Using the elastic $\gamma^* p \rightarrow J/\psi p$ cross-section to probe the small x gluon density

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Outline

- 1 Introduction
- 2 Theory
- 3 Results
- 4 Summary



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Introduction.



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Motivation

> Deep inelastic scattering data (e.g. from HERA) poorly constrains the gluon density at low scales ($\mu^2 < 6 \text{ GeV}^2$) and small $x (\sim 10^{-4})$



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- > Deep inelastic scattering data (e.g. from HERA) poorly constrains the gluon density at low scales ($\mu^2 < 6 \text{ GeV}^2$) and small $x (\sim 10^{-4})$
- Exclusive processes such as exclusive heavy vector meson production should be a sensitive probe of the gluon in this domain
- > Can such processes constrain the gluon?



Photoproduction of J/ψ



$$> \gamma^* p \rightarrow J/\psi p$$

- Photoproduction of a cc̄ state (Dipole model)
- Colorless two gluon (pomeron) exchange with the proton → Squared dependency of the gluon distribution in the proton
- > J/ψ in the final state



Situation so far



Elastic J/w photoproduction

> So far: Few PDFs able to describe new data, especially NLO not Quite some theoretical interest in this >



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Situation so far



Elastic J/ ψ photoproduction

- > So far: Few PDFs able to describe new data, especially NLO not
- > Quite some theoretical interest in this
- > Two papers published, while we were here

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Theory.



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Leading logarithmic order theory

Differential cross-section

$$\frac{\frac{\mathrm{d}\,\sigma}{\mathrm{d}\,t}(\gamma^* p \to J/\psi p)\Big|_{t=0} = \frac{\Gamma_{ee}M_{J/\psi}^3 \pi^3}{48\alpha} \left[\frac{\alpha_s(\bar{Q}^2)}{\bar{Q}^4} xg(x,\bar{Q}^2)\right]^2 \left(1 + \frac{Q^2}{M_{J/\psi}^2}\right)^2 a$$

^aS.P. Jones, A.D. Martin, T. Teubner et. al.



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Leading logarithmic order theory

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$$\xrightarrow{\ ^3\text{S.P. Jones, A.D. Martin, T. Teubner et. al.}$$

$$ar{Q}^2 = (Q^2 + M_{J/\psi}^2)/4, \quad x = (Q^2 + M_{J/\psi}^2)/(W^2 + Q^2).$$

 Q^2 : virtuality of the photon, $W_{\gamma p}$: center-of-mass energy.



Leading logarithmic order theory

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$$\frac{1}{\mathrm{a}^3\mathrm{S.P. Jones, A.D. Martin, T. Teubner et. al.}}$$

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For t integration assume $\sigma \propto \exp(-Bt)$, with B the slope parameter:

$$B(W) = (b_0 + 4\alpha' \ln(W/90 \text{GeV})) \text{GeV}^{-2}$$



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Aim

calculate cross-section with different LO gluon densities



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Aim

> calculate cross-section with different LO gluon densities

Gluon densities

- S.P. Jones LO (Theory)
- HERAPDF-like LO parametrisation
- > CTEQ6I (LO, except for α_s)



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Results.



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Using theoretical LO prediction



$$g(x,\mu^2) = N x^{-\lambda} \quad \rightarrow \quad \sigma(W_{\gamma p}) = N' W^{-\delta}, \quad \delta = 4\lambda$$

DESY

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Using HERAPDFLO at μ^2 =2.4 GeV²



Using $k = 0.56 \pm 0.01$ for normalization

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Probing small x gluons



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Using HERAPDFLO at $\mu^2=3$ GeV²

Elastic J/ ψ photoproduction



Using $k = 0.91 \pm 0.01$ for normalization





Comparison with CTEQ6I



Using k = 0.41 for normalization



General

- > Factorization scale and therefore normalization not exactly known.
- The W_{γp}-dependency of the cross-section yields the x-dependency of the gluon density, down to low x.



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HERAPDF LO parametrization

- > J/ψ cross-section not a straight line anymore in LO.
- > Good agreement with low $W_{\gamma p}$ solutions from LHCb data and with data from fixed target experiments, when choosing different scale.
- > Also good agreement with high $W_{\gamma p}$ (low x) solutions from LHCb.
- Better agreement than CTEQ6I.



Thank you for your attention.



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Backup.



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Correction factors

Skewed gluon distribution

> It is more likely, that one gluon has a significantly lower imes

> This can be estimated by multiplying with:

$$R_g = rac{2^{2\lambda+3}}{\sqrt{\pi}} rac{\Gamma(\lambda+rac{5}{2})}{\Gamma(\lambda+4)}$$

Real Part

- The gluon PDF correpsonds only to the imaginary part of the amplitude
- > Using: $\frac{\Re A}{\Im A} \approx \frac{\pi}{2}\lambda$
- > And therefore multiplying with:

$$1 + \tan^2\left(\frac{\pi}{2}\lambda\right)$$

with x axis





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Comparison between NLO and LO PDF



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