#### PANIC 2014 Particles and Nuclei International Conference Hamburg University, August 25-29, 2014

#### **Parity Violating Electron Scattering**



# Recent Results and Future Prospects

#### Krishna Kumar

UMass, Amherst & Stony Brook U.



Historical Overview Electroweak Measurement of the Neutron **Skin of a Heavy Nucleus** ★ PREX and CREX at Jefferson Lab Precision Measurements of Weak Charges ★ Future Program with the JLab 12 GeV Upgrade ★ (Followup talks on 6 GeV program by P. King and V. Sulkosky)

Parity-Violating Electron Scattering

# Historical Perspective





Fourier transform of charge distribution



## **Weak Interactions**

Neutron & nuclear  $\beta$  Decay



charge and flavor-changing

Fermi Theory for weak interactions Universal strength: coupling constant G<sub>F</sub>

"Effective" low energy theory that explains many observed properties of radioactive nuclear decays

### Observed NOT to be invariant under parity transformations Weak Interactions

Neutron & nuclear  $\beta$  Decay



charge and flavor-changing

parity transformation (reflection)

$$x, y, z \rightarrow -x, -y, -z$$



Fermi Theory for weak interactions Universal strength: coupling constant G<sub>F</sub>

"Effective" low energy theory that explains many observed properties of radioactive nuclear decays

Parity-Violating Electron Scattering

## Observed NOT to be invariant under parity transformations Weak Interactions

Neutron & nuclear  $\beta$  Decay



Fermi Theory for weak interactions Universal strength: coupling constant G<sub>F</sub>

"Effective" low energy theory that explains many observed properties of radioactive nuclear decays



## **Observed NOT** to be invariant under parity transformations **Weak Interactions**



Universal strength: coupling constant GF "Effective" low energy theory that explains many

observed properties of radioactive nuclear decays



Magnetic

field

#### Zel'dovich speculation: Is Electron Scattering Parity-Violating? Parity Violation Signature JETP 36, pp 964-66 (1959)

Parity-Violating Electron Scattering

















## Glashow, Weinberg and Salam: SU(2)<sub>L</sub>XU(1)<sub>Y</sub> Weak Interaction Theory

**The Z boson incorporated** One free parameter: weak mixing angle  $\theta_W$ 

	Left-	Right-
γ Charge	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$
W Charge	$T = \pm \frac{1}{2}$	zero
Z Charge	$T-q\sin^2\theta_W$	$-q\sin^2\theta_W$

Glashow, Weinberg and Salam: SU(2) <sub>L</sub> X U(1) <sub>Y</sub>				
M	<b>leak</b> I	ntera	ction Theory	
The Z boson incorporated One free parameter: weak mixing angle $\theta_W$				
	Left-	Right-	μ	
γ Charge	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	N M	
W Charge	$T = \pm \frac{1}{2}$	zero	μW	
Z Charge	$T-q\sin^2\theta_W$	$-q\sin^2\theta_W$	Charged Current	

Glashow, Weinberg and Salam: SU(2) <sub>L</sub> XU(1) <sub>Y</sub> Weak Interaction Theory					
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	γ Charge	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	v n	V all
	W Charge	$T = \pm \frac{1}{2}$	zero	μ W	μ Ζ΄
	Z Charge	$T - q\sin^2\theta_W$	$-q\sin^2\theta_W$	Charged Current	Neutral Current
-					

Glashow, Weinberg and Salam: SU(2) <sub>L</sub> XU(1) <sub>Y</sub> Weak Interaction Theory						
The Z b	oson incorpo	prated On	ne free parameter: weak	$x$ mixing angle $\theta_W$		
	Left-	Right-		μ		
γ Charge	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$v_{\mu}$	VI 70 Y		
W Charge	$T = \pm \frac{1}{2}$	zero	μ			
Z Charge	$T - q\sin^2\theta_W$	$-q\sin^2\theta_W$	Charged Current	Neutral Current		
<b>Z</b> Charge $T - q\sin^2 \theta_W - q\sin^2 \theta_W$ $f = \int_{Z_0} f \frac{g}{\cos \theta_W} Z_\mu \bar{f} \gamma^\mu (T_{3f} - 2Q_f \sin^2 \theta_W - T_{3f} \gamma_5) f,  T_{3f} = \pm 1/2$						

Glashow, Weinberg and Salam: SU(2) <sub>L</sub> X U(1) <sub>Y</sub> Weak Interaction Theory						
The Z b	<b>The Z boson incorporated</b> One free parameter: weak mixing angle $\theta_W$					
	Left-	Right-	μ	Ϋμ		
γ Charge	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	$0,\pm 1,\pm \frac{1}{3},\pm \frac{2}{3}$	V. Wt	V. 70 V		
W Charge	$T = \pm \frac{1}{2}$	zero	μ	μΖ		
Z Charge	$T - q\sin^2\theta_W$	$-q\sin^2\theta_W$	Charged Current	Neutral Current		
$f = \begin{bmatrix} 1 & q \sin \theta_W \\ \hline q & q \sin \theta_W \end{bmatrix} = \frac{q \sin \theta_W}{T_{3f} - 2Q_f \sin^2 \theta_W - T_{3f} \gamma_5} f,  T_{3f} = \pm 1/2$						

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E122 at SLAC (1978): PV Deep Inelastic Scattering **Anatomy of a Parity Experiment**  $\int_{\gamma, Z'} \sigma \alpha |A_{\gamma} + A_{weak}|^{2} = A_{PV} = \frac{\sigma - \sigma}{\sigma + \sigma}$ Need few x 10<sup>11</sup> events  $\int_{\gamma, Z'} Count at \sim 100 \text{ kHz}$   $\int_{\gamma, Z'} \delta(A_{PV}) \sim few ppm$ longitudinally polarized e

# E122 at SLAC (1978): PV Deep Inelastic Scattering **Anatomy of a Parity Experiment** $\int_{\text{Negitudinally polarized target}} \left[ \sigma \alpha | A_{\gamma} + A_{\text{weak}} |^2 \right]^2 = \left[ \sigma \phi - \sigma_{\phi} -$



•Optical pumping of a GaAs wafer: "black magic" chemical treatment to boost quantum efficiency

•Rapid helicity reversal: polarization sign flips ~ 100 Hz to minimize the impact of drifts

Accelerator •Helicity-correlated beam motion: under sign flip, beam stability at the micron level





## Anatomy of a Parity Experiment



# Anatomy of a Parity Experiment



## Anatomy of a Parity Experiment




## Anatomy of a Parity Experiment



## Anatomy of a Parity Experiment



•Parity Violation in Weak Neutral Current Interactions • $sin^2\theta_w = 0.224 \pm 0.020$ : same as in neutrino scattering

$$A_{PV} \sim 10^{-4}$$
  
$$\delta(A_{PV}) \sim 10^{-5}$$

Glashow, Weinberg, Salam Nobel Prize awarded in 1979

Parity-Violating Electron Scattering

**Z**\*

e-

e-

Krishna Kumar, August 28 2014



photocathodes, polarimetry, high power cryotargets, nanometer beam stability, precision beam diagnostics, low noise electronics, radiation hard detectors

# The Neutron Skin of a Heavy Nucleus

Pb-Radius EXperimentEW Probe of Neutron Densities
$$\mathcal{V}$$
 $\mathcal{M}^{EM} = \frac{4\pi\alpha}{Q^2} F_p(Q^2)$  $M_{PV}^{NC} = \frac{G_p}{\sqrt{2}} [(1 - 4\sin^2\theta_w)F_p(Q^2) - F_n(Q^2)]$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^n_W \sim -1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^n_W \sim -1$ 

Pb-Radius EXperimentEW Probe of Neutron Densities
$$\mathcal{M}^{EM} = \frac{4\pi\alpha}{Q^2} F_p(Q^2)$$
 $M_{PV}^{NC} = \frac{G_F}{\sqrt{2}} [(1 - 4\sin^2\theta_w)F_p(Q^2) - F_n(Q^2)]]$  $A_{PV} \approx \frac{G_FQ^2}{4\pi\alpha\sqrt{2}} \frac{F_n(Q^2)}{F_p(Q^2)}$  $Q^p_{EM} \sim 1$  $Q^p_{EM} \sim 1$  $Q^n_{EM} \sim 0$  $Q^n_W \sim -1$  $Q^p_W \sim 1 - 4sin^2\theta_W$  $\frac{Proton neutron}{Electric charge 1 0}$  $\frac{Weak charge ~0.08 -1}{V}$ 

Pb-Radius EXperiment EW Probe of Neutron Densities				
$M^{EM} = \frac{4\pi\alpha}{Q^2} F_p(Q^2) \qquad M^{NC}_{PV} = \frac{G_F}{\sqrt{2}} \Big[ \Big( 1 - 4\sin^2\theta_W \Big) F_p(Q^2) - F_n(Q^2) \Big] \\ A_{PV} \approx \frac{G_F Q^2}{4\pi\alpha\sqrt{2}} \frac{F_n(Q^2)}{F_n(Q^2)} \qquad Q^p_{EM} \sim 1 \qquad Q^n_{EM} \sim 0 \qquad Q^n_W \sim -1  Q^p_W \sim 1 - 4\sin^2\theta_W$				
$A_{\rm PV} \sim 0.6 \rm ppm$			proton	neutron
$Q^2 \sim 0.01 \text{ GeV}^2$ $\longrightarrow$ Rate ~ 1 C	GHz	Electric charge	1	0
$5^{\circ}$ scattering angle $\delta(A_{PV}) \sim 2$	20 ppb!	Weak charge	~0.08	-1





Parity-Violating Electron Scattering

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## PREX at JLab

1 GeV electron beam, 50-70 μA high polarization, ~89% helicity reversal at 120 Hz





0.5 mm isotopically pure <sup>208</sup>Pb target 5° scattered electrons Q<sup>2</sup> =0.0088 GeV<sup>2</sup>/c<sup>2</sup> new thin quartz detectors

### $A_{PV} = 0.656 \ ppm \pm 0.060(stat) \pm 0.014(syst)$ PRL 108 (2012) 112502 **PREX at JLab**

1 GeV electron beam, 50-70 μA high polarization, ~89% helicity reversal at 120 Hz





0.5 mm isotopically pure <sup>208</sup>Pb target 5° scattered electrons Q<sup>2</sup> =0.0088 GeV<sup>2</sup>/c<sup>2</sup> new thin quartz detectors

## $A_{PV} = 0.656 \ ppm \pm 0.060(stat) \pm 0.014(syst)$ PRL 108 (2012) 112502 **PREX at JLab**

1 GeV electron beam, 50-70 µA

high polarization, ~89%

helicity reversal at 120 Hz

 $\bar{q} = 0.475 \text{ fm}^{-1}$ 

 $F_W(\bar{q}) = 0.204 \pm 0.028(\exp) \pm 0.001(\text{model}) \text{ fm}$ 

 $R_n = 5.751 \pm 0.175(\text{exp})$  $\pm 0.026(\text{model})$  $\pm 0.005(\text{strange}) \text{ fm}$ 



0.5 mm isotopically pure <sup>208</sup>Pb target 5° scattered electrons Q<sup>2</sup> =0.0088 GeV<sup>2</sup>/c<sup>2</sup> new thin quartz detectors

Establishes the existence of a neutron skin in a heavy nucleus at 95% C.L. with an electroweak probe

### Presented to JLab PAC in 2011 and 2013: Approved with strong endorsement **PREX-II and CREX**

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$$\frac{\delta(A_{PV})/A_{PV} \sim 3\%}{\delta(R_n)/R_n \sim 1\%}$$

$$\delta(R_n) \sim \pm 0.06 \, fm$$

Full precision in 25 additional PAC days

#### PREX-II likely to run in late 2016



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# Precision Electroweak Physics







# Weak Mixing Angle at 1-Loop

For electroweak interactions, 3 input parameters needed:

- 1. Rb-87 mass + Ry constant
- 2. The muon lifetime
- 3. The Z line shape

 $\alpha_{QED} \ G_F \ M_Z$ 

Muon decay Z production

**Weak Neutral Current interactions** 



4th and 5th best

measured parameters:

M<sub>W</sub> and  $sin^2\theta_W$ 







Parity-Violating Electron Scattering

## **PV Electron-Electron Scattering**

 $\underbrace{e^{-}}_{W_{e^{-}}} \xrightarrow{e^{-}}_{e^{-}} \underbrace{e^{-}}_{e^{-}} \underbrace{e^{-}} \underbrace{e^{-}}_{e^{-}} \underbrace{e^{-}}_{e^{-}} \underbrace{e^{-}}$ 

## **PV Electron-Electron Scattering**



electron target:

 $\mathbf{Q}_{\mathbf{W}} = \mathbf{1} - \mathbf{4}\sin^2\theta_{\mathbf{W}}$ 

 $\frac{\delta(Q_W)}{Q_W} \sim 10\% \Longrightarrow \frac{\delta(\sin^2 \theta_W)}{\sin^2 \theta_W} \sim 0.5\%$ 









### $A_{PV}$ = (-131 ± 14 ± 10) x 10-9 Tree-level prediction: ~ 270 ppb Low Q<sup>2</sup>: 3 Best Measurements







Parity-Violating Electron Scatterin

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# **Unique Spectrometer Concept**




# **MOLLER Status**

Director's Review chaired by C. Prescott: strong, positive endorsement

### **Technical Challenges**

- ~ 150 GHz scattered electron rate
  - Design to flip Pockels cell ~ 2 kHz
  - 80 ppm pulse-to-pulse statistical fluctuations

#### 1 nm control of beam centroid on target

- Improved methods of "slow helicity reversal"
- > 10 gm/cm<sup>2</sup> liquid hydrogen target
  - 1.5 m: ~ 5 kW @ 85 μA
- Full Azimuthal acceptance with  $\theta_{\text{lab}}$  ~ 5 mrad
  - novel two-toroid spectrometer
  - radiation hard, highly segmented integrating detectors

#### Robust and Redundant 04% beam polarimetry

 Pursue both Compton and Atomic Hydrogen techniques

#### • MOLLER Collaboration

- ~ 100 authors, ~ 30 institutions
- Expertise from SAMPLE, A4, HAPPEX, G0, PREX, Qweak, E158
- 4th generation JLab parity experiment



- 20M\$ proposal to DoE NP
- 2-3 years construction
- 2-3 years running
- Science review scheduled



 $\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\overline{e}\gamma^{\mu}\gamma_5 e(C_{1u}\overline{u}\gamma_{\mu}u + C_{1d}\overline{d}\gamma_{\mu}d)]$  $+\overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_{5}u+C_{2d}\overline{d}\gamma_{\mu}\gamma_{5}d)]$ 



$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\overline{e} \gamma^{\mu} \gamma_5 e(C_{1u} \overline{u} \gamma_{\mu} u + C_{1d} \overline{d} \gamma_{\mu} d) + \overline{e} \gamma^{\mu} e(C_{2u} \overline{u} \gamma_{\mu} \gamma_5 u + C_{2d} \overline{d} \gamma_{\mu} \gamma_5 d)]$$

**A**<sub>PV</sub> in elastic e-p scattering:



For a <sup>1</sup>H target, nucleon structure contribution well-constrained from measurements

$$\rightarrow 0) = -\frac{G_F}{4\pi\alpha\sqrt{2}} \left[ Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$

$$Q_{weak}^{p} = 2C_{1u} + C_{1d} \propto 1 - 4\sin^{2}\vartheta_{W}$$

Parity-Violating Electron Scattering



$$\mathcal{L}^{PV} = \frac{G_F}{\sqrt{2}} [\overline{e} \gamma^{\mu} \gamma_5 e(C_{1u} \overline{u} \gamma_{\mu} u + C_{1d} \overline{d} \gamma_{\mu} d) + \overline{e} \gamma^{\mu} e(C_{2u} \overline{u} \gamma_{\mu} \gamma_5 u + C_{2d} \overline{d} \gamma_{\mu} \gamma_5 d)]$$

**A**<sub>PV</sub> in elastic e-p scattering:

e(p,s)

Ν

P(k,s)

**Z**\*

e¬(p',s',

For a <sup>1</sup>H target, nucleon structure contribution well-constrained from measurements

$$A(Q^2 \rightarrow 0) = -\frac{G_F}{4\pi\alpha\sqrt{2}} \left[ Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$

$$Q_{weak}^{p} = 2C_{1u} + C_{1d} \propto 1 - 4\sin^{2}\vartheta_{W}$$

$$\begin{array}{rcl} C_{1u} &=& -\frac{1}{2} + \frac{4}{3} \, \sin^2 \theta_W &\approx & -0.19 \\ C_{1d} &=& \frac{1}{2} - \frac{2}{3} \, \sin^2 \theta_W &\approx & 0.35 \\ C_{2u} &=& -\frac{1}{2} + 2 \, \sin^2 \theta_W &\approx & -0.04 \\ C_{2d} &=& \frac{1}{2} - 2 \, \sin^2 \theta_W &\approx & 0.04 \end{array}$$



 $\mathcal{L}_{f_1f_2} =$ 

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_{\mu} f_{1i} \bar{f}_{2j} \gamma_{\mu} f_{2j}$ 





**A**<sub>PV</sub> in elastic e-p scattering:

 $e \neg (p,s)$ 

Ν

P(k,s)

**Z**\*

For a <sup>1</sup>H target, nucleon structure contribution well-constrained from measurements

$$A(Q^2 \rightarrow 0) = -\frac{G_F}{4\pi\alpha\sqrt{2}} \left[ Q^2 Q_{weak}^p + Q^4 B(Q^2) \right]$$

$$Q_{weak}^{p} = 2C_{1u} + C_{1d} \propto 1 - 4\sin^{2}\vartheta_{W}$$



 $C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$ 

new physics  $f_1 
eq f_1$ 

 $\mathcal{L}_{f_1f_2} =$ 

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_\mu f_{1i} \bar{f}_{2j} \gamma_\mu f_{2j}$ 

PV elastic e-p scattering, Atomic parity violation

Parity-Violating Electron Scattering



$$\mathcal{C}^{PV} = \frac{G_F}{\sqrt{2}} [\overline{e}\gamma^{\mu}\gamma_5 e(C_{1u}\overline{u}\gamma_{\mu}u + C_{1d}\overline{d}\gamma_{\mu}d) + \overline{e}\gamma^{\mu}e(C_{2u}\overline{u}\gamma_{\mu}\gamma_5u + C_{2d}\overline{d}\gamma_{\mu}\gamma_5d)]$$

A<sub>PV</sub> in elastic e-p scattering:

 $C_{1u} = -\frac{1}{2} + \frac{4}{3} \sin^2 \theta_W \approx -0.19$ 

 $C_{2d} = \frac{1}{2} - 2 \sin^2 \theta_W \approx 0.04$ 

 $A(Q^2 \rightarrow 0)$ 

 $C_{1q} \propto (g_{RR}^{eq})^2 + (g_{RL}^{eq})^2 - (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \Longrightarrow$ 

 $e \neg (p,s)$ 

Ν

P(k,s)

**Z**\*

For a <sup>1</sup>H target, nucleon structure contribution well-constrained from measurements

$$= -\frac{G_F}{4\pi\alpha\sqrt{2}} \left[ Q^2 Q^p_{weak} + Q^4 B(Q^2) \right]$$

 $Q_{weak}^{p} = 2C_{1u} + C_{1d} \propto 1 - 4\sin^{2}\vartheta_{W}^{Q}$  Qweak at JLab: talk by P. King

new physics  $\begin{array}{rcl} C_{1d} &=& \frac{1}{2} - \frac{2}{3} \sin^2 \theta_W &\approx & 0.35 \\ C_{2u} &=& -\frac{1}{2} + 2 \sin^2 \theta_W &\approx & -0.04 \end{array}$  $f_1 \rightarrow f_2$ 

 $\mathcal{L}_{f_1 f_2} =$ 

 $\sum_{i,j=L,R} \frac{(g_{ij}^{12})^2}{\Lambda_{ij}^2} \bar{f}_{1i} \gamma_{\mu} f_{1i} \bar{f}_{2j} \gamma_{\mu} f_{2j}$ 

#### PV elastic e-p scattering, Atomic parity violation

Parity-Violating Electron Scattering

### Weak Charge and Neutron Skin at Mainz

#### Future: MESA/P2 at Mainz

New ERL complex will also support a highcurrent extracted beam suitable for a PV measurement of proton weak charge

A<sub>PV</sub> = -20 ppb to 2.1% (0.4ppb)
 δ(sin<sup>2</sup>θ<sub>W</sub>) = 0.2%

- Funding approved from DFG
- Development starting now
- Planned running 2017-2020



### Weak Charge and Neutron Skin at Mainz



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Parity-Violating Electron Scattering

Davoudiasl, Lee, Marciano



**Dark Photons:** Beyond kinetic mixing; introduce mass mixing with the Z<sup>0</sup>

 $\epsilon_Z = \frac{m_{Z_d}}{M_Z} \delta$ 

 $1 \times 10^{-4}$ 5×10 arXiv:1203.2947v2  $a_{\mu}$ KLOE a<sub>e</sub>  $1 \times 10^{-5}$ plaineMoller  $5 \times 10^{-6}$ 87 E774 MAMI SAAPE APV Combined  $1 \times 10^{-6}$  $5 \times 10^{-7}$ [For  $\delta^2 = 10^{-6}$  $1 \times 10^{-7}$ 5 10 50 100 500 1000  $m_{\rm Zd}$  [MeV]

**Dark Photons:** Beyond kinetic mixing; introduce mass mixing with the Z<sup>0</sup>

$$\epsilon_Z = \frac{m_{Z_d}}{M_Z} \delta$$

Davoudiasl, Lee, Marciano

 Potentially Observable Effects (for δ≥10<sup>-3</sup>) APV & Polarized Electron Scattering at low <Q> BR(K→πZ<sub>d</sub>)≈ 4x10<sup>-4</sup>δ<sup>2</sup> BR(B→KZ<sub>d</sub>)≈0.1δ<sup>2</sup>

#### δ<sup>2</sup> roughly probed to10<sup>-6</sup>

<u>Davoudiasl, Lee, Marciano</u>



 $K \rightarrow \pi Z_d \rightarrow \pi + "missing energy"$ ε and δ effects could partially cancel! Dark Photons:Beyond kinetic mixing; $\epsilon_Z = \frac{m_{Z_d}}{M_Z} \delta$ introduce mass mixing $\omega$ with the  $Z^0$ • Potentially Observable Effects (for  $\delta \geq 10^{-3}$ )

APV & Polarized Electron Scattering at low <Q> BR(K $\rightarrow \pi Z_d$ )  $\approx 4x10^{-4}\delta^2$  BR(B $\rightarrow KZ_d$ )  $\approx 0.1\delta^2$ 

#### δ<sup>2</sup> roughly probed to10<sup>-6</sup>



**Dark Photons:** 

with the  $Z^0$ 

Davoudiasl, Lee, Marciano

 $\epsilon_Z = \frac{m_{Z_d}}{M_Z} \delta$ 



• Potentially Observable Effects (for  $\delta \ge 10^{-3}$ ) APV & Polarized Electron Scattering at low <Q> BR(K→ $\pi$ Z<sub>d</sub>) ≈ 4x10<sup>-4</sup>δ<sup>2</sup> BR(B→KZ<sub>d</sub>) ≈0.1δ<sup>2</sup> δ<sup>2</sup> roughly probed to10<sup>-6</sup> 0.242  $\nu$ -DIS 0.240



Parity-Violating Electron Scattering

$$\begin{array}{c} \textbf{PUPUIS at JLab}\\ \textbf{Deuterium Target}\\ \overrightarrow{P} = \frac{G_F Q^2}{2\sqrt{2}\pi \alpha} \Big[ g_A \frac{F_1^{*Z}}{F_1^{*}} + g_V \frac{f(y)}{2} \frac{F_3^{*Z}}{F_1^{*}} \Big]\\ \textbf{Q}^* \gg 1 \ GeV^2, W^2 \gg 4 \ GeV^2 \end{array}$$



$$C_{2q} \propto (g_{RR}^{eq})^2 - (g_{RL}^{eq})^2 + (g_{LR}^{eq})^2 - (g_{LL}^{eq})^2 \implies PV \text{ deep inelastic scattering} \\ PVDIS at JLab 6 GeV: \\ talk by V. Sulkosky \\ Deuterium Target$$

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$$PV = \frac{G_F Q^2}{2\sqrt{2\pi\alpha}} \left[g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_1^{\gamma Z}}{F_1^{\gamma}}\right] \\ Q^{2 > 1 GeV^2, W^2 > 4 GeV^2} \\ A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \left[a(x) + f(y)b(x)\right] \\ b(x): function of \\ C_{2i}'s$$









## **SOLID Sensitivity**





## Summary

- Parity-Violating Electron Scattering
  - ★ Technical progress has enabled unprecedented precision
  - ★ flagship experiments at electron accelerators
  - ★ Nuclear/Nucleon Physics
    - Neutron RMS radii of heavy nuclei: JLab (PREX, CREX) & MESA
    - valence quark structure of protons and neutrons: SOLID
  - ★ Electroweak Physics
    - Qweak, 6 GeV PVDIS, MOLLER, SOLID (JLab) and P2 (MESA)
    - Search for new dynamics at the multi-TeV scale and the 100 MeV scale
    - precision measurements of the weak mixing angle at various  $Q^2$  values

#### A rich experimental program over the next 5 to 10 years at Jefferson Laboratory and the new Mainz MESA facility





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- Precision beam position monitoring
- Active calibration of detector slopes



• Active calibration of detector slopes



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Parity-Violating Electron Scattering





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# **MOLLER Design**

3-D Layout of all 14 coils



Acceptance collimator

Vacuum Tank

Concept

**Field Profile** 

Full Azimuthal Acceptance
 Warm copper coils

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Water cooling

20 10

Optics are being fine-tuned

-130 -120 -110 100

- Reduce backgrounds
- Optimize asymmetry
- Symmetric forward/backward

**Detector plane distribution** 

Moller and ep electrons (GHz/cm<sup>2</sup>)

0.09

0.08

0.07

0.06

0.05

0.04

0.03

0.02

0.01

- **Collimator optimization**
- Position sensitivity study
- Engineering work
  - Native CAD model
  - Water-cooling
  - Support structure
  - Force calculations 3



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## **SOLID Sensitivity**

 $\chi$ 

 $\mathcal{N}Z, \gamma$ 



Does Supersymmetry provide a candidate for dark matter?

•B and/or L need not be conserved: neutralino decay

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## **SOLID Sensitivity**



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#### Leptophobic Z'

 $\sim Z, \gamma$ 

Virtually all GUT models predict new Z's <u>arXiv:1203.1102v1</u>
 LHC reach ~ 5 TeV, but.... Buckley and Ramsey-Musolf
 Little sensitivity if Z' doesnt couple to leptons
 Leptophobic Z' as light as 120 GeV could have escaped detection

Since electron vertex must be vector, the Z' cannot couple to the  $C_{1q}$ 's if there is no electron coupling: can only affect  $C_{2q}$ 's

### **The Three Best Measurements**

#### Atomic Parity Violation $\sin^2 \theta_W(m_Z)_{\overline{\rm MS}} = 0.2283(20)$ ★ The 6S - 7S transition in <sup>133</sup>Cs atom $\langle Q \rangle \simeq 2.4 \,\,\mathrm{MeV}$ Neutrino Deep Inelastic Scattering $\sin^2 \theta_W(m_Z)_{\overline{MS}} = 0.2356(16)$ ★ The NuTeV Experiment $\langle Q \rangle \simeq 5 \text{ GeV}$ Parity-Violating Møller Scattering $\sin^2 \theta_W(m_Z)_{\overline{MS}} = 0.2329(13)$ ★ The E158 Experiment $\langle Q \rangle \simeq 160 \text{ MeV}$
**Fundamental Symmetries & Neutrinos** (also HEP Intensity Frontier) Compelling arguments for "New Dynamics" in the Early Universe A comprehensive search to understand the origin of matter requires: The Large Hadron Collider, astrophysical observations as well as Lower Energy: Q<sup>2</sup> << M<sub>Z</sub><sup>2</sup>

Nuclear/Atomic systems address several topics; unique & complementary

- Neutrino mass and mixing  $0\nu\beta\beta$  decay,  $\theta_{13}$ ,  $\beta$  decay, long baseline neutrino expts...
- Rare or Forbidden Processes EDMs, charged LFV, 0vββ decay...
- Dark Matter Searches direct detection, dark photon searches...
- Precision Electroweak Measurements: (g-2)μ, charged & neutral current amplitudes

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#### **Experimental Facilities/Initiatives/Programs**

- Neutrons: Lifetime, Asymmetries (LANSCE, NIST, SNS...)
- Underground Detectors: Dark Matter, Double-Beta Decay
- Nuclei: Precision Weak Decays, Atomic Parity Violation, EDMs (MSU, ANL, TAMU, Tabletop...)
- Muons, Kaons, Pions: Lifetime, Branching ratios, Michel parameters, g-2, EDMs (BNL, PSI, TRIUMF, FNAL, J-PARC...)
- Electron Beams: Weak neutral current couplings, precision weak mixing angle, dark photons (JLab, Mainz)

# EW & Hadron Physics Interplay

**MOLLER Inelastic backgrounds** On LH<sub>2</sub>



- ★ Inelastic e-p scattering in diffractive region ( $Q^2 << 1 \text{ GeV}^2$ , W<sup>2</sup> > 2 GeV<sup>2</sup>) pollutes the Møller peak
- Box diagram uncertainties
  - \* Proton weak charge modified; inelastic intermediate states
- Parton dynamics in nucleons and nuclei
  - ★ Higher twist effects
  - \* charge symmetry violation in the nucleon



\* "EMC" style effects: quark pdfs modified in nuclei

## **Charge Symmetry Violation**

Parton-level charge symmetry assumed in deriving <sup>2</sup>H A<sub>PV</sub>

### **Charge Symmetry Violation**

 $\delta u(x) = u^{p}(x) - d^{n}(x)$  $\delta d(x) = d^{p}(x) - u^{n}(x)$ 

u,d quark mass difference
electromagnetic effects



$$R_{CSV} = \frac{\delta A_{PV}(x)}{A_{PV}(x)} = 0.28 \frac{\delta u(x) - \delta d(x)}{u(x) + d(x)}$$

• Direct observation of partonlevel CSV would be very exciting!

• Important implications for high energy collider pdfs

• Could explain significant portion of the NuTeV anomaly

### Elastic Electron-Proton Scattering P2 at Mainz



E <sub>Beam</sub>	200 MeV
Q²/θ <sub>e</sub>	0.0048 GeV <sup>2</sup> /20°
Time/current/target	10000h/150µA/60cm
Aphys	-20.25 ppb
ΔA <sub>tot</sub>	0.34 ppb (1.7 %)
ΔA <sub>stat</sub>	0.25 ppb
ΔA <sub>sys</sub>	0.19 ppb (0.9%)
Polarization	(85 ± 0.5) %
Rate	0.44 10 <sup>12</sup> Hz
$\Delta sin^2 \theta_{W stat}$	2.8 10-4
$\Delta sin^2 \theta_{W tot}$	3.6 10 <sup>-4</sup> (0.15 %)

 Funding approval from DFG
 R&D in progress
 Aim to run from 2017-20 Technically challenging: great synergy with JLab program
 Recent joint beam test of integrating quartz detectors successful

Parity-Violating Electron Scattering



Parity-Violating Electron Scattering

Krishna Kumar, August 28 2014



Parity-Violating Electron Scattering

## **New Physics Sensitivity**

$$\begin{split} Q_W^e &= -0.0435(9)[1+0.25\,T-0.34\,S+0.7\,X(Q^2)+7m_Z^2/m_{Z_\chi}^2] \\ Q_W^p &= 0.0707(9)[1+0.15\,T-0.21\,S+0.43\,X(Q^2)+4.3m_Z^2/m_{Z_\chi}^2] \\ Q_W(^{12}C) &= -5.510(5)[1-0.003\,T+0.016\,S-0.033\,X(Q^2)-m_Z^2/m_{Z_\chi}^2] \\ Q_W(^{133}Cs) &= -73.24(5)[1+0.0\,T+0.011\,S-0.023\,X(Q^2)-0.9m_Z^2/m_{Z_\chi}^2] \\ \text{Kumar, Mantry, Marciano & Souder, arXiv:1302.6263} \end{split}$$

Oblique Corrections (vacuum polarization)

- Contact Interactions
- Heavy Z's
- Light Z's

**X** Parameter (Q dependence of  $\sin^2\theta_W$ )

### **PV Deep Inelastic Scattering** off the simplest isoscalar nucleus and at high Bjorken x



$$A_{PV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[ g_A \frac{F_1^{\gamma Z}}{F_1^{\gamma}} + g_V \frac{f(y)}{2} \frac{F_3^{\gamma Z}}{F_1^{\gamma}} \right]$$

 $Q^{2} \gg 1 \text{ GeV}^{2}, W^{2} \gg 4 \text{ GeV}^{2}$  $A_{PV} = \frac{G_{F}Q^{2}}{\sqrt{2}\pi\alpha} \left[a(x) + f(y)b(x)\right]$ 

$$\begin{split} x &\equiv x_{Bjorken} \\ y &\equiv 1 - E'/E \\ Y &= \frac{1 - (1 - y)^2}{1 + (1 - y)^2 - y^2 \frac{R}{R + 1}} \end{split}$$

 $R(x,Q^2) = \sigma^l / \sigma^r \approx 0.2$ 

$$\begin{split} A_{\rm iso} &= \frac{\sigma^l - \sigma^r}{\sigma^l + \sigma^r} & \text{At high x, } A_{\rm iso} \text{ becomes independent of pdfs, x \& W,} \\ &= \frac{\sigma^l - \sigma^r}{\sigma^l + \sigma^r} & \text{with well-defined SM prediction for } Q^2 \text{ and y} \\ &= -\left(\frac{3G_FQ^2}{\pi\alpha2\sqrt{2}}\right) \frac{2C_{1u} - C_{1d}\left(1 + R_s\right) + Y\left(2C_{2u} - C_{2d}\right)R_v}{5 + R_s} \end{split}$$

$$\begin{array}{lll}
R_s(x) &=& \frac{2S(x)}{U(x) + D(x)} \xrightarrow{\text{Large } x} 0 \\
R_v(x) &=& \frac{u_v(x) + d_v(x)}{U(x) + D(x)} \xrightarrow{\text{Large } x} 1
\end{array}$$

#### **Interplay with QCD**

- Parton distributions (u, d, s, c)
- Charge Symmetry Violation (CSV)
- Higher Twist (HT)
- Nuclear Effects (EMC)

#### Krishna S. Kumar

Low Energy Standard Model Tests with Parity-Violating Electron Scattering

### A Novel "EMC" Effect

Consider PVDIS on a heavy nucleus

Cloet, Bentz, Thomas, arXiv 0901.3559

- Neutron or proton excess in nuclei leads to a isovector-vector mean field (ρ exchange)
- shifts quark distributions: "apparent" CSV
- Isovector EMC effect: explain additional 2/3 of NuTeV anomaly
- new insight into medium modification of quark distributions



# **A Special HT Effect**

The observation of Higher Twist in PV-DIS would be exciting direct evidence for diquarks

following the approach of Bjorken, PRD 18, 3239 (78), Wolfenstein, NPB146, 477 (78)

Isospin decomposition before using PDF's

$$A_{PV} = \frac{G_F Q^2}{\sqrt{2\pi\alpha}} \left[ a(x) + f(y)b(x) \right]$$

**Higher-Twist valence quark-quark correlation** 

Zero in quark-parton model

 $\langle VV \rangle - \langle SS \rangle = \langle (V-S)(V+S) \rangle \propto l_{\mu\nu} \int \langle D | \overline{u}(x)\gamma^{\mu}u(x)\overline{d}(0)\gamma^{\nu}d(0) \rangle e^{iq \times t} d^4x$ 



(c) type diagram is the only operator that can contribute to a(x) higher twist: theoretically very interesting!

 $\sigma_L$  contributions cancel

Use v data for small b(x) term.

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