

Renormalization Group approaches to quantum gravity

Astrid Eichhorn

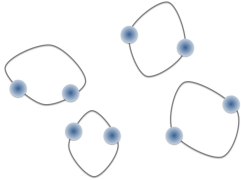
Perimeter Institute, Waterloo

DESY, Theory Seminar, 31.3.2014

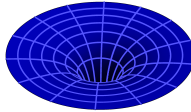


Quantum fields and gravitational fields

quantum fields:

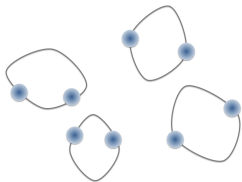


gravity:

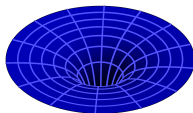


Quantum fields and gravitational fields

quantum fields:

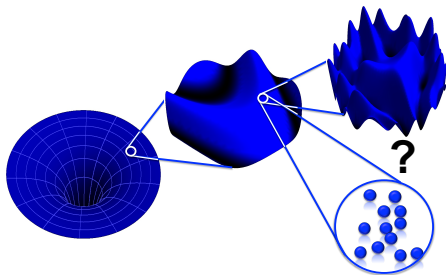


gravity:



→ quantum gravity:

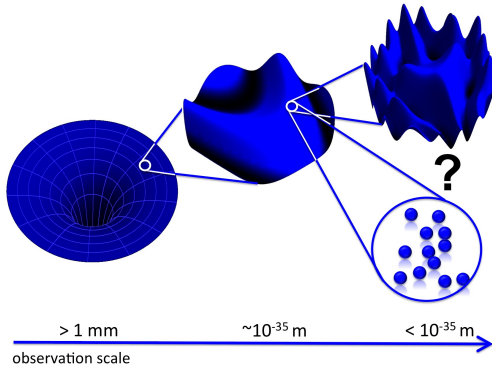
spacetime fluctuations
at the Planck scale



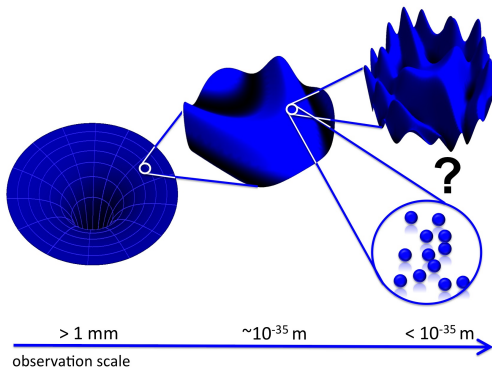
quantum theory of gravity in the path-integral framework:

Goal: $\int_{\text{spacetimes}} e^{iS}$

The Renormalization Group in quantum gravity



The Renormalization Group in quantum gravity

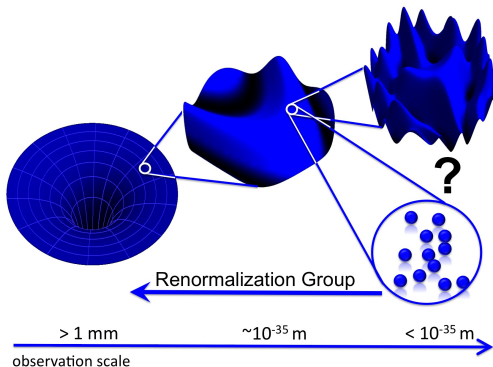


effective action for effective degrees of freedom at k :

$$e^{-\Gamma_k[\phi]} = \int_{p>k} \mathcal{D}\varphi e^{-S[\varphi]}$$



The Renormalization Group in quantum gravity

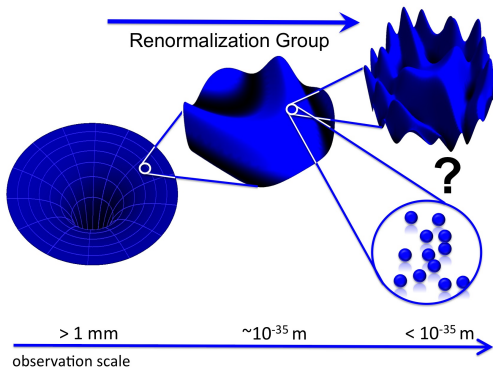


effective action for effective degrees of freedom at k :

$$e^{-\Gamma_k[\phi]} = \int_{p>k} \mathcal{D}\varphi e^{-S[\varphi]}$$



The Renormalization Group in quantum gravity

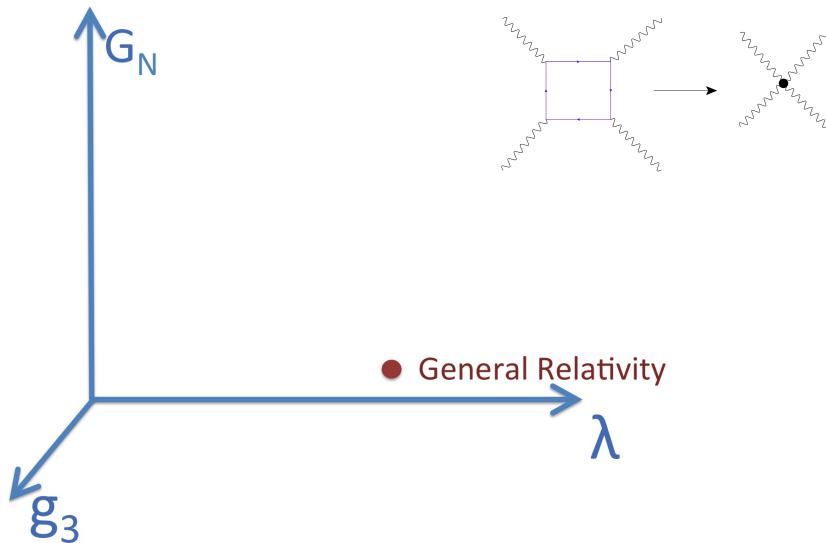


effective action for effective degrees of freedom at k :

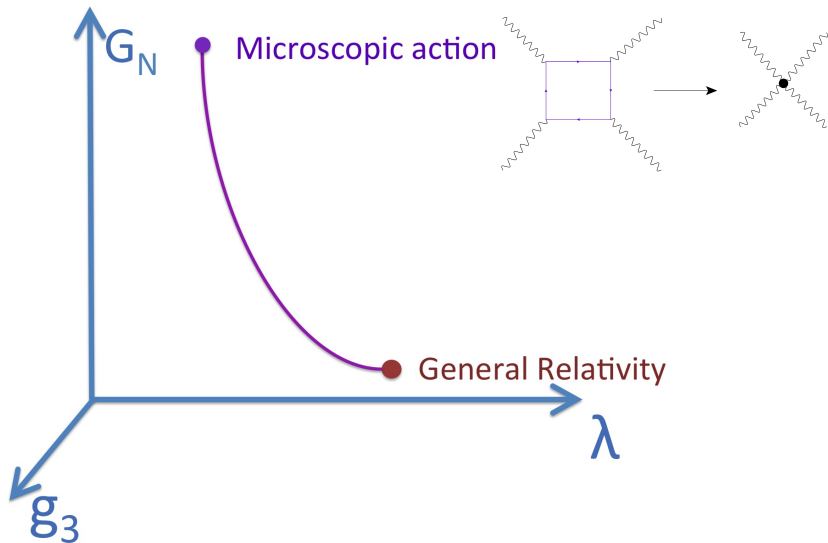
$$e^{-\Gamma_k[\phi]} = \int_{p>k} \mathcal{D}\varphi e^{-S[\varphi]}$$



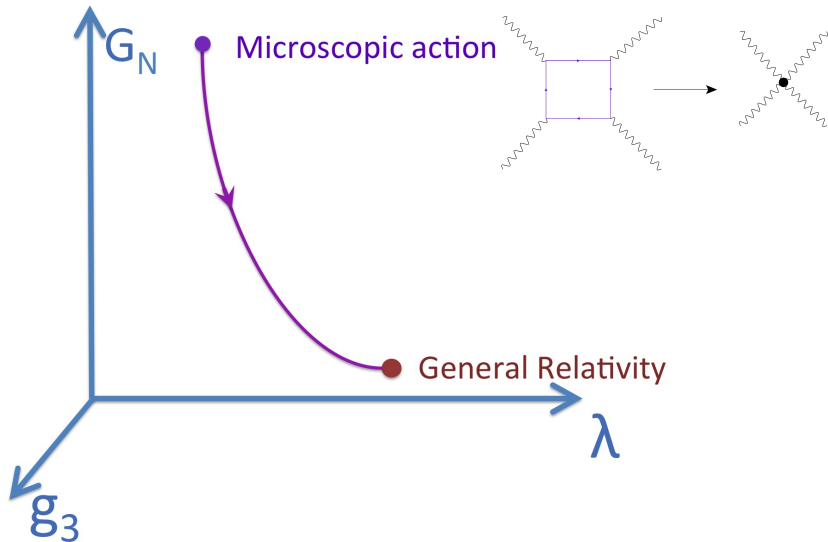
The Renormalization Group in quantum gravity



The Renormalization Group in quantum gravity

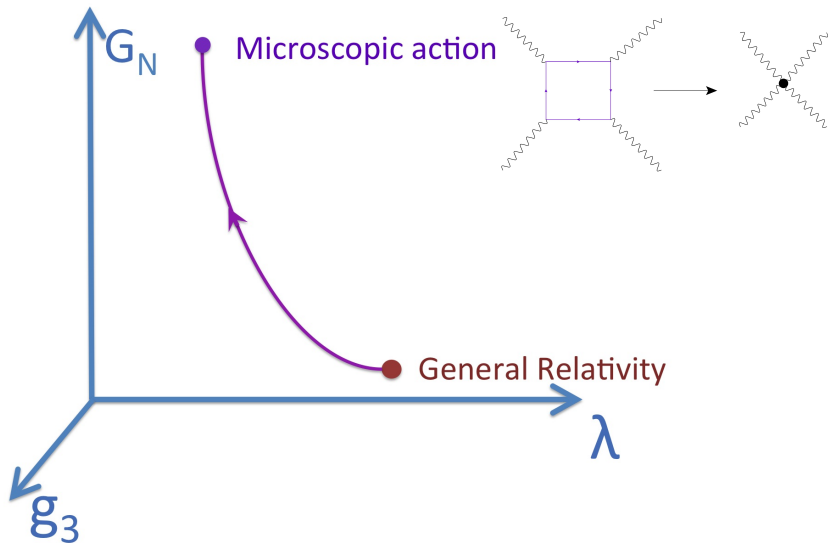


The Renormalization Group in quantum gravity



test continuum limit of discrete models (matrix/ tensor models)

The Renormalization Group in quantum gravity

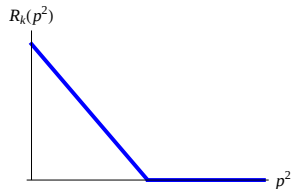


find out whether quantum gravity is asymptotically safe

Functional Renormalization Group

include high- energy quantum fluctuations first:

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int_p \varphi(p) R_k(p) \varphi(-p)}$$



Functional Renormalization Group

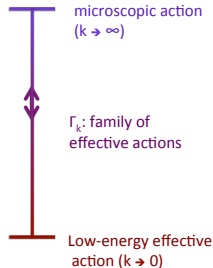
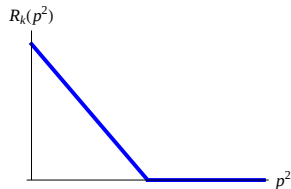
include high- energy quantum fluctuations first:

$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int_p \varphi(p) R_k(p) \varphi(-p)}$$

scale-dependent action:

Γ_k contains effect of quantum fluctuations above
momentum scale k

$$\begin{aligned} \Gamma_{k \rightarrow 0} &= \Gamma & \Gamma_{k \rightarrow \Lambda \rightarrow \infty} &\rightarrow S \\ \Gamma &= \Gamma_{\text{EH}} & S: &\text{prediction!} \end{aligned}$$



Functional Renormalization Group

include high- energy quantum fluctuations first:

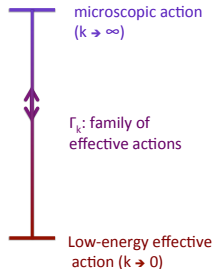
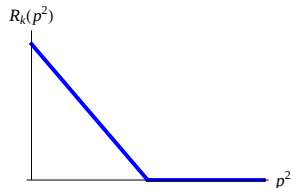
$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int_p \varphi(p) R_k(p) \varphi(-p)}$$

scale-dependent action:

Γ_k contains effect of quantum fluctuations above
momentum scale k

$$\Gamma_{k \rightarrow 0} = \Gamma \quad \Gamma_{k \rightarrow \Lambda \rightarrow \infty} \rightarrow S$$

$$\Gamma = \Gamma_{\text{EH}} \quad S: \text{prediction!}$$



Wetterich equation [1993]:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$


Functional Renormalization Group

include high- energy quantum fluctuations first:

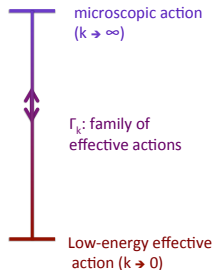
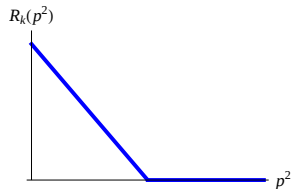
$$e^{-\Gamma_k[\phi]} = \int \mathcal{D}\varphi e^{-S[\varphi] - \frac{1}{2} \int_p \varphi(p) R_k(p) \varphi(-p)}$$

scale-dependent action:

Γ_k contains effect of quantum fluctuations above
momentum scale k

$$\Gamma_{k \rightarrow 0} = \Gamma \quad \Gamma_{k \rightarrow \Lambda \rightarrow \infty} \rightarrow S$$

$$\Gamma = \Gamma_{\text{EH}} \quad S: \text{prediction!}$$



Wetterich equation [1993]:

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k =$$


$$\Gamma_k = \sum_i g_i(k) \mathcal{O}_i$$

$\rightarrow$ non-perturbative $\beta_g = k \partial_k g(k) \rightarrow$ Quantum Gravity [Reuter, 1996]

Functional Renormalization Group

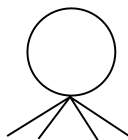
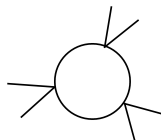
$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k: \text{ exact one-loop equation}$$

Functional Renormalization Group

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k: \text{ exact one-loop equation}$$

$$\Gamma_k = \sum_i \bar{g}_i(k) \mathcal{O}_i[g_{\mu\nu}] \rightarrow \text{infinitely many coupled equations}$$

example: $\lambda\phi^4$ theory:

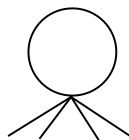
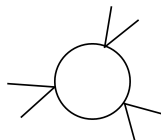


Functional Renormalization Group

$$\partial_k \Gamma_k = \frac{1}{2} \text{STr} \left(\Gamma_k^{(2)} + R_k \right)^{-1} \partial_k R_k: \text{ exact one-loop equation}$$

$$\Gamma_k = \sum_i \bar{g}_i(k) \mathcal{O}_i[g_{\mu\nu}] \rightarrow \text{infinitely many coupled equations}$$

example: $\lambda\phi^4$ theory:



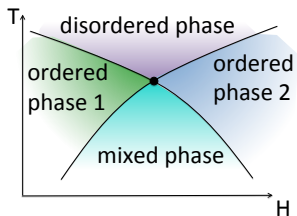
in practice:

- 1 truncated ansatz for $\Gamma_k = \sum_i^I \bar{g}_i(k) \mathcal{O}_i \rightarrow$ evaluate β functions for $g_i(k)$
- 2 enlarge truncation for Γ_k
- 3 Convergence of results?

FRG in different physical systems: Does it work?

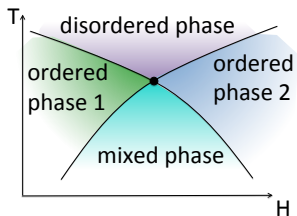
FRG in different physical systems: Does it work?

multicritical points in condensed matter systems
with competing orders



FRG in different physical systems: Does it work?

multicritical points in condensed matter systems
with competing orders



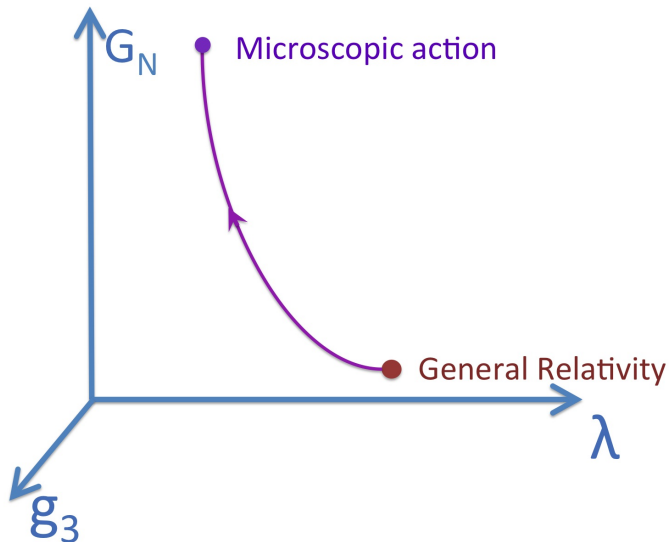
scaling in the vicinity of critical point: can
be determined from Renormalization Group
fixed-point scaling

results: scaling exponent for $O(2) \oplus Z_2$:

$$y_{2,2} = 1.790 \quad [\text{AE, Mesterházy, Scherer, Phys. Rev. E 88 (2013) 042141}];$$

$$y_{2,2} = 1.7906(3) \quad (\text{Monte Carlo}) \quad [\text{Hasenbusch, Vicari (2007)}]$$

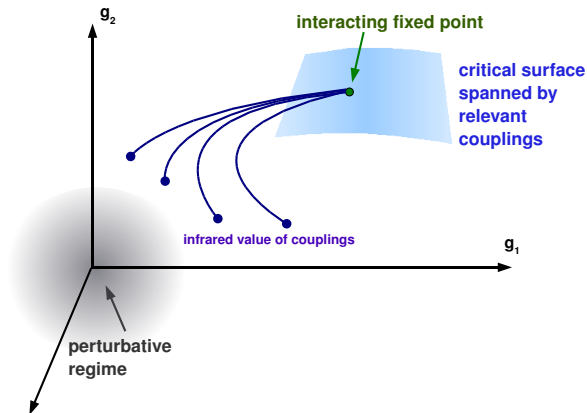
The Renormalization Group in quantum gravity



find out whether quantum gravity is asymptotically safe

Asymptotic safety

Fundamental QFT:
 $\int_{p>k} \mathcal{D}\varphi e^{-S}$
couplings stay finite
for $k \rightarrow \infty$
scale-free!

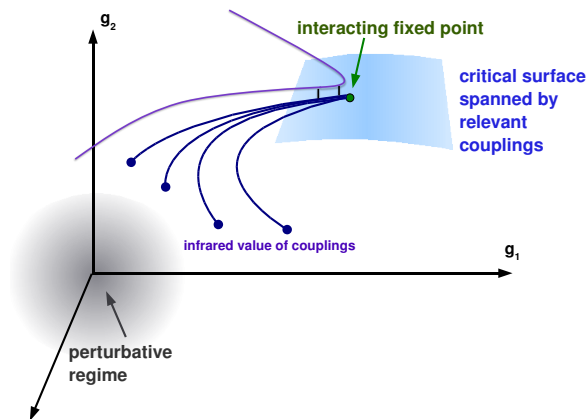


Fundamental theory:

Running dimensionless couplings approach fixed point towards UV

Asymptotic safety

Fundamental QFT:
 $\int_{p>k} \mathcal{D}\varphi e^{-S}$
couplings stay finite
for $k \rightarrow \infty$
scale-free!



Fundamental theory:

Running dimensionless couplings approach fixed point towards UV

Predictive theory:

Finite number of relevant (UV-attractive) couplings:

$$g_i(k) = g_* + c \left(\frac{k}{k_0} \right)^{-\theta_i} \quad \theta_i = d_i + \eta_i$$

Asymptotically Safe Quantum Gravity: Evidence

$$\Gamma_{k \text{ EH}} = \frac{-1}{16\pi G_N(k)} \int \sqrt{g} (R - 2\bar{\lambda}(k))$$

$$G = G_N k^2 \text{ and } \lambda = \bar{\lambda}/k^2$$

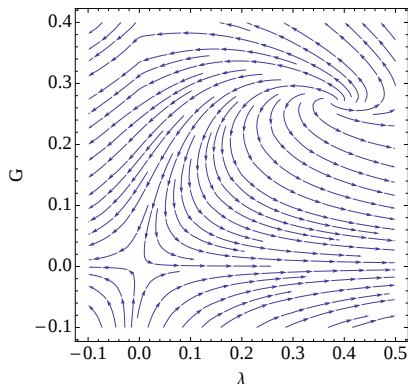
fixed point in dimensionless couplings $\rightarrow$ scale-free regime

Asymptotically Safe Quantum Gravity: Evidence

$$\Gamma_{k\text{EH}} = \frac{-1}{16\pi G_N(k)} \int \sqrt{g}(R - 2\bar{\lambda}(k))$$

$$G = G_N k^2 \text{ and } \lambda = \bar{\lambda}/k^2$$

fixed point in dimensionless couplings $\rightarrow$ scale-free regime



Compatibility with observations:
 G and λ relevant directions

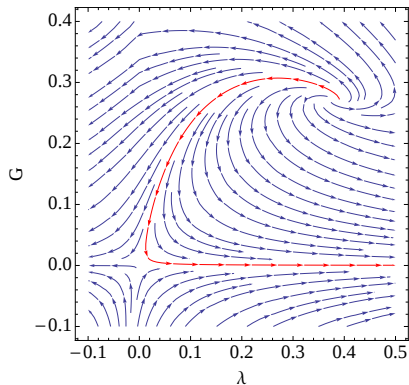
Semiclassical gravity?

Asymptotically Safe Quantum Gravity: Evidence

$$\Gamma_{k\text{EH}} = \frac{-1}{16\pi G_N(k)} \int (R - 2\bar{\lambda}(k))$$

$$G = G_N k^2 \text{ and } \lambda = \bar{\lambda}/k^2$$

fixed point in dimensionless couplings $\rightarrow$ scale-free regime

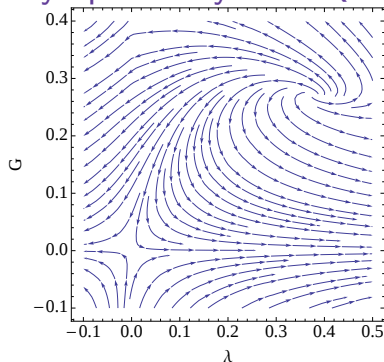


Compatibility with observations:
 G and λ relevant directions

Semiclassical gravity?

trajectory with $G_N \rightarrow \text{const}$
and $\bar{\lambda} \rightarrow \text{const}$ and measured
values in infrared

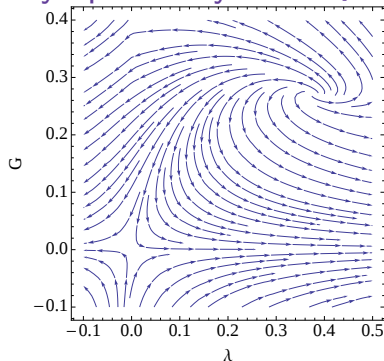
Asymptotically Safe Quantum Gravity: Evidence



$$\Gamma_{k\text{EH}} = \frac{-1}{16\pi G_N(k)} \int \sqrt{g} (R - 2\bar{\lambda}(k))$$

fixed-point action: *prediction*

Asymptotically Safe Quantum Gravity: Evidence



$$\Gamma_{k\text{EH}} = \frac{-1}{16\pi G_N(k)} \int \sqrt{g} (R - 2\bar{\lambda}(k))$$

fixed-point action: *prediction*

$$\Gamma_k = \Gamma_{k\text{EH}} + \Gamma_{\text{gauge-fixing}} + \Gamma_{\text{ghost}} + \int \sqrt{g} (f(R) + R_{\mu\nu} R^{\mu\nu} + \dots)$$

E. Manrique, M. Reuter, F. Saueressig (2009, 2010);

I. Donkin, J. Pawłowski (2012);

A. Codello, G. D'Odorico, C. Pagani (2013)

A.E., H. Gies, M. Scherer (2009), A.E., H. Gies (2010),

A.E. (2013)

A. Codello, R. Percacci, C. Rahmede (2008);

D. Benedetti, F. Caravelli (2012);

K. Falls, D. Litim, K. Nikolakopoulos (2013);

J. Dietz, T. Morris (2013);

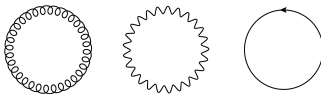
M. Demmel, F. Saueressig, O. Zanusso (2014)

D. Benedetti, P. Machado, F. Saueressig (2009)

Does matter matter?



quantum fluctuations of all fields drive Renormalization Group flow:



Quantum gravity and matter

Quantum gravity and matter

Quantum gravity will have an effect on matter:

$$\beta_{g_{\text{matter}}} = f(g_{\text{matter}}) + c_1 G g_{\text{matter}} + c_2 G^2 + \dots$$

Quantum gravity and matter

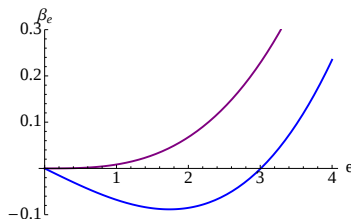
Quantum gravity will have an effect on matter:

$$\beta_{g_{\text{matter}}} = f(g_{\text{matter}}) + c_1 G g_{\text{matter}} + c_2 G^2 + \dots$$

- possible UV fixed point for Higgs sector [Zanusso, Zambelli, Vacca, Percacci (2009); A.E. (2012)]
- possible UV fixed point for QED [Harst, Reuter (2011); Doná, A.E., Percacci (2013)]
(cf. [Robinson, Wilczek (2005)])

$$\beta_e = \frac{e^3}{12\pi^2} N_D - \frac{3}{8\pi} e (G - 4G\lambda^2) + \dots$$

Can predict e at low energies?



Quantum gravity and matter

Quantum gravity and matter

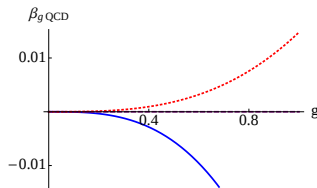
Matter will have an effect on quantum gravity

Quantum gravity and matter

Matter will have an effect on quantum gravity

analogy: Quantum Chromodynamics:

Asymptotic freedom only for $N_f < 16.5$



Quantum gravity and matter

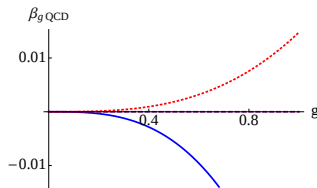
Matter will have an effect on quantum gravity

analogy: Quantum Chromodynamics:

Asymptotic freedom only for $N_f < 16.5$

Possibility to test quantum gravity models!

Is quantum gravity compatible with Standard Model matter?



Matter matters

with P. Doná, R. Percacci (2013): Truncation of the effective action:

$$\Gamma_k = \Gamma_{k \text{ Einstein-Hilbert}} + \Gamma_{k \text{ matter}} \quad \text{with minimally coupled matter:}$$

Matter matters

with P. Doná, R. Percacci (2013): Truncation of the effective action:

$\Gamma_k = \Gamma_{k \text{ Einstein-Hilbert}} + \Gamma_{k \text{ matter}}$ with minimally coupled matter:

$$N_S \text{ scalars: } S_S = \frac{Z_S}{2} \int d^d x \sqrt{g} g^{\mu\nu} \sum_{i=1}^{N_S} \partial_\mu \phi^i \partial_\nu \phi^i$$

$$N_D \text{ Dirac fermions } S_D = iZ_D \int d^d x \sqrt{g} \sum_{i=1}^{N_D} \bar{\psi}^i \not{D} \psi^i$$

N_V Abelian vector bosons:

$$S_V = \frac{Z_V}{4} \int d^d x \sqrt{g} \sum_{i=1}^{N_F} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa}^i F_{\nu\lambda}^i + \frac{Z_V}{2\xi} \int d^d x \sqrt{g} \sum_{i=1}^{N_F} (\bar{g}^{\mu\nu} \bar{D}_\mu A_\nu^i)^2$$

Matter matters

with P. Doná, R. Percacci (2013): Truncation of the effective action:

$\Gamma_k = \Gamma_{k \text{ Einstein-Hilbert}} + \Gamma_{k \text{ matter}}$ with minimally coupled matter:

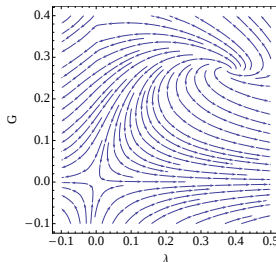
$$N_S \text{ scalars: } S_S = \frac{Z_S}{2} \int d^d x \sqrt{g} g^{\mu\nu} \sum_{i=1}^{N_S} \partial_\mu \phi^i \partial_\nu \phi^i$$

$$N_D \text{ Dirac fermions } S_D = iZ_D \int d^d x \sqrt{g} \sum_{i=1}^{N_D} \bar{\psi}^i \not{\nabla} \psi^i$$

N_V Abelian vector bosons:

$$S_V = \frac{Z_V}{4} \int d^d x \sqrt{g} \sum_{i=1}^{N_F} g^{\mu\nu} g^{\kappa\lambda} F_{\mu\kappa}^i F_{\nu\lambda}^i + \frac{Z_V}{2\xi} \int d^d x \sqrt{g} \sum_{i=1}^{N_F} (\bar{g}^{\mu\nu} \bar{D}_\mu A_\nu^i)^2$$

$$\rightarrow \beta_G, \beta_\lambda, \eta_h, \eta_c, \\ \eta_F, \eta_D, \eta_S$$

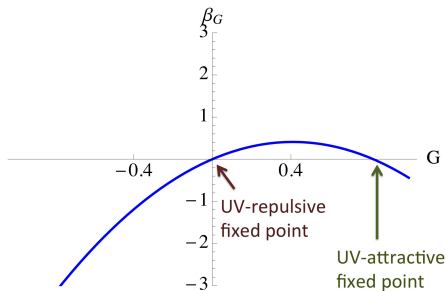


???

Perturbative analysis

(neglect graviton and matter wave function renormalizations)

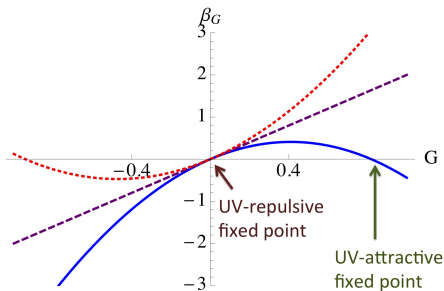
$$\beta_G = 2G + \frac{G^2}{6\pi} (-46),$$



Perturbative analysis

(neglect graviton and matter wave function renormalizations)

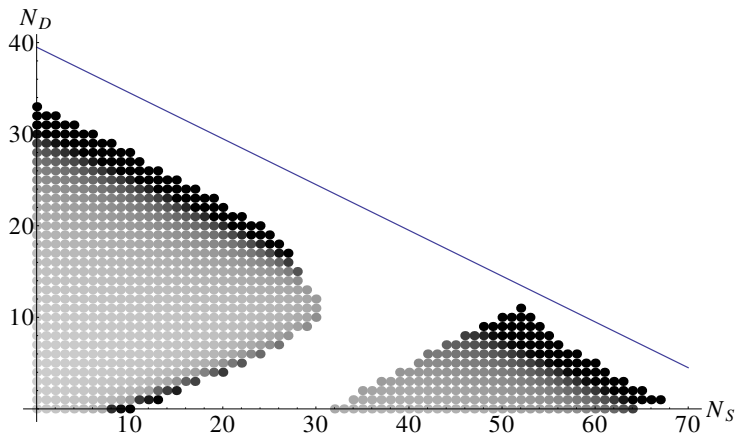
$$\beta_G = 2G + \frac{G^2}{6\pi} (N_S + 2N_D - 4N_V - 46),$$



→ for a given number of vectors N_V , there is an upper limit on the number of scalars N_S and Dirac fermions N_D !

Matter matters in asymptotically safe quantum gravity!

Full analysis for $N_V = 12$



upper limit on N_D and N_S

Standard Model: $N_V = 12$, $N_D = 45/2$, $N_S = 4$: compatible with gravitational fixed point

Specific matter models

Standard Model: ($N_S = 4$, $N_D = 45/2$, $N_V = 12$) ✓

Specific matter models

Standard Model: ($N_S = 4$, $N_D = 45/2$, $N_V = 12$) ✓

→ right-handed neutrinos? ✓

→ dark matter scalar? ✓

→ axion? ✓

Specific matter models

Standard Model: ($N_S = 4$, $N_D = 45/2$, $N_V = 12$) ✓

→ right-handed neutrinos? ✓

→ dark matter scalar? ✓

→ axion? ✓

supersymmetric extension (MSSM: $N_S = 49$, $N_D = 61/2$, $N_V = 12$) ✗

Specific matter models

Standard Model: ($N_S = 4$, $N_D = 45/2$, $N_V = 12$) ✓

→ right-handed neutrinos? ✓

→ dark matter scalar? ✓

→ axion? ✓

supersymmetric extension (MSSM: $N_S = 49$, $N_D = 61/2$, $N_V = 12$) ✗

GUT (SO(10): $N_S = 97$, $N_D = 24$, $N_V = 45$) ✗

Specific matter models

Standard Model: ($N_S = 4$, $N_D = 45/2$, $N_V = 12$) ✓

→ right-handed neutrinos? ✓

→ dark matter scalar? ✓

→ axion? ✓

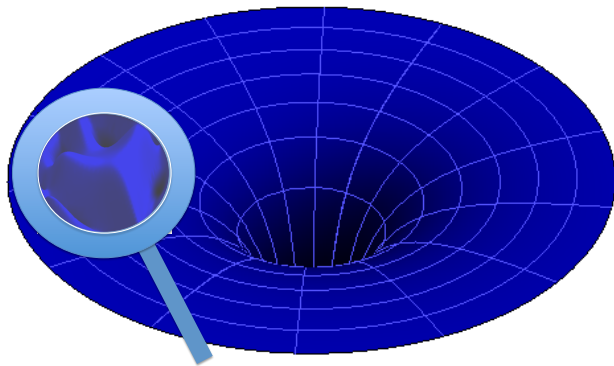
supersymmetric extension (MSSM: $N_S = 49$, $N_D = 61/2$, $N_V = 12$) ✗

GUT (SO(10): $N_S = 97$, $N_D = 24$, $N_V = 45$) ✗

Prediction:

Only specific models with restricted matter content are compatible with Asymptotically Safe Quantum Gravity within our truncation!

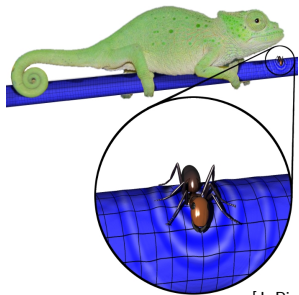
Tests of quantum gravity



Does testing quantum gravity require galaxy-size accelerators?

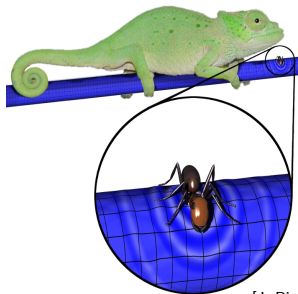
Possibly could test Asymptotically Safe Quantum Gravity at LHC, 14 TeV:
Look for Beyond-Standard-Model particle physics

Extra dimensions



[J. Pivarski]

Extra dimensions



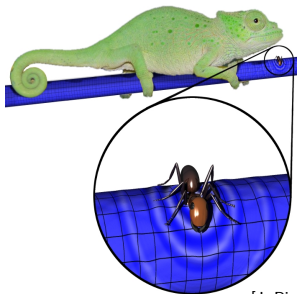
[J. Pivarski]

Extra dimensions in asymptotic safety?

pure-gravity fixed point exists in $d \geq 4$

(Einstein-Hilbert [Fischer, Litim, 2006] and higher derivatives [Ohta, Percacci, 2013])

Extra dimensions

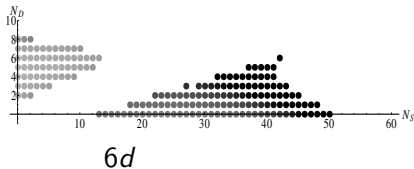
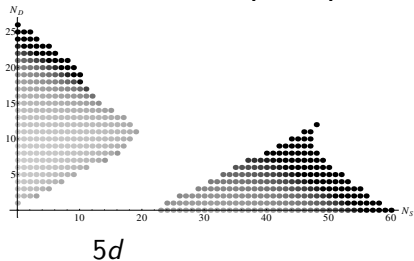


[J. Pivarski]

Extra dimensions in asymptotic safety?

pure-gravity fixed point exists in $d \geq 4$

(Einstein-Hilbert [Fischer, Litim, 2006] and higher derivatives [Ohta, Percacci, 2013])



Asymptotic safety: Summary

Continuum quantum field theory for gravity $\rightarrow$ interacting UV fixed point

$\Rightarrow$ nonperturbative tool: functional Renormalization Group

Asymptotic safety: Summary

Continuum quantum field theory for gravity $\rightarrow$ interacting UV fixed point

$\Rightarrow$ nonperturbative tool: functional Renormalization Group

Status:

- Pure gravity fixed point persists in extended truncations
- Coupling to Standard Model matter consistent (not all BSM-scenarios compatible!)

Asymptotic safety: Summary

Continuum quantum field theory for gravity $\rightarrow$ interacting UV fixed point

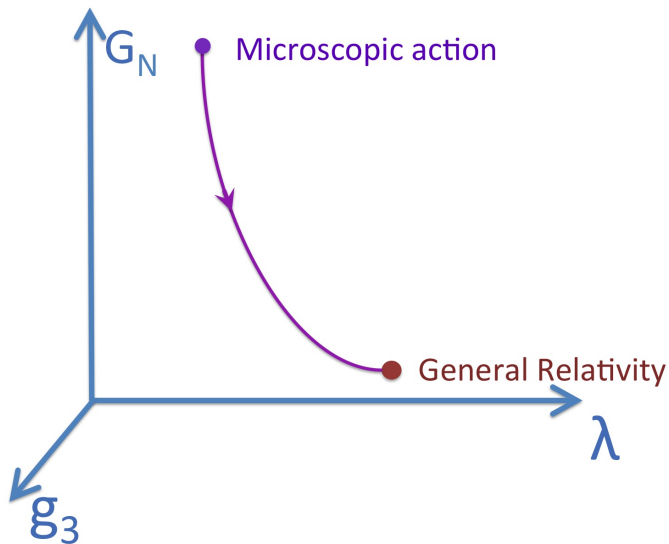
$\Rightarrow$ nonperturbative tool: functional Renormalization Group

Status:

- Pure gravity fixed point persists in extended truncations
- Coupling to Standard Model matter consistent (not all BSM-scenarios compatible!)

New test of quantum gravity at LHC and future accelerators could be possible!

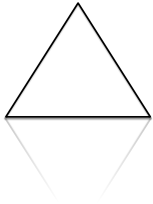
The Renormalization Group in quantum gravity



test continuum limit of discrete models (matrix/ tensor models)

Building a spacetime

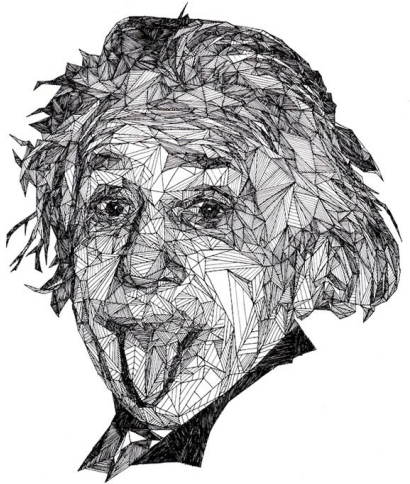
building blocks of spacetime:



2d



3d

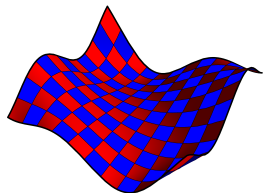


[by Josh Bryan]

Building a spacetime

$$\int \mathcal{D}g_{\mu\nu} e^{-S[g_{\mu\nu}]} \longrightarrow \sum_{\mathcal{T}} e^{-S[\mathcal{T}]}$$

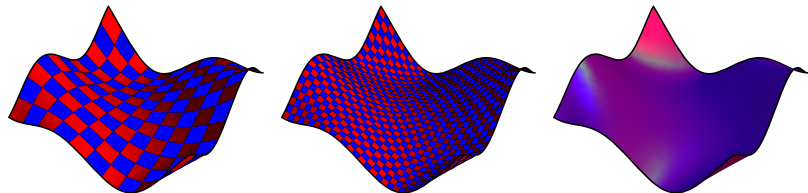
superposition of quantum spacetimes $\rightarrow$ sum over triangulations



Building a spacetime

$$\int \mathcal{D}g_{\mu\nu} e^{-S[g_{\mu\nu}]} \longrightarrow \sum_{\mathcal{T}} e^{-S[\mathcal{T}]}$$

superposition of quantum spacetimes $\rightarrow$ sum over triangulations



discretization: “unphysical” tool

Main challenge: What is the continuum limit?

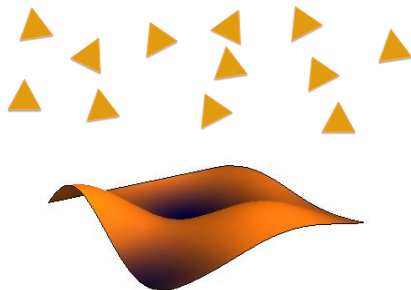
(e.g. Causal Dynamical Triangulations program [Ambjorn, Jurkiewicz, Loll],

Tensor Track [Bonzom, Ben Geloun, Freidel, Gurau, Oriti, Rivasseau])

Building a spacetime

Continuum limit from fixed point of Renormalization Group

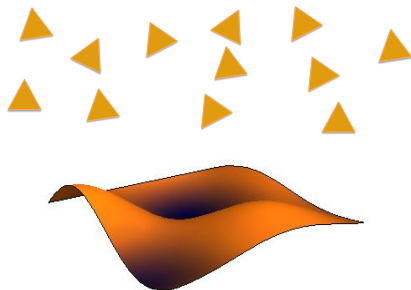
“condensation” of building blocks



Building a spacetime

Continuum limit from fixed point of Renormalization Group

“condensation” of building blocks



two-dimensional gravity: exact results known $\rightarrow$ testing ground for Functional RG methods!

Two-dimensional quantum gravity

gravity in two dimensions: $\frac{1}{4\pi} \int d^2x \sqrt{g} R = \chi \rightarrow$ action is topological

Euler character $\chi = 2 - 2h$ with h number of handles



quantum gravity nontrivial: $Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-\beta A + \gamma \chi},$

Two-dimensional quantum gravity

gravity in two dimensions: $\frac{1}{4\pi} \int d^2x \sqrt{g} R = \chi \rightarrow$ action is topological

Euler character $\chi = 2 - 2h$ with h number of handles



quantum gravity nontrivial: $Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-\beta A + \gamma \chi},$

How to evaluate partition function? $\rightarrow$ use triangulation!

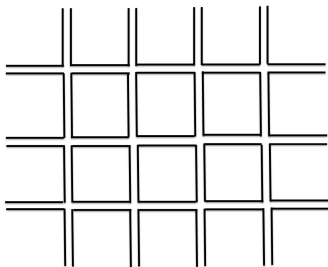
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\mathcal{T}} e^{-S[\mathcal{T}]},$$

reformulate as hermitian matrix model: ϕ_{ij} $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr} \phi^2 + g \text{tr} \phi^4} \right)$$

Feynman diagrams:



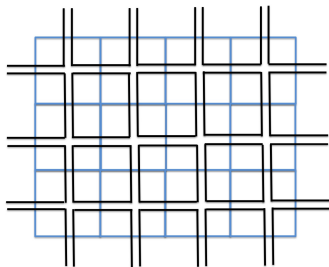
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\mathcal{T}} e^{-S[\mathcal{T}]},$$

reformulate as hermitian matrix model: ϕ_{ij} $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr} \phi^2 + g \text{tr} \phi^4} \right)$$

Feynman diagrams:
dual is “squarulation”



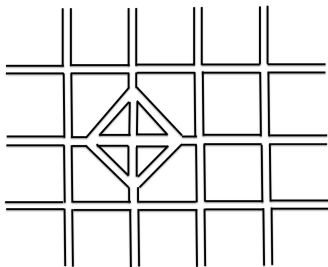
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\mathcal{T}} e^{-S[\mathcal{T}]},$$

reformulate as hermitian matrix model: ϕ_{ij} $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 + g\text{tr}\phi^4} \right)$$

Feynman diagrams:
dual is “squarulation”



Matrix model for two-dimensional quantum gravity

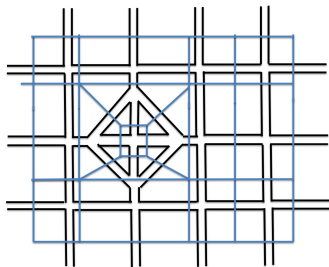
$$Z_{grav} \sim \sum_{\mathcal{T}} e^{-S[\mathcal{T}]},$$

reformulate as hermitian matrix model: ϕ_{ij} $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr} \phi^2 + g \text{tr} \phi^4} \right)$$

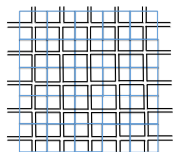
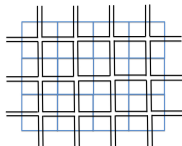
Feynman diagrams:
dual is “squarulation”

Feynman diagram expansion
yields sum over all “squarulations”



Matrix model for two-dimensional quantum gravity

continuum limit:
$$N^2 Z_{grav} = \ln \left(\int d\phi e^{-\text{tr}\phi^2 + g \text{tr}\phi^4} \right)$$

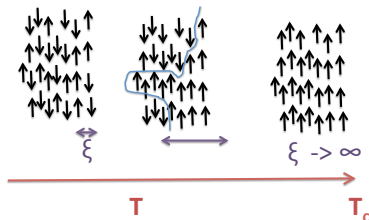


$$g \rightarrow g_c$$

$$\langle A \rangle \rightarrow \infty$$

(cf. phase transition)

→ universal behavior!



Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h=0 \text{ dominates}$$

Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h=0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_{str}/2} = \text{const.}, \quad \gamma_{str} = -1/2.$$



Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h=0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_{str}/2} = \text{const.}, \quad \gamma_{str} = -1/2.$$



$$\text{here: } g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_{str})}$$

Large N limit

$$Z_{\text{grav}} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h=0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_{\text{str}}/2} = \text{const.}, \quad \gamma_{\text{str}} = -1/2.$$



$\rightarrow$ looks like fixed-point behavior: $g(k) = g_* + c \left(\frac{k}{k_0} \right)^{-\theta}$

$$\text{here: } g(N) = g_c + \left(\frac{N}{c} \right)^{-2/(2-\gamma_{\text{str}})}$$

Large N limit

$$Z_{\text{grav}} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_{\text{str}}/2} = \text{const.}, \quad \gamma_{\text{str}} = -1/2.$$



$\rightarrow$ looks like fixed-point behavior: $g(k) = g_* + c \left(\frac{k}{k_0} \right)^{-\theta}$

$$\text{here: } g(N) = g_c + \left(\frac{N}{c} \right)^{-2/(2-\gamma_{\text{str}})}$$

double-scaling limit = interacting fixed point

$k \leftrightarrow N$ [Brézin, Zinn-Justin, 1992]

Large N limit

$$Z_{\text{grav}} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_{\text{str}}/2} = \text{const.}, \quad \gamma_{\text{str}} = -1/2.$$



$\rightarrow$ looks like fixed-point behavior: $g(k) = g_* + c \left(\frac{k}{k_0} \right)^{-\theta}$

$$\text{here: } g(N) = g_c + \left(\frac{N}{c} \right)^{-2/(2-\gamma_{\text{str}})}$$

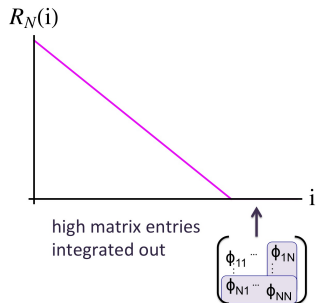
double-scaling limit = interacting fixed point $k \leftrightarrow N$ [Brézin, Zinn-Justin, 1992]

main advantage: can extend to $d=4$ tensor models!

Renormalization Group in matrix models

$$N\partial_N\Gamma_N = \frac{1}{2}\text{tr}(\Gamma_N^2 + R_N)^{-1} N\partial_N R_N$$

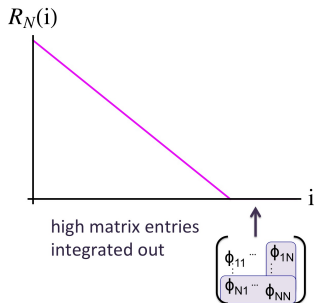
A.E., T. Koslowski PRD 88, 084016 (2013))



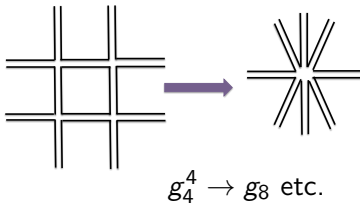
Renormalization Group in matrix models

$$N\partial_N\Gamma_N = \frac{1}{2}\text{tr}(\Gamma_N^2 + R_N)^{-1} N\partial_N R_N$$

A.E., T. Koslowski PRD 88, 084016 (2013))



theory space:



→ all random polygonizations generated! $\Gamma_N = \frac{Z}{2}\text{tr}\phi^2 + \sum_i \frac{g_i}{i}\text{tr}\phi^i + \dots$

Results: Double scaling limit

Truncation: $Z, g_4, g_6, \dots, g_{14}$ [A.E., T. Koslowski, PRD 88, 084016 (2013)]

fixed point with one relevant direction ($g \rightarrow g_c$ tuning)

RG works qualitatively in matrix models!

Results: Double scaling limit

Truncation: $Z, g_4, g_6, \dots, g_{14}$ [A.E., T. Koslowski, PRD 88, 084016 (2013)]

fixed point with one relevant direction ($g \rightarrow g_c$ tuning)

RG works qualitatively in matrix models!

critical exponent

$\theta \approx 1+$ (cf. exact value $\theta = 4/5$)

benchmark: perturbative RG (explicit integration)

by Brézin, Zinn-Justin: $\theta = 1$ (1992)

Results: Double scaling limit

Truncation: $Z, g_4, g_6, \dots, g_{14}$ [A.E., T. Koslowski, PRD 88, 084016 (2013)]

fixed point with one relevant direction ($g \rightarrow g_c$ tuning)

RG works qualitatively in matrix models!

critical exponent

$\theta \approx 1+$ (cf. exact value $\theta = 4/5$)

benchmark: perturbative RG (explicit integration)

by Brézin, Zinn-Justin: $\theta = 1$ (1992)

What precision to expect? Wilson-Fisher fixed point in scalar ϕ^4 in $d = 3$:

leading order approximation $\sim 20\%$

optimization tools, extended truncations give quantitative precision

→ route to quantitative precision in matrix models clear

Summary and Outlook: RG in matrix models

Matrix models: Gravitational path-integral from continuum limit in discretized setting

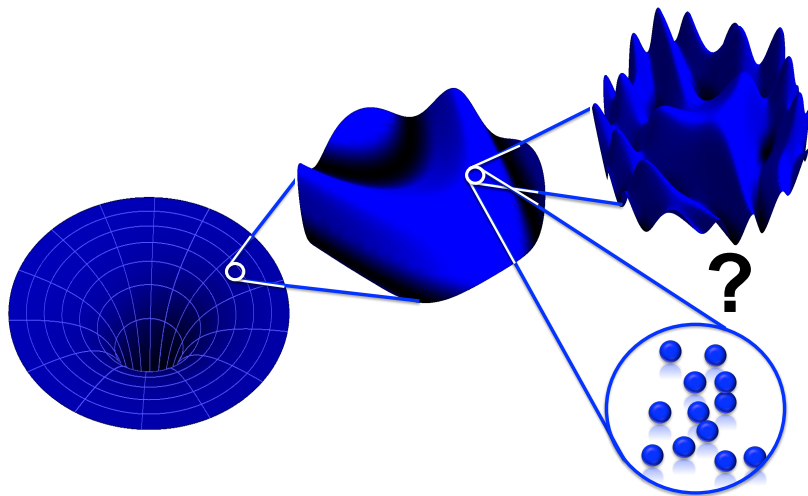
Double scaling limit accessible as fixed point of RG

Qualitative results ✓

Outlook: Development of optimization methods, extended truncations → towards quantitative precision!

Extension to higher dimensions $d = 4$
("recipe" in PRD 88, 084016 (2013))

The Renormalization Group in quantum gravity



Thank you for your attention!