# Measurement of the Higgs boson mass with the ATLAS detector

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### Measurement of the mass of the Higgs boson

- Measurement based on 2011 and 2012 LHC pp collision data corresponding to 4.5 fb<sup>-1</sup> at  $\sqrt{s} = 7$  TeV and 20.3 fb<sup>-1</sup> at  $\sqrt{s} = 8$  TeV.
- Mass obtained from a simultaneous fit to the  $m_{\gamma\gamma}$  and  $m_{4\ell}$  ( $\ell = e, \mu$ ) spectra.



• Original mass measurement (Phys. Lett. B 726 (2013) 88)

 $m_H = \left[ 125.5 \pm 0.2 (\text{stat})^{+0.5}_{-0.6} (\text{sys}) \right] \text{ GeV}$ 

o minated by the systematic uncertainties on the  $\gamma$ , e,  $\mu$  energy scales. • Topic of this talk:  $m_H$  measurement with highly reduced energy scale uncertainties due to improved calibration procedures.

### Inner detector and e.m. calorimeter for $e/\gamma$ detection



### Electron and photon reconstruction



- $E_{e/\gamma}$  = sum of the energies of the calorimeter cells associated to the  $e/\gamma$  corrected for energy loss due to absorption in the passive material and leakage outside the cluster.
- Previous calibration of the energy measurement:
  - 1. Gain of the individual amplifiers determined periodically with test pulses.
  - 2. Simulation and test-beam based corrections.
  - 3. Energy scale correction derived from  $Z \rightarrow e^+e^-$  decays.
- A more advanced calibration strategy has been adopted for the updated Higgs mass measurement (see next slide).

# Refined calibration of the $e/\gamma$ energy scale



- 1. Cell energy calibration with test pulses
- 2. Intercalibration of the different calorimeter layers
  - No muon energy loss before the ECAL.
  - $\Rightarrow$  Intercalibration of layers 1 to 3 with muons from *Z* decays.
  - Relative calibration of the presampler with electrons as a function of the longitudinal shower development in the ECAL.
- 3. Determination of the material in front of the EM calorimeter
  - Measurement of the material between the presampler and the first layer with unconverted photons as a function of the longitudinal shower development.
  - Integral material in front of the presampler is extracted from the difference of electron and unconverted photon longitudinal shower profiles.
- 4. Global calorimeter energy scale adjustment with  $Z \rightarrow e^+e^-$  events

### Energy scale uncertainties

### Checks of the $e/\gamma$ energy scale

- $J/\psi \rightarrow e^+e^-$  probes the electron energy scale at low  $E_{\rm T} \sim 7 \dots 35$  GeV.
- $Z \rightarrow \ell^+ \ell^- \gamma$  probes the photon energy scale for  $E_{\rm T} \sim 30$  GeV.



#### Total energy scale uncertainties



#### Main source of the scale uncertainties

- Non-linearity of the E measurement at cell level:  $\sim 0.1\%.$
- Relative calibration of the different calorimeter layers:  $\sim 0.1$ %.
- Material in front of the calorimeter:  $0.1 \dots 0.3\%$ .

### Muon reconstruction



### Energy/momentum measurements

• Calo- and segment tagged muons:

 $p_{muon} = p_{ID}.$ 

Stand-alone muons:

 $p_{muon} = p_{MS} + E_{loss}.$ 

Combined mouns:

 $p_{muon} = \text{Combination}(p_{ID}, p_{stand-alone}).$ 

⇒ Calibrating the momentum scale only with  $Z \rightarrow \mu^+ \mu^-$  data does not reduce the uncertainty of the energy loss correction.

 $\Rightarrow$  New calibration uses  $Z \rightarrow \mu^+ \mu^-$  and  $J/\psi \rightarrow \mu^+ \mu^-$  decays.



• Use of  $Z \to \mu^+ \mu^-$  and  $J/\psi \to \mu^+ \mu^-$  decays in the calibration procedure reduces the muon momentum scale uncertainty from  $\sim 0.1\%$  to  $\sim 0.01\%$ .

# $H\to\gamma\gamma$ selection and event categories



### Requirements

- 2 isolated high  $p_{\rm T}$  photons:
  - $|\eta| < 2.37, |\eta| \notin [1.37, 1.56]$  $E_{\rm T}^{(1)} > 0.35m_{\gamma\gamma}$  $E_{\rm T}^{(2)} > 0.25m_{\gamma\gamma}$
- Main background: γγ continuum production. Significantly smaller contributions from γj and jj.

To improve the  $m_H$  measurement accuracy 10 event categories are used.

| Both photons uncoverted                               |                                   |                            |                            | ≥1 photon converted        |                            |                                   |   |
|---|-----------------------------------|----------------------------|----------------------------|----------------------------|----------------------------|-----------------------------------|---|
| "central"   | "transition"                      | "rest"                     |                            | "central"                  |                            | "transition"                      | "rest"  |
| ( ŋ <sub>1,2</sub>  <0.75)                            | (1.3< η <sub>1 or 2</sub>  <1.75) |                            |                            | ( η <sub>1,2</sub>  <0.75) |                            | (1.3< η <sub>1 or 2</sub>  <1.75) |   |
| P <sub>T</sub> (γγ)<70 GeV P <sub>T</sub> (γγ)>70 GeV |                                   | P <sub>T</sub> (γγ)<70 GeV | P <sub>τ</sub> (γγ)>70 GeV | P <sub>T</sub> (γγ)<70 GeV | P <sub>T</sub> (γγ)>70 GeV |                                   | $P_{T}(\gamma\gamma)$ <70 GeV $P_{T}(\gamma\gamma)$ >70 GeV |

The categories differ from each other in  $\frac{S}{B}$ ,  $\sigma(m_{\gamma\gamma})$ , systematic uncertainties.

#### Invariant mass resolution

$$m_{\gamma\gamma} = 2E_{\gamma_1}E_{\gamma_2}\left[1 - \cos(\theta(\gamma_1, \gamma_2))\right]$$

Thanks to the longitudinal segmentation of the electromagnetic calorimeter  $z_{\gamma_{1,2}}$  is known with 15 mm precision.

 $\Rightarrow \sigma (m_{\gamma\gamma}) \text{ is dominated by the } \gamma \text{ energy resolution!} \\\Rightarrow \sigma (m_{\gamma\gamma}) \sim 1.7 \text{ GeV at } m_{\gamma\gamma} \sim 125 \text{ GeV.}$ 

Fit function for the simultaneous fit of the  $m_{\gamma\gamma}$  spectra of the 10 categories

- Linear combination of signal and background functions.
- Signal: Crystal Ball with category dependent resolutions.
- **Background**: Category dependent analytic functions.
- Free parameters:
  - Signal strength  $\mu_{\gamma\gamma}$ ;
  - Higgs boson mass  $m_H$ ;
  - Background function parameters;
  - Nuisance parameters for systematic uncertainties.

# Measured Higgs boson mass in the $H \rightarrow \gamma \gamma$ channel



New result

 $m_{H} = [125.98 \pm 0.42 ({\rm stat}) \pm 0.28 ({\rm syst})] ~{\rm GeV}$ 

with  $\mu_{\gamma\gamma} = 1.29 \pm 0.30$ .

### Old result

$$m_{H} = [126.8 \pm 0.2 ({\rm stat}) \pm 0.7 ({\rm syst})] ~~{\rm GeV}$$
 with  $\mu_{\gamma\gamma} = 1.55.$ 

### Comments

- Shift of  $m_H$  to lower value than before due to new calibration.
- Increased statistical error on  $m_H$  compared to before due to reduced signal strength.
- Reduced systematic uncertainty on  $m_H$  compared to before thanks to reduced photon energy scale uncertainty.

# $H \rightarrow ZZ^* \rightarrow 4\ell$ selection and event categories



### Requirements

- 2 pairs of isolated same-flavour, opposite-sign leptons (e or μ) from a common vertex.
- Invariant mass requirement:

 $50 \; \mathrm{GeV} < m_{12} < 106 \; \mathrm{GeV}$ 

- To improve the m<sub>4ℓ</sub> resolution: FSR recovery and Z mass constraint.
- Main background:  $ZZ^*$  continuum production. Small contributions from Z + jand  $t\bar{t}$ .

To improve the  $m_H$  measurement accuracy 4 event categories are used:  $4\mu$ , 4e,  $2\mu 2e$ ,  $2e2\mu$ . The categories differ from each other in  $\frac{S}{B}$ ,  $\sigma(m_{4\ell})$ , systematic uncertainties.

# Mass measurement procedure in the $H \to ZZ^* \to 4\ell$ channel

Multivariate discriminant to distinguish  $ZZ^*$  background from  $H \rightarrow ZZ^*$  signal: "BDT<sub>ZZ</sub> output".



### Cross checks:

Fit function for the simultaneous fit of the BDT $_{ZZ}$  output versus  $m_{4\ell}$  spectra of the 4 categories

- Linear combination of a 2D signal and a 2D background distribution.
- The 2D signal and ZZ distributions are taken from MC samples, the reducible background shape is data-driven.

### • Fit parameters:

- Signal strength  $\mu_{4\ell}$ ;
- Higgs boson mass  $m_H$ ;
- Nuisance parameters for systematic uncertainties.
- ${\scriptstyle \circ}$  Simultaneous fits of the 1D  $m_\ell$  distributions like in the past.
- Simultaneous fits of the 2D distributions taken into account per-event  $m_{4\ell}$  resolution values.

# Measured Higgs boson mass in the $H \to ZZ^* \to 4\ell$ channel



 Individual mass measurements in the 4 categories lead to compatible results.

#### New result

 $m_H = [124.51 \pm 0.52 (\text{stat}) \pm 0.06 (\text{syst})] \text{ GeV}$  with  $\mu_{4\ell} = 1.66^{+0.45}_{-0.38}$ .

Old result

$$m_H = \begin{bmatrix} 124.3^{+0.6}_{-0.5}(\text{stat})^{+0.5}_{-0.3}(\text{syst}) \end{bmatrix} \text{ GeV}$$
 with  $\mu_{4\ell} = 1.43^{+0.40}_{-0.35}$ .

#### Comments

- Shift of  $m_H$  to slightly higher value than before due to new electron energy calibration.
- Increased signal strength compared to before due to increased  $m_H$  value.
- Reduced systematic uncertainty on  $m_H$  compared to before thanks to reduced lepton energy/momentum scale uncertainty.



New result

$$m_H = [125.36 \pm 0.37(\text{stat}) \pm 0.18(\text{syst})]$$
 GeV

$$m_H = \left[ 125.49 \pm 0.24 (\text{stat})^{+0.50}_{-0.58} (\text{syst}) \right] \text{ GeV}$$

#### Comments

- Shift of  $m_H$  to slightly smaller value than before due to the lower  $m_H$  value in the  $H \rightarrow \gamma \gamma$  channel.
- Increased statistical error on  $m_H$  compared to before due to increased statistical uncertainty of  $m_H$  in the  $H \rightarrow \gamma \gamma$  channel.
- Compatibility between the  $m_H$  measurements in the  $H \to \gamma \gamma$  and  $H \to ZZ^* \to 4\ell$  channels is  $2.0\sigma$  compared to  $2.5\sigma$  before.

### BONUS SLIDE



- The  $H \to 4\ell$  data were used to set an upper limit on the natural width  $\Gamma_H$  of the Higgs boson.
- $\Gamma_H$  was added as a free parameter to the 3D fit while fixing  $m_H$  to the measured value.

 $\Gamma_H < 2.6~{\rm GeV}$  at 95% CL

Expected limits:

- $\Gamma_H < 6.2$  GeV for  $\mu_{4\ell} = 1$ .
- $\Gamma_H < 3.5 \text{ GeV}$  for  $\mu_{4\ell} = 1.66$ .