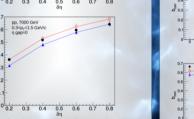
• In case of hard process, the transverse momentum of a cluster of strings: $p_i^* = \sum^n p_{exi}^*$, where $p_{exi}^* = \frac{1}{dx^2} + \frac{1}{dx^2} + p_n^*$. Every parton can interact with anther one only once (contrary to Glauber supposition of constant nucleon cross section). • Parameters of the model are constrained from the data on total inelastic cross-section and multiplicity [2, 11]. We used r_{exi} =0.2*Im* (in case with string fusion), and $p_i = 0.2$ *GeV/c* in case with the hard process or $p_i = 0.4$ *GeV/c* without hard process.

Transverse momentum distribution

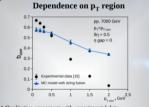


 Inclusion on nate process is necessary in order to reproduce the transverse momentum spectra of charged particles in pp collisions.
 Reasonable good description of transverse momentum spectra of charged particles in the MC model with string fusion and hard process included.





The general trends (growth of b_{corr} with energy and pseudorapidity windows width) are well described by the model
 The data does not allow to distinguish between fusion/no fusion

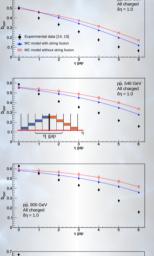


Qualitative agreement with experimental data
 Decrease of the correlation coefficient on the lower bound of
the p_r range due to lack of phase space and less multiplicity in
the window

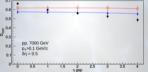
- Sensitivity to the details of the \mathbf{p}_{T} spectra

Summary and conclusions

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- The model with string fusion is better supported by the data, while the case without string fusion is disfavored



Dependence of b_{corr} on pseudorapidity gap



The model describes the behavior of b_{corr} in wide energy range
At small η gap the data is contaminated by short-range effects
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