

## Shift Symmetry

What is the Origin of the Shift. Sym.?

Tatsu Yanagida

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## SHIFT SYMMETRY:

K. Y. Y. (2000)

Inflaton  $\phi \rightarrow \phi + iC$

$$K = \frac{1}{2} (\phi + \phi^*)^2 + g^i S$$

$$W = m S \phi$$



$$V = \frac{1}{2} m^2 |\phi|^2$$

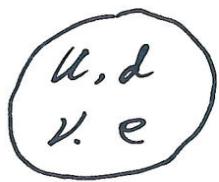
Chaotic Inflation

Linde (1983)

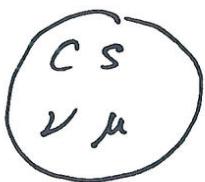
Why we have the shift symm. ?

Motivated by Particle Physics.

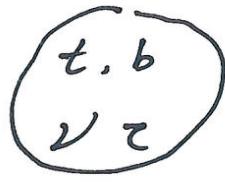
Why do we have three families?



$$\underline{5^*+10}$$



$$\underline{5^*+10}$$



$$\underline{5^*+10}$$

of  $SU(5)_{GUT}$

$$m_{u, d, e} \ll m_{c, s, \mu} \ll m_{t, b, \tau} \quad ?$$

c.f. QCD

$$m_\pi^2 \ll m_{K, \eta}^2 \ll m_N^2$$

$\pi, K, \eta$  are Nambu-Goldstone bosons !!

Quarks/Leptons are Nambu-Goldstone  
chiral multiplets ?

Bachmäder, Peccei, Yanazida (1982)

$$\Phi(x, \theta) = \varphi(x) + \psi(x) \cdot \theta$$

↓  
 Quarks/Leptons

Given  $G/H$ , we can determine # of families.

But,

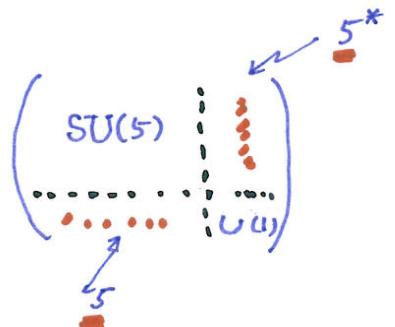
What is  $G/H$  ?

$G \rightarrow H : NG \overset{?}{\Phi} \leftarrow \text{the Broken Generators}$

### Simple Exercise

- $SU(6)/SU(5) \times U(1)$ : Broken generator is  $X(5) + X^*(5)$

$$\therefore NG \Phi(x) \simeq \underline{5}^*$$



- $SO(10)/SU(5) \times U(1)$

$$SO(10) = \underbrace{24}_{SU(5)} + \underbrace{10 + 10^*}_{U(1)} + 1$$

Broken generator ( $10 + 10^*$ )

NG chiral multiplet  $\Phi = \underline{10}$

We need  $G/H$  which contains

Both of  $\underline{5}^* + \underline{10}$ .



$E_6$  ?

E<sub>6</sub>

$$\text{Generators } 78 = 45 + \underbrace{16 + 16^*}_{\text{32 broken generators}} + 1$$

↙ SO(10)      ↘ U(1)

$E_6/SO(10) \times U(1)$  :

NG chiral multiplet  $\Phi(\underline{16})$

$$\underline{16} \text{ of } SO_{10} = \underline{5^*} + \underline{10} + 1$$

↑  
right-handed 2

Exceptional Group is also interesting, since  
it has the maximal group  $E_8$  !!!



We may have the prediction of  
# of families .

$E_7 / SU(5) \times U(1)^3$ :

$$NG \text{ } \Xi = \underbrace{3 \times (5^* + 10)}_{\text{three families}} + 5 + 1's$$

Kugo, Yanagida (1983)

$E_8 / SU(5) \times U(1)^4$

$$\begin{aligned} NG \text{ } \Xi = & \underbrace{3 \times (5^* + 10)}_{\text{Vector-like matters}} + 5 + 1's \\ & + \underbrace{(5^* + 10) + (5 + 10^*)}_{\text{Vector-like matters}} \end{aligned}$$

$N(\#)$  of Families  $\leq \underline{3}$  !!!

$E_7 / SU(5) \times U(1)^3$  :

Nonlinear Sigma Model Anomaly

But it is canceled by introducing a matter multiplet  $\Psi(\underline{5}^*)$ .

Yasui (1986)

$$N6 \oplus = 3 \times (\underline{5}^* + \underline{10}) + \underline{1}'_s + \underline{5} \oplus \Psi(\underline{5}^*)$$

If  $E_7$  is exact, all Yukawa couplings and Masses are zero !!!

Let introduce explicit breaking:

$E_7 \xrightarrow{\text{e.b.}} E_6$  : breaking parameter.

- $\underline{5} + \Psi(\underline{5}^*)$  has a mass term :  $M \underline{5} \cdot \Psi(\underline{5}^*)$
- $t, b, \tau$  get Yukawa Couplings

$$m_t, m_b, m_\tau \neq 0$$

$E_6 \xrightarrow{\text{e.b.}} SO(10)$

- $c, s, \mu$  get Yukawa Couplings

$$m_c, m_s, m_\mu \neq 0$$

$SO(10) \xrightarrow{\text{e.b.}} SU(5)$  :

- $u, d, e$  get Yukawa Couplings

$$m_u, m_d, m_e \neq 0$$

Mass hierarchy :

$$m_{t, b, \tau} \gg m_{c, s, \mu} \gg m_{u, d, e}$$

C.f. QCD

$$\frac{SU(3)_L \times SU(3)_R}{SU(3)_V} : \pi, K, \eta \quad \text{massless NG bosons}$$

Explicit Breaking parameter  $m_S \neq 0$

$$SU(3)_L \times SU(3)_R \xrightarrow{\text{e.b.}} SU(2)_L \times SU(2)_R$$

- $K, \eta$  get masses.

Explicit Breaking  $m_{u, d} \neq 0$

$$SU(2)_L \times SU(2)_R \xrightarrow{\text{e.b.}} SU(1)_V$$

- $\pi$  get masses.

$$m_{K, \eta}^2 \gg m_\pi^2$$

Kähler Potential is given by  $G/H$  geometry.

$$K(\bar{\varphi}_i, \bar{\varphi}_i^*) = \underbrace{G/H}_{\text{Kähler manifold.}} \text{ Kähler Potential}$$

$E_7$  transformation

$$K(\bar{\varphi}_i, \bar{\varphi}_i^*) \rightarrow K(\bar{\varphi}_i, \bar{\varphi}_i^*) + H(\bar{\varphi}_i) + H^*(\bar{\varphi}_i^*)$$

$$\delta \int d^4\theta K(\bar{\varphi}_i, \bar{\varphi}_i^*) - \int d^4\theta [H(\bar{\varphi}) + H(\bar{\varphi}^*)] = 0$$

$K(\bar{\varphi}, \bar{\varphi}^*)$  is invariant of  $E_7$ .

But.

A Big Problem in Supergravity !!!

In Supergravity :  $\int d^4\theta K(\bar{\Xi}, \bar{\Xi}^*)$

$$\delta K = \int d^4\theta [H(\bar{\Xi}) + H^*(\bar{\Xi}^*)] \neq 0$$

$E_7$  invariance is maximally broken !!!

Witten, Bagger (1983)

A solution was found !!

Kugo, Yanagida (2010)  
Komargostri, Seiberg (2010)

$$K(\bar{\Xi}, \bar{\Xi}^*)_{\text{new}} = K(\bar{\Xi}, \bar{\Xi}^*) + Z + Z^*$$

$$\left\{ f_8 = M_{Pl} = 1 \right\}$$

$$\left\{ \begin{array}{l} K(\bar{\Xi}, \bar{\Xi}^*) \xrightarrow{E_7} K(\bar{\Xi}, \bar{\Xi}^*) + H(\bar{\Xi}) + H^*(\bar{\Xi}^*) \\ Z \xrightarrow{E_7} Z - H(\bar{\Xi}) \end{array} \right.$$

$K(\bar{\Xi}, \bar{\Xi}^*)_{\text{new}}$  is completely invariant.

The new field  $Z$  has a shift symmetry.

$$Z \rightarrow Z + iC !!$$

$Z$  can be the Inflaton  $\phi$ .

A simple exercise:

Consider  $CP^1 \simeq SU(2)/U(1)$  : Kähler Manifold.

$$\begin{cases} \text{unbroken generator } T \\ \text{broken generator } X, X^+ \end{cases}$$

NG multiplet  $\Xi$ :

$$\begin{cases} [T, \Xi] = +i\Xi \\ [X, \Xi] = \Xi^2 \\ [X^+, \Xi] = 1 \end{cases} \quad \begin{matrix} \text{nonlinear realization} \\ \text{of } SU(2) \end{matrix}$$

Kähler Potential

$$K = \log(1 + \Xi\Xi^*)$$

$$\delta K = [X, K] = \frac{\Xi^2\Xi^* + \Xi \cdot 1}{1 + \Xi\Xi^*} = \Xi$$

$$K \rightarrow K + \Xi$$

Introduce  $Z$  to cancel  $\Xi$  term.

$$\delta Z = [X, Z] = -\Xi, \quad [X^+, Z] = 0$$

$$[X, K + Z + Z^*] = 0$$

"  
Knew

$$\delta K_{\text{new}} = 0 !$$

c.f. Witten-Bagger Theory

$$\left\{ \begin{array}{l} K = f_\phi^2 \left( \log \left( 1 + \frac{\phi \phi^*}{f_\phi^2} \right) \right) \\ W = 0 ! \end{array} \right.$$

If  $f_\phi^2 = 2n M_{pl}^2$ , the theory is consistent.

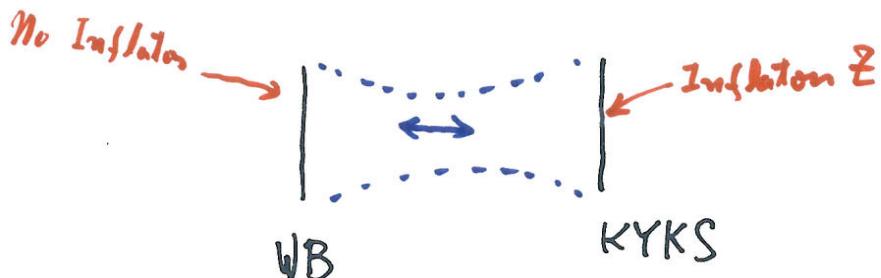
$f_\phi \gg M_{pl}$  :  $n = 1, 2, 3 \dots$

May be consistent with  $\Sigma_0 \gg M_{pl}$ .  
the large initial value  
for the Chaotic Inflation.

KYKS Model :

$$\left\{ \begin{array}{l} K = F(K_{new}) = F(f_\phi^2 \log \left( 1 + \frac{\phi \phi^*}{f_\phi^2} \right) + \frac{\Sigma + \Sigma^*}{f_\phi}) \\ W \neq 0 \end{array} \right.$$

$f_\phi$  is free.



$f_\phi \gg M_{pl}$

Hollerman, Kehagias, Yano, jso  
(2014)

## Explicit Breaking of $E_7$

$$E_7 \xrightarrow{\text{ob.}} E_6 \xrightarrow{\text{ob.}} E_5(SO(10)) \xrightarrow{\text{ob.}} E_8(SU(5))$$

$SO(10)$  must be a good symmetry.

Let consider  $SO(10)/SU(5) \times U(1)$  manifold.

N.G. multiplet  $\Phi = \underline{10}$

The up quark mass is given by

$$W = y_u \cdot \underline{10} \cdot \underline{10} \cdot \langle H \rangle$$

If  $SO(10)$  is exact symmetry,  $y_u = 0 \rightarrow m_u = 0$ .

$y_u \simeq SO(10)$  explicit breaking parameter

$$\simeq \underline{10^{-5}} \quad \text{for } m_u \simeq \text{a few MeV.}$$

Now,

$$\leftarrow G/H = SO(10)/SU(5) \times U(1)$$

$$K_{\text{new}} = K(\phi, \phi^*) + \Sigma + \Sigma^*$$

Introduce the explicit breaking of  $SO(10)$

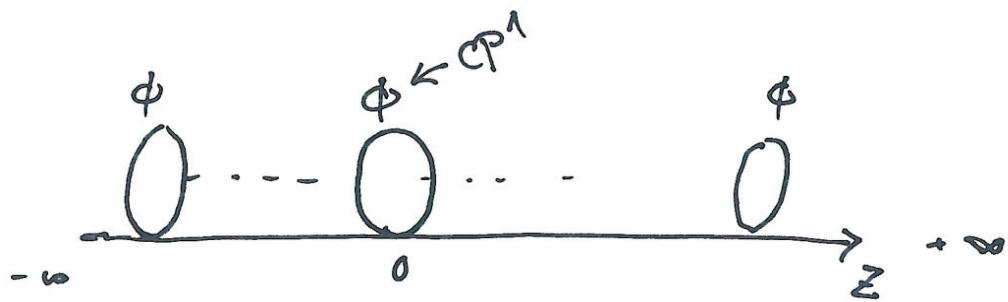
$$W = y_u \underline{10} \cdot \underline{10} \cdot \langle H \rangle + m S \Sigma$$

$$\nearrow 10^{-5}$$

$$\left\{ 10^{-5} \simeq 10^{13} \text{ GeV} !!! \right.$$

Good for Inflaton mass.

Now Our Manifold is Non compact !!!



In  $\Sigma$  can be  $-\infty \longleftrightarrow +\infty$ .

$$|\Sigma_0| \gg M_{pl}$$

We may live naturally the large initial value for the Inflaton  $\Sigma$ .

What is the Prediction ?

$$\left\{ \begin{array}{l} m_{3/2} \simeq m_{\text{sgen.}} \simeq 100 \text{ TeV} \\ m_{\chi} \simeq 125 \text{ GeV} \end{array} \right. !!$$

Consider  $\mathbb{C}\mathbb{P}^1$  as an example :

$$\begin{aligned}
 K &= K(\log(1+\phi^*\phi) + Z^*Z) \\
 &= \log(1+\phi^*\phi) + Z + Z^* \\
 &\quad + \frac{1}{2} [\log(1+\phi^*\phi) + Z + Z^*]^2 + \dots
 \end{aligned}$$

needed for the kinetic term of  $\phi$   
 needed for the kinetic term of  $Z$

$$\text{Kinflation} = \underline{Z + Z^*} + \frac{1}{2} (Z + Z^*)^2 + \dots$$


  
 $e^Z W$

$$\text{Kinflation} = \frac{1}{2} (Z + Z^*)^2$$

$$W = e^Z \bar{W}$$

$$N = g_{TR} g_L H + \dots + \frac{1}{2} M N N - \dots$$

$$Z \xrightarrow{\text{decay}} N + N : g t H : \dots$$

$$T_R \approx 10^9 - 10^{10} \text{ GeV}$$

Good for thermal leptogenesis.

Let consider Dynamical SUSY Breaking.

Dynamical Scale  $\Lambda$

Composite hadrons mass  $\sim \Lambda$

$$\text{SUSY breaking } f \simeq \alpha \cdot \frac{\Lambda^2}{4\pi} \quad \alpha \simeq 1 \sim 0.1$$

$$W = W_{\text{standard model}} + W_{\text{hidden}}$$

$$W_{\text{hidden}} = M_{\text{composite}} \cdot \Psi \bar{\Psi} + \dots$$

$$M_{\text{composite}} \simeq \Lambda$$

$$\text{If } \Lambda \lesssim \frac{1}{2} m_{\text{inflaton}} \simeq \frac{1}{2} \cdot 10^{13} \text{ GeV},$$

The inflaton decays mainly to hidden  
hadrons.



$$\text{If } \Lambda \lesssim \frac{1}{2} M_{\text{inflation}} \approx \frac{1}{2} \cdot 10^{13} \text{ GeV},$$

The inflation decay produces too many LSP !!

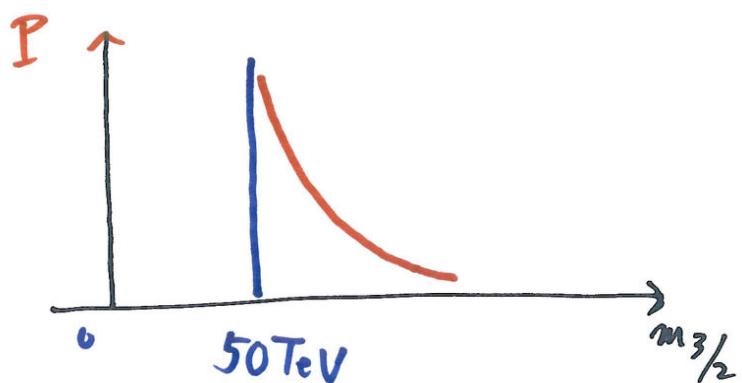
$$\Sigma_{\text{LSP}} \gg \Sigma_{\text{observation}}$$



$$\Lambda \gtrsim \frac{1}{2} \cdot 10^{13} \text{ GeV}$$

$$F \gtrsim 0.25 \times 10^{25} \text{ GeV} \times (\alpha)$$

$$m_{3/2} \gtrsim \alpha \cdot 500 \text{ TeV} \quad \alpha(1 - 0.1)$$



Hariyaya, T.T.Y.  
(2014)

If SUSY breaking is low-energy biased,

$$m_{3/2} \simeq 100 \text{ TeV} !!!$$

$m_{\chi}$  is predicted as  $m_{\chi} \simeq \underline{125 \text{ GeV}}$ .

Okada, Yamagishi, Yanagisawa  
(1992).

## Conclusions

We have a very serious problem in Particle Physics.

$\left\{ \begin{array}{l} \text{LHC found Higgs boson of mass } \approx 125 \text{ GeV.} \\ \text{LHC found no SUSY particle } \lesssim 1 \text{ TeV} \end{array} \right.$

These observations suggest the high scale SUSY  
 $m_0 \gtrsim 10 \text{ TeV} !!!$

Why so high?

If SUSY is the solution to the  
naturalness problem, we expect  
 $m_0 \lesssim 1 \text{ TeV}$ .

Without a convincing answer to the question  
I can not believe SUSY.

Now, the Chaotic Inflation gives us an answer !

$$\Lambda_{\text{dyn}} \geq \frac{1}{2} M_{\text{inflaton}} \approx \frac{1}{2} \cdot 10^{13} \text{ GeV}$$



$$m_{3/2} \gtrsim 100 \text{ TeV}$$

If  $M_{\text{inf}} \approx 10^{12} \text{ GeV}$ , we get

$$m_{3/2} \approx 500 \text{ GeV}. \quad F \propto \Lambda^2 \approx m^2$$

$M_{\text{inf}} \approx 10^{13} \text{ GeV}$  is crucial !!!

The pure gravity mediation predicts

$$\left\{ \begin{array}{l} m_{\text{glino}} \gtrsim 3 \text{ TeV} \\ m_{\text{Wino}} \approx \text{a few TeV} \end{array} \right.$$

Ibe, Yanagisawa  
(2011)

Wino is the DM.

Ibe, Matsunaga, Yanagisawa  
(2012)

If R parity is broken, Wino can decay through

$$W_R = \lambda L \bar{L} E$$
$$\lambda \ll 1. \quad (\lambda \approx 10^{-19})$$

The wino decay explains the recent observation of the cosmic positron anomaly by AMS02-2014.

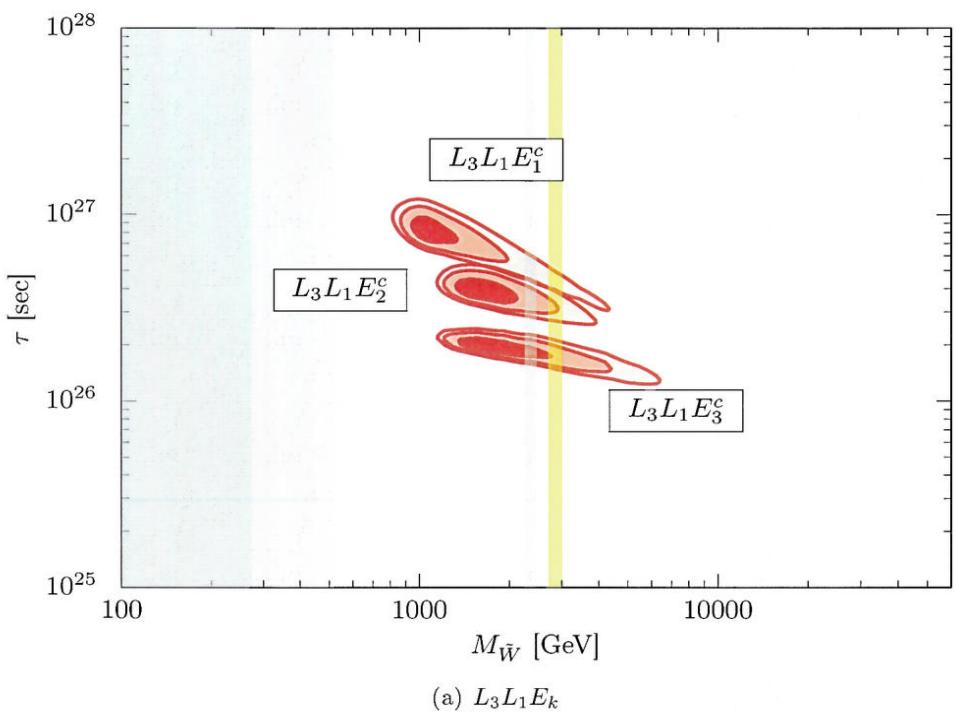
We derive the mass of the Wino :

$$\underline{m_W \simeq 1 - 3 \text{ TeV}} \quad !!!$$

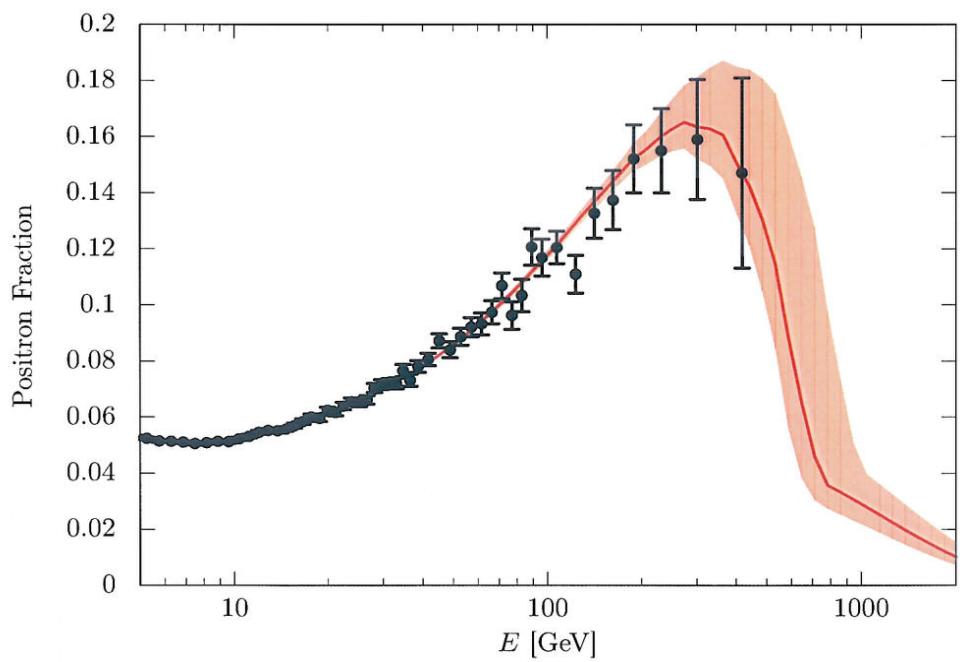
See figures

Ibe, Matsui, Shirai  
(now)

Consistent with the prediction of  
the "Pare gravity mediation".



(a)  $L_3 L_1 E_k$



(b)  $L_3 L_1 E_2$  positron fraction