

NON-RELATIVISTIC MAJORANA NEUTRINOS IN A THERMAL BATH AND LEPTOGENESIS

Simone Biondini

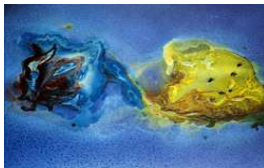
in collaboration with N. Brambilla, M. A. Escobedo and A. Vairo

T30f - Technische Universität München

PANIC 2014 - Hamburg University



Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)



- 1 MOTIVATION AND INTRODUCTION
- 2 NEUTRINO THERMAL WIDTH
- 3 CP ASYMMETRY AT FINITE TEMPERATURE
- 4 CONCLUSIONS

BARYON ASYMMETRY AND LEPTOGENESIS

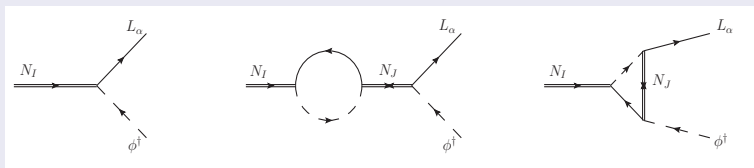
BARYON ASYMMETRY IN THE UNIVERSE

- our universe is strongly matter-antimatter asymmetric

$$\eta_B = \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.20 \pm 0.15) \times 10^{-10}$$

[WMAP collaboration (2011)]

- within SM: $\eta_B \sim 10^{-18}$
- Baryogenesis via Leptogenesis: $\Delta L \rightarrow \Delta B$ [M. Fukugida and T. Yanagida (1986)]



- Three Sakharov conditions: $\Delta L \neq 0$, C and CP violation, out-of-equilibrium
- Massive Majorana neutrinos play the relevant role

STRONG WASH-OUT

DECAY PARAMETER FOR THE WASH-OUT

$$K = \frac{\Gamma_1^{(T=0)}}{H(T = M_1)} = \frac{M_1 (F^\dagger F)_{11}}{8\pi 1.66 \sqrt{g^*} \frac{M_1^2}{M_{Pl}}} = \frac{\tilde{m}_1}{m_*}$$

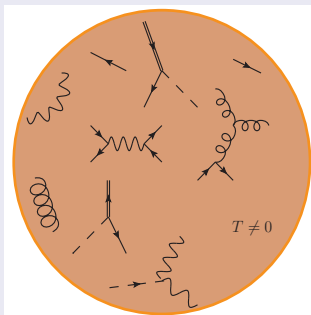
where we may define

- $\tilde{m} = (F^\dagger F)_{11} \frac{v^2}{M_1}$, effective neutrino mass
- $m_* = 8\pi 1.66 \sqrt{g^*} \frac{v^2}{M_{Pl}} \simeq 1.1 \times 10^{-3} \text{ eV}$ [E. Nardi, Y. Nir and S. Davidson]
- Because $\tilde{m} \simeq m_{\text{sol}} = \sqrt{\Delta m_{\text{sol}}^2} \simeq 9 \times 10^{-3} \Rightarrow K > 1$

- 1) Majorana neutrinos is coupled until $T \ll M_1$
- 2) The final lepton asymmetry is produced in a non-relativistic regime

NON-RELATIVISTIC NEUTRINOS IN MEDIUM

INTERACTION AT $T \neq 0$



1) Thermal production rate

A. Salvio, P. Lodone and A. Strumia (2011)

M. Laine and Y. Schroder (2012)

$$\Gamma(T=0) \rightarrow \Gamma(T)$$

1) CP asymmetry at finite T

M. Garry, A. Hohenegger and A. Kartavtsev

(2010)

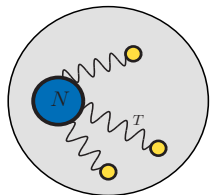
$$\epsilon = \frac{\Gamma_{\ell} - \Gamma_{\bar{\ell}}}{\Gamma_{\ell} + \Gamma_{\bar{\ell}}} \rightarrow \epsilon(T)$$

CALCULATIONS AT FINITE TEMPERATURE ARE DIFFICULT AND TRICKY

However...

- the hierarchy of scales $M \gg T$ has been not fully exploited
- we want to provide an effective field theory (EFT) approach

HEAVY MAJORANA NEUTRINOS



$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{\psi} (i\not{\partial} - M) \psi - F_f \bar{L}_f \tilde{\phi} P_R \psi - F_f \bar{\psi} P_L \tilde{\phi}^\dagger L_f$$

- $p^\mu = Mv^\mu + k^\mu$ with $|k| \sim T \ll M$
- What are the low-energy degrees of freedom? What is the \mathcal{L}_{EFT} ?

- In a reference frame where the heavy particle is at rest up to to $k \ll M$

$$\mathcal{L}_{\text{EFT}} = N^\dagger i\partial_0 N + \sum_i c_n \left(\frac{\mu}{M} \right) \frac{\mathcal{O}(\mu, T)_n}{M^{d_n-4}} + \mathcal{L}_{\text{light}}$$

OBSERVATIONS

- In the heavy particle sector: expansion in $1/M$
- Contribution of higher order operators are counted in power of T/M
- The **Wilson coefficients** may be computed setting $T=0$, in vacuum

EFT AT WORK: MATCHING

MATCHING COMPUTATION AT $T=0$

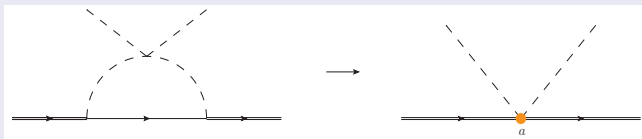
- 1) The low-energy Lagrangian is organized as follows:

$$\mathcal{L}_{\text{EFT}} = N^\dagger \left(i\partial_0 - i\frac{\Gamma_0}{2} \right) N + \frac{\mathcal{L}^{(1)}}{M} + \frac{\mathcal{L}^{(2)}}{M^2} + \frac{\mathcal{L}^{(3)}}{M^3} + \mathcal{O}\left(\frac{1}{M^4}\right)$$

- 2) By dimensional analysis the thermal corrections scale as

$$\delta\Gamma_{(1)} \propto \frac{T^2}{M}, \quad \delta\Gamma_{(2)} \propto \frac{T^3}{M^2}, \quad \delta\Gamma_{(3)} \propto \frac{T^4}{M^3}$$

- By symmetry, the leading dimension 5 operator is: $\mathcal{L}^{(1)} = a N^\dagger N \phi^\dagger \phi$



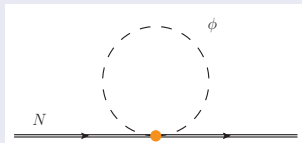
EFT AT WORK: THERMAL WIDTH

- In the $v^\mu = (1, \vec{0})$ frame, the Majorana neutrino propagator has the form

$$\left(\frac{1+\gamma_0}{2}\right)^{\alpha\beta} \frac{iZ}{k^0 - E + i\Gamma/2} = \left(\frac{1+\gamma_0}{2}\right)^{\alpha\beta} Z \left[\frac{i}{k^0 + i\epsilon} - (iE + \frac{\Gamma}{2}) \left(\frac{i}{k^0 + i\epsilon}\right)^2 + \dots \right]$$

THERMAL WIDTH IN THE EFT

- The self-energy diagram we can build is a tadpole in the EFT



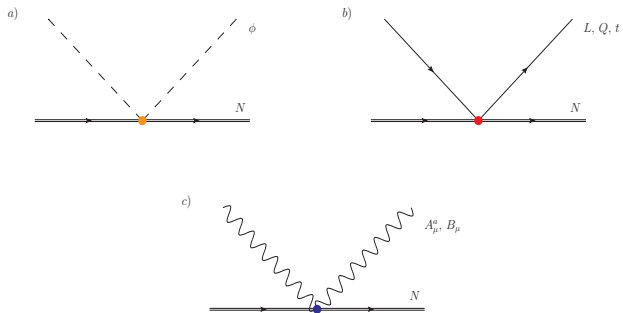
$$\Gamma_\phi = 2 \frac{\text{Im } a}{M} \langle \phi^\dagger(0) \phi(0) \rangle_T = -\frac{\lambda |F|^2 M}{8\pi} \left(\frac{T}{M}\right)^2$$

- Higgs propagator in the Real Time Formalism (RTF)

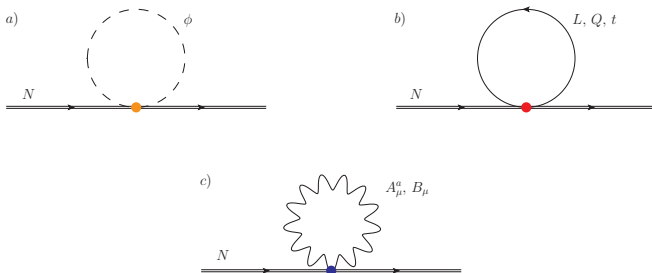
$$i\Delta_{11}(x-y) = \int \frac{d^4k}{(2\pi)^4} \left[\frac{i}{k^2 + i\epsilon} + (2\pi) n_B(|k_0|) \delta(k^2) \right] e^{-ik \cdot (x-y)}$$

SUB-LEADING OPERATORS

- Higher order operators in the effective Lagrangian
- Higgs, fermions (quarks and leptons) and gauge bosons effective vertices



- Thermal tadpoles induced by the sub-leading operators



EFT FOR MAJORANA FERMIONS WORKS:

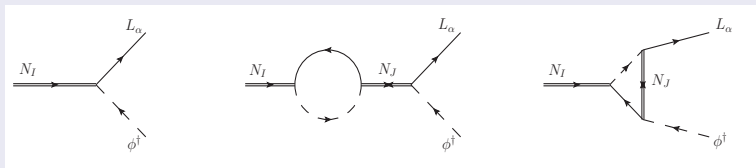
$$\Gamma_T = \frac{|F|^2 M}{8\pi} \left\{ -\lambda \left(\frac{T}{M}\right)^2 - \frac{\pi^2}{80} \left(\frac{T}{M}\right)^4 (3g^2 + g'^2) - \frac{7\pi^2}{60} \left(\frac{T}{M}\right)^4 |\lambda_t|^2 + \mathcal{O}\left(\frac{T}{M}\right)^6 \right\}$$

M. Laine and Y. Schroeder (2012)

CP ASYMMETRY

- Thermal corrections to the CP asymmetry defined as follows

$$\epsilon = \sum_{l,f} \frac{\Gamma(N_I \rightarrow l_f) - \Gamma(N_I \rightarrow \bar{l}_f)}{\Gamma(N_I \rightarrow l_f) + \Gamma(N_I \rightarrow \bar{l}_f)} \Rightarrow \boxed{\eta_B \sim \epsilon}$$



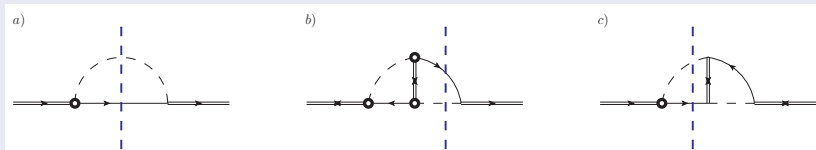
- We consider two Majorana neutrinos N_1 and N_2 , $\epsilon \sim \text{Im} [(F_1 F_2^*)^2]$

$$M_1 \ll M_2$$

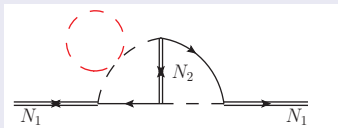
- We focus on the direct CP asymmetry (vertex diagram)

AT $T=0$, IN VACUUM ALREADY TWO LOOPS

- Diagrams relevant for the direct asymmetry: $\epsilon_0 = -\frac{\text{Im}[(F_1 F_2^*)^2]}{16\pi |F_1|^2} \left(\frac{M_1}{M_2}\right) + \dots$



IN MEDIUM THERMAL PARTICLES ENTER THE GAME

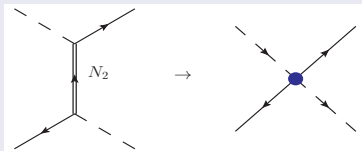


- three loops at finite temperature
→ **very hard task**
- there are hierarchies of scales!

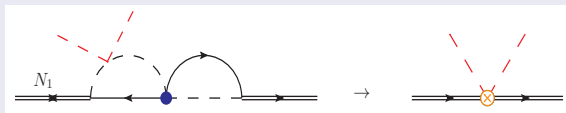
$$M_2 \gg M_1 \gg T$$

WE CAN USE THE EFT APPROACH

- 1) Integrate out the energy modes of order
- M_2



- 2) Integrate out the energy modes of order
- M_1



RESULT:

$$\epsilon_T = -\frac{\text{Im} [(F_1 F_2^*)^2]}{16\pi |F_1|^2} \frac{3\lambda}{2} \left(\frac{M_1}{M_2}\right) \left(\frac{T}{M_1}\right)^2 + \dots$$

CONCLUSIONS

- Non-relativistic regime may be relevant for leptogenesis
- Interactions occur in medium

- Formulation of an EFT to deal with Majorana neutrinos in medium
- Test the EFT: neutrino thermal width up to $\mathcal{O}\left(\frac{T}{M}\right)^4$

2-loops ($T \neq 0$) \rightarrow 1-loop ($T=0$) + 1-loop (thermal tadpoles)

- Address the CP asymmetry in the EFT formalism
- Matching calculation (cutting rules at $T = 0$) and thermal tadpoles ($T \neq 0$)
- In the non-relativistic regime thermal corrections are under control
- To do list: self-energy diagram, resonant case $M_1 \simeq M_2$