

# Summer Student Presentation.

Simulation studies for the total decay width of the Higgs Boson at  
the International Linear Collider

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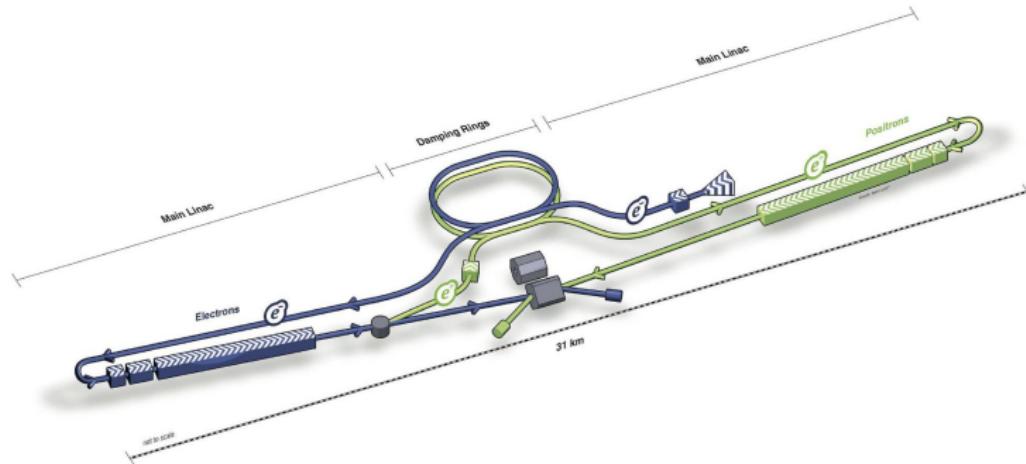
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**2** Theory

**3** Determination of  $\Gamma_H^{tot}$  Precision

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# International Linear Collider.



- The International Linear Collider will be 31 km long.
- Runs at  $\sqrt{s} = 500\text{GeV}$  with possible upgrade to  $1\text{TeV}$
- Two detectors: ILD and SiD
- Designed to make precision measurements
- Produces spin polarised beams



# Introduction.

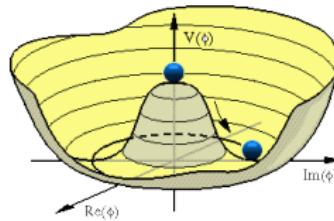
**Aim:** Explore how accurately the  $\Gamma_H^{tot}$  can be measured at the ILC

- Interested in full simulation at  $\sqrt{s} = 500 \text{ GeV}$  and  $\mathcal{L} = 500 \text{ fb}^{-1}$
- Use the WW fusion as signal process  $e^+ e^- \rightarrow H \nu_e \bar{\nu}_e$

Previously a fast simulation was carried out for  $m_H = 120 \text{ GeV}$  and  $\sqrt{s} = 500 \text{ GeV}/350 \text{ GeV}$

*Niels Meyers: Higgs-Bosons At TESLA: Studies On Production In  
WW-Fusion And Total Decay Width*

# Higgs Boson.



The Higgs Mechanism results in *spontaneous symmetry breaking*.  
For real scalar field  $\phi$ ,

$$\mathcal{L} = \frac{1}{2}(\partial_\mu\phi)^2 - \frac{1}{2}\mu^2\phi^2 - \frac{1}{4}\lambda\phi^4$$

when  $\mu^2 < 0$  we get non-zero minima

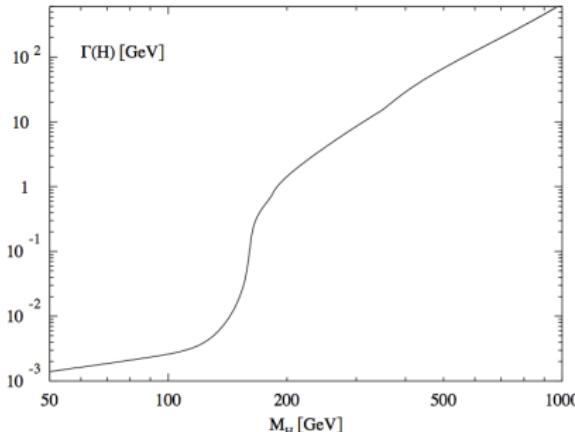
$$\phi = \pm\sqrt{-\frac{\mu^2}{\lambda}}$$

This gives *vacuum expectation value*.

Field  $\phi$  can be written in terms of  $h$ , real scalar field: Higgs Boson



# Decay Width.



- Decay width of Higgs much for small masses is smaller than experiment resolution
  - $\Gamma_H^{tot} \sim 5$  MeV
  - Must be determined indirectly

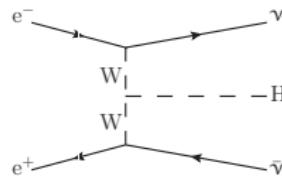


# Higgs Production.

- Dominant method of Higgs formation at linear collider:

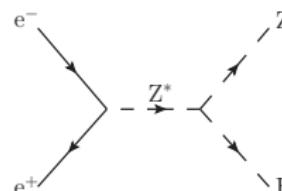
- WW Fusion

$$e^+ e^- \rightarrow \nu_e \bar{\nu}_e H$$



- Higgs-Strahlung

$$e^+ e^- \rightarrow Z^* \rightarrow ZH$$

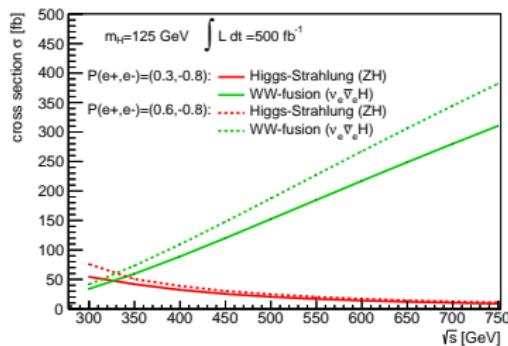


- ZZ fusion also possible but suppressed by weak mixing angle



# Cross Sections.

I have calculated the cross section for WW-fusion events and Higgs-Strahlung using *Whizard* as a function of  $\sqrt{s}$  for two different beam polarisations:



- WW fusion dominates at  $\sqrt{s} = 500$  GeV, Higgs Strahlung dominates at lower energy



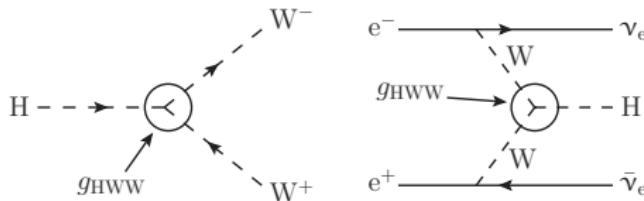
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- Measurement of  $\sigma_{wwfusion}$  allows  $\Gamma_H^{tot}$  to be determined:

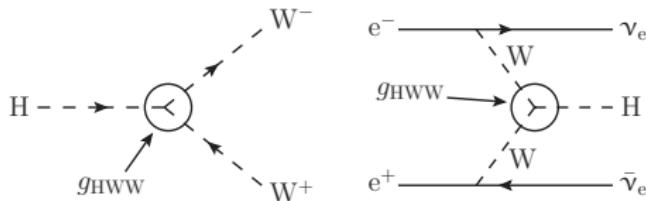


$$\Gamma(H \rightarrow WW) \propto g_{HWW}^2 \quad \sigma_{wwfusion} \propto g_{HWW}^2$$



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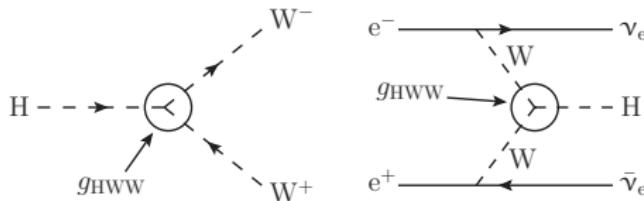
Therefore

$$\Gamma_H^{tot} = \frac{\Gamma(H \rightarrow WW)}{BR(H \rightarrow WW)} \propto \frac{\sigma_{WWfusion}}{BR(H \rightarrow WW)}$$



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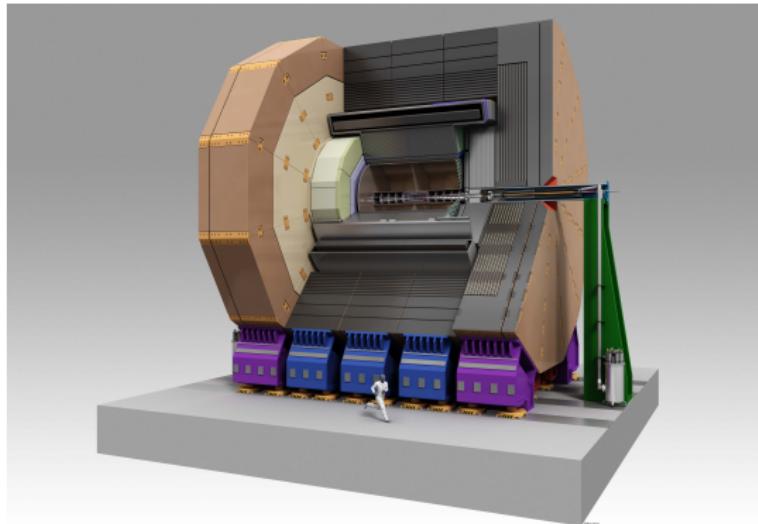
$$\Gamma_H^{tot} = \frac{\Gamma(H \rightarrow WW)}{BR(H \rightarrow WW)} \propto \frac{\sigma_{WWfusion}}{BR(H \rightarrow WW)}$$

- For  $m_H = 125$  GeV the decay  $H \rightarrow b\bar{b}$  is largest

$$\sigma_{WWfusion} = \frac{N'_{WW}}{\epsilon \cdot \mathcal{L}}$$



# International Large Detector (ILD).



## Full simulation of ILD

- $\sqrt{s} = 500 \text{ GeV}$  and  $\mathcal{L} = 500 \text{ fb}^{-1}$
- Simulations generate Monte Carlo data sets
- Takes into account detector design

# Signal and Background Processes.



**Signal:**

$$e^+ e^- \rightarrow H \nu_e \bar{\nu}_e \rightarrow b \bar{b} \nu_e \bar{\nu}_e$$

- Total number of WW fusion events is  $N_{wwfusion} = 7.6 \times 10^4$ .



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**Background:**

Type	Decay	$N_{events}$
Semileptonic	$W^+ W^- \rightarrow \nu_\ell \ell^\pm q \bar{q}$	2,785,120
Hadronic	$W^+ W^- \rightarrow q \bar{q} q \bar{q}$	2,245,500
Semileptonic	$ZZ \rightarrow \ell \bar{\ell} q \bar{q}$	182,999
Hadronic	$ZZ \rightarrow q \bar{q} q \bar{q}$	203,310
Semileptonic	$W^\mp e^\pm \nu_e \rightarrow e^\pm \nu_e q \bar{q}$	2,283,520
Semileptonic	$Z e^+ e^- \rightarrow q \bar{q} e^- e^+$	603,845
Hadronic	$Z \rightarrow q \bar{q}$	9,805,180
Semileptonic	$Z \nu \bar{\nu} \rightarrow q \bar{q} \nu \bar{\nu}$	279,408

- Total no. of background events is  $N_{Background} = 1.8 \times 10^7$ .
- Total no. of Higgs Strahlung events is  $N_{HiggsStrahlung} = 1.0 \times 10^4$ .



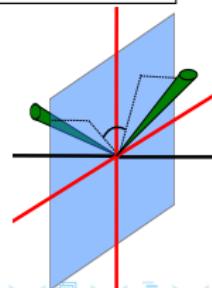
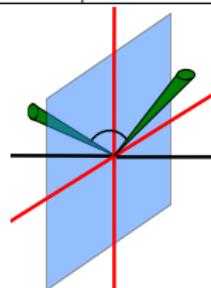
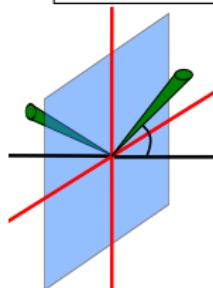
# Event Selection and Cuts.

Name	Cut
Isolated lepton removal	No. isolated lep $\leq 1$
Visible mass	$105 \text{ GeV} \leq m_{vis} \leq 135 \text{ GeV}$
Visible energy	$105 \text{ GeV} \leq E_{vis} \leq 255 \text{ GeV}$
Visible $P_t$	$5 \text{ GeV} \leq \sum P_T \leq 200 \text{ GeV}$
Polar angle of jet	$ \cos \theta_{jet}  \leq 0.9$
Angle between jets	$\cos \alpha \leq 0.2$
Acoplanarity	$\text{Acop} \geq 10^\circ$
Durham, $Y_{12}$ (minus)	$0.2 \leq Y_{12} \leq 0.8$
b-tagging	$\text{b-tag} \geq 0.5$
Number of tracks	$10 \leq N \leq 60$

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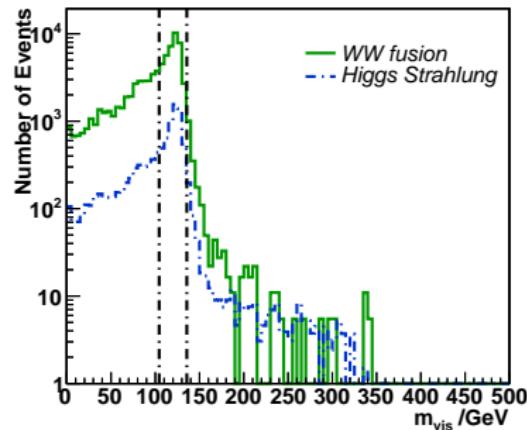
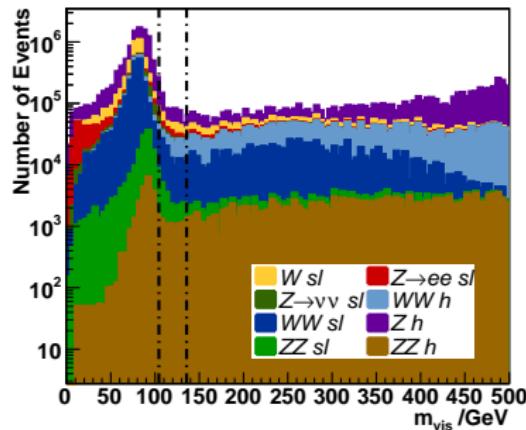


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# Background and Signal Distributions.



- The distribution of the visible mass,  $m_{vis}$ , before any cuts are applied

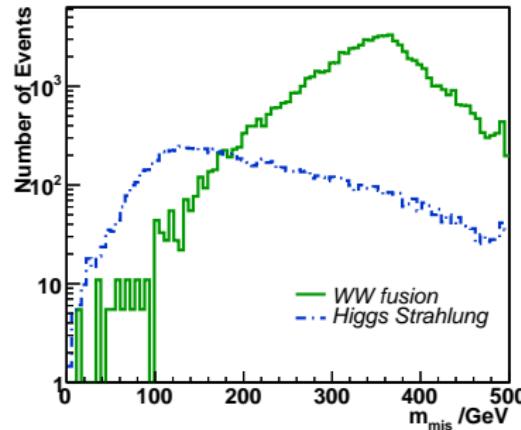
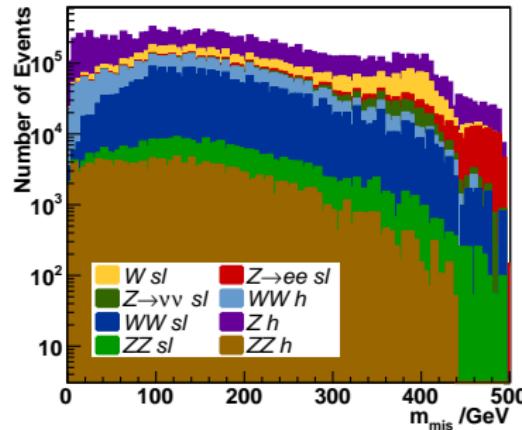
$$m_{vis} = \sqrt{E_{vis}^2 - |\mathbf{p}_{vis}|^2}$$

**Background:**  $m_{vis} \approx m_Z$  or  $m_W$

**WW fusion/Higgs Strahlung:**  $m_{vis} \approx m_H$



# Missing Mass Distribution.



$$m_{mis} = \sqrt{E_{mis}^2 - |\mathbf{p}_{mis}|^2}$$

$$\mathbf{p}_{mis} = -\mathbf{p}_{vis} \text{ and } E_{mis} = \sqrt{s} - E_{vis}$$

We expect different characteristics for the  $\nu\bar{\nu}$  invariant mass:

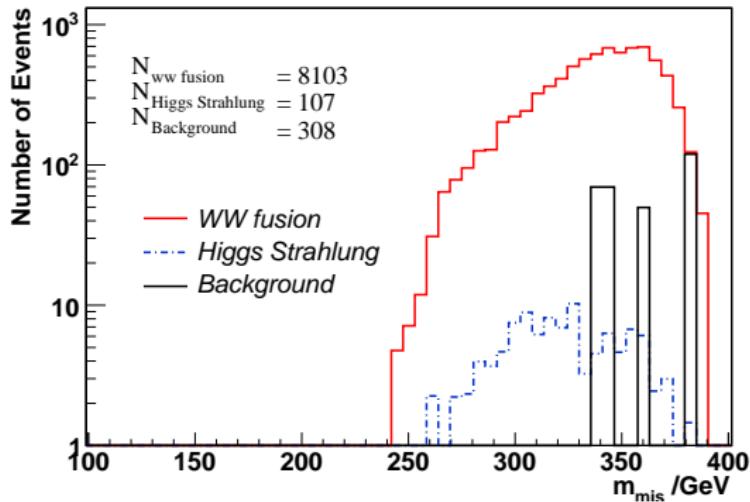
**Higgs Strahlung:**  $m_{mis} \approx m_Z$

**WW fusion:** Higher peak



# Missing Mass After Cuts.

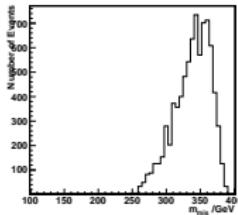
After cuts have been applied:



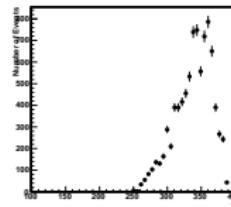
- Background affected by low statistics
- Low  $m_{mis}$  missing due to  $Y_{12}$  cut



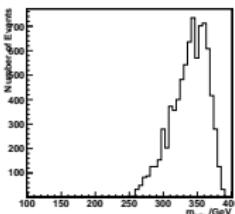
# Missing Mass Fit.



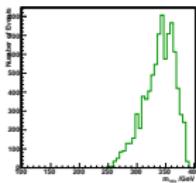
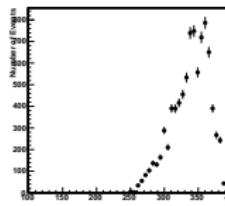
Toy MC →



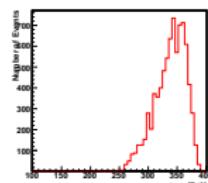
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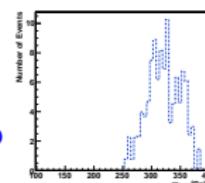
Toy MC →



==



+ B



+

Fit

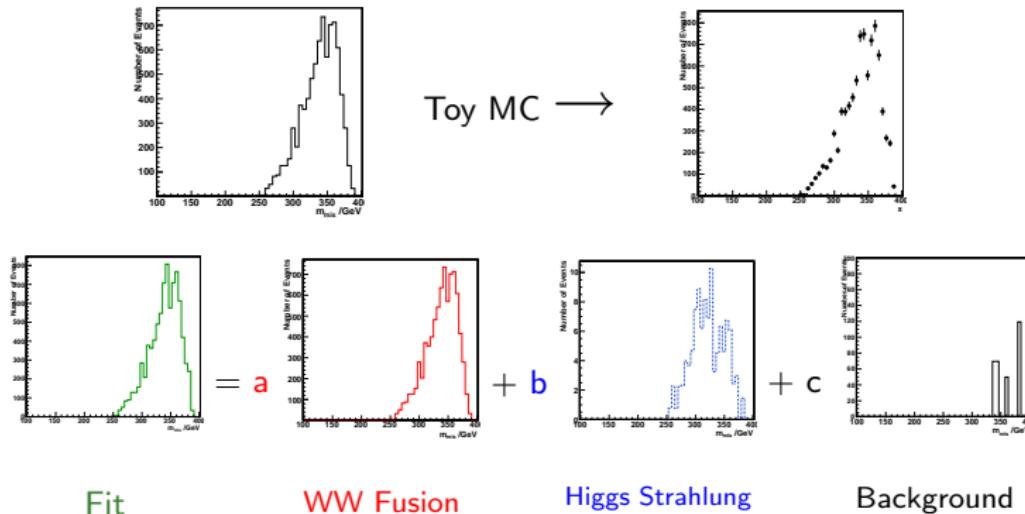
WW Fusion

## Higgs Strahlung

## Background



# Missing Mass Fit.

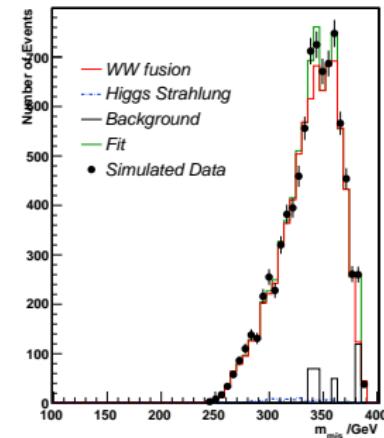
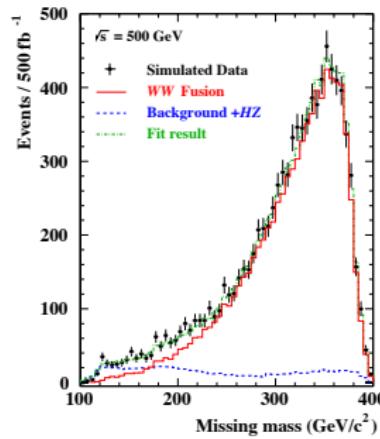
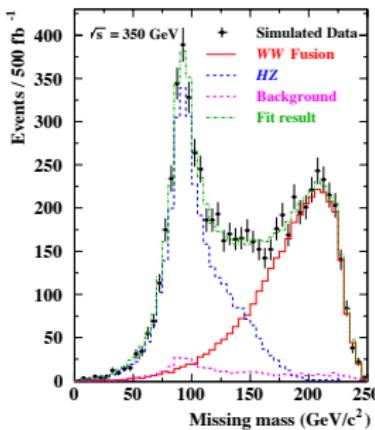


Adjust co-efficients **a** **b** and **c** such that the  $\chi^2$  function is minimised:

$$\chi^2 = \sum_i^{Nbins} \left( \frac{N_{MCdata,i} - N_{pred,i}}{\sigma_{MCdata,i}} \right)^2$$



# Comparison.



$N_{WW\text{Fusion}}$	$3421 \pm 75$	$8181 \pm 98$	$8049 \pm 160$
$N_{\text{HiggsStrahlung}}$	$3192 \pm 83$	$229 \pm 21$	$101 \pm 124$
$N_{\text{Background}}$	$363 \pm 88$	$348 \pm 38$	$340 \pm 39$

\*first two from *Niels Meyers: Higgs-Bosons At TESLA (2000)*



# Obtaining Uncertainty.

The overall uncertainty on the total width can be calculated using propagation of errors

- The fit gives us the value  $\left( \frac{\Delta N'_{WW}}{N'_{WW}} \right) = 1.99\%$

$$\left( \frac{\Delta \sigma_{WWfusion}}{\sigma_{WWfusion}} \right) = \sqrt{\left( \frac{\Delta N'_{WW}}{N'_{WW}} \right)^2 + \left( \frac{\Delta BR(H \rightarrow b\bar{b})}{BR(H \rightarrow b\bar{b})} \right)^2} = 2.48\%$$



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- The uncertainties for the branching ratios come from previous studies
- *Fast simulation found*  $\left( \frac{\Delta N'_{WW}}{N'_{WW}} \right) = 1.2\% \text{ for } m_H = 120 \text{ GeV and}$   
 $\left( \frac{\Delta N'_{WW}}{N'_{WW}} \right) = 2.4\% \text{ for } m_H = 130 \text{ GeV}$



# Summary.

- Conducted full simulation of  $\Gamma_H^{tot}$  at  $\sqrt{s} = 500$  GeV
- Used selection cuts to keep WW fusion events
- Generated toy MC data and fitted data
- Propagated uncertainty using branching ratio uncertainties
- Calculated overall uncertainty to be  $\left( \frac{\Delta \Gamma_H^{tot}}{\Gamma_H^{tot}} \right) = 2.8\%$
- Result consistent with previous analysis