

Yang-Baxter Equation

from SUSY Gauge Theories

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IAS



KAVLI
IPMU

[IGST 2014, Jul/18, DESY]

Happy 10th anniversary, IGST!



mostly based on

M.Y. 1307.1128

D. Xie + M.Y. 1207.0811

M.Y. 1203.5784

Y. Terashima + M.Y. 1203.579

earlier works by

V. V. Bazhanov - S. M. Sergeev 1006.0651, 1106.5874

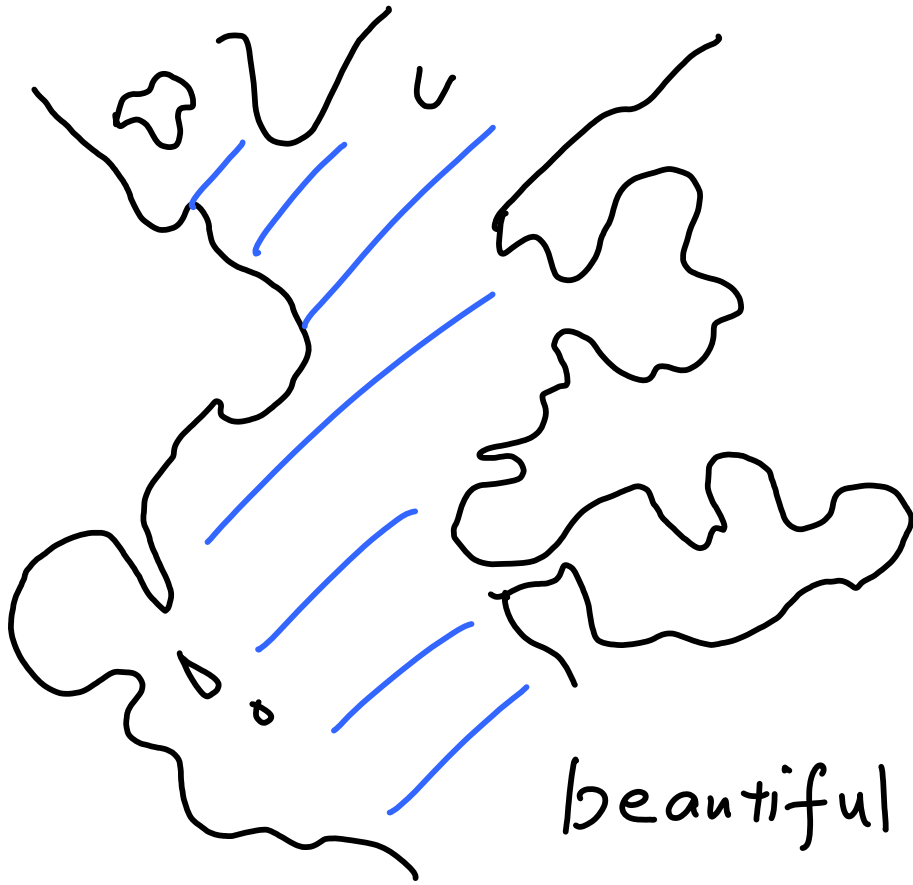
V. P. Spiridonov 1011.3798

yesterday

T. Bargheer, Y. Huang & F. Loebbert + M.Y. 1407.4449

Introduction

ocean of integrability



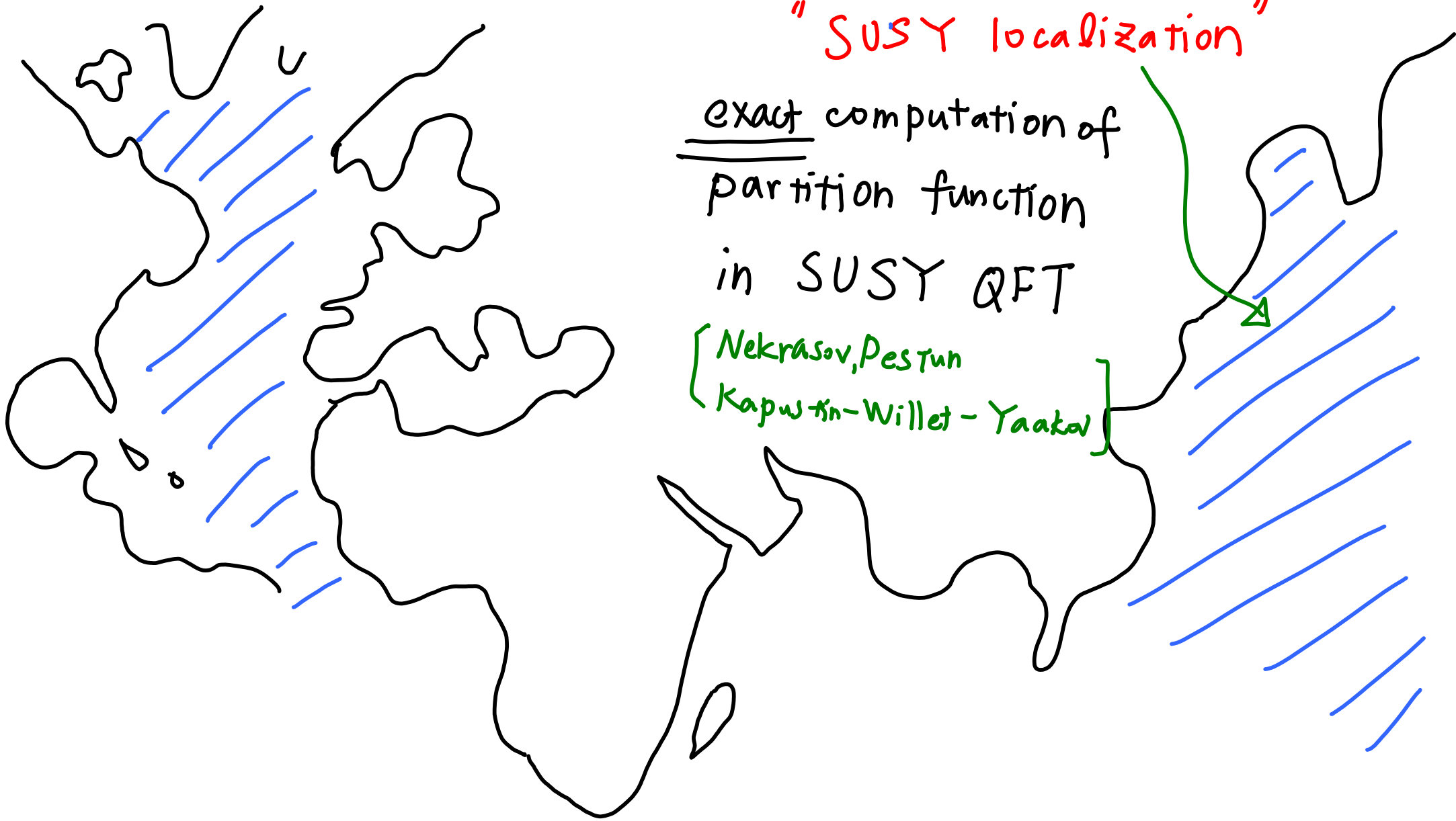
beautiful developments in
understanding 4d $N=4$, 3d ABJM,

& their gravity duals

"SUSY localization"

exact computation of
partition function
in SUSY QFT

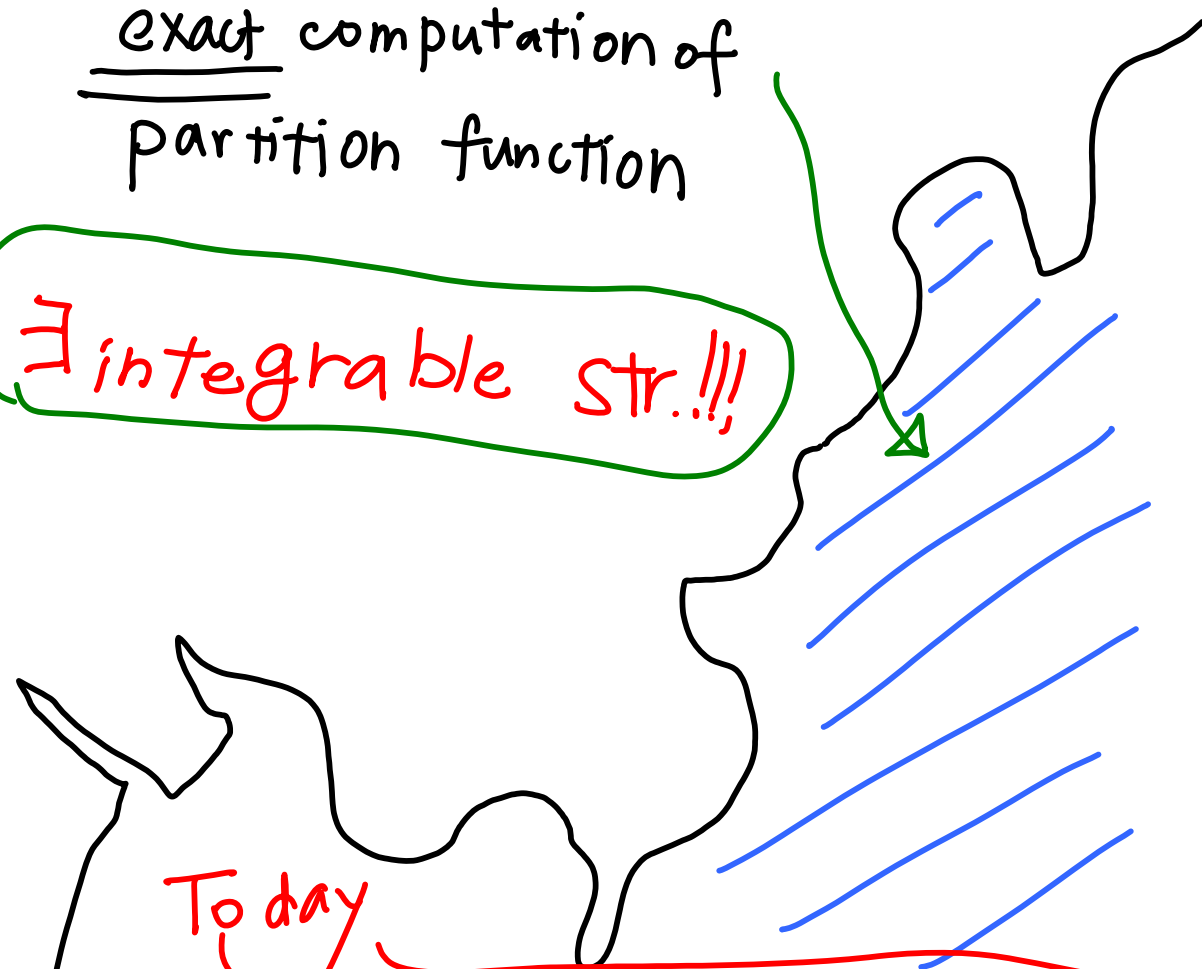
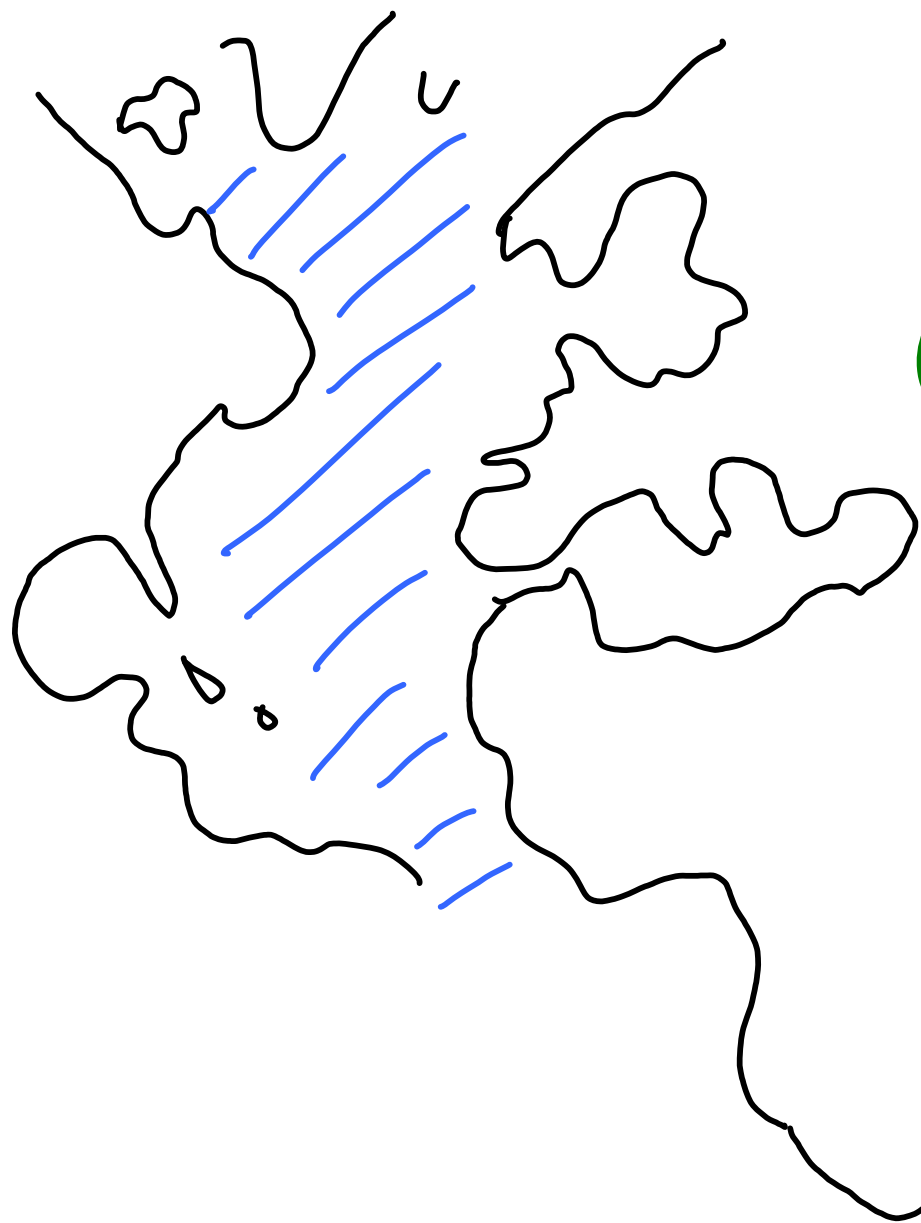
[Nekrasov, Pestun
Kapustin-Willet-Yaakov]



"SUSY localization"

exact computation of
partition function

\exists integrable str.!!!



To day

New integrable model
from 4d $N=1$ quiver
gauge theory

"SUSY localization"

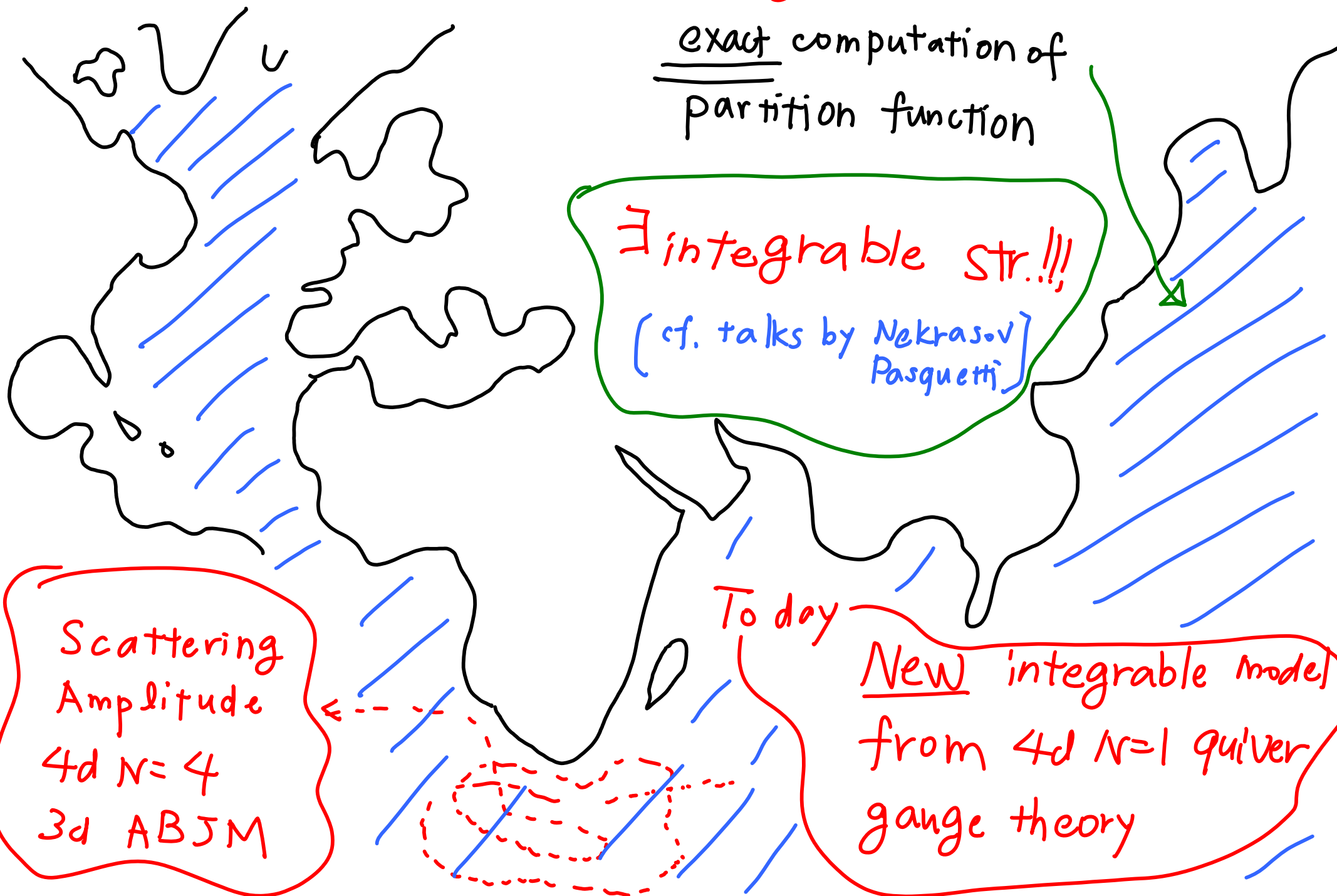
exact computation of
partition function

\exists integrable str.!!!
(cf. talks by Nekrasov
Pasquetti)

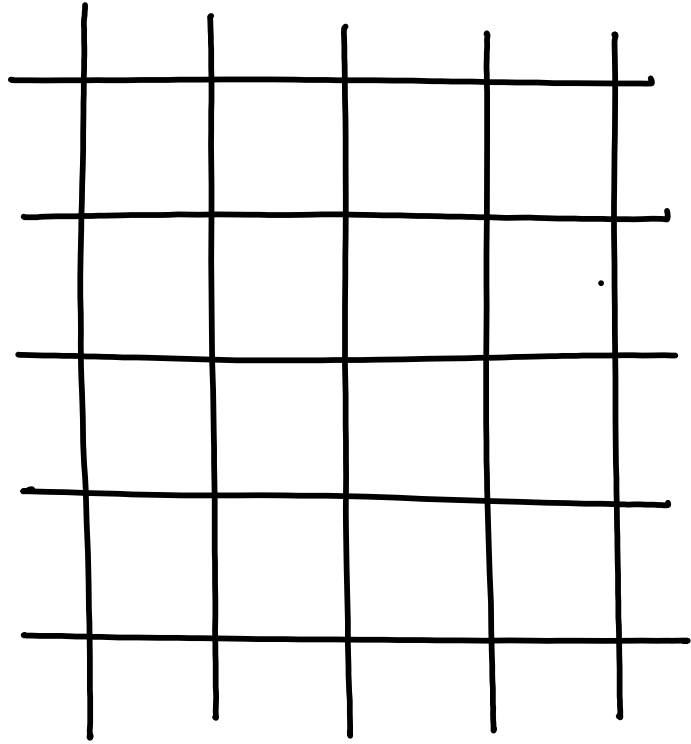
Scattering
Amplitude
4d $N=4$
3d ABJM

To day

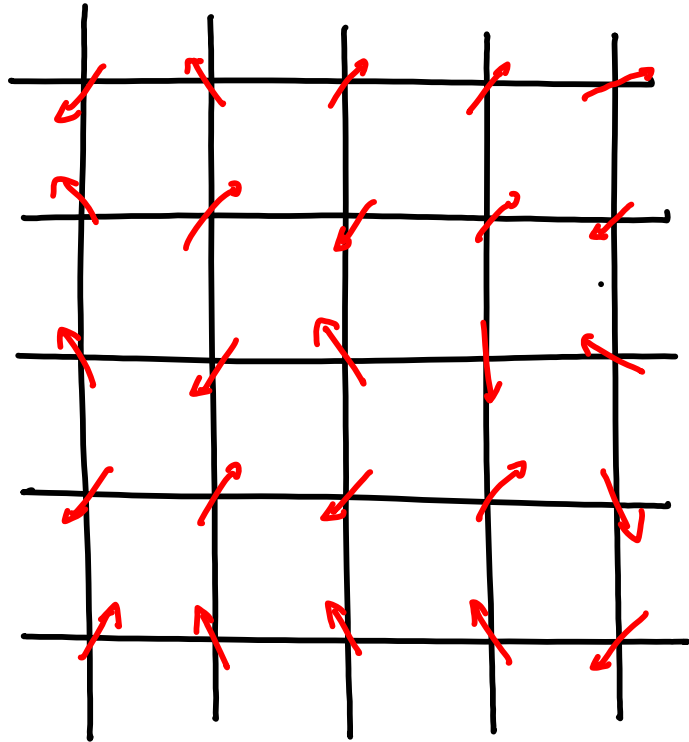
New integrable model
from 4d $N=1$ quiver
gauge theory



Gänge / YBE



stat-mech



- Spin \vec{s}_v at vertex $v \in V$

- energy

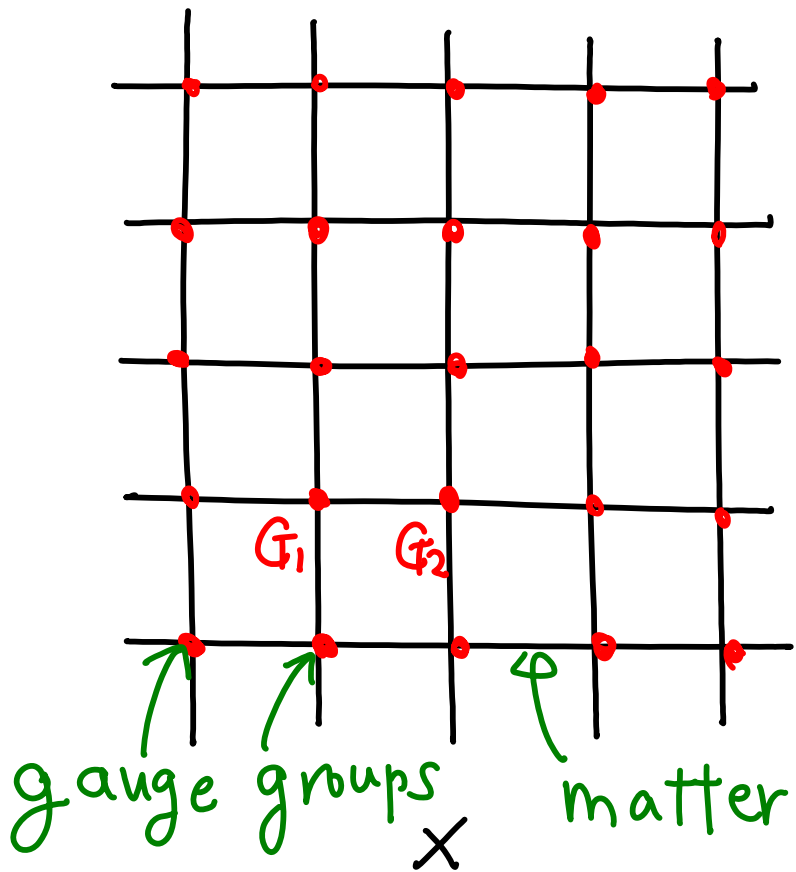
$$\mathcal{E} = \sum_{v \in V} \mathcal{E}_v(\vec{s}_v) + \sum_{e \in E} \mathcal{E}_e(\vec{s}_v, \vec{s}_{v'})$$

↑ edge e between v & v'

- partition function

$$\mathcal{Z} = \sum_{\{\vec{s}_v\}} e^{-\beta \mathcal{E}(\{\vec{s}_v\})}$$

(quiver) gauge theory



$\mathbb{R}^{d,1}$

- gauge group G_v at vertex $v \in V$
 $A_\mu^v(x)$

- matter charged under $G_v \times G_{v'}$ for an edge e between v, v'
 $\Phi_e(x)$

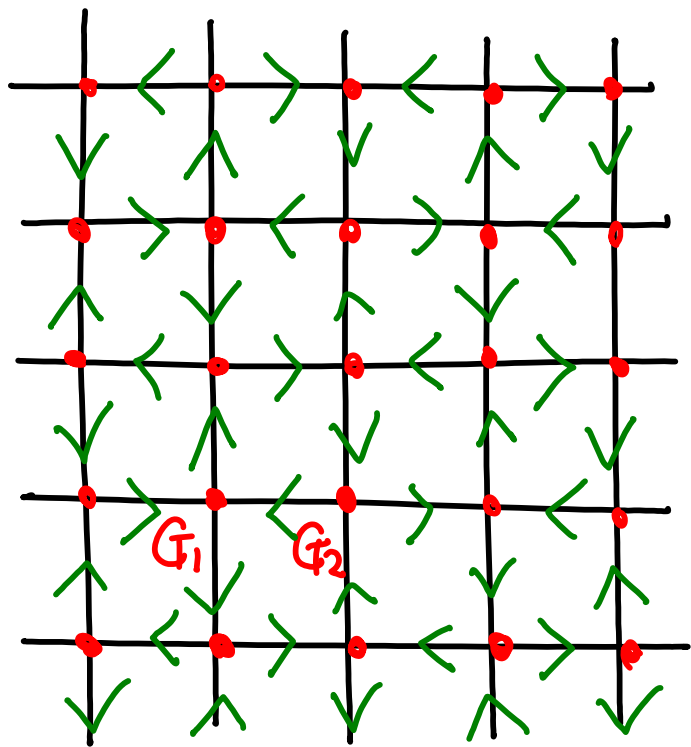
- Lagrangian

$$\mathcal{L} = \sum_{v \in V} \mathcal{L}_v(A_\mu^v) + \sum_{e \in E} \mathcal{L}(\Phi_e, A_\mu^v, A_\mu^{v'})$$

- partition function

$$Z = \int \prod_{v \in V} \pi \delta A_\mu^v \prod_{e \in E} \pi \delta \Phi_e e^{i\mathcal{L}}$$

4d $N=1$
 quiver gauge theory

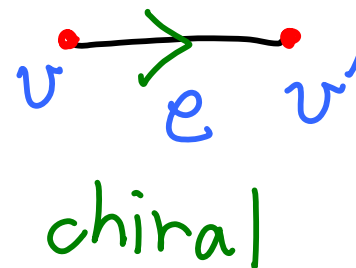


x

$\mathbb{R}^{3,1}$

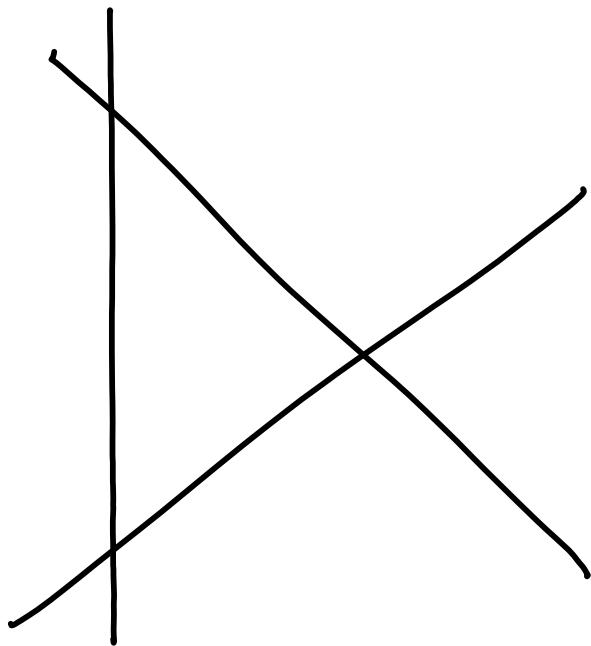
For concreteness we choose

- $G_v = SU(N_v)$ vector multiplet
 at vertex $v \in V$
 mostly $N \equiv N_v$
- bifundamental chiral multiplet
 $(N_v, \overline{N_{v'}})$ under $SU(N_v) \times su(N_{v'})$

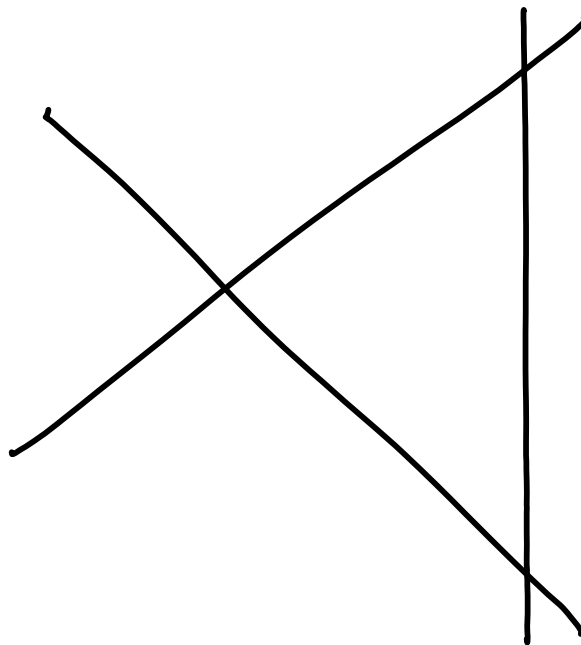


integrability?

YBE

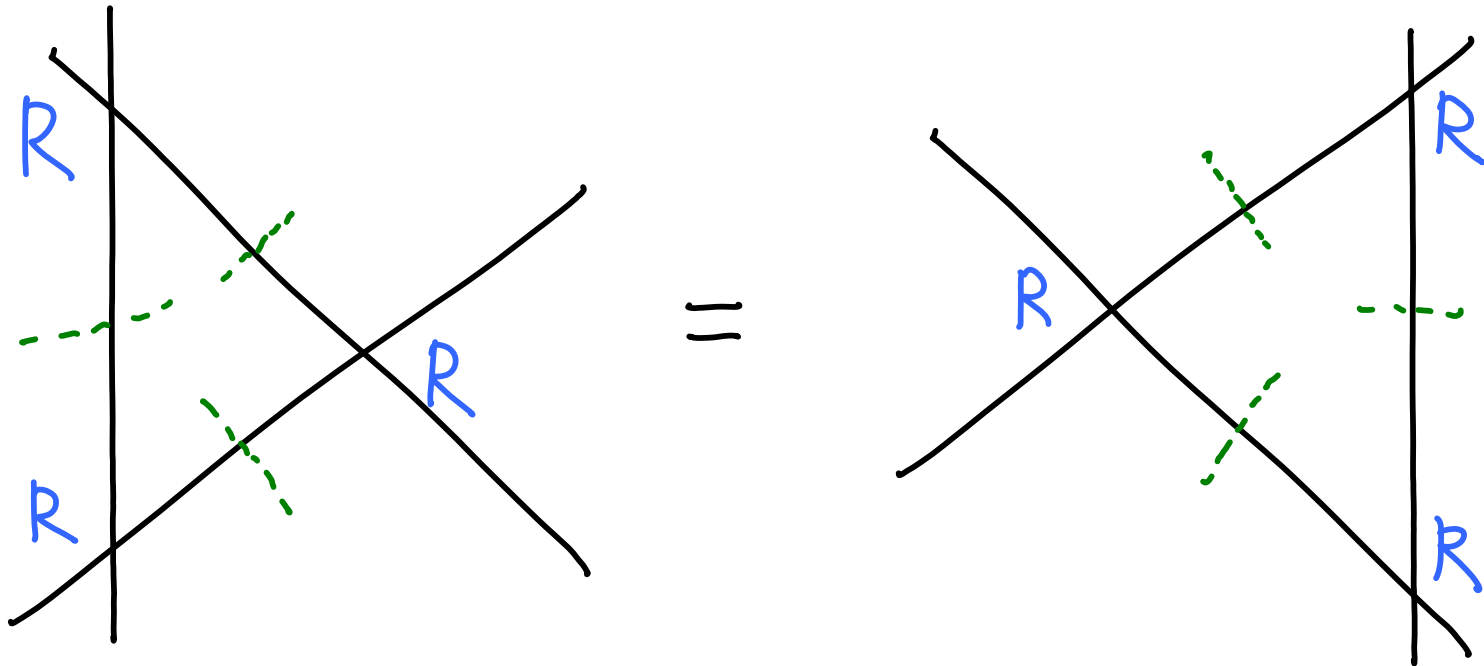


=



integrability?

YBE



$$\begin{array}{c}
 d \\
 | \\
 a \text{ --- } c \\
 | \\
 b
 \end{array}
 = R_{a \ b}^{\ d \ c}$$

$$\begin{array}{c}
 d \qquad g \\
 | \qquad | \\
 a \text{ --- } c \text{ --- } f \\
 | \qquad | \\
 b \qquad e
 \end{array}
 = \sum_c R_{a \ b}^{\ d \ c} R_{c \ e}^{\ g \ f}$$

$$= (R \circ R)_{a \ b}^{\ d \ g}$$

$$R: V_1 \otimes V_2 \rightarrow V_3 \otimes V_4$$

lift to gauge theory?

spin

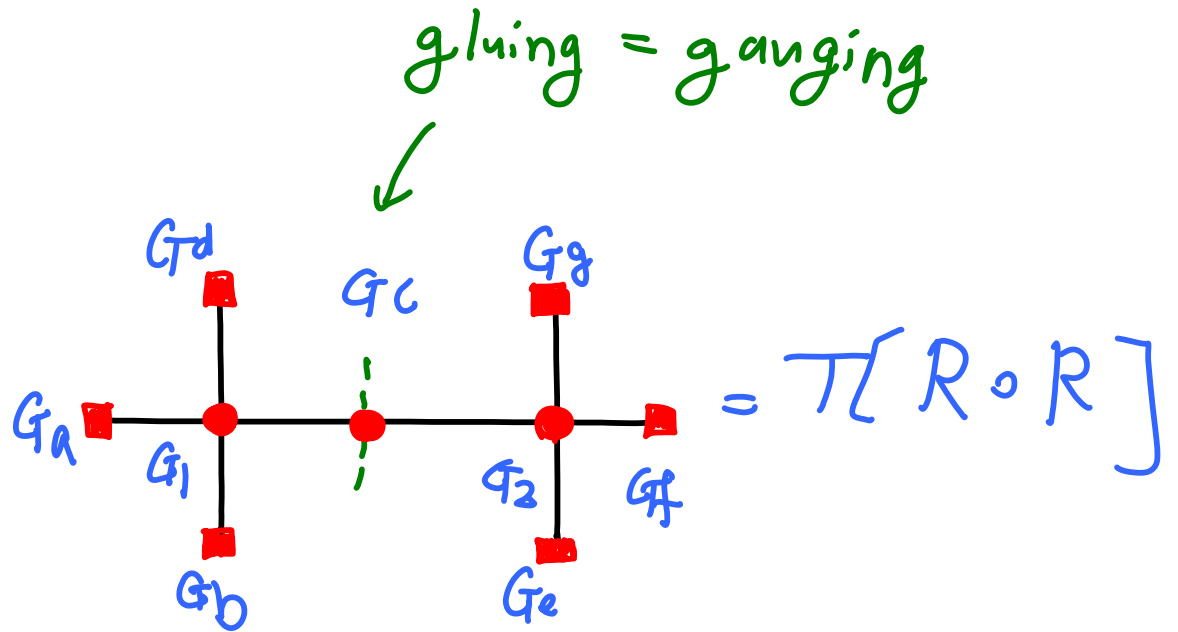
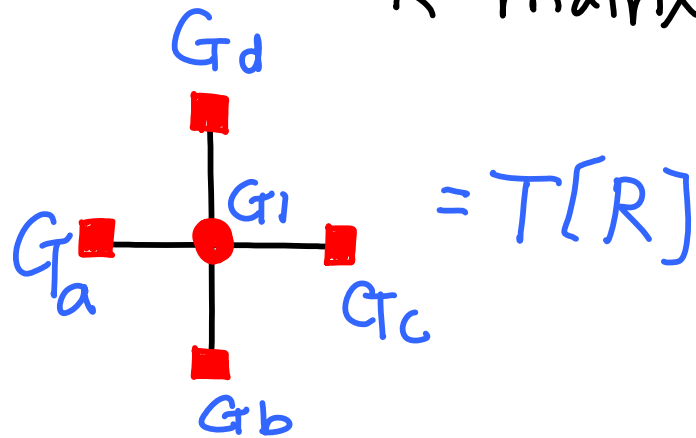
V



gauge group (vect mult.)

G

" the theory for the " R-matrix "

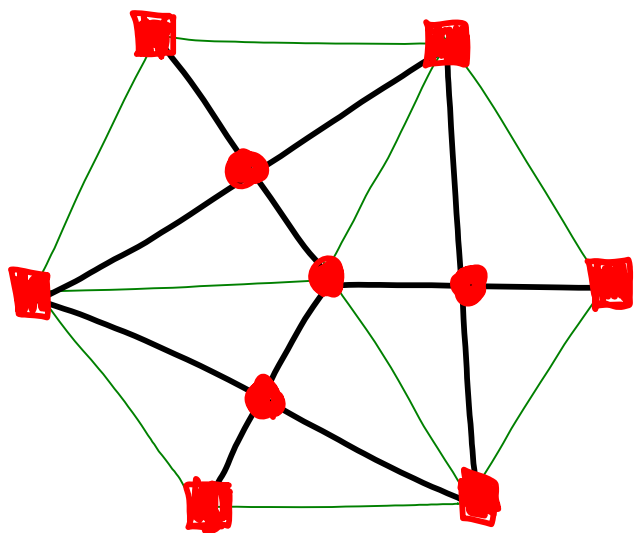
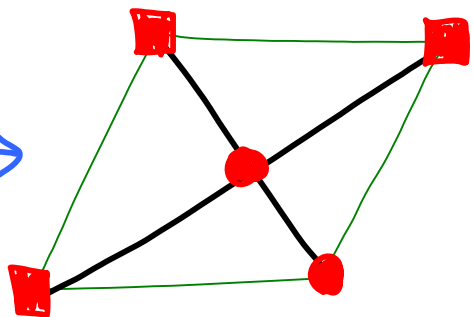


● = gauge sym ; dynamical ; sum in \mathbb{Z}

■ = flavor sym ; non-dynamical ; fixed parameter in \mathbb{Z}

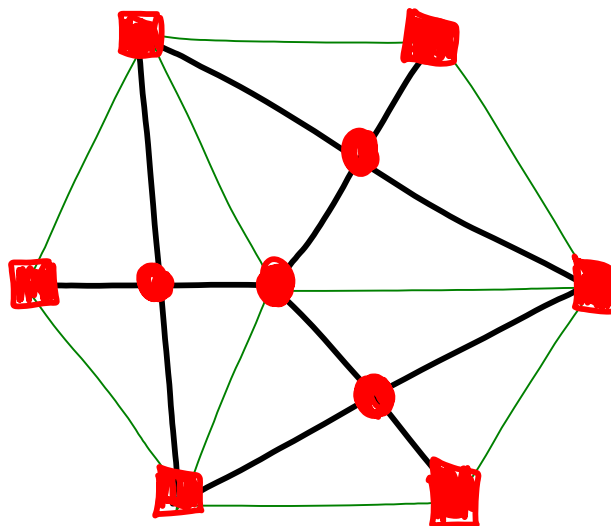
YBE as a duality?

R →



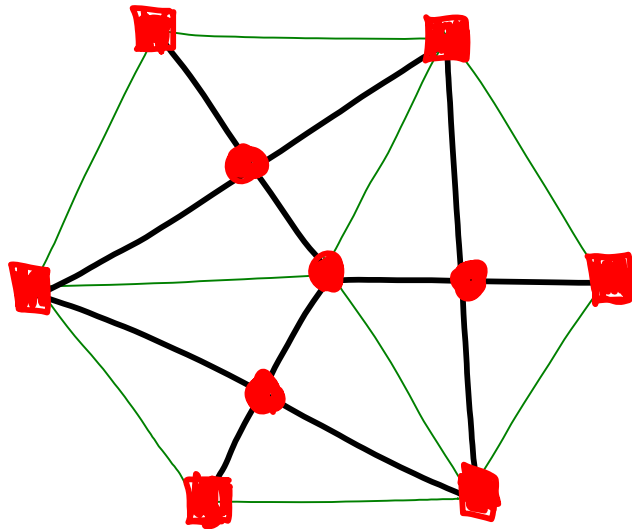
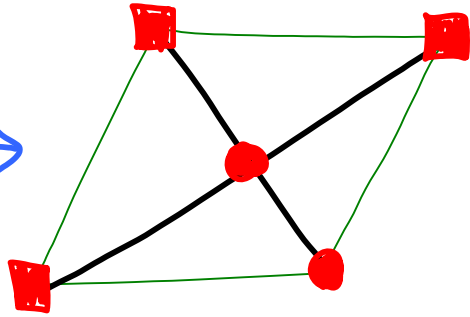
??
=

dual?



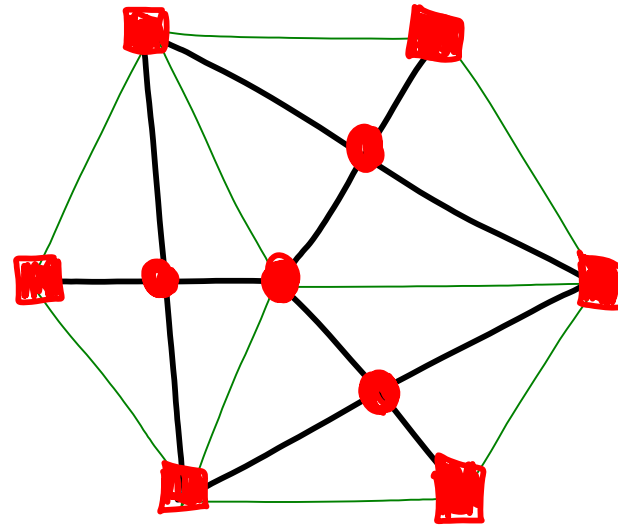
YBE as a duality?

R →



??
=

dual?

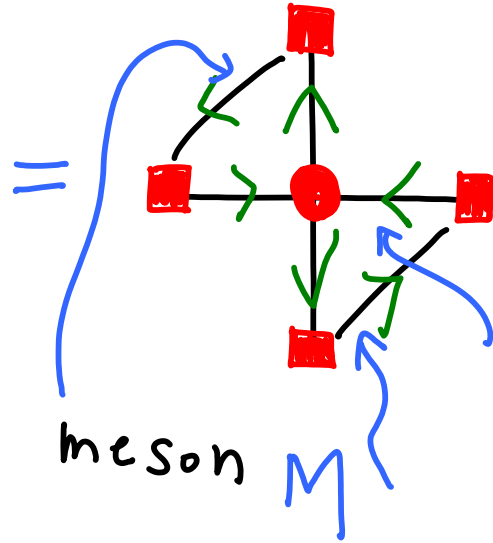
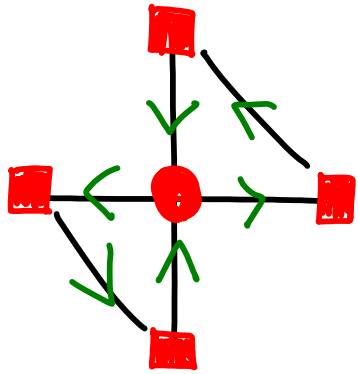


unfortunately there is no such known duality
a bit too naive so far...

the story works with a little modification,

with the help of **4d Seiberg duality**

[for $SU(N)$ $N_f = 2N$]

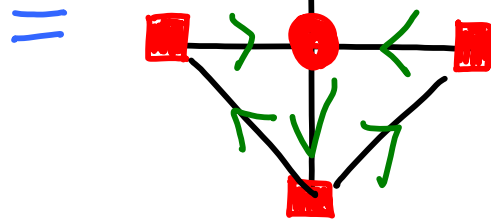
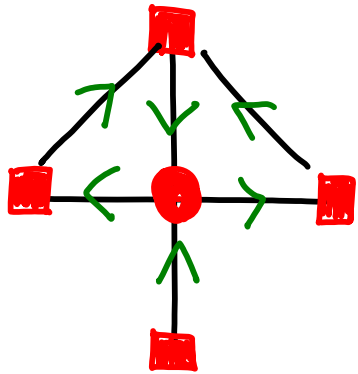


$= T[R]$

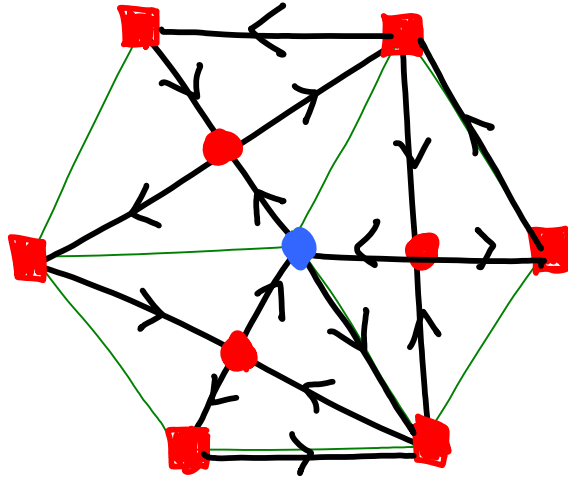
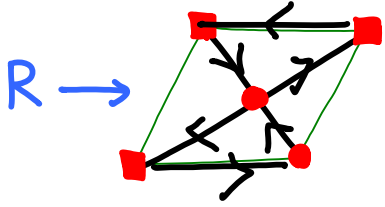
closed loop:

superpotential

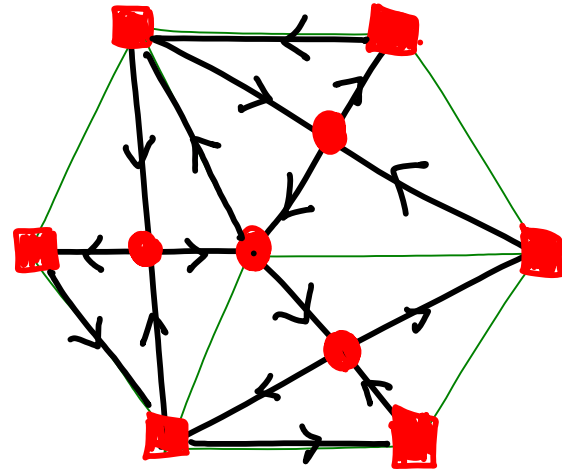
$$W = M q \tilde{q}$$



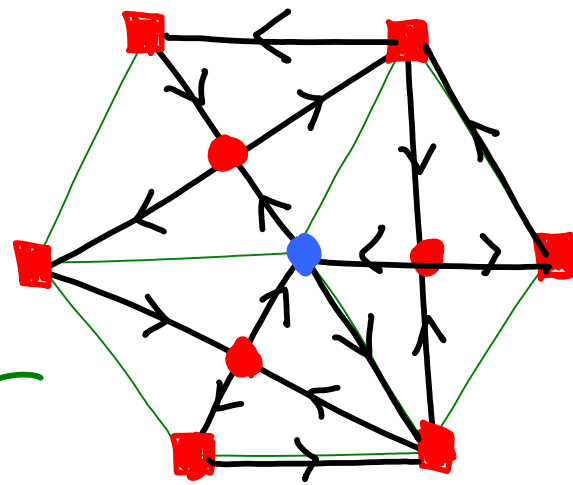
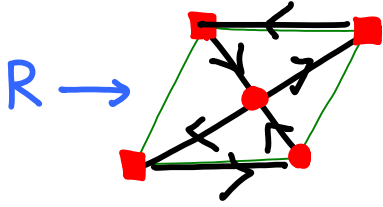
YBE \rightsquigarrow duality?



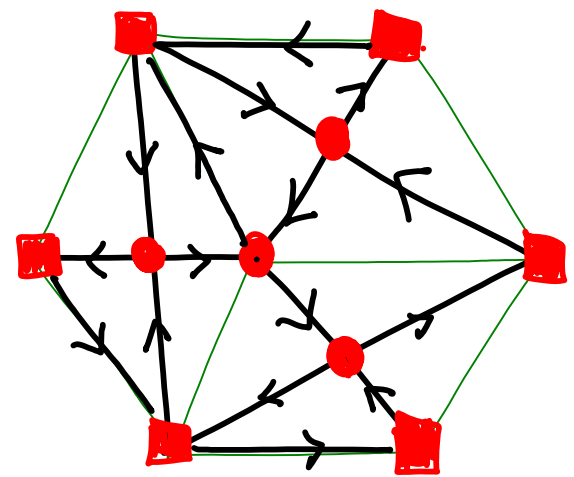
??
=
dual?



YBE \rightsquigarrow duality?

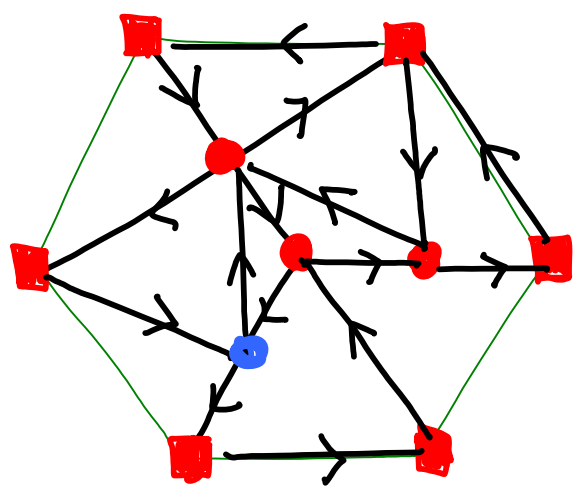


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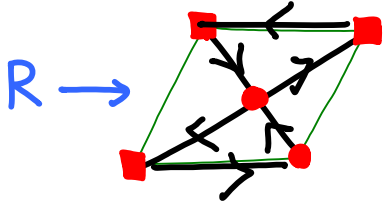
dual?

SD

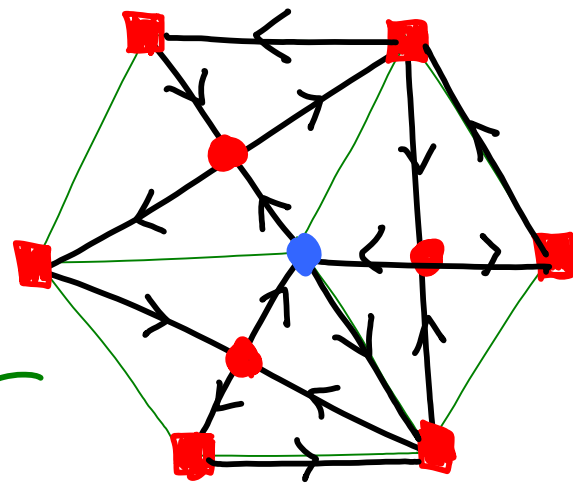


YBE \rightsquigarrow "YB duality" = (Seiberg duality)⁴ !!!

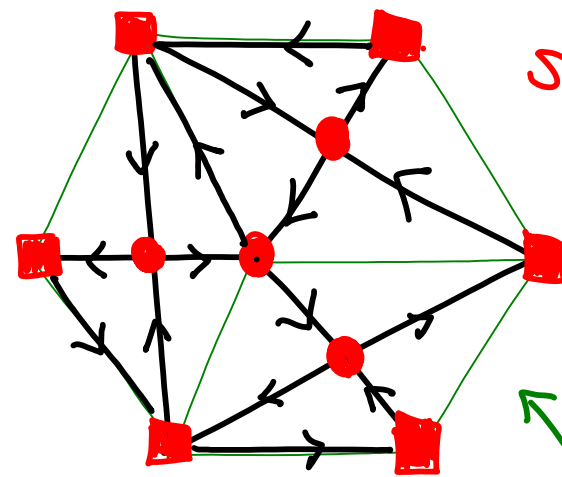
[cf. Baxter ('86)
Bazhanov-Baxter ('92)]



R \rightarrow



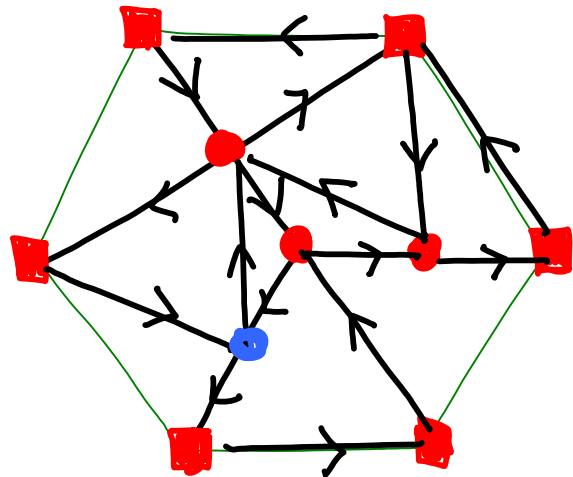
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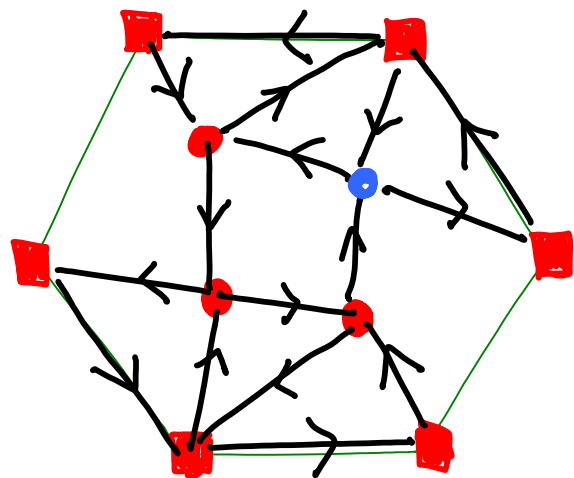
star-star
relation
 \downarrow
YBE

SD

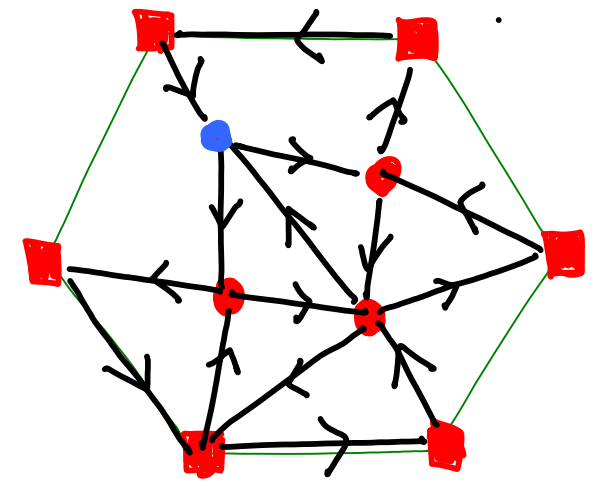
SD



SD



SD



back to integrable models

$$T[R_{12} \circ R_{13} \circ R_{23}] \overset{\text{dual}}{\longleftrightarrow} T[R_{23} \circ R_{13} \circ R_{12}] \quad \text{YB duality}$$

$$\rightsquigarrow \mathbb{Z}[T[R_{12} \circ R_{13} \circ R_{23}]] = \mathbb{Z}[T[R_{23} \circ R_{13} \circ R_{12}]]$$

back to integrable models

$$T[R_{12} \circ R_{13} \circ R_{23}] \overset{\text{dual}}{\longleftrightarrow} T[R_{23} \circ R_{13} \circ R_{12}] \quad \text{YB duality}$$

gauging \rightsquigarrow $Z[T[R_{12} \circ R_{13} \circ R_{23}]] = Z[T[R_{23} \circ R_{13} \circ R_{12}]]$ YBE

$$Z[\pi[R_{12}]] \cdot Z[T[R_{13}]] \cdot Z[T[R_{23}]] \quad Z[T[R_{23}]] \cdot Z[T[R_{13}]] \cdot Z[T[R_{12}]]$$

sum

$$Z = \int_{v+V} \prod_{e \in E} \Theta_{A_\mu^v} \prod_{e \in E} \Theta_{\Phi_e} e^{i\mathcal{L}}$$

localization

$$\mathcal{L} = \sum_{v \in V} \mathcal{L}_v(A_\mu^v) + \sum_{e \in E} \mathcal{L}(\Phi_e, A_\mu^v, A_\mu^{v'})$$

$$\begin{pmatrix} A_\mu(x) \rightarrow \sigma \\ \Phi_e \rightarrow 0 \end{pmatrix}$$

$$Z = \int_{v+V} \prod_{v \in V} d\sigma_v e^{i\mathcal{E}}$$

"spin"

$$\mathcal{E} = \sum_{v \in V} \mathcal{E}(\sigma_v) + \sum_{e \in E} \mathcal{E}(\sigma_v, \sigma_{v'})$$

$$Z = Z [S' \times S^3 / \mathbb{Z}_r] \quad [\text{Benini-Nishioka-Y (11)}]$$

$$= Z [N, r; [P, q; \vec{R}_i]] \quad (\text{R-charge}) = (\text{spectral parameter})$$

rk of $G = \text{SU}(N)$ \rightarrow N
 order of S^3 / \mathbb{Z}_r \rightarrow r
 P \rightarrow fugacity
 q \rightarrow "quantum parameter"
 \vec{R}_i \rightarrow (R-charge) = (spectral parameter)

"Spin"
 \vec{S}_V = holonomy of the gauge field along $S' \times S^3 / \mathbb{Z}_r$
 $= \vec{z}_V \in \text{U}(1)^{N-1}$ & $\vec{m}_V \in (\mathbb{Z}_r)^{N-1}$

The resulting model is one of the most general solutions to YBE known in the literature

$$\begin{aligned}
Z &= \sum_{s_v (v \in V)} \left(\prod_{v \in V} \mathbb{S}^v(s) \right) \left(\prod_{e \in E} \mathbb{W}^e(R; s) \right) \\
&= \sum_{\underbrace{m_v}_{\text{wavy}} (v \in V)} \int_{\underbrace{|z_{v,m}|=1}_{\text{wavy}}} \prod_{v \in V} \prod_{i=1}^{N-1} \frac{dz_{v,i}}{2\pi\sqrt{-1}z_{v,i}} \left(\prod_{v \in V} \underbrace{\mathbb{S}^v(z, m)}_{\text{wavy}} \right) \left(\prod_{e \in E} \underbrace{\mathbb{W}^e(R; z, m)}_{\text{wavy}} \right) ,
\end{aligned}$$

$$\underbrace{\mathbb{S}^v(s)}_{\text{wavy}} = \mathbb{S}^v(s_v) = \left(\prod_{a=0}^{r-1} \frac{1}{n_{v,a}!} \right) \underbrace{\mathcal{S}_0^v(s_v)}_{\text{wavy}} (\Gamma_{r,0}(1; p, q))^{-(N-1)} \prod_{i \neq j} \Gamma_{r, \llbracket m_i - m_j \rrbracket} \left(\frac{z_{v,i}}{z_{v,j}}; p, q \right)^{-1} ,$$

$$\mathcal{S}_0^v(s_v) = (pq)^{-\frac{1}{4r}} \sum_{i \neq j} \llbracket m_{v,i} - m_{v,j} \rrbracket \llbracket -m_{v,i} + m_{v,j} \rrbracket .$$

$$\underbrace{\Gamma_{r, \llbracket m \rrbracket}(x; p, q)}_{\text{wavy}} = \Gamma(xp^{\llbracket m \rrbracket}; pq, p^r) \Gamma(xq^{r - \llbracket m \rrbracket}; pq, q^r) ,$$

$$\underbrace{\Gamma(x; p, q)}_{\text{wavy}} = \prod_{n_1, n_2 \geq 0} \frac{1 - x^{-1}p^{n_1+1}q^{n_2+1}}{1 - xp^{n_1}q^{n_2}} .$$

$$\begin{aligned}
\mathbb{W}^e(R; s) &= \mathbb{W}_{R_e}^e(s_{t(e)}, s_{h(e)}) \\
&= \mathcal{W}_0^e(s_{t(e)}, s_{h(e)}) \prod_{1 \leq i, j \leq N} \Gamma_{r, \llbracket m_{t(e), i} - m_{h(e), j} \rrbracket} \left((pq)^{\frac{R_e}{2}} \frac{z_{t(e), i}}{z_{h(e), j}}; p, q \right),
\end{aligned}$$

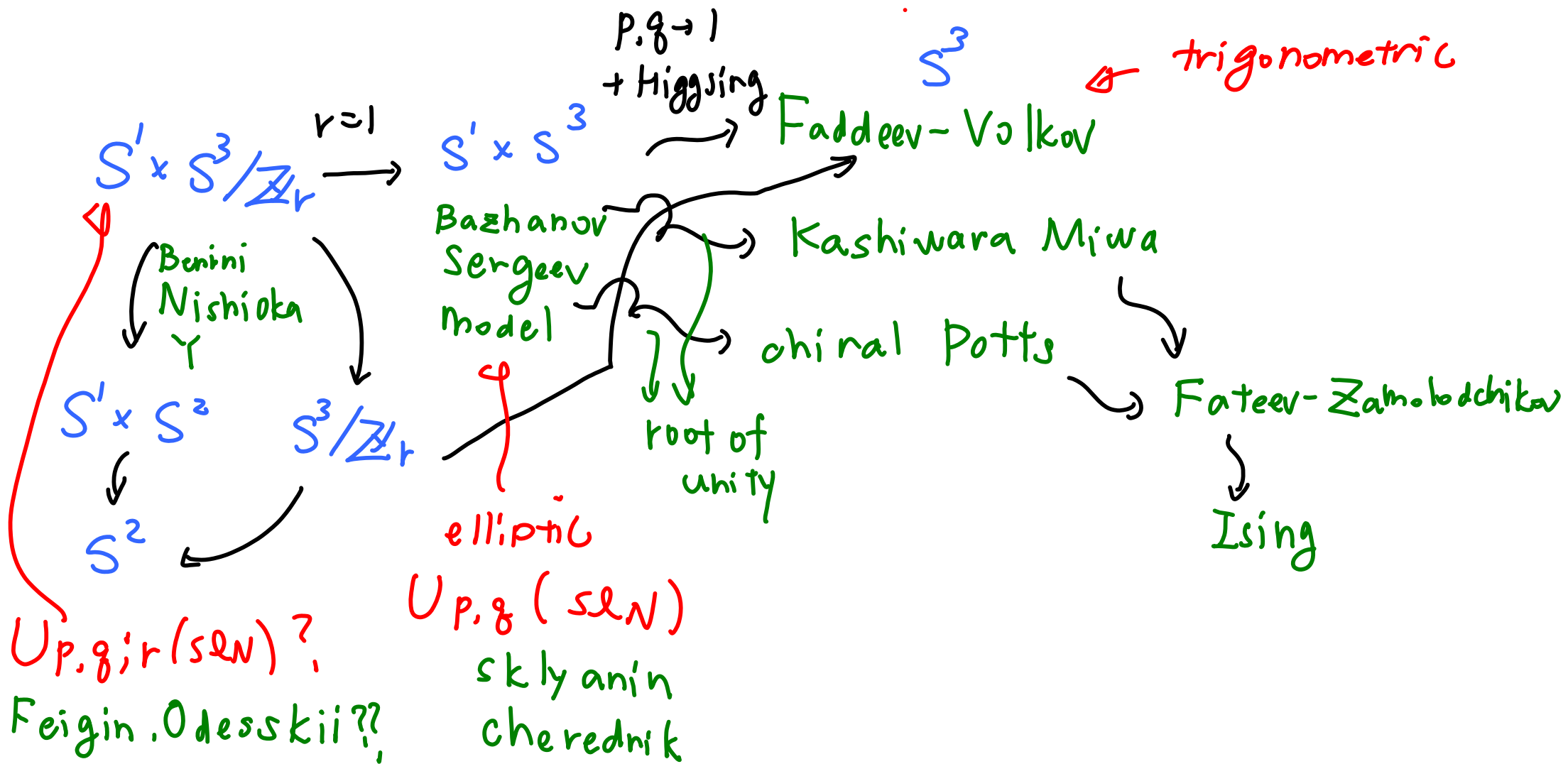
$$\begin{aligned}
\mathcal{W}_0^e(s_{t(e)}, s_{h(e)}) &= \prod_{1 \leq i, j \leq N} \left[(pq)^{\frac{1}{4r} \llbracket m_{t(e), i} - m_{h(e), j} \rrbracket} \llbracket -m_{t(e), i} + m_{h(e), j} \rrbracket (1 - R_e) \right. \\
&\quad \times \left(\frac{p}{q} \right)^{-\frac{1}{12r} \llbracket m_{t(e), i} - m_{h(e), j} \rrbracket} \llbracket -m_{t(e), i} + m_{h(e), j} \rrbracket (2 \llbracket m_{t(e), i} - m_{h(e), j} \rrbracket - r) \\
&\quad \left. \times \left(\frac{z_{t(e), i}}{z_{h(e), j}} \right)^{-\frac{1}{2r} \llbracket m_{t(e), i} - m_{h(e), j} \rrbracket} \llbracket -m_{t(e), i} + m_{h(e), j} \rrbracket \right].
\end{aligned}$$

$$\Gamma_{r, \llbracket m \rrbracket}(x; p, q) = \Gamma(xp^{\llbracket m \rrbracket}; pq, p^r) \Gamma(xq^{r - \llbracket m \rrbracket}; pq, q^r),$$

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Comments

The model has parameters $N, r; p, q$
 We can consider specializations



Witten, 1990 "gauge theories, vertex models & quantum groups"

There are several obvious areas for further investigation. In terms of statistical mechanics, one compelling question is to understand the origin of the spectral parameter (and the elliptic modulus) in IRF and vertex models; this is essential for explaining the origin of integrability.

✓ Spectral parameter = R-charge

Another question, which may or may not be related, is to understand the spin models formulated only rather recently [24] in which the spectral parameter is not an abelian variable but is a point on a Riemann surface of genus greater than one.

If possible, one would also like to understand vertex models and quantum groups directly at physical values of q , without the less than appealing analytic continuation that is used in this paper.

Perhaps related to this, one would like if possible to see the origin of quantum groups and not just quantum Lie algebras.

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chiral Potts ✓ root of unity reduction. gauge theory meaning?

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? We found $U_{p,q,r}(g)$

See however [Costello ('13)]

Gauge/YBE from 6d

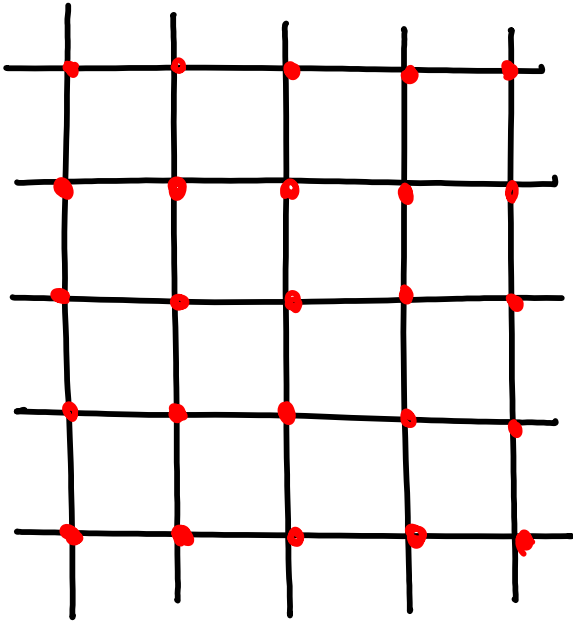
(Similar to AGT)

6d origin:
D5/NS5 system

fixed boundary condition
[Heckman-Vafa-Xie-MY]
(12)

periodic boundary
condition

2d spin
system



$\times \mathbb{R}^{3,1}$
4d gauge
theory

[Imamura
Imamura-Kimura-Isono-Y
(07)]

Gauge / YBE

directly gives integrable model 😊

$$I_{S^1 \times S^3 / \mathbb{Z}_n} \left[\begin{array}{l} 4d \\ N=1 \end{array} \right] = \mathbb{Z} \left[\text{integrable model} \right]$$

~~||~~

Gauge / Bethe [Nekrasov Shatashvili]

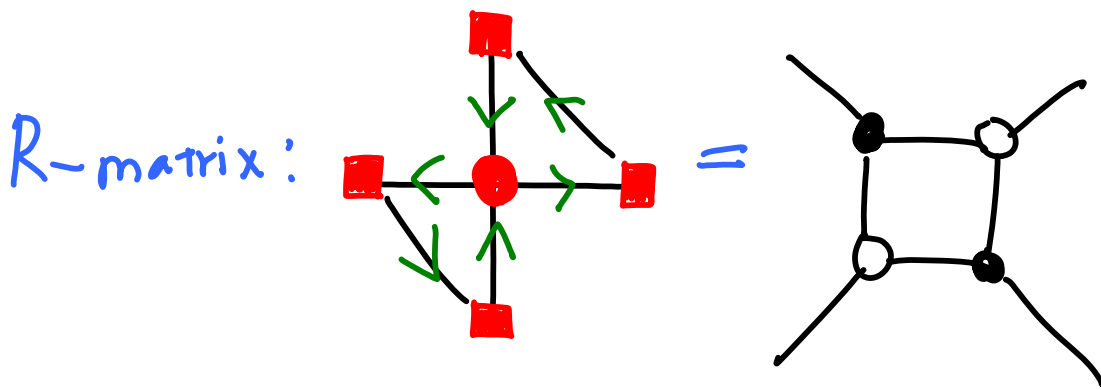
Yang-Yang function

$$W_{\text{twisted}} \left[\begin{array}{l} 2d \\ N=(2,2) \end{array} \right] = Y \left[\begin{array}{l} \text{integrable} \\ \text{model} \end{array} \right]$$

(~~*~~ in both cases we have "integrability in theory space")

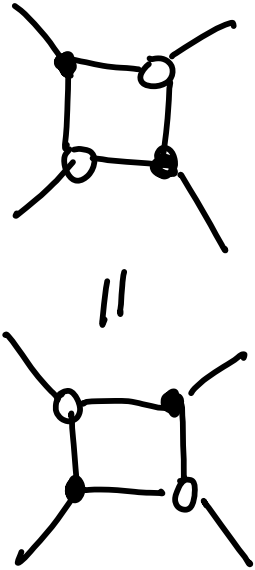
Scattering Amplitude

quivers \leftrightarrow bipartite graph \rightsquigarrow 4d $N=1$ gauge theory

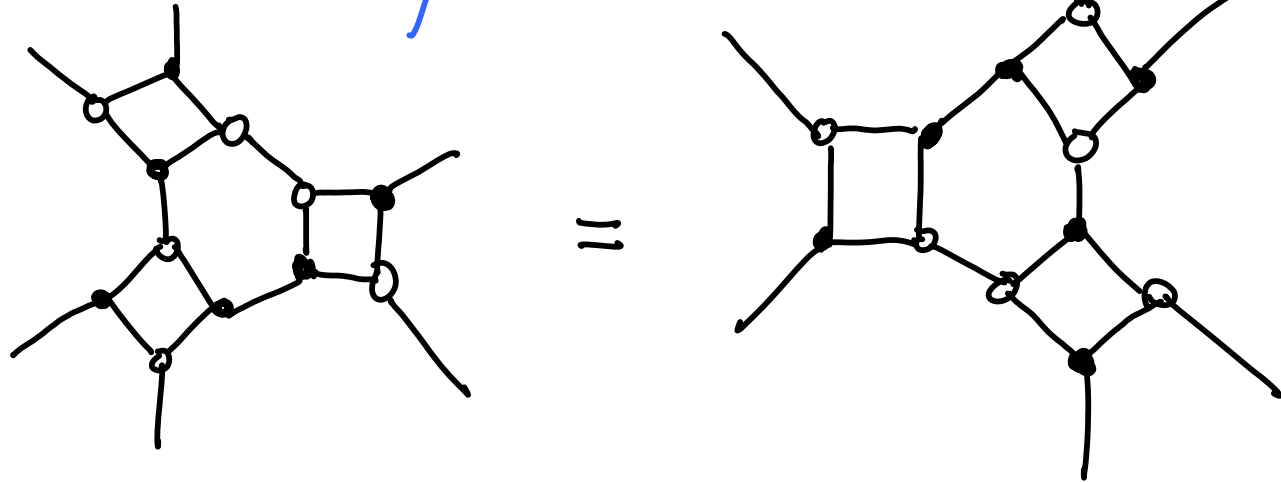


Hanany-Kennaway, ...
 "brane tilings"
 Xie-Yamazaki, Franco ('12)

Seiberg duality



YB-duality



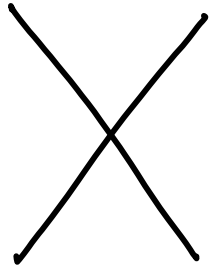
On-shell diagram for 4d $N=4$

[Arkani-Hamed, Bourjaily, Goncharov, Postnikov, Trnka ('12)]
 [also Ferro-Lukowski-Meneghelli-Plefka-Staudacher]

The connection with YBE is simpler for

3d ABJM theory

building block: 4 pt [ABGPT, Huang-Wen('13)]



orthogonal Grassmannian
OG(2,4)

$$A_4(\mathbb{Z}) = \frac{\delta^3(P)\delta^3(Q)}{\langle 12 \rangle^{1+\mathbb{Z}} \langle 3 \rangle^{1-\mathbb{Z}}} = \int \frac{d^{2 \times 4} C}{|GL(2)|} \frac{\delta^3(C \cdot C^T) \delta^4(C \cdot \Lambda)}{M_1^{1+\mathbb{Z}} M_2^{1-\mathbb{Z}}}$$

invariant under Yangian $Y[OSp(6|4)]$

with evaluation parameters $(\mathbb{Z}, 0, \mathbb{Z}, 0)$

[Bargheer - Loebbert - Huang - Y]
(yesterday)

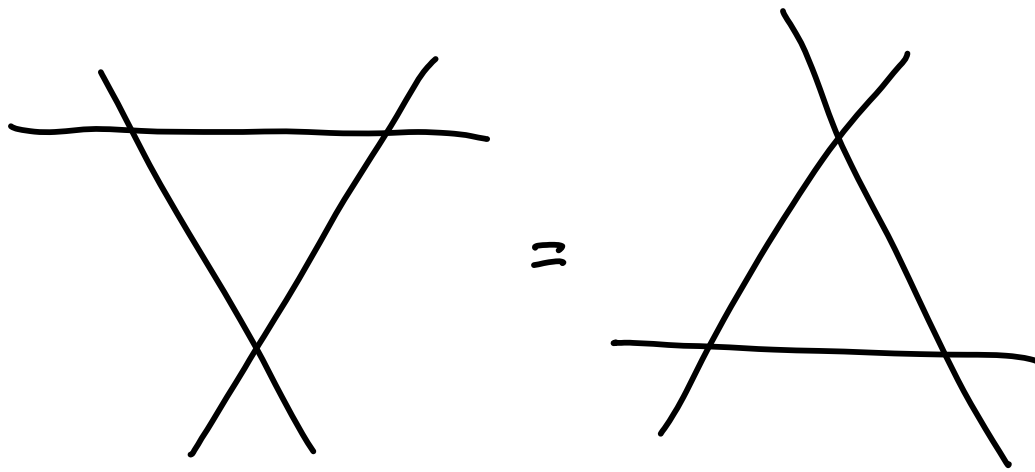
R-matrix: 4-pt

$$R(z) = A_4(z)$$

YBE: 6-pt / triangle move \leftarrow uniqueness of

top cell of

$$OG(3,6)$$



The Grassmannian integral $R(z)$

is an R-matrix for $Y[OSp(6|4)]$;

satisfy YBE / RLL = LLR relation

Conclusion

- YB -duality = (Seiberg duality)⁴ $\xrightarrow{\mathbb{Z}}$ YBE
(4d quiver gauge theory) (IM)

$$I_{S^1 \times S^3 / \mathbb{Z}_r} [4d, N=1] = \mathbb{Z} [\text{integrable model}]$$

- New, one of the most general solutions to YBE
- mysterious parallel w/ scattering amplitude

Q: We now understand R-matrix/YBE, but
can we lift the rest to SUSY gauge
theory?

[e.g. $RLL = LLR$ relation as a duality?
BAE?]

Q: Integrable structure in *space of QFTs*?

New formulation of QFTs?

