

Measurement of the top quark mass in topologies enhanced with single top-quarks produced in the t -channel at $\sqrt{s} = 8$ TeV using the ATLAS experiment

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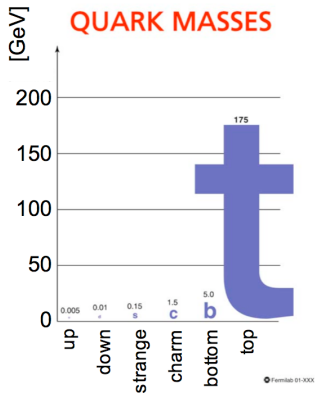


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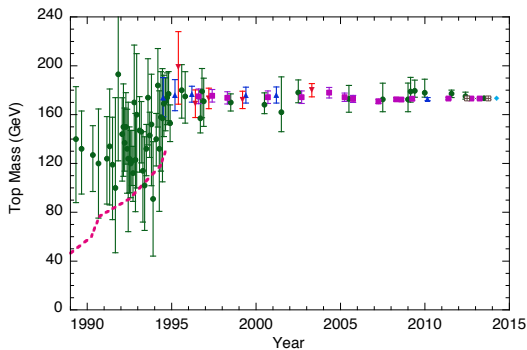


Measurement of the top quark mass in topologies enhanced with single top-quarks produced in the t -channel at $\sqrt{s} = 8$ TeV using the ATLAS experiment

1. The top quark and its mass
2. Basic event selection and neural network based event selection
3. Measurement of the top quark mass
 - More details in ATLAS-CONF-2014-055



[1]



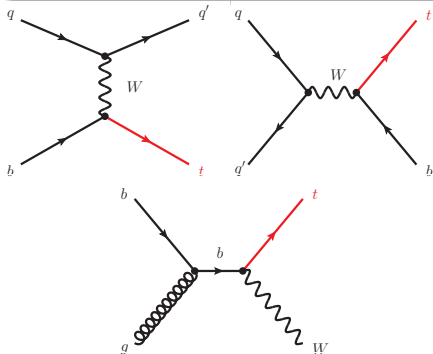
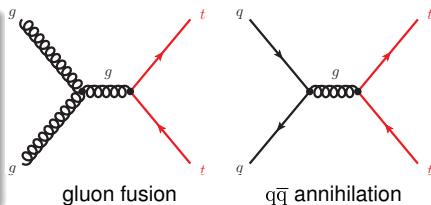
[2]

- The top quark has a much higher mass than all other known elementary particles
- First indirect mass measurements in the early 1990s
- Current world combination from 2014 yields [3]:

$$m_{\text{top}} = (173.34 \pm 0.27 (\text{stat.}) \pm 0.71 (\text{syst.})) \text{ GeV}$$

[1] http://www-d0.fnal.gov/Run2Physics/top/public/fall06/singletop/plain_english_summary.html[2] <http://lutece.fnal.gov/TTS/> [3] ATLAS-CONF-2014-008, CDF-NOTE-11071, CMS-PAS-TOP-13-014, D0-NOTE-6416

- Discovery in $t\bar{t}$ pair production in 1995 at the Tevatron
- $t\bar{t}$ pair production is always mediated by the strong interaction
- All previous mass measurements performed in $t\bar{t}$ final states



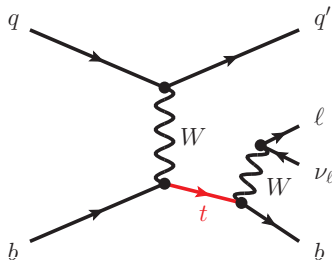
- First observation of single top quarks in 2009 in t -s-channel at the Tevatron
- Evidence of associated production with a W -boson observed in 2012 at the LHC
- Single top quarks are always produced by the weak interaction

Measuring m_{top} in single-top t -channel topologies

Goal: First measurement of the top quark mass in topologies enhanced with single top-quarks produced in the t -channel

- Event topology is different and has never been studied to measure m_{top}
 - ▶ Production via the weak interaction at different typical energy scale
 - ▶ Different final state without ambiguities in the jet-parton assignment
 - ▶ Complementary sensitivity of systematic uncertainties
- Orthogonal phase-space

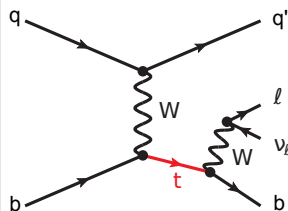
→ Very good prospects for future combinations



- Electron or muon with $p_T > 25$ GeV and $|\eta| < 2.5$
- Exactly two jets with $p_T > 30$ GeV and $|\eta| < 4.5$
 - ▶ $p_T > 35$ GeV (if $2.75 < |\eta^{\text{jet}}| < 3.5$)
- $E_T^{\text{miss}} > 30$ GeV
- $m_T(W) = \sqrt{2p_T^\ell E_T^{\text{miss}} [1 - \cos \Delta\varphi(\ell, E_T^{\text{miss}})]} > 50$ GeV

Additional cut reducing multijet background

- $p_T^\ell > 40$ GeV $\left(1 - \frac{\pi - |\Delta\varphi(\text{jet}_1, \ell)|}{\pi - 1}\right)$
- Reject leptons with low p_T and large opening angle with leading jet (e.g. from dijet production with fake lepton)

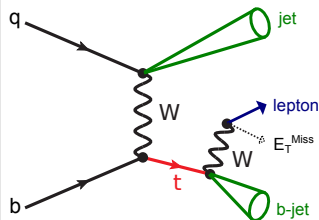


- Signal region:
 - ▶ Exactly one jet is tagged by MV1c at $\varepsilon_b^{t\bar{t}} = 50\%$
- Control region enriched with main background (W +jets)
 - ▶ Veto against signal region (No jet is tagged by MV1c at $\varepsilon_b^{t\bar{t}} = 50\%$)
 - ▶ Exactly one jet is tagged by MV1 at $\varepsilon_b^{t\bar{t}} = 80\%$

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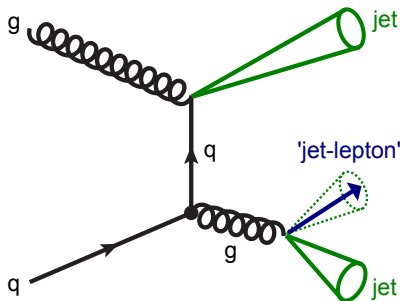


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- Two different models to mimic the shape of the multijet background

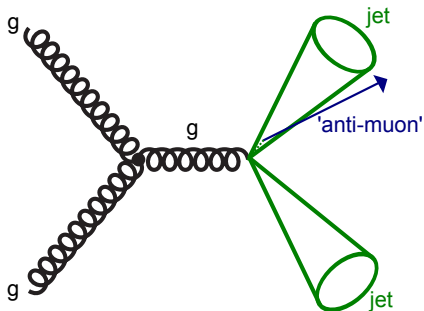
Jet-lepton model

- Requirement of a selected lepton is replaced by a lepton-like jet
- Jet-lepton has to fulfil the same kinematic requirements as a true lepton

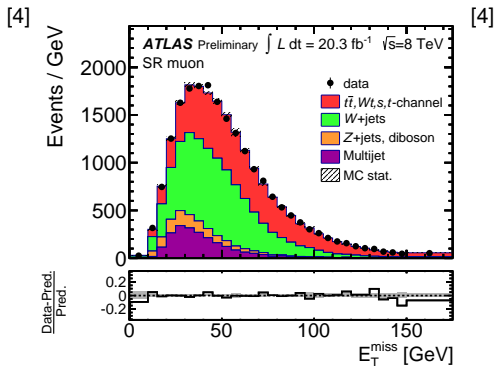
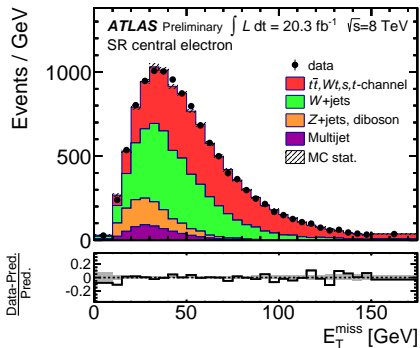


Anti-muon model

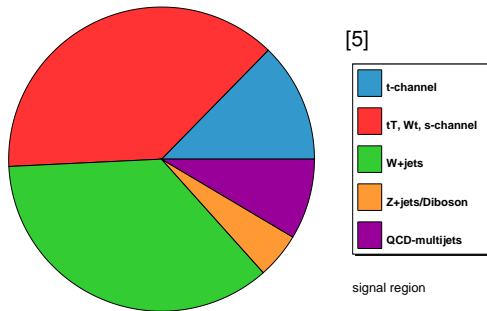
- The anti-muon template is taken from a sample highly enriched with fake muons
- Done by relaxing or inverting some of the muon identification cuts



- Both models provide a shape of the multijet background but no normalisation
- Normalisation is estimated using a binned likelihood fit in the full E_T^{miss} distribution in data
 - W +jets, Z +jets, Diboson and top component from simulation fitted with Gaussian constraints
 - Multijet rate is obtained from the fit

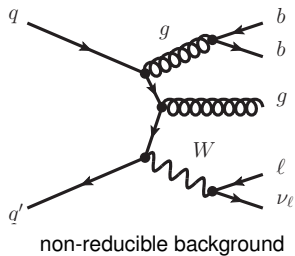
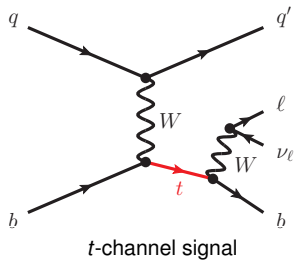


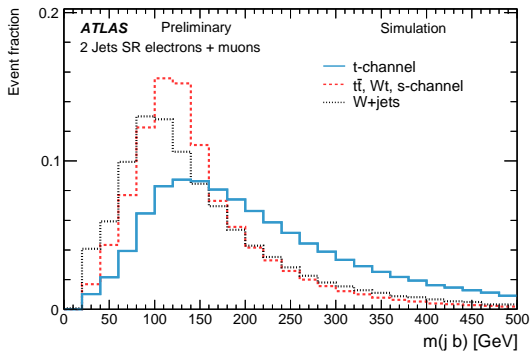
[4] ATLAS-CONF-2014-007



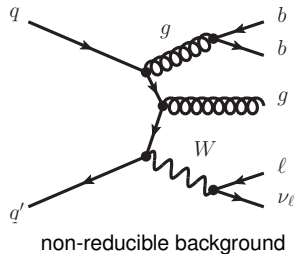
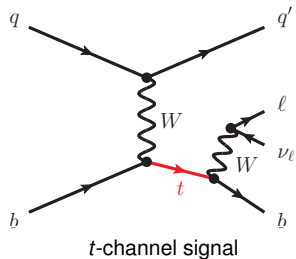
- About 143000 events selected from the full dataset
- Caveat of the basic event selection is a high fraction of non-reducible background

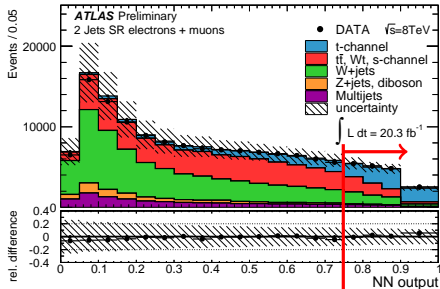
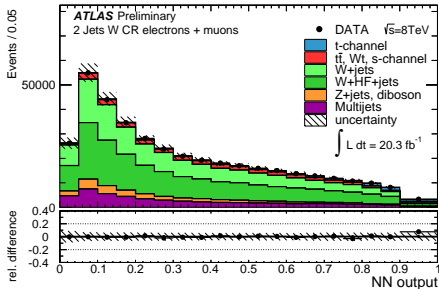
[5] from Table 2 in ATLAS-CONF-2014-007



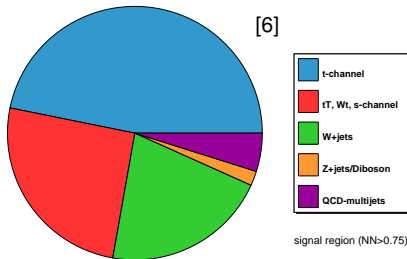


- t -channel process can be fully reconstructed
- Exploit full event topology to distinguish signal from background events
- E.g. invariant mass of the two jets ($m(jb)$)
→ from gluon splitting in W +jets

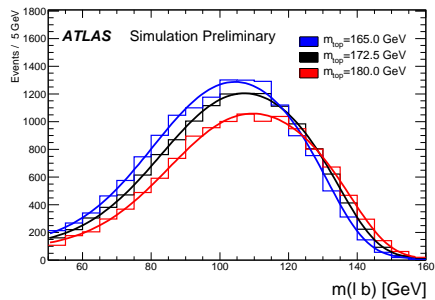
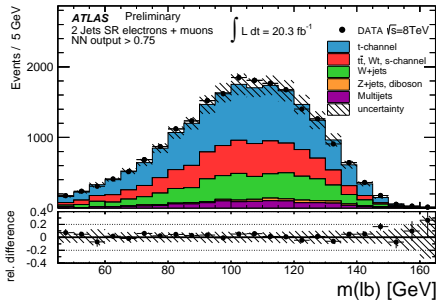




Variable	loss of total correlation (%)
$m(\ell\nu b)$	38
$m(jb)$	31
$m(\ell b)$	18
$ \eta(j) $	14
$\eta(\ell\nu)$	13
$H_T(\ell, \text{jets}, E_T^{\text{miss}})$	10
E_T^{miss}	7
$m_T(W)$	7
$\cos\theta(\ell, j)_{\ell\nu b, r.f.}$	6
$p_T(W)$	3
$\eta(\ell\nu b)$	2
$\Delta R(\ell, \ell\nu b)$	1



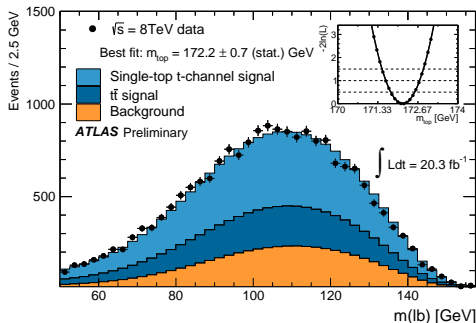
[6] from Table 2 in ATLAS-CONF-2014-007



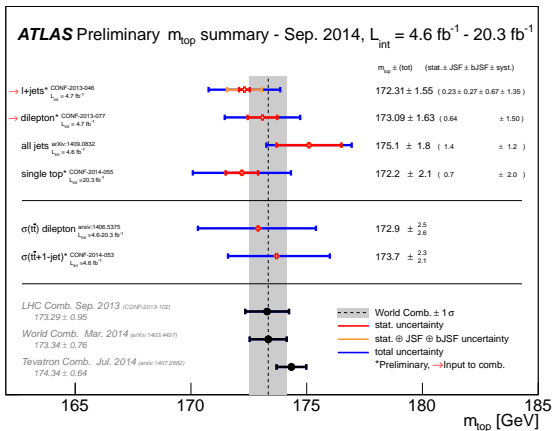
- The invariant mass $m(lb)$ of the charged lepton and b -tagged jet is employed as an estimator of the top quark mass using a template method
 - ▶ Same technique with the same variable has already been used e.g. in the dileptonic $t\bar{t}$ decay channel (ATLAS-CONF-2013-077)
- In mass measurement all processes involving top-quarks are treated as signal
- Normalised templates for the signal and background are used in a binned maximum likelihood fit to data with:

$$\mathcal{L} = \prod_{\text{bins } i} P\left(m(lb)_i^{\text{data}} | \lambda_i(N, f, s_i(m_{\text{top}}), b_i)\right) \cdot G\left(f | f_{\text{bkg}}, \sigma_{f_{\text{bkg}}}\right) \quad (1)$$

	Value [GeV]
Measured value	172.2
Statistical uncertainty	0.7
Jet energy scale	1.5
Jet energy resolution	< 0.1
Jet vertex fraction	< 0.1
Flavour tagging efficiency	0.3
Electron uncertainties	0.3
Muon uncertainties	0.1
Missing transverse momentum	0.2
W +jets normalisation	0.4
W +jets shape	0.3
Z +jets/diboson normalisation	0.2
Multijet normalisation	0.2
Multijet shape	0.3
Top normalisation	0.2
t -channel generator	< 0.1
t-channel hadronisation	0.7
t -channel colour reconnection	0.3
t -channel underlying event	< 0.1
$t\bar{t}, Wt$, and s -channel generator	0.2
$t\bar{t}$ hadronisation	< 0.1
$t\bar{t}$ colour reconnection	0.2
$t\bar{t}$ underlying event	0.1
$t\bar{t}$ ISR/FSR	0.2
Proton PDF	< 0.1
Simulation sample statistics	0.3
Total systematic uncertainty	2.0
Total uncertainty	2.1

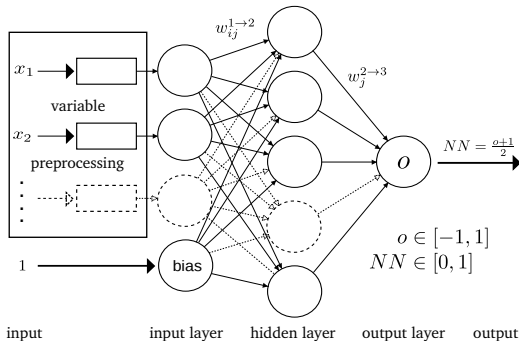


- Measurement is dominated by systematic uncertainties (mainly JES)
- More details about this analysis are given in ATLAS-CONF-2014-055

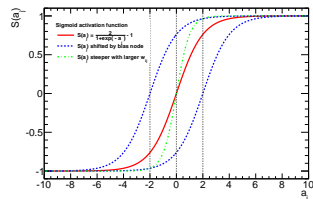
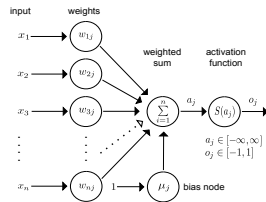


- Result is in good agreement with other ATLAS measurements in different $t\bar{t}$ final states
- The presented analysis is the first top mass measurement in this complementary topology
 → Many ideas to further improve the analysis
- The prospects to include this result in future top mass combinations are very good

Backup



- At each node the weighted sum of the input variables is passed through a sigmoid activation function.



- An artificial neural network is inspired by the human brain with neurons and synapses.
- During the training the weights (synapses) between neurons are modified such that the prediction of the network is as close as possible to the true expectation.
- Whole neural network framework and variable preprocessing is implemented in the Neurobayes package.

- The goal of the training procedure is to minimise the difference between the predicted output of the network, $o_K \in [-1, 1]$, and the true target value, $T_K \in [-1, 1]$ for event k .
- This is quantified by the entropy loss function which is defined based on the sum of n events in the training sample

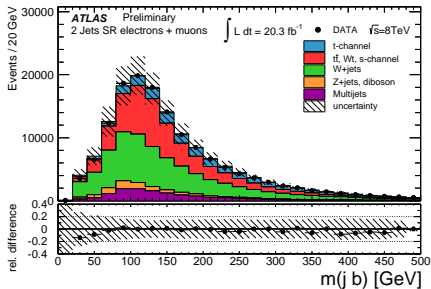
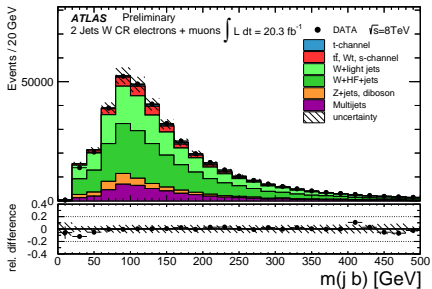
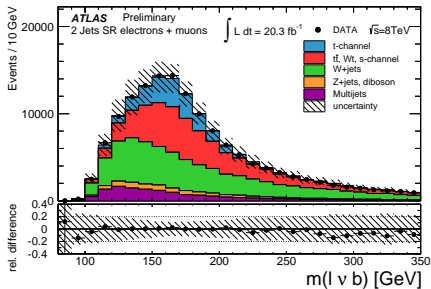
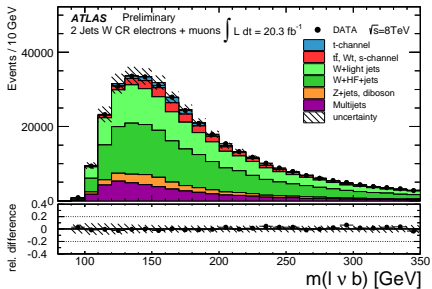
$$E_D = \sum_{k=1}^n \log \left[\frac{1}{2} (1 + T_K \cdot o_K + \varepsilon) \right].$$

- The network is trained by propagating E_D backwards from output to input layer. The change of the i input weights of neuron j is defined as

$$\begin{aligned} \Delta w_{ij} &= -\eta \frac{\partial E_D}{\partial w_{ij}} = -\eta \frac{\partial E_D}{\partial o_j} \frac{\partial o_j}{\partial a_j} \frac{\partial a_j}{\partial w_{ij}} \\ &= -\eta \frac{\partial E_D}{\partial o_j} S'(a_j) x_i = -\eta \delta_j x_i. \end{aligned}$$

- The error signal δ_j of neuron j depends on the layer

$$\begin{aligned} \delta_j &= S'(a_j) \cdot \frac{\partial E_D}{\partial o_j} = S'(a_j) \cdot \frac{T_j}{1 + T_j o_j + \varepsilon} \quad (\text{if } j \text{ is output neuron}), \\ \delta_j &= S'(a_j) \cdot \sum_k \delta_k w_{jk} \quad (\text{if } j \text{ is hidden neuron}). \end{aligned}$$



Process	SR	SR ($NN > 0.75$)
t -channel	18100 ± 1800	9100 ± 1300
$t\bar{t}$, Wt , s -channel	54200 ± 4300	4940 ± 600
W +jets	51000 ± 28000	4090 ± 2200
Z +jets, diboson	6900 ± 1700	360 ± 90
Multijet	12200 ± 6100	950 ± 480
Total expectation	142000 ± 29000	19470 ± 2700
Data	143332	19833

[7]

- About 75% of the obtained sample is from top processes and roughly 50% from single top t -channel production.

[7] ATLAS-CONF-2014-055

- JES uncertainty is dominated by η -intercalibration and modelling NP1
- JES uncertainty is the dominant systematic uncertainty of the method

[8]

	Δm_{top} [GeV]
<i>b</i> -jet energy scale	0.4
Modelling1	0.9
Modelling2	< 0.1
Modelling3	0.2
Modelling4	< 0.1
Eta intercalibration (modelling)	0.9
Statistical1	0.1
Statistical2	< 0.1
Statistical3	< 0.1
Eta intercalibration (statistical)	0.2
Detector1	0.4
Detector2	< 0.1
Detector3	< 0.1
Mixed1	< 0.1
Mixed2	< 0.1
Pile-up offset (μ term)	0.2
Pile-up offset (NPV term)	0.2
Pile-up (p_T term)	< 0.1
Pile-up (ρ topology)	0.1
Single particle high p_T	< 0.1
Flavour composition	0.2
Flavour response	< 0.1
Jet energy scale	1.5

[8] ATLAS-CONF-2014-055